

# Implementation of poloidal density variation in collision operator

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These notes are concerned with the implementation of the `poloidalVariationInCollisionOperator` option in SFINCS version 3.

## $\Phi_1$ implementation in SFINCS

In a previous set of notes ( *$\Phi_1$  implementation*), poloidal density variation to the lowest order was implemented by linearizing around a poloidally varying Maxwellian  $f_0$ , rather than a flux-function Maxwellian  $f_M$ .

Poloidal variations can occur to lowest order if the potential varies on a flux-surface, in which case the lowest order distribution function becomes

$$f_0(\psi, \theta, \zeta) = f_M(\psi) e^{-Ze\Phi_1(\theta, \zeta)/T}. \quad (1)$$

[We will in the future suppress the  $\psi$  dependence, as SFINCS only solves for a single flux-surface.]

In the  $\Phi_1$  *implementation* notes, the equations were modified to linearize around this  $f_0$  rather than  $f_M$ . However, the linearization in the collision operator was not treated, which is the subject of these notes.

## $\Phi_1$ implementation in the linearized collision operator

The linearized collision operator can be written as

$$C_{ab}^{L:f_0} = C_{ab}\{f_{a0}, f_{b0}\} + C_{ab}\{f_{a1}, f_{b0}\} + C_{ab}\{f_{a0}, f_{b1}\}, \quad (2)$$

where we use the superscript  $L : f_0$  to indicate that the operator has been linearized around  $f_0$  rather than  $f_M$ .

As  $f_0$  and  $f_M$  have the same velocity space-structure, the terms in the above linearization can easily be expressed in terms of the terms in the linearization around  $f_M$ :

$$C_{ab}^{L:f_0} = C_{ab}\{f_{aM}, f_{bM}\} e^{-(Z_a/T_a + Z_b/T_b)e\Phi_1(\theta, \zeta)} + C_{ab}\{f_{a1}, f_{b0}\} e^{-Z_b e\Phi_1(\theta, \zeta)/T_b} + C_{ab}\{f_{aM}, f_{b1}\} e^{-Z_a e\Phi_1(\theta, \zeta)/T_a}. \quad (3)$$

Effectively, from the explicit expressions for  $C^{L:f_M}$ , one can obtain  $C^{L:f_0}$  by substituting the density of each species by:

$$n_a \rightarrow n_a e^{-\frac{Z_a e \Phi_1(\theta, \zeta)}{T_a}}. \quad (4)$$

## Changes to the Jacobian

As the old collision operator only depend on  $f_1$  (neglecting the temperature equilibration term  $C_{ab}\{f_{aM}, f_{bM}\}$  for the moment), its contribution to the Jacobian was

$$J^{L:f_M} = \frac{\delta C^{L:f_M}}{\delta f_1}. \quad (5)$$

Furthermore, as the linearized collision operator is linear in  $f_1$ , we have

$$\frac{\delta C^{L:f_M}}{\delta f_1} = C^{L:f_M}[\cdot], \quad (6)$$

where the notation  $C^{L:f_M}[\cdot]$  implies that the operator does not act on anything, i.e. the matrix used to represent  $C^{L:f_M}[\cdot]$  enters as is into the Jacobian.

In the system linearized around  $f_0$ , we have essentially the same  $f_1$  dependence, if we substitute the density as in (4)

$$\frac{\delta C^{L:f_0}}{\delta f_1} = C^{L:f_M}[\cdot] \left\{ n_a \rightarrow n_a e^{-\frac{Z_a e \Phi_1(\theta, \zeta)}{T_a}} \right\}, \quad (7)$$

however, as  $\Phi_1$  is an unknown, we get new non-zero elements in the Jacobian due to  $\frac{\delta C^{L:f_0}}{\delta \Phi_1}$ . As  $C^{L:f_0}$  only depends on  $\Phi_1$  through  $n_a e^{-\frac{Z_a e \Phi_1(\theta, \zeta)}{T_a}}$ , we can also get the new expressions through as a density substitution

$$n_a e^{-\frac{Z_a e \Phi_1}{T_a}} \rightarrow -\frac{Z_a e}{T_a} n_a e^{-\frac{Z_a e \Phi_1}{T_a}}. \quad (8)$$

Thus

$$\frac{\delta C^{L:f_0}}{\delta \Phi_1} = C^{L:f_M}[f_1] \left\{ n_a \rightarrow -\frac{Z_a e}{T_a} n_a e^{-\frac{Z_a e \Phi_1}{T_a}} \right\}. \quad (9)$$

Note that the collision operator here acts on  $f_1$ .

In addition, if we include temperature equilibration, we get an additional contribution to the Jacobian

$$\frac{\delta C^{L:f_0}}{\delta \Phi_1} = -e \frac{Z_a T_b + Z_b T_a}{T_a T_b} e^{-\frac{Z_a T_b + Z_b T_a}{T_a T_b} e \Phi_1} C_{ab}\{f_{aM}, f_{bM}\}. \quad (10)$$

This term, like all the other in this document, is merely an extra factor to an expression already calculated in the code.

## Changes to the code

The collision operator is linearized around  $f_0$  if `poloidalVariationInCollisionOperator` is set to true in the input file. This switch adds the extra factor  $e^{-\frac{Z_a e \Phi_1}{T_a}}$  to the densities in the collision operator as calculated in the code, and reuses the calculations to evaluate  $\frac{\delta C^{L:f_0}}{\delta \Phi_1}$ . If `includeTemperatureEquilibrationTerm` is set to true, the contribution from this term is also included in the Jacobian.