

Lee 2-Symmetries as Extended Defect Ops

$$S = \int_M d^d x \mathcal{L}$$

$$\mathcal{L}[\phi] \quad \phi \mapsto \phi + \alpha \delta\phi$$

$$S \mapsto S + \delta S$$

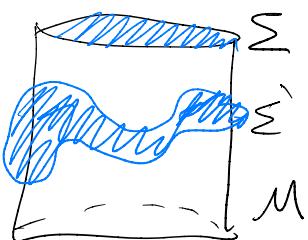
$$\delta S = \int d^d x \partial_\mu J^\mu$$

$$\delta S = \int d^d x \partial_\mu \left(\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi)} \delta\phi + \mathcal{L} \delta x^\mu \right)$$

$$\partial_\mu j^\mu = 0$$

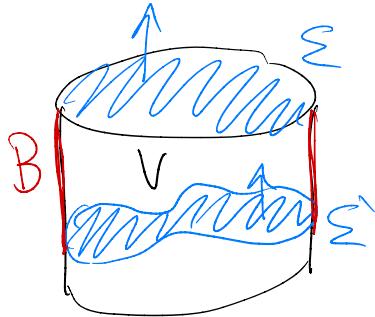
$$j^\mu = \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi)} \delta\phi + \mathcal{L} \delta x^\mu - J^\mu$$

$$\partial_\mu j^\mu = 0 \quad d * j = 0$$



$$Q = \int_{\Sigma} d^d x j^0$$

$$Q(\Sigma) = \int_{\Sigma} * j$$



$$\partial V = \Sigma - B - \Sigma'$$

$$\int_V d * j = 0$$

||

$$\int_{\partial V} * j = \int_{\Sigma} * j - \int_{\Sigma'} * j - \int_B * j$$

$$Q(\Sigma) = Q(\Sigma') + Q(B)$$

$$Q(\Sigma) = Q(\Sigma')$$

Homology classes of the surface

$$U_g = e^{i\lambda \underline{Q}(\underline{\varepsilon})}$$

Quantisation $Q(\underline{\varepsilon}) \rightarrow$ charge operator.

$\lambda \rightarrow$ lie Algebra valued group action a .

$$g = e^{i\lambda} \quad g \in G \quad \lambda \in \mathfrak{g}$$

$$\langle U_{g_1}(\underline{\varepsilon}) \ U_{g_2}(\underline{\varepsilon}) \rangle = \langle U_{g_1 g_2}(\underline{\varepsilon}) \rangle$$

$$\rightarrow U_{g_1}(\underline{\varepsilon}) \ U_{g_2}(\underline{\varepsilon}) = e^{i\lambda_1 Q(\underline{\varepsilon})} e^{i\lambda_2 Q(\underline{\varepsilon})}$$

$$e^X e^Y = e^Z \quad Z = X + Y + \frac{1}{2}[X, Y] + \dots$$

$$\partial_\mu j^\mu_a(x) j^\nu_b(y) = f_{ab}{}^c j_c(x) \delta^{(d)}(x-y)$$



$$\langle d \star j \rangle = 0$$

action of sym A.

$$\Theta_R(x) \mapsto R(g) \Theta_R(x)$$

$$d \star \langle j(x) \Theta_R(y) \rangle = \delta^{(d)}(x-y) \langle \overbrace{R(T^a)}^{\sim} \Theta_R(y) \rangle$$

$$\downarrow$$

$$\langle d \star j(x) \Theta_R(y) \rangle = \langle \delta^{(d)}(x-y) R(T^a) \Theta_R(y) \rangle$$

$$U_g(\Sigma) = U_g(\Sigma')$$

$$\underline{U_g(\Sigma)} = e^{i\lambda Q(\Sigma)} = e^{i\lambda Q(\Sigma')} = \underline{\underline{U_g(\Sigma')}}$$



$$U_g U_{g_2} = U_{g \oplus g_2}$$

Unitary operators

$$\begin{aligned} U_g(\varepsilon) U_{g_1}(\varepsilon) &= e^{i\lambda Q(\varepsilon)} e^{-i\lambda Q(\varepsilon)} \\ &= e^{i(\lambda - \lambda) Q(\varepsilon)} \\ &= e^0 = 1 \end{aligned}$$

$$\begin{aligned} g^{-1} &= (g)^{-1} \\ &= (e^{i\lambda})^{-1} \\ &= e^{-i\lambda} \end{aligned}$$

$$U_{\tilde{g}}(\varepsilon) = (U_g(\varepsilon))^{-1}$$

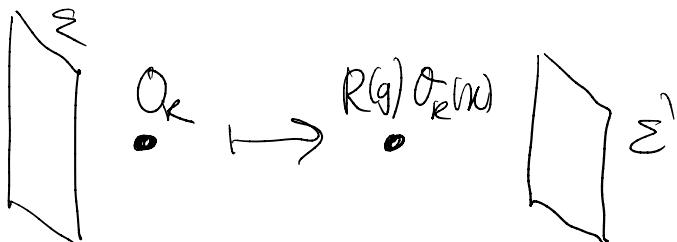
$j^{(i)}$ → 1 form currents

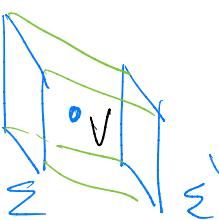
$Q^{(0)}$ → 0 form charges

$$U_g(\varepsilon) = e^{i\lambda Q(\varepsilon)} \rightarrow \text{act on } \mathcal{O}_R(x)$$

Claim

$$\langle U_g(\varepsilon) \mathcal{O}_R(x) \rangle = \langle R(g) \mathcal{O}_R(x) U_{g'}(\varepsilon') \rangle$$



$$U_g(\varepsilon) \mathcal{O}_R(x) U_{g^{-1}}(\varepsilon') = e^{i\lambda Q(\varepsilon)} \mathcal{O}_R(x) e^{-i\lambda Q(\varepsilon')} \\ = e^{i\lambda Q(\varepsilon)} e^{-i\lambda Q(\varepsilon')} \mathcal{O}_R(x) \quad \xrightarrow{\text{L} \circ \gamma} \\ = e^{i\lambda(Q(\varepsilon) - Q(\varepsilon'))} \mathcal{O}_R(x)$$


$$Q(\varepsilon) - Q(\varepsilon') = \int_{\varepsilon} * j - \int_{\varepsilon'} * j + \underbrace{\int_B * j}_{=0} \\ = \int_{\varepsilon - \varepsilon' + B} * j \\ = \int_V d * j$$

$$U_g(\varepsilon) \mathcal{O}_R(x) U_{g^{-1}}(\varepsilon) = e^{i\lambda \int_V d * j} \mathcal{O}_R(x) \\ = \sum_{k=0}^{\infty} \frac{(i\lambda)^k}{k!} \left(\int_V d * j \right)^n \mathcal{O}_R(x).$$

$$d * \langle j(y) \mathcal{O}_R(x) \rangle = * \delta^d(x-y) R(T^a) \mathcal{O}_R(x)$$

$$\lambda = \lambda^a T^a$$

$$= \sum_{k=0}^{\infty} \frac{(i\lambda^a)^k}{k!} \left(\int_V dy \delta^{(d)}(x-y) R(T^a) \right)^k \phi_R(x)$$

$$= \sum_{k=0}^{\infty} \frac{(i\lambda^a R(T^a))^k}{k!} \underbrace{\left(\int_V dy \delta^{(d)}(x-y) \right)^k}_{1} \phi_R(x).$$

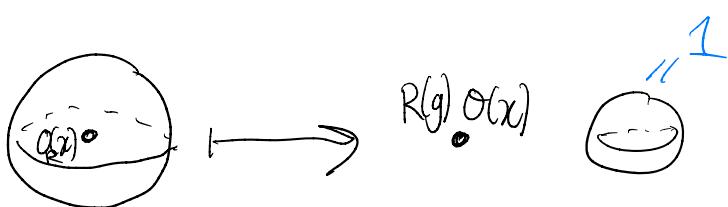
$$= \sum_{k=0}^{\infty} \frac{(i\lambda^a R(T^a))^k}{k!} \phi_R(x)$$

$$\leftarrow e^{\frac{i\lambda^a R(T^a)}{R(a)}} \phi_R(x)$$

$$= R(g) \phi_R(x)$$

$$U_g(\varepsilon) \phi_R(x) = R(g) \phi_R(x) U_g(\varepsilon)$$

Corollary



$$U_g(S^{d-1}) \phi_R(x) = R(g) \phi_R(x)$$

$G = U(1) \rightarrow$ electromagnetism.

$$\mathcal{O}_q(x) \mapsto e^{i\lambda q} \mathcal{O}_q(x)$$

Aharanov-Bohm effect.

Generalised Non- \mathfrak{cte} symmetries.

Lagrangian



Noether currents \rightarrow Noether charges.



$$U_g(\varepsilon) = e^{i\lambda Q(\varepsilon)}$$

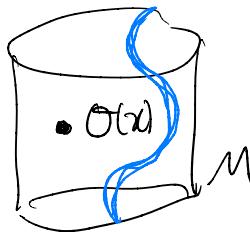
↔ Discrete groups.
Non-Lagrange Theories.

0-form Sym: $Q^{(0)} \quad j^{(1)} \quad \Sigma^{(d-1)}$

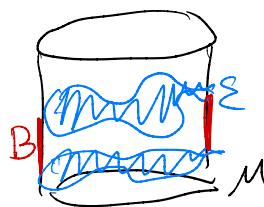
↓ 1-form Sym: $Q \quad j^{(2)} \quad \Sigma^{(d-2)}$

↪ Abelian.

Local operators



$$V_g(\varepsilon) = e^{i \int_{\underline{\varepsilon}}^* j}$$



$$Q(\varepsilon) = Q(\varepsilon') + Q(B)$$

$$e^{i\lambda_1 Q(\varepsilon)} e^{i\lambda_2 Q(\varepsilon)} = e^Z$$

$$Z = \lambda_1 Q + \lambda_2 Q + [\lambda_1 Q(\varepsilon), \lambda_2 Q(\varepsilon)] + \dots$$

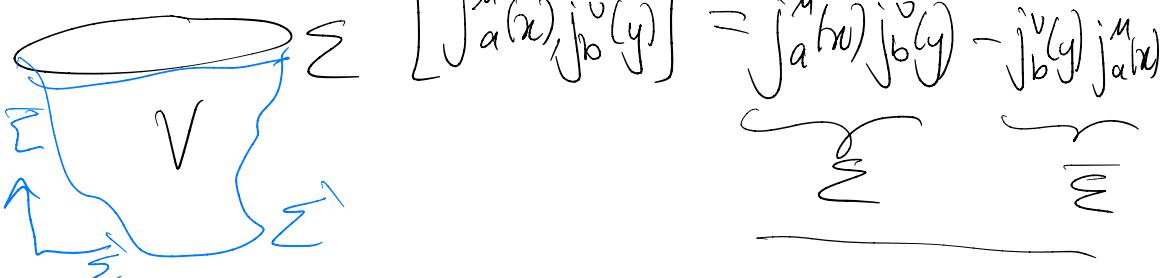
$$e^{i\lambda_1} e^{i\lambda_2} = e^{i\lambda_3} \quad Q^a \rightarrow \text{Generators}$$

$$\begin{aligned} [\lambda_1^a Q(\varepsilon), \lambda_2^b Q(\varepsilon)] &= \lambda_1^a \lambda_2^b [Q^a(\varepsilon), Q^b(\varepsilon)] \\ &= \lambda_1^a \lambda_2^b f^{ab}_c Q^c(\varepsilon) \\ &= \lambda^c Q^c(\varepsilon) \end{aligned}$$

$$[\mathbb{Q}^a(z), \mathbb{Q}^b(z)] = \int_{\Sigma} \int_{\Sigma} [\star j^a, \star j^b]$$

$$\partial_\mu j_a^\mu(x) j_b^\nu(y) = f_{ab}{}^c j_c^\nu(x) \delta^{(d)}(x-y).$$

$$\int_{\Sigma} \int_{\Sigma} n_v n_\mu [j_a^\mu(x), j_b^\nu(y)] d^{d-1}x d^{d-1}y$$



$\int_{\Sigma} \int_{\Sigma} n_v n_\mu [j_a^\mu(x), j_b^\nu(y)] = j_a^\mu(x) j_b^\nu(y) - j_b^\nu(y) j_a^\mu(x)$

The diagram shows a cylinder with a top surface labeled Σ and a bottom surface labeled $\bar{\Sigma}$. The volume inside is labeled V . Blue arrows indicate the flow of vectors from the top surface down through the volume to the bottom surface.

$$\int_{\Sigma} \int_{\Sigma-\bar{\Sigma}} n_v n_\mu j_a^\mu(x) j_b^\nu(y) d^{d-1}x d^{d-1}y$$

$$= \int_{\Sigma} \int_{\Sigma} n_v n_\mu j_a^\mu j_b^\nu d^{d-1}x d^{d-1}y - \int_{\Sigma} \int_{\Sigma} n_v n_\mu j_b^\nu j_a^\mu d^{d-1}x d^{d-1}y$$

$$= \int_{\Sigma} \int_{\Sigma-\bar{\Sigma}} n_v n_\mu j_a^\mu j_b^\nu d^{d-1}x d^{d-1}y$$

$$= \int_{\Sigma} \int_V n_v \partial_m \left(j_a^m j_b^v \right) d^{d-1}x d^{d-1}y \quad >$$

$$= \int_{\Sigma} \int_V n_v f_{ab}^c j_c^v(x) \delta^d(x-y)$$

$$= \int_{\Sigma} n_v f_{ab}^c j_c^v(y)$$

$$= f_{ab}^c Q^c(\Sigma)$$

$$\therefore [Q^a(\varepsilon), Q^b(\varepsilon)] = \langle f_{ab}^c Q^c(\varepsilon) \rangle$$

$$\lambda Q = \lambda^a Q^a \quad g = e^{i\lambda}$$

Then $U_{g_1}(\varepsilon) U_{g_2}(\varepsilon) = U_{g_1 g_2}(\varepsilon)$