

HOMEWORK 8

4.1. a) For $f(x, y, z)$, according to Clairaut's Theorem:

$$[x, y]f = (xy - yx)f = 0 = [y, z]f = [z, x]f.$$

$$\rightarrow [\hat{p}_x, \hat{p}_y]f = (\hat{p}_x \hat{p}_y - \hat{p}_y \hat{p}_x)f = \left[\left(-i\hbar \frac{\partial}{\partial x} \right) \left(-i\hbar \frac{\partial}{\partial y} \right) - \left(-i\hbar \frac{\partial}{\partial y} \right) \left(-i\hbar \frac{\partial}{\partial x} \right) \right] f \\ = \left(-\hbar^2 \frac{\partial^2}{\partial x \partial y} + \hbar^2 \frac{\partial^2}{\partial y \partial x} \right) f = -\hbar^2 \frac{\partial^2 f}{\partial x \partial y} + \hbar^2 \frac{\partial^2 f}{\partial y \partial x} = 0$$

$$\rightarrow [\hat{p}_y, \hat{p}_z] = [\hat{p}_x, \hat{p}_z] = 0$$

$$\rightarrow [\hat{x}, \hat{p}_x]f = (\hat{x} \hat{p}_x - \hat{p}_x \hat{x})f = \left[x \left(-i\hbar \frac{\partial}{\partial x} \right) - \left(-i\hbar \frac{\partial}{\partial x} \right) x \right] f = -i\hbar x \frac{\partial f}{\partial x} + i\hbar \frac{\partial}{\partial x} (xf) \\ = -i\hbar x \frac{\partial f}{\partial x} + i\hbar x \frac{\partial f}{\partial x} + i\hbar f = i\hbar f$$

$$\rightarrow [\hat{y}, \hat{p}_y]f = [\hat{z}, \hat{p}_z]f = i\hbar f$$

$$\rightarrow [\hat{x}, \hat{p}_y]f = (\hat{x} \hat{p}_y - \hat{p}_y \hat{x})f = \left[x \left(-i\hbar \frac{\partial}{\partial y} \right) - \left(-i\hbar \frac{\partial}{\partial y} \right) x \right] f = -i\hbar x \frac{\partial f}{\partial y} + i\hbar \frac{\partial}{\partial y} (xf) \\ = -i\hbar x \frac{\partial f}{\partial y} + i\hbar x \frac{\partial f}{\partial y} = 0$$

$$\rightarrow [\hat{x}, \hat{p}_z]f = [\hat{y}, \hat{p}_x]f = [\hat{y}, \hat{p}_z]f = [\hat{z}, \hat{p}_x]f = [\hat{z}, \hat{p}_y]f = 0$$

$$\rightarrow [\hat{x}, \hat{y}] = [\hat{y}, \hat{z}] = [\hat{y}, \hat{x}] = [\hat{z}, \hat{y}] = [\hat{x}, \hat{z}] = [\hat{z}, \hat{x}] = 0$$

$$\rightarrow [\hat{p}_x, \hat{p}_y] = [\hat{p}_y, \hat{p}_z] = [\hat{p}_x, \hat{p}_z] = [\hat{p}_z, \hat{p}_y] = [\hat{p}_z, \hat{p}_x] = [\hat{p}_x, \hat{p}_z] = 0$$

$$\rightarrow [\hat{y}, \hat{p}_x] = [\hat{p}_x, \hat{y}] = [\hat{z}, \hat{p}_y] = [\hat{p}_y, \hat{z}] = [\hat{x}, \hat{p}_z] = [\hat{p}_z, \hat{x}] = [\hat{y}, \hat{p}_z] = [\hat{p}_z, \hat{y}] \\ = [\hat{z}, \hat{p}_x] = [\hat{p}_x, \hat{z}] = [\hat{p}_x, \hat{z}] = [\hat{p}_y, \hat{z}] = 0$$

$$\rightarrow [\hat{p}_x, \hat{x}] = [\hat{p}_y, \hat{y}] = [\hat{p}_z, \hat{z}] = -[\hat{x}, \hat{p}_x] = -[\hat{y}, \hat{p}_y] = -[\hat{z}, \hat{p}_z] = -i\hbar.$$

b) $i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi + V(x, y, z, t) \psi(x, y, z, t)$

$$\Leftrightarrow \frac{\partial \psi}{\partial t} = \frac{i\hbar}{2m} \nabla^2 \psi - \frac{i}{\hbar} V \psi \quad \Leftrightarrow \frac{\partial \psi^*}{\partial t} = -\frac{i\hbar}{2m} \nabla^2 \psi^* + \frac{i}{\hbar} V \psi^*$$

$$\rightarrow \frac{d}{dt} \langle x_j \rangle = \frac{d}{dt} \langle \psi | \hat{x}_j | \psi \rangle = \frac{d}{dt} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \psi^*(x, y, z, t) (x_j) \psi(x, y, z, t) dx dy dz$$

$$= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} x_j \left[\left(\frac{-i\hbar}{2m} \nabla^2 \psi^* + \frac{i}{\hbar} V \psi^* \right) \psi + \psi^* \left(\frac{i\hbar}{2m} \nabla^2 \psi - \frac{i}{\hbar} V \psi \right) \right] dx dy dz$$

$$= \frac{i\hbar}{2m} \left(\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} x_j \psi^* \nabla^2 \psi dx dy dz - \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} x_j \psi \nabla^2 \psi^* dx dy dz \right)$$

$$= \frac{i\hbar}{2m} \left[- \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \nabla(x_j \psi^*) \cdot \nabla \psi dx dy dz + \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \nabla(x_j \psi) \cdot \nabla \psi^* dx dy dz \right]$$

(I don't wanna do this...)

$$= \frac{1}{m} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \psi^* \left(-i\hbar \frac{\partial}{\partial x_j} \right) \psi dx dy dz = \frac{1}{m} \langle \psi | \hat{p}_j | \psi \rangle = \frac{1}{m} \langle p_j \rangle = \frac{d}{dt} \langle \hat{x}_j \rangle = \frac{1}{m} \langle \hat{p}_j \rangle$$

$$\begin{aligned}
 \frac{d}{dt} \langle \vec{p} \rangle &= \frac{d}{dt} \langle \Psi | \vec{p} | \Psi \rangle = \frac{d}{dt} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \Psi^*(x, y, z, t) (-i\hbar \nabla) \Psi(x, y, z, t) dx dy dz \\
 &= -i\hbar \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \left[\left(\frac{-\hbar^2}{2m} \nabla^2 \Psi^* - \frac{i}{\hbar} \Psi^* \right) \nabla \Psi + \Psi^* \nabla \left(\frac{\hbar^2}{2m} \nabla^2 \Psi - \frac{i}{\hbar} \Psi \right) \right] dx dy dz \\
 &= -i\hbar \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \left[\frac{-\hbar^2}{2m} \nabla^2 \Psi^* \nabla \Psi + \frac{\hbar^2}{2m} \Psi^* \nabla^2 (\nabla \Psi) - \Psi^* \nabla \nabla \Psi \right] dx dy dz \\
 &= \frac{\hbar^2}{2m} \cdot 0 - \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \Psi^* \nabla \nabla \Psi dx dy dz = \langle \Psi | -\nabla \nabla | \Psi \rangle = \langle -\nabla \nabla \rangle
 \end{aligned}$$

$$c) \frac{\partial^2 \sigma_x^2}{\partial t^2} = \left(\frac{1}{2} \langle [A, B] \rangle \right)^2$$

$$\begin{aligned}
 [x, \hat{p}_x] &= [y, \hat{p}_y] = [z, \hat{p}_z] = [x, \hat{p}_y] = [\hat{p}_x, y] = [\hat{p}_x, z] = [\hat{p}_y, x] = [\hat{p}_y, z] = 0 \\
 \Rightarrow \partial_x \partial_y &\neq 0, \partial_y \partial_z \neq 0, \partial_x \partial_z \neq 0, \partial_{p_x} \partial_{p_y} \neq 0, \partial_{p_y} \partial_{p_z} \neq 0, \partial_{p_x} \partial_{p_z} \neq 0 \\
 [x, \hat{p}_y] &= [x, \hat{p}_z] = [y, \hat{p}_x] = [y, \hat{p}_z] = [z, \hat{p}_x] = [z, \hat{p}_y] = 0 \\
 \Rightarrow \partial_x \partial_{p_y} &\neq 0, \partial_x \partial_{p_z} \neq 0, \partial_y \partial_{p_x} \neq 0, \partial_y \partial_{p_z} \neq 0, \partial_z \partial_{p_x} \neq 0, \partial_z \partial_{p_y} \neq 0 \\
 [x, \hat{p}_x] &= [y, \hat{p}_y] = [z, \hat{p}_z] = i\hbar \\
 \Rightarrow \partial_x \partial_{p_x} &\neq \frac{\hbar}{2}, \partial_y \partial_{p_y} \neq \frac{\hbar}{2}, \partial_z \partial_{p_z} \neq \frac{\hbar}{2}
 \end{aligned}$$

$$4.2. V(x, y, z) = \begin{cases} 0, & x, y, z \text{ all between } 0 \text{ and } a \\ \infty, & \text{otherwise} \end{cases}$$

$$a) i\hbar \frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m} \left(\frac{\partial^2 \Psi}{\partial x^2} + \frac{\partial^2 \Psi}{\partial y^2} + \frac{\partial^2 \Psi}{\partial z^2} \right) + V\Psi, \quad -\infty < x, y, z < \infty, t > 0$$

$$\Psi(x, y, z, 0) = \Psi_0(x, y, z)$$

$$b) i\hbar \frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m} \left(\frac{\partial^2 \Psi}{\partial x^2} + \frac{\partial^2 \Psi}{\partial y^2} + \frac{\partial^2 \Psi}{\partial z^2} \right) + \infty \Psi, \quad x, y, z \notin (0, a)$$

outside cube

$$i\hbar \frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m} \left(\frac{\partial^2 \Psi}{\partial x^2} + \frac{\partial^2 \Psi}{\partial y^2} + \frac{\partial^2 \Psi}{\partial z^2} \right), \quad 0 < x, y, z < a, t > 0$$

$$\Psi(0, y, z, t) = 0, \Psi(x, 0, z, t) = 0, \Psi(x, y, 0, t) = 0, \Psi(x, y, z, 0) = \Psi_0(x, y, z)$$

$$\Psi(a, y, z, t) = 0, \Psi(x, a, z, t) = 0, \Psi(x, y, a, t) = 0$$

For $\Psi(x, y, z, t) = X(x) Y(y) Z(z) T(t)$ and apply boundary conditions:

$$+ X(0) = X(a) = Y(0) = Y(a) = Z(0) = Z(a) = 0$$

Apply Separation of Variables:

$$i\hbar \frac{T'(t)}{T(t)} = -\frac{\hbar^2}{2m} \left[\frac{X''(x)}{X(x)} + \frac{Y''(y)}{Y(y)} + \frac{Z''(z)}{Z(z)} \right] = E \text{ (const)}$$

$$\left. \begin{aligned} \psi T'(t) &= E (\text{const}) \\ \frac{T(t)}{X(x)} &= F (\text{const}) \\ \frac{Y(y)}{Z(z)} &= G (\text{const}) \\ -\frac{2mE}{\hbar^2} - F - \frac{Z''(z)}{Z(z)} &= G \end{aligned} \right\} \begin{aligned} X(x) &= C_1 e^{\mu x} + C_2 e^{-\mu x} \quad (F = \mu^2) \Rightarrow X' = 0 \\ X(x) &= C_3 x + C_4 \quad (F = 0) \Rightarrow X'' = -\gamma^2 X \\ X(x) &= C_5 \cos \gamma x + C_6 \sin \gamma x \quad (F = -\gamma^2) \\ \text{if } F = -\gamma^2 = -\frac{T^2 \pi^2}{a^2}, T = 1, 2, 3, \dots \Rightarrow X(x) &= C_6 \sin \frac{T \pi x}{a} \end{aligned}$$

$$\Rightarrow \text{For } G = -\frac{k^2 \pi^2}{a^2} \quad (k=1, 2, 3) \Rightarrow Y(y) = C_7 \sin \frac{k \pi y}{a}$$

$$\Rightarrow Z(z) = C_8 \cos \sqrt{\frac{2mE}{\hbar^2} - \frac{T^2 \pi^2}{a^2} - \frac{k^2 \pi^2}{a^2}} z + C_9 \sin \sqrt{\frac{2mE}{\hbar^2} - \frac{T^2 \pi^2}{a^2} - \frac{k^2 \pi^2}{a^2}} z$$

$$\Rightarrow Z(z) = \frac{2mE}{\hbar^2} - \frac{T^2 \pi^2}{a^2} - \frac{k^2 \pi^2}{a^2} = \frac{\ell^2 \pi^2}{a^2} \quad (\ell = 0, \pm 1, \pm 2, \dots)$$

$$\Rightarrow E_{Tkl} = \frac{\pi^2 \hbar^2}{2ma^2} (T^2 + k^2 + \ell^2)$$

$$\Rightarrow \text{Stationary states: } \Psi_{Tkl}(x, y, z, t) = X_T(x) Y_k(y) Z_\ell(z) T_{Tkl}(t) \\ = A \sin \frac{T \pi x}{a} \sin \frac{k \pi y}{a} \sin \frac{\ell \pi z}{a} e^{-i \pi^2 \hbar / 2ma^2 (T^2 + k^2 + \ell^2) t}$$

$$\text{Normalized to find } A = \pm \left(\frac{2}{a} \right)^{3/2}$$

$$\Rightarrow \Psi_{Tkl}(x, y, z, t) = \left(\frac{2}{a} \right)^{3/2} \sin \frac{T \pi x}{a} \sin \frac{k \pi y}{a} \sin \frac{\ell \pi z}{a} e^{-i \pi^2 \hbar / 2ma^2 (T^2 + k^2 + \ell^2) t}$$

According to Principle of Superposition:

$$\Psi(x, y, z, t) = \sum_{T=1}^{\infty} \sum_{k=1}^{\infty} \sum_{\ell=1}^{\infty} c_{Tkl} \Psi_{Tkl}(x, y, z, t) \quad (\text{do math...})$$

$$\Rightarrow c_{Tkl} = \left(\frac{2}{a} \right)^{3/2} \int_0^a \int_0^a \int_0^a \Psi_0(x, y, z) \sin \frac{T \pi x}{a} \sin \frac{k \pi y}{a} \sin \frac{\ell \pi z}{a} dx dy dz$$

NRG Degen $J^2 + k^2 + \ell^2$

$$b) \begin{array}{lll} E_{111} & 1 & 3 = E_1 \end{array}$$

$$E_{211/112/112} \quad 3 \quad 6 = E_2$$

$$E_{311/212/22} \quad 3 \quad 9 = E_3$$

$$E_{311/131/113} \quad 3 \quad 11 = E_4$$

$$E_{222} \quad 1 \quad 12 = E_5$$

$$E_{312/321/132/123/213/231} \quad 6 \quad 14 = E_6$$

c) For E_{14} : $E_{555/511/51/115} \Rightarrow \text{degen} = 4$. We have the same set of numbers for E_{14} (333) similar to E_2 (111) & E_3 (222) but unlike those 2 where $\text{degen} = 1$, E_{14} has a $\text{degen} = 4$ by having other set of numbers (511, 151, 115).

4.5. a) $i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi + V(r) \psi(r, t) = -\frac{\hbar^2}{2m} \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \psi}{\partial r} \right) + V(r) \psi(r, t)$
 For $\psi(r, t) = \psi(r) T(t)$

$$\Rightarrow i\hbar \frac{\partial [\psi(r) T(t)]}{\partial t} = -\frac{\hbar^2}{2m} \frac{1}{r^2} \frac{\partial}{\partial r} \left[r^2 \frac{\partial}{\partial r} [\psi(r) T(t)] \right] + V(r) [\psi(r) T(t)]$$

$$\Leftrightarrow i\hbar \psi(r) T'(t) = -\frac{\hbar^2}{2m} \frac{T(t)}{r^2} \frac{d}{dr} \left(r^2 \frac{d\psi}{dr} \right) + V(r) [\psi(r) T(t)]$$

Separating variables:

$$\Rightarrow i\hbar \frac{T'(t)}{T(t)} = E$$

$$-\frac{\hbar^2}{2m} \frac{1}{r^2 \psi(r)} \frac{d}{dr} \left(r^2 \frac{d\psi}{dr} \right) + V(r) = E \quad \Rightarrow V(r) = E + \frac{\hbar^2}{2m r^2 \psi(r)} \frac{d}{dr} \left(r^2 \frac{d\psi}{dr} \right)$$

If $\psi(r) = Ae^{-r/a} \Rightarrow V(r) = E + \frac{\hbar^2}{2m r^2 \psi(r)} \frac{d}{dr} \left(r^2 \frac{d\psi}{dr} \right)$

$$= E + \frac{\hbar^2}{2m r^2 (Ae^{-r/a})} \frac{d}{dr} \left[r^2 \frac{d}{dr} (Ae^{-r/a}) \right] = E - \frac{\hbar^2}{2ma^2 e^{-r/a}} \left(2e^{-r/a} - \frac{r^2}{a} e^{-r/a} \right)$$

$$= E - \frac{\hbar^2}{2ma^2} \left(2 - \frac{r}{a} \right)$$

$V(r) \rightarrow 0$ as $r \rightarrow \infty \Rightarrow \lim_{r \rightarrow \infty} V(r) = E - \frac{\hbar^2}{2ma^2} (-1) = 0 \Rightarrow E = \frac{-\hbar^2}{2ma^2}$
 $\Rightarrow V(r) = -\frac{\hbar^2}{2ma^2}$

b) If $\psi(r) = Ae^{-r^2/a^2} \Rightarrow V(r) = E + \frac{\hbar^2}{2m r^2 \psi(r)} \frac{d}{dr} \left(r^2 \frac{d\psi}{dr} \right)$

$$= E + \frac{\hbar^2}{2m r^2 (Ae^{-r^2/a^2})} \frac{d}{dr} \left[r^2 \frac{d}{dr} (Ae^{-r^2/a^2}) \right] = E - \frac{\hbar^2}{ma^2 r^2 e^{-r^2/a^2}} \frac{d}{dr} (r^3 e^{-r^2/a^2})$$

$$= E - \frac{\hbar^2}{ma^2} \left(3 - 2 \frac{r^2}{a^2} \right)$$

$V(0) = 0 \Rightarrow E - \frac{\hbar^2}{ma^2} \cdot 3 = 0 \Rightarrow E = \frac{3\hbar^2}{ma^2} \Rightarrow V(r) = \frac{2\hbar^2}{ma^4} r^2$

4.7. Eqn 4.32: $Y_l^m(\theta, \phi) = \sqrt{\frac{(2l+1)(l-m)!}{4\pi(l+m)!}} e^{im\phi} P_l^m \cos \theta$

c) $Y_3^2(\theta, \phi) = \sqrt{\frac{(2 \cdot 3 + 1)(3-2)!}{4\pi(3+2)!}} e^{2i\phi} P_3^2 \cos \theta$

$$= \sqrt{\frac{7}{4\pi}} \frac{1}{5!} e^{i\phi} \left[\frac{(-1)^2}{2^3 3!} (1-x^2)^{3/2} \frac{d^{3+2}}{dx^{3+2}} (x^2-1)^3 \right] \Big|_{x=\cos\theta}$$

$$= \sqrt{\frac{7}{480\pi}} e^{2i\phi} \left[\frac{1}{48} (1-x^2) \frac{d^5}{dx^5} (x^6 - 3x^4 + 3x^2 - 1) \right] \Big|_{x=\cos\theta}$$

$$= \frac{1}{4} \sqrt{\frac{7}{30\pi}} e^{2i\phi} \left[\frac{1}{48} (1-x^2)(720x) \right] \Big|_{x=\cos\theta} = \frac{15}{4} \sqrt{\frac{7}{30\pi}} e^{2i\phi} (1-\cos^2\theta) \cos\theta$$

$$\Rightarrow Y_3^2(\theta, \phi) = \frac{1}{4} \sqrt{\frac{105}{2\pi}} e^{2i\phi} \cos\theta \sin^2\theta$$

$$\text{Eqn 4.18: } \sin\theta \frac{\partial}{\partial\theta} \left(\sin\theta \frac{\partial Y}{\partial\theta} \right) + \frac{\partial^2 Y}{\partial\phi^2} = -l(l+1) \sin^2\theta Y, \quad l=3$$

Satisfy check.

$$\Rightarrow \sin\theta \frac{\partial}{\partial\theta} \left(\sin\theta \frac{\partial Y_3^2}{\partial\theta} \right) + \frac{\partial^2 Y_3^2}{\partial\phi^2} = \sin\theta \frac{\partial}{\partial\theta} \left[\sin\theta \frac{\partial}{\partial\theta} \left(\frac{1}{4} \sqrt{\frac{105}{2\pi}} e^{2i\phi} \cos\theta \sin^2\theta \right) \right] + \frac{\partial^2}{\partial\phi^2} \left(\frac{1}{4} \sqrt{\frac{105}{2\pi}} e^{2i\phi} \cos\theta \sin^2\theta \right)$$

$$= \frac{1}{4} \sqrt{\frac{105}{2\pi}} \left(4e^{2i\phi} \cos\theta \sin^2\theta - 4e^{2i\phi} \cos^3\theta \sin^2\theta - 12e^{2i\phi} \cos\theta \sin^4\theta \right)$$

$$= -(12\sin^2\theta) \left(\frac{1}{4} \sqrt{\frac{105}{2\pi}} e^{2i\phi} \cos\theta \sin^2\theta \right) = -3(3+1) \sin^2\theta Y_3^2(\theta, \phi) \quad (\checkmark)$$

$$Y_l^l(\theta, \phi) = \sqrt{\frac{(2l+1)(l-1)!}{4\pi(l+l)!}} e^{il\phi} P_l^l(\cos\theta) = \sqrt{\frac{(2l+1)!}{4\pi(2l)!}} e^{il\phi} \left[\frac{(-1)^l (1-x^2)^{l/2}}{2^l l!} \frac{d^{2l}}{dx^{2l}} (x^2-1)^l \right] \Big|_{x=\cos\theta}$$

$$= \sqrt{\frac{(2l+1)!}{4\pi(2l)!}} e^{il\phi} \left[\frac{(-1)^l}{2^l l!} (1-x^2)^{l/2} \frac{d^{2l}}{dx^{2l}} (x^2-1)^l \right] \Big|_{x=\cos\theta}$$

$$l=0 \Rightarrow \frac{d^0}{dx^0} (x^2-1)^0 = 1, \quad l=1 \Rightarrow \frac{d^2}{dx^2} (x^2-1)^1 = 2, \quad l=2 \Rightarrow \frac{d^4}{dx^4} (x^2-1)^2 = 24, \quad l=3 \Rightarrow \frac{d^6}{dx^6} (x^2-1)^3 = 720$$

$$\Rightarrow \frac{d^{2l}}{dx^{2l}} (x^2-1)^l = (2l)! \Rightarrow Y_l^l(\theta, \phi) = \sqrt{\frac{(2l+1)!}{4\pi(2l)!}} e^{il\phi} \left[\frac{(-1)^l (1-x^2)^{l/2}}{2^l l!} (2l)! \right] \Big|_{x=\cos\theta}$$

Satisfy check.

$$= \frac{(-1)^l}{2^l l!} \sqrt{\frac{(2l+1)!}{4\pi}} e^{il\phi} \sin^l\theta$$

$$\sin\theta \frac{\partial}{\partial\theta} \left(\sin\theta \frac{\partial Y_l^l}{\partial\theta} \right) + \frac{\partial^2 Y_l^l}{\partial\phi^2} = \sin\theta \frac{\partial}{\partial\theta} \left[\sin\theta \frac{\partial}{\partial\theta} \left(\frac{(-1)^l}{2^l l!} \sqrt{\frac{(2l+1)!}{4\pi}} e^{il\phi} \sin^l\theta \right) \right] + \frac{\partial^2}{\partial\phi^2} \left[\frac{(-1)^l}{2^l l!} \sqrt{\frac{(2l+1)!}{4\pi}} e^{il\phi} \sin^l\theta \right]$$

$$= \frac{(-1)^l}{2^l l!} \sqrt{\frac{(2l+1)!}{4\pi}} \sin^l\theta \frac{d^2}{d\phi^2} (e^{il\phi}) = -l(l+1) \sin^2\theta \left[\frac{(-1)^l}{2^l l!} \sqrt{\frac{(2l+1)!}{4\pi}} e^{il\phi} \sin^l\theta \right]$$

$$= -l(l+1) \sin^2\theta Y_l^l(\theta, \phi)$$