

NUCLEAR DEPENDENCE OF PROTON-INDUCED DRELL-YAN DIMUON
PRODUCTION AT 120GEV AT SEAQUEST

BY

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DISSERTATION

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Abstract

A measurement of the atomic mass dependence of Drell-Yan dimuons produced by 120GeV protons is presented here. Over 88,000 dimuon pairs with dimuon mass $M \geq 4.2\text{ GeV}$ were recorded at SeaQuest from targets 1H , 2H , C , Fe , and W . The ratio of dimuon yield per nucleon for heavy nuclei versus 2H , $R = Y_A/Y_{^2H}$, is sensitive to modifications in the anti-quark sea in nuclei for the case of proton-induced Drell-Yan. No nuclear modifications are observed over the target quark momentum fraction range $0.1 \leq x_t \leq 0.45$. These results are compared to predictions of models of the EMC effect.

Dedicated to my grandfather, Ted, who took me to school every day.

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List of Abbreviations

ACNET	Accelerator control system network
AD	Accelerator division (of Fermilab)
BIM	Beam intensity monitor
BNC	Bayonet NeillConcelman signal cable connector standard
BOS	Beginning of spill
BPM	Beam profile monitor
CEBAF	Continuous Electron Beam Accelerator Facility
CODA	CODA On-line Data Acquisition
CS	Collins-Soper reference frame
CTEQ	Coordinated theoretical-experimental project on QCD
DAQ	Data acquisition
DC	Drift chamber
DIS	Deep-inelastic scattering
DY	Drell-Yan
EOS	End of spill
EOF	End of file
EPICS	Experimental physics and industrial control system
EVIO	CODA event input/output
FEE	Front-end electronics
FNAL	Fermi National Accelerator Laboratory (Fermilab)
FPGA	Field programmable gate arrays
IC	Ion chamber
LED	Light-emitting diode
LINAC	Fermilab Linear Accelerator
LO	Leading order

MI	(Fermilab) Main Injector
MOSFET	Metaloxidesemiconductor field-effect transistor
NDF	Neutral density filter
NLO	Next-to-leading order
NM	Neutrino-muon beam line
NNLO	Next-to-next-to-leading order
PCB	Printed circuit board
PDF	Parton distribution function
PID	Particle identification
PMT	Photomultiplier tube
RF	Radio frequency
RFQ	Radio frequency quadrupole
ROC	Readout controller
SEM	Secondary emission monitor
SHV	Secure high voltage connector standard
SLAC	Stanford Linear Accelerator Center
SRC	Short range correlations
TS	Trigger supervisor
QCD	Quantum chromodynamics
QED	Quantum electrodynamics
QIE	Charge (Q) integrator and encoder
RDBMS	Relational database management system
SQL	Structured querying language
SWIC	Segmented wire ion chamber
VME	Versa Module Europa

Preface

Preface here talking about what is contained in this thesis. Emphasize that the thesis contains the summary of my personal efforts on the SeaQuest experiment, and describe the layout of the paper. Mention that the Apparatus, PMT, and DB sections are standalone chapters that are not required in order to understand the rest of the chapters.

Chapter 1

Theoretical Background

The topic of this paper is the exploration of the sub-structure of the *nucleon*, a particle that makes up an atom's nucleus which can be either a proton (p) or a neutron (n). By sub-structure, I refer to the nucleon's composition and the momentum distributions carried by the nucleon's constituent particles, quarks and gluons.

While a complete review of the history and physics behind nucleon structure and its investigative probes is beyond the scope of this paper, a brief overview of Drell-Yan, parton distribution functions, and nuclear structure phenomenology will help in understanding concepts and terminology relevant to this and later chapters.

1.1 Introduction

The first indication that the proton may have some internal structure was in a 1933 experiment by Estermann *et al.* measuring the magnetic moment of the proton [41]. Since the proton was thought to be a point-like Dirac particle, it's magnetic moment (μ_p) was expected to be $\mu_p = \frac{e}{2m_p} = 1\text{n.m.}$, or one *nuclear magneton*. The experiment resulted in a value of 2.5 n.m., leading many to reconsider the notion that the proton is indeed point-like.

Around the same time, Hideki Yukawa is credited for establishing the first theory of a *strong force*, a force binding together nucleons in a nucleus against the sizable *Coulomb* repulsion of protons against each other. The force was theorized to be mediated by the exchange of particles called *mesons*, and its range was limited to nuclei-scale distances, seeing as it's not observed at larger distances. Based on the size of the nucleus, Yukawa estimated the mass of the intermediating particles to be approximately $2 \times 10^2 m_e \approx 100\text{MeV}$, where m_e is the electron mass. The following year, Anderson *et al.* discovered the muon (μ) at around this mass [10], which confused many as it did not seem to partake in strong interactions. Eventually, by 1947, the meson theory was validated by the discovery of the *pion* by Powell *et al.* [90], and Yukawa was awarded a Nobel Prize for his theory in 1949. While the pion turned out to be just another composite particle, its

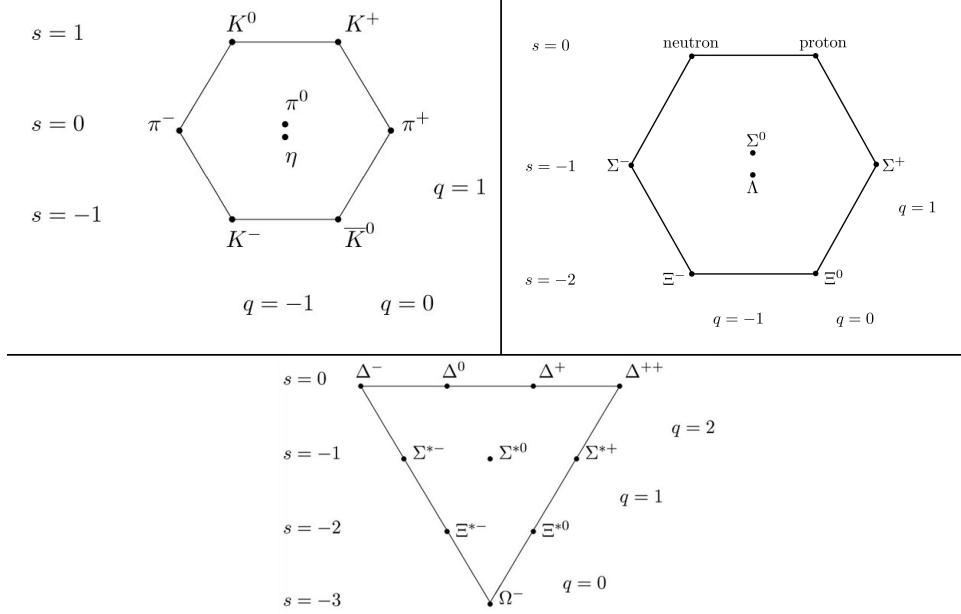


Figure 1.1: The octets and decuplet of Gell-Mann’s “Eightfold Way”. Top left: the eight particles of the meson octet; top right: the spin- $\frac{1}{2}$ baryon octet; bottom: the spin- $\frac{3}{2}$ baryon decuplet.

discovery was a watershed moment in particle physics that led to a cascading series of discoveries.

Near the end of the 1960’s, SLAC began the first set of experiments to use an electron beam to investigate the nucleon in what is now known as Deep Inelastic Scattering (DIS). What they discovered was that the functions used to describe nuclear structure were largely independent of momentum transfer (Q^2) above a certain threshold. This so-called “scaling” behavior lent credence to the theory that nucleons were composed of point-like particles. This discovery coupled with the proliferation of new particles (the “particle zoo”) led Murray Gell-Mann and Yuval Ne’eman to construct a framework that would make some sense of it all. Gell-Mann had organized the host of mesons and baryons discovered into a geometric order named the “Eightfold Way”¹ as depicted in Figure 1.1. The underlying explanation for all of this, as described independently by Gell-Man and Ne’eman, was to characterize these many particles as the several combinations of three *flavors* of constituent particles, named *quarks* [67]. The flavors, *u*, *d*, and *s* were named *up*, *down*, and *strange*. Here, mesons were predicted to be composite particles of integer spin containing two quarks and baryons be half-integer spin particles composed of three quarks. This model led Gell-Mann to predict the existence of the Ω^- particle, along with its *strangeness*, charge, and mass. The discovery of this particle [17] earned Gell-Mann the 1969 Nobel Prize for his work on the quark model.

Though this model was a breakthrough in the understanding of fundamental particles, it did not comprehensively describe all experimental data. For example, when accounting for the total momentum of a

¹Gell-Mann here makes a rather poetic allusion to the Buddhist “Noble Eightfold Path”

nucleon, it was found that only $\sim 50\%$ of the momentum was being carried by the quarks. It was not until the theory of quantum chromodynamics (QCD) came along that this and several other mysteries could be explained. QCD described the mechanism of a 3-fold *color* charge and the gauge bosons, gluons, that intermediate the strong force between quarks (and other gluons). In the particular case of the missing momentum, this was found to reside in the intermediating gluons [16], which carry only color charge, to which the electroweak probes used were insensitive.

Many nuclear probes and methods have been used to characterize the distributions and characteristics of nuclear partons, including the aforementioned DIS process. As I will describe later in this section, much has been learned about parton probability distribution functions with this and other processes, but it is the process of hadron-hadron di-lepton production that is of primary focus in this paper, which allows for an alternative approach to investigate the structure and characteristics of the quarks inside of a nucleon, and perhaps shed some light on some ongoing mysteries in nuclear physics.

1.2 Di-lepton production in nucleon-nucleon collisions

Studying the production of pairs of leptons resulting from hadron-hadron collisions has proven to be a powerful tool in probing nucleon structure and parton distribution functions (PDFs). Figure 1.2 shows the mass spectrum of the muon mode of di-lepton production in proton-nucleus collisions measured at the E-866/NuSea experiment at Fermilab [60]. The resonances of the J/Ψ , Ψ' , Υ , Υ' , and Υ'' can be seen atop a smooth, continuous distribution which decreases with mass. The process responsible for this distribution of $\mu^+\mu^-$ pairs is the Drell-Yan process [34], which is illustrated in the Feynman diagram in Fig. 1.3. This process is characterized by a quark annihilating with an anti-quark to form a virtual time-like photon which then decays into a pair of leptons. The study of this process can lend some insight into nucleon structure due to the fact that the mass and momentum of the di-lepton pair directly reflects the momentum distributions of the interacting quarks and anti-quarks. The most comprehensive and precise knowledge and models of these momentum distributions come primarily from deep-inelastic lepton scattering (DIS) and the Drell-Yan (DY) process.

One particular focus of study for the Drell-Yan process is its nucleon number-dependent (A -dependent) behavior of its cross sections. This is of particular interest due to a phenomenon known as the EMC effect, in which the European Muon Collaboration (EMC) discovered in 1983 that parton distribution functions become modified when in the presence of the nuclear medium. This A -dependent behavior was not expected for hard scattering processes as the modest (maximum 8.8 MeV) binding energy of a nucleus on a nucleon was

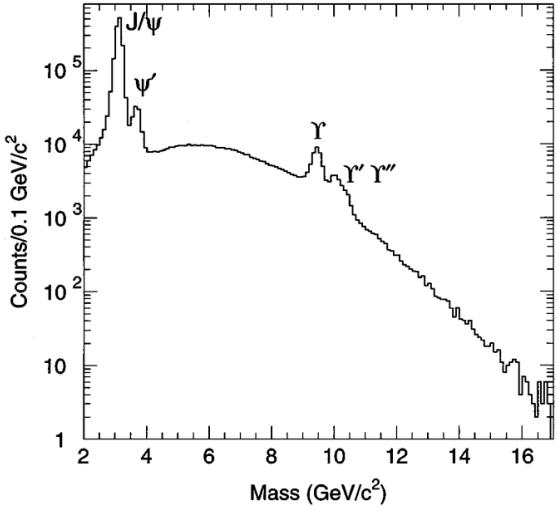


Figure 1.2: Dimuon mass spectrum from E-866 $p + h$ collisions at 800 GeV/c [60].

not thought to have a significant effect on quark momentum distributions within a nucleon ($938\text{ Mev}/c^2$). Many experiments since have confirmed, extended, and precisely characterized the A-dependent behavior observed. Among the many hard scattering probes for investigating this phenomenon, the Drell-Yan process is uniquely capable of isolating the effect on the anti-quark distributions to a high degree of accuracy. The study of the effects of the nuclear medium on nucleon anti-quark distributions is the secondary research goal of the SeaQuest experiment and is the main focus of this thesis.

1.2.1 The Drell-Yan Process

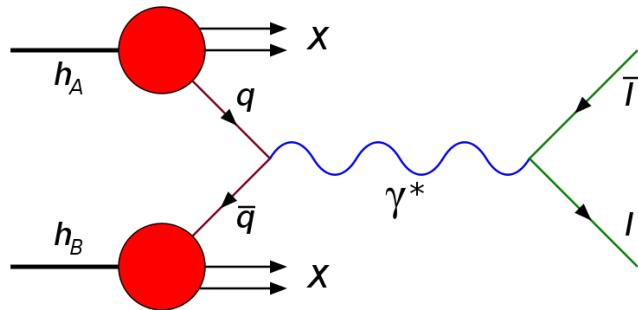


Figure 1.3: The Drell-Yan process, an s channel interaction consisting of the annihilation of a quark with an anti-quark.

The study of continuum di-lepton production in hadron collisions,

$$h_A + h_B \rightarrow l^+l^- + X \quad (1.1)$$

can be used to study hadronic structure in a way that is complementary to the study of deep-inelastic scattering,

$$l + h \rightarrow l' + X. \quad (1.2)$$

In 1970, S. Drell and T.M. Yan were the first to suggest that, at high $Q^2 (= M_{l^+l^-}^2) \geq 16\text{GeV}^2$, the quarks inside the hadrons h_A and h_B can be considered free fermions in the instantaneous moment that they interact. The Drell-Yan model addresses the dominant subprocess here,

$$q_A + \bar{q}_B \rightarrow \gamma^* \rightarrow l^+l^- \quad (1.3)$$

as an electromagnetic annihilation process. In this high- Q^2 kinematic space, the final state of the hadrons that contain these quarks becomes irrelevant. By the energy-time uncertainty principle [1] of $\Delta E \Delta t \sim \hbar$, the timescales under consideration are $< 10^{-25}\text{s}$, and the corresponding distances are $< 10^{-17}\text{m}$, where the size of a nucleon is $\sim 10^{-15}\text{m}$ [89]. At these short space-time scales, electromagnetic annihilation dominates this quark-quark interaction.

This is further reinforced by the widely accepted non-Abelian gauge field theory of Quantum Chromodynamics (QCD) which describes the interactions between quarks and gluons. QCD provides a theoretical justification for treating Drell-Yan processes as events isolated from the rest of the hadron states, and it does so through the concept of the running of the coupling constant, or *asymptotic freedom* [20]. This characteristic feature of QCD is described by the decrease of the strength of the strong force coupling constant as the space-time scale of the interaction approaches zero (i.e. as $Q^2 \rightarrow \infty$). The strong coupling constant can be expressed as a function of Q^2 :

$$\alpha_s(Q^2) = \frac{1}{\beta_0 \ln(Q^2/\Lambda^2)} \quad (1.4)$$

where

$$\beta_0 = \frac{12\pi}{33 - n_f} \quad (1.5)$$

Here, Λ is the QCD scale parameter that depends on the number of quark flavors, n_f , and the renormalization scheme, measured to be around $\sim 217\text{MeV}$. Experimental measurements of the running of the coupling

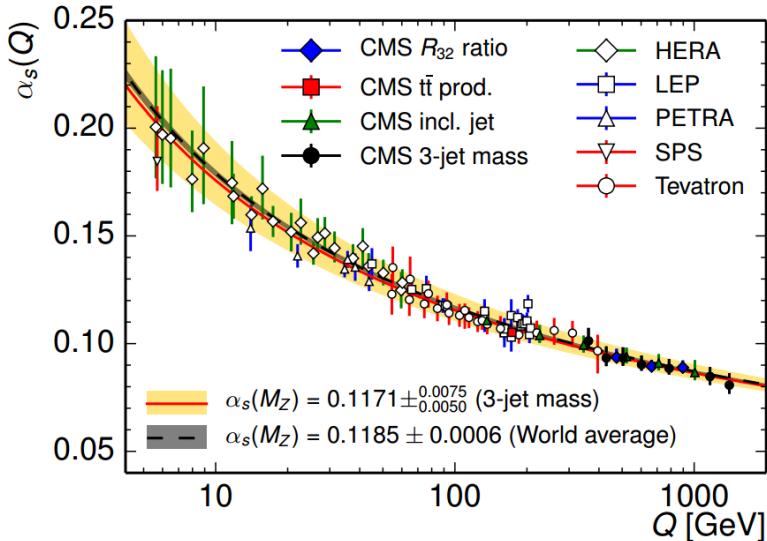


Figure 1.4: A compilation of measurements of the running of the strong coupling constant, $\alpha_s(Q^2)$ [72].

constant as measured by many different processes can be seen in Figure 1.4.

It is useful to have some intuition as to the reason for these “running” of the coupling constants. The phenomenon responsible are termed “screening” and “anti-screening” (or “camouflaging”) [91]. This straightforward in the case of the electromagnetic force which can be seen depicted in the left pane of Figure 1.5. Whether it is in some kind of dielectric medium or in an isolated vacuum (where quantum fluctuations lead to virtual e^+e^- pairs), the charges surrounding an individual charged particle are oriented in such a way as to decrease the overall observed charge. In the realm of QED, the higher energies you use to probe a particle, the closer you actually get to seeing the “actual” charge of the particle, and as such, the EM coupling constant increases at higher energies.

The same is seen in the case of QCD with its non-binary charge scheme. It’s more difficult to conceptualize, but the same screening effect exists in the case of a particle with color charge. The middle pane of Fig. 1.5 is a depiction of screening analogous to the familiar EM screening. The difference between QCD and QED with the behavior of the coupling constants is attributable to the fact that the force-carrying boson of QCD, the gluon, carries with it its own color charge. In the QED picture, the charge at the center is continuously emitting and absorbing photons as it interacts with the surrounding charges, but its charge does not change. With the color-charged particle, gluons are also being continuously emitted and absorbed, but an effect of that is a change in color charge. The result ends up being the phenomenon of camouflaging where the closer you are to a color-charged particle, the less that you can “see” the *net* charge, and the actual charge of the particle is obfuscated. The net charge of the particle, statistically, becomes neutral,

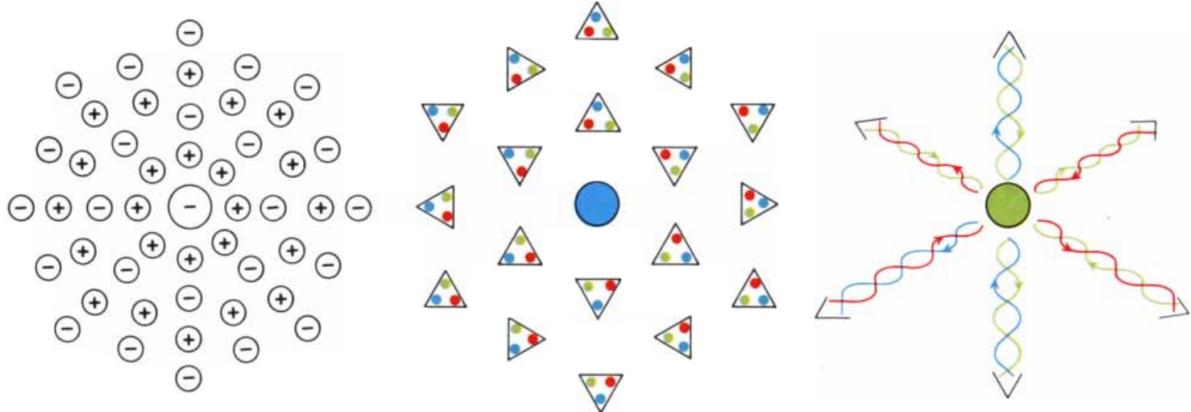


Figure 1.5: (Left) classic electromagnetic screening of a free charge in a dielectric or even a vacuum. (Center) The analogous screening in the three-charge QCD regime. (Right) The color camouflaging due to the constant absorption/emission of colored gluons [91].

causing the strong force to become increasingly feeble at smaller scales (compared to the strength of QED interactions). This effect competes with and surpasses the color screening that occurs and is the root cause of the asymptotic freedom in QCD. On the flip side, the farther you get from the color charge, the more you are exposed to the full true net charge, and the stronger the coupling becomes. This increase in the coupling is the QCD confinement that is observed. An example if in the right pane of Fig. 1.5 where the quark emits many gluons, thus changing its color from Blue to Green, but when summing over all colors emitted and absorbed, the net color is still Blue.

1.2.2 Drell-Yan Kinematics

In the center-of-mass frame, the Drell-Yan process can be broken down into three stages with three sets of kinematics (refer to Fig. 1.3). Beginning with the quarks, x is defined as the fraction of the hadron's momentum carried by the interacting quark or antiquark. Conventionally in fixed-target experiments, subscripts are assigned as x_1 and x_2 , which refer to the quark/antiquark from the beam and the antiquark/quark from the target, respectively. This x is called the Bjorken x , and is well-known in DIS processes to have a value of

$$x = -q^2/2p \cdot q \quad (1.6)$$

where p and q are the 4-momenta of the hadron and the photon. In DY, the same momentum fraction x is used to describe part of the fundamental basis of the interaction. In this case, q is the 4-momentum of the virtual photon and p is the 4-momentum of the hadron that contains the quark in question.

The next stage of the process is the virtual photon. This photon's properties are effectively equivalent

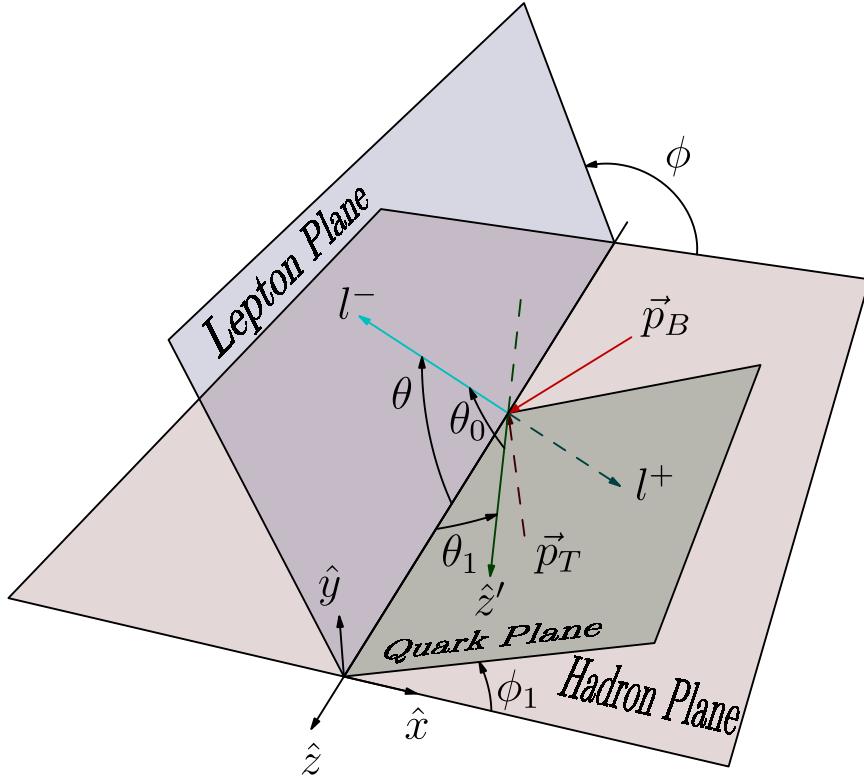


Figure 1.6: A depiction of the Collins-Soper reference frame by Evan McClellan (UIUC). The CS variables are best understood by identifying the three defined planes and then referring to the specific angles between them.

to the ‘dimuon’ or ‘dilepton’, which are the terms more commonly used in referring to kinematics. The first of its relevant kinematics is its mass M_{γ^*} , which represents its energy and virtuality. The value x_F , or Feynman-x, is the fraction of the maximum possible longitudinal momentum carried by the virtual photon in the beam direction. The transverse momentum, p_T , and the azimuthal production angle, ϕ_{γ^*} are the remaining kinematics associated with the virtual photon.

The third stage regards the pair of leptons produced. In the frame of the virtual photon, there is a polar and azimuthal decay angle, θ_μ and ϕ_μ , respectively, for each of the decay muons. It becomes impossible, however, to reconstruct these variables, as the individual transverse momenta of the quarks are unknown, and thus the quark-antiquark annihilation axis is unknown. This is remedied by shifting the process into the Collins-Soper (CS) reference frame [28] which orients the reference axis to be parallel to the bisector of the angle between the interacting hadrons in the rest frame of the muon pair. A depiction of the Collins-Soper frame can be found in Fig. 1.6. In total, this brings a total of eight kinematic variables, summarized in Table 1.1.

Experimentally, six independent variables are measured, which form a basis by which all eight can be

known. This is due to the fact that two pairs of variables, (M_{γ^*}, x_F) and (x_1, x_2) are correlated by the following, ignoring quark masses and considering $p_\perp \ll p_\ell$:

$$x_F \equiv \frac{p_\ell}{\sqrt{s}/2} \approx x_1 - x_2 \quad (1.7)$$

$$M_{\gamma^*}^2 \equiv E^2 - p_\ell^2 = sx_1 x_2 \quad (1.8)$$

$$E = \frac{1}{2}(x_1 + x_2)\sqrt{s} \quad (1.9)$$

$$p_\ell = \frac{1}{2}(x_1 - x_2)\sqrt{s} \quad (1.10)$$

In this frame, the longitudinal momenta of the quarks are $x_1\sqrt{s}/2$ and $-x_2\sqrt{s}/2$, with \sqrt{s} being the center of mass energy of the hadronic collision. So, by measuring the 3-momenta of the μ^+ and μ^- , the quantities $(M_{\gamma^*}, x_F, p_T, \phi_{\gamma^*}, \theta_\mu, \phi_\mu)$, which is sufficient to calculate the remaining (x_1, x_2) in our approximation.

Variable	Description
$x_{1/2}$	Momentum fraction of the beam/target quark
M_{γ^*}	Mass of the virtual photon (dimuon)
x_F	Fraction of the max. possible p_ℓ carried by virt. photon
p_T	Transverse momentum carried by the virt. photon
θ_μ, ϕ_μ	Polar and azimuthal decay angle of one of the muons, in the CS ref. frame
α	The fine structure constant
$K(x_1, x_2)$	High-order QCD correction term
\sqrt{s}	Center of mass energy of the hadronic collision
$\sqrt{\hat{s}}$	Center of mass energy of the $q\bar{q}$ collision
Q^2	Four-momentum of the intermediate time-like photon, squared
$q_i^{t/b}(x)$	The quark number density in the nucleon of the target/beam

Table 1.1: Kinematic variables relevant to the Drell-Yan process .

1.2.3 Cross-Section

As the participating quarks are asymptotically free within each hadron, there will be no correlations between the probability distributions of the annihilating particles, and the process is independent of the distributions. As a result, the cross section of the Drell-Yan process can be reduced to a function of the electromagnetic annihilation process and the quark probability distribution functions. With these components and some QCD considerations, we can construct it piece by piece.

The first step is to begin with the known hard scattering cross section of $\epsilon + \bar{\epsilon} \rightarrow l^+l^-$, where ϵ is an

arbitrary particle. This cross section[58] is given by

$$\sigma(\epsilon\bar{\epsilon} \rightarrow l^+l^-) = \frac{4\pi\alpha^2}{3M_{\gamma^*}^2} e_f^2 \quad (1.11)$$

where α is the electromagnetic fine structure constant, e_f is the charge of the particle, and M_{γ^*} is the dilepton mass. With this, we add the QCD consideration that only $q\bar{q}$ of opposite color can annihilate with each other into a colorless virtual photon. Possible combinations are $R\bar{R}$, $B\bar{B}$, and $G\bar{G}$ out of 3×3 possible cases. As such, an overall factor of $\frac{1}{3}$ is added to this cross section. Finally, factoring in the conservation of flavor (there can only be $u\bar{u}$, $d\bar{d}$, etc. combinations) and the quark structure of hadrons A and B, we use the product ($q_f^A(x_1)\bar{q}_f^B(x_2)$) of the quark probability distributions for finding quarks of the same flavor-antiflavor combination in the two hadrons. It must also be considered that the quark or antiquark may be found in either hadron A or hadron B. The product of these three factors leads us to the Drell-Yan cross section[33]

$$\frac{d^2\sigma}{dx_1 dx_2} = \frac{1}{3} \frac{4\pi\alpha^2}{3M_{\gamma^*}^2} \sum_f e_f^2 [q_f^A(x_1)\bar{q}_f^B(x_2) + \bar{q}_f^A(x_1)q_f^B(x_2)] \quad (1.12)$$

$$= \frac{4\pi\alpha^2}{9sx_1x_2} \sum_f e_f^2 [q_f^A(x_1)\bar{q}_f^B(x_2) + \bar{q}_f^A(x_1)q_f^B(x_2)] \quad (1.13)$$

where the sum is summing over flavors of quarks ($f \in \{u, d, s, \dots\}$). This can be evaluated in terms of the measurables M_{γ^*} and x_F via Equations 1.7 and 1.8

$$M_{\gamma^*}^2 \frac{d^2\sigma}{dM_{\gamma^*}^2 dx_F} = \frac{1}{3} \frac{4\pi\alpha^2}{3M_{\gamma^*}^2} \frac{x_1 x_2}{x_1 + x_2} \sum_f e_f^2 [q_f^A(x_1)\bar{q}_f^B(x_2) + \bar{q}_f^A(x_1)q_f^B(x_2)] \quad (1.14)$$

where x_1 and x_2 can be expressed as

$$x_1 = \frac{1}{2} \left[\sqrt{x_F^2 + 4\tau} + x_F \right], \quad x_2 = \frac{1}{2} \left[\sqrt{x_F^2 + 4\tau} - x_F \right], \quad \tau = \frac{M_{\gamma^*}^2}{s} \quad (1.15)$$

The cross section can also be represented by dimensionless variables in its scaling form,

$$s \frac{d^2\sigma}{d\sqrt{\tau} dy} = \frac{1}{3} \frac{4\pi\alpha^2}{3} \sum_f e_f^2 [q_f^A(x_1)\bar{q}_f^B(x_2) + \bar{q}_f^A(x_1)q_f^B(x_2)] \quad (1.16)$$

where we introduce the rapidity term y in describing x_1 and x_2 ,

$$y = \frac{1}{2} \ln \frac{E + p_\ell}{E - p_\ell} = \frac{1}{2} \ln \frac{x_1}{x_2}, \quad x_1 = \sqrt{\tau} e^y, \quad x_2 = \sqrt{\tau} e^{-y} \quad (1.17)$$

It should be noted that with collider experiments, it is conventional to refer to the hadrons A and B in terms of the beam and target hadrons, respectively, and as such, x_1 refers to the the quark in the beam hadron and x_2 refers to the quark in the target hadron.

In each of these equations 1.13, 1.14, and 1.16, the cross sections can be factored into two parts: one subprocess cross section and one part that has only a dependence on the parton distribution functions. They are independent of each other because one of the staples of the quark parton model is that the PDFs ($q(x)$ and $\bar{q}(x)$) are independent of the process by which they are probed. This feature is commonly referred to as the “universality” of the PDFs.

1.3 Nucleon Structure

There are a few distributions, functions, terminologies, sum rules, and integrals that must be established and explained in order to speak the language of nucleon structure. Here, some of these are explained so that they may be used freely later. Also included are brief sections about how nucleon structure was tested by the Drell-Yan process, and how the QCD understanding of the nucleon was used to improve the understanding of the Drell-Yan process.

1.3.1 PDFs, Structure Functions, and the Quark Parton Model

In the context of hadron-hadron collisions, one can interpret the PDFs as follows: consider two hadrons A and B colliding; a parton of type a ($a \in \{u, d, s, g, \dots\}$) comes from A and carries with it a fraction of A’s momentum (x_A). The same goes for hadron B; a parton of type b comes from B and carries momentum fraction x_B . Now, the probability of finding the discussed parton from A at momentum fraction x_A is given by $q_{a/A}(x_A)$. Likewise, the probability of finding the discussed parton from B at x_B is $q_{b/B}(x_B)$.

In general, experimental data measuring F_2 is used in conjunction with some general constraints in order to arrive at the calculated PDFs. One constraint is to consider the known number of valence quarks in a nucleon. Since the function $q_f(x)$ can be interpreted to be the number density of quark flavor f as a function of momentum fraction x , the value $q_f(x)dx$ represents the number of quarks with flavor f with fractional momentum in the range of $[x, x + dx]$. Therefore, the known number of valence quarks of flavor f (N_f) provides the following condition.

$$\int_0^1 dx[q_f(x) - \bar{q}_f(x)] = N_f \quad (1.18)$$

Another condition considers the value $xq(x)dx$, which is the total momentum fraction of the hadron of the number of quarks $q(x)dx$. This provides a logical constraint is the summation of the momentum fractions

of all partons must add up to the full nucleon momentum:

$$\sum_{q,\bar{q},g} \int_0^1 dx [xq(x)] = 1 \quad (1.19)$$

It should also be noted that the quark PDFs can be split with respect to their spin alignment (+) or anti-alignment (-) with the nucleon's spin

$$q(x) = q^+(x) + q^-(x) \quad (1.20)$$

and the polarized (or helicity) PDF can be defined as

$$\Delta q(x) = q^+(x) - q^-(x) \quad (1.21)$$

though, due to the fact that SeaQuest uses neither a polarized $\Delta q(x)$. There are two standing proposals, however, to extend the SeaQuest spectrometer for an additional period of data taking, augmenting it with a polarized proton beam (FNAL E-1027) and/or a polarized target (FNAL E-1039). The polarized beam proposal adds two beam-polarizing apparatuses named *Siberian Snakes* to provide sufficiently polarized protons. The polarized target modification calls for the construction and installation of a new cryogenic liquid target system. Both proposals have been approved by the FNAL review committee and are likely to be acted upon near the end of the SeaQuest contract.

The structure functions F_1 , F_2 , and g_1 of the nucleon are interpreted in the QPM as the charge-weighted sums over the quark flavors of the corresponding PDFs:

$$F_1(x) = \frac{1}{2} \sum_f e_f^2 q_f(x) \quad (1.22)$$

$$F_2(x) = \sum_f e_f^2 x q_f(x) \quad (1.23)$$

$$g_1(x) = \frac{1}{2} \sum_f e_f^2 \Delta q_f(x) \quad (1.24)$$

where here f represents all flavors of of quarks and anti-quarks. We see above the relation between $F_1(x)$ and $F_2(x)$ known as the Callan-Gross relation ($F_2 = 2xF_1$). This relation relies on the partons having spin- $\frac{1}{2}$ and no transverse momentum in the infinite momentum frame. The full expression is

$$F_2 = 2xF_1 \frac{1+R}{1+2M_N x/\nu} \quad (1.25)$$

where M_N is the nucleon mass and $R = \sigma_L/\sigma_T$ is the ratio of cross sections for absorbing a longitudinal to that for a transverse photon. Experimentally, R is small[31] ($\lesssim 0.1$) for $x \gtrsim 0.1$ and for $Q^2 \gtrsim 5 \text{ GeV}^2$. Also, when assuming the infinite momentum frame, ν , the energy transferred by the photon is assumed to trend towards infinity. As a result, this relation reduces to the commonly-used Callan-Gross relation.

Due to the complex nature of lattice QCD² simulations, the parton distributions $q_f(x)$ and $\bar{q}_f(x)$ within the nucleus are determined empirically, with only a few rules based on theory. The primary probe used to measure them has been deep inelastic scattering (DIS), examining the inclusive jet production in order to get a clearer picture of nucleon structure. Data from many different experiments are combined to extract the unpolarized PDFs. The earlier parametrizations of the PDFs first relied on fits to the measured structure functions solely from DIS experiments which are primarily sensitive to the light quark distributions and are unable to distinguish between quarks and antiquarks of the same flavor. More modern parametrizations use many different processes to extract supplementary and complementary information that can be incorporated. Lepton-charge asymmetry observed in W^\pm production provides additional light quark distribution information. Jet production and photon measurements are also used to set constraints on gluon distributions. It should be noted that Drell-Yan dilepton production from such experiments as E605 and E866 that have contributed constraints on the light anti-quark distributions from the nucleon sea.

Current experimental determinations of these PDFs compiled by CTEQ (Coordinated Theoretical-Experimental Project on QCD) and TEA and are illustrated in Figures 1.7a and 1.7b. The CTEQ-TEA calculated these PDFs from data on inclusive, high-momentum transfer processes, for which perturbative QCD is expected to be reliable. In the case of deep inelastic lepton scattering, only data with $Q > 2 \text{ GeV}$ is used. Data in this region are expected to be relatively free of non-perturbative effects, such as higher twists or nuclear corrections. Thus, there is no need to introduce phenomenological models for nonperturbative corrections beyond the leading-twist perturbative contributions[36].

Looking to Figure 1.7 we see that at $x > 0.1$, u and d quarks dominate \bar{u} and \bar{d} quarks. The SeaQuest spectrometer (described in detail in the next chapter) is a “forward spectrometer”, which means that, for PDF-relevant purposes, the high- x_F ($x_F > 0$) kinematic phase space is explored, which translates to high x_1 and low x_2 . With this consideration, the Drell-Yan process measured at SeaQuest is probabilistically dominated by a quark from the beam annihilating with an anti-quark from the target. Further, all anti-quarks that exist in the target nucleons must come from what are called the *sea quarks*, or the virtual $q\bar{q}$ pairs that arise from gluons splitting.

Other interesting observations and measurements that have gone into these distributions are the obser-

²An approach to solving quantum chromodynamics situations by simulating quarks and gluons on a set of lattice points in spacetime and modeling their well-defined interactions.

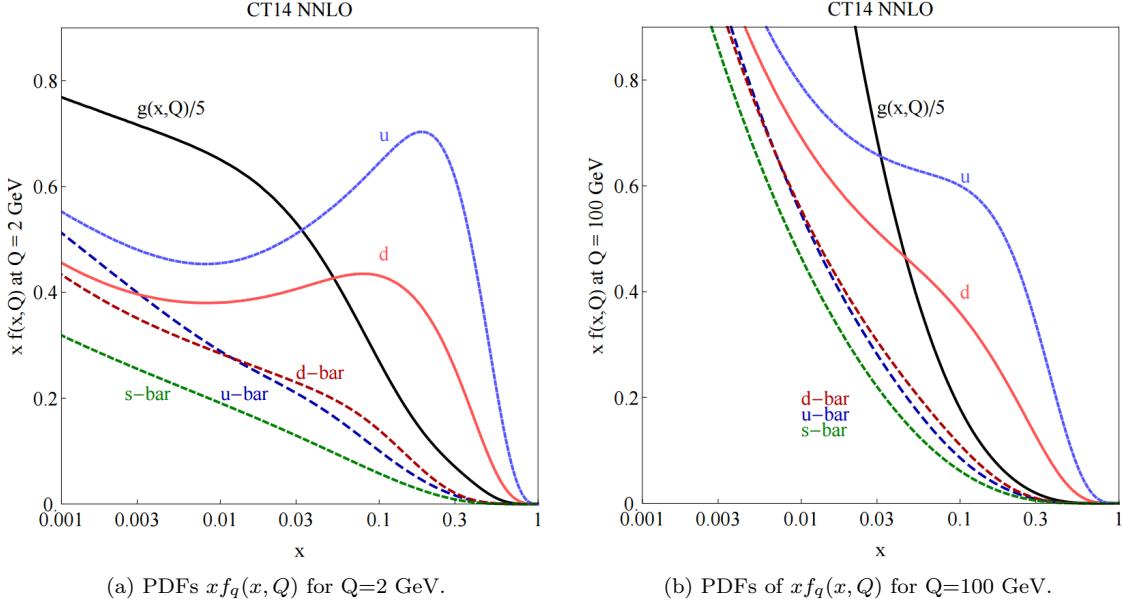


Figure 1.7: Parton distribution functions for quarks and gluons and their x and Q dependence as calculated to NNLO by the CTEQ-TEA (CT) global analysis group[36].

vations of momentum contributions from all quarks and antiquarks and the momentum fraction from just antiquarks. The momentum sum

$$x_{tot} = \int_0^1 dx [F_2(x)] \simeq \int_0^1 dx [x(q(x) + \bar{q}(x))] \quad (1.26)$$

showed that x_{tot} for quarks is 0.512 ± 0.018 [54]. The rest was found to be attributed to the gluons, though they could not be detected directly by an electron beam in DIS since they carry no electric charge. The momentum fraction from just anti-quarks comes from clever use of another structure function, F_3 .

$$F_3^{\nu N}(x) = q(x) - \bar{q}(x) - 2s(x) + 2c(x) \quad (1.27)$$

$$F_3^{\bar{\nu} N}(x) = q(x) - \bar{q}(x) + 2s(x) - 2c(x) \quad (1.28)$$

$$F_3(x) = \frac{1}{2}[F_3^{\nu N}(x) + F_3^{\bar{\nu} N}(x)] = q(x) - \bar{q}(x) = u_v(x) + d_v(x) \quad (1.29)$$

Here, the v subscript denotes valence quark distributions, $q/\bar{q}(x)$ are

$$q(x) = u(x) + d(x) + s(x) + c(x) \quad (1.30)$$

$$\bar{q}(x) = \bar{u}(x) + \bar{d}(x) + \bar{s}(x) + \bar{c}(x) \quad (1.31)$$

and the individual $F_3^{\nu/\bar{\nu}N}$ structure functions are measured from νN and $\bar{\nu} N$ scattering and then combined. With F_3 , there is the Gross-Llewellyn Smith (GSL) sum rule[55]

$$\int_0^1 F_3(x) = 3 \quad (1.32)$$

which is to say that there are three valence quarks. The new information gained here with F_3 though is the measurement of $\int dx[xF_3(x)] = 0.341 \pm 0.036$, which gives the momentum fraction carried by the valence quarks. Put this together with the results of Eq. 1.26 and one can derive the conclusion that the anti-quarks in the nucleon possess between 13% and 17% of the nucleon's momentum[42].

The last interesting sum rule to discuss is the Gottfried sum rule which is focused on flavor asymmetries between the proton and its isospin partner, the neutron.

$$S_G = \int_0^1 dx \left[\frac{F_2^p - F_2^n}{x} \right] \quad (1.33)$$

$$= \int_0^1 dx \left[\sum_f e_f^2 [q_f^p(x) + \bar{q}_f^p(x) - q_f^n(x) - \bar{q}_f^n(x)] \right] \quad (1.34)$$

We assume charge symmetry insofar as $u^p(x) = d^n(x)$ and $\bar{u}^n(x) = \bar{d}^p(x)$, yielding

$$S_G = \int_0^1 dx \left[\frac{1}{3}(u(x) + \bar{u}(x) - d(x) - \bar{d}(x)) \right] \quad (1.35)$$

where these PDFs refer to the distributions in the proton. By some reorganization and grouping, this can be re-expressed in a more familiar way.

$$S_G = \int_0^1 dx \left[\frac{1}{3}[u(x) - \bar{u}(x)] \right] - \int_0^1 dx \left[\frac{1}{3}[d(x) - \bar{d}(x)] \right] - \int_0^1 dx \left[\frac{2}{3}[\bar{d}(x) - \bar{u}(x)] \right] \quad (1.36)$$

The first two integrals here are the definitions of the valence quark sum rule mentioned in Eq. 1.18, and this is known to be 2 and 1 for the u and d valence quarks of the proton, respectively. As a result, we arrive at the common representation of the Gottfried Sum Rule (GSR) [52].

$$S_G = \frac{1}{3} - \int_0^1 dx \left[\frac{2}{3}[\bar{d}(x) - \bar{u}(x)] \right] \quad (1.37)$$

If one were to assume that, within a proton, $\int dx[\bar{u}(x)] = \int dx[\bar{d}(x)]$, then this sum reduces to $\frac{1}{3}$. More on this sum and its measured violation in Section 1.4.1.

1.3.2 Drell-Yan Tests of the QPM

The Drell-Yan process can be described within the frameworks of QPM and QCD, so it provides a good testing ground for these models. Even further, the very clean leptonic DY signature offers a good vantage point by which to confirm the parton picture experimentally. The testing process began with establishing the validity of the quark parton model via Drell-Yan. Once this is done, the following step would be to study the deviation between measurements and what is expected. Here we step through some tests of the QPM that Drell-Yan was applied to in order to validate.

Point-like quarks

If the quarks in the QPM are point-like and without any substructure, then the DY cross section (expressed in Eq. 1.16) across Q^2 should exhibit scaling behavior, i.e. the parton distributions measured should be Q^2 -invariant. The experiments E-288 and E-605 at Fermilab compared the cross section yields across four different beam energies corresponding to \sqrt{s} values between 19.4 and 38.8 GeV. For each setting, the cross section was binned in x_1 and x_2 and compared across different beam settings (this is the same as binning in $Q^2 = M_{\gamma^*} = sx_1x_2$).

If what was being seen was a quark, and the quark was point-like, then the cross section would be Q^2 -independent. If there was substructure to the quark, then at higher Q^2 , sub-quarks would strike or annihilate with other sub-quarks, and one would observe the cross section increasing with Q^2 as new interaction phase space becomes available. The measurements of the experiments exhibited no Q^2 -dependence [66, 23] and lent support to the notion that what was being struck was indeed a point-like particle with no sub-structure.

Quark charge

One of the predictions of QPM was that the up (down) quarks possessed $e_u = +2/3$ ($e_d = -1/3$) times the elementary charge. The use of π^+ and π^- beams to induce Drell-Yan interactions was used in the Fermilab E-444 and the OMEGA experiments to measure the Drell-Yan cross section ratios of π^+ to π^- , which should approach a value of $\frac{e_d^2}{e_u^2}$. The results indicated a value of $\frac{1}{4}$ [56, 65], which is consistent with at least the magnitudes of the charges of the quarks.

Spin-1/2 quarks

If the quarks that annihilate are in fact spin- $\frac{1}{2}$ particles, then the virtual photon that the quarks annihilate to must be predominantly transversely polarized. This, in turn, should define the angular distribution of the

leptons that are produced. As such, a transversely polarized photon should result in a dilepton distribution that goes as $dN/d\theta = 1 + \cos^2 \theta$ in the rest frame [70].

Several experiments sought to test this prediction, but the values of the measurement where $\cos \theta \approx \pm 1$ had a lack of data, allowing for a sizable amount of variation. Regardless, the data collected by CIP, NA3, ABCS, CHFMNP, Omega, and Goliath all fit the angular distribution to the function $f(\theta) = 1 + \lambda \cos^2 \theta$ and arrived at results for λ consistent with $\lambda = 1$, though each had relatively sizable uncertainties [70]. This is in agreement with what the simple Drell-Yan model predicts.

1.3.3 QCD-Improved Drell-Yan

According to the quark-parton picture, the simple Drell-Yan model predicts a certain magnitude of the cross-section and a certain measurement of p_T . Deviations from these predictions indicated a need to accommodate for higher-order effects that arise due to QCD effects.

It was determined over the course of many experiments that the shape of the cross section was predicted very well while the scale of the measurement was off by a non-negligible amount. The discrepancy was studied thoroughly in the form of an arbitrary correction constant, K , known as the K -factor:

$$\frac{d\sigma^{DY}}{dQ^2} \Big|_{exp} = K \cdot \frac{d\sigma^{DY}}{dQ^2} \Big|_{LO} \quad (1.38)$$

where LO stands for leading order (NLO, NNLO stand for next-to- and next-to-next-to-leading order). Relatively wide ranges of values of K were measured [56], from 1.6 to 3.1, with the combined value converging to around 2.

The expected measurement of p_T was also surprising, as the simple Drell-Yan predicted a simple, clean annihilation that carried no other p_T than the intrinsic p_T caused by nucleon confinement and Fermi motion, which should be on the order of 400 MeV (this value depends on the proton dimensions only). It turned out that not only did the measurement of p_T not match this constant value, but the Omega collaboration showed that the value had a dependence on the \sqrt{s} in the form of

$$\hat{p}_T^p = 0.45 + 0.025\sqrt{s} \quad ; \quad \hat{p}_T^\pi = 0.54 + 0.029\sqrt{s} \quad (1.39)$$

for proton-induced pion-induced DY, respectively [70].

These discrepancies can be understood when one considers gluon interactions that happen on very short time scales. When a quark participates in a DY interaction, it may emit a gluon before encountering the other quark. The effects of time dilation in the moving nucleon essentially result in this (*soft*) gluon happening

“long before” the hard interaction occurs. This will not alter the factorized structure of the Drell-Yan cross section. It is, however, the case that gluon emission and absorption can occur at timescales on the order $\Delta t \sim 1/Q$ before the actual annihilation (i.e. *hard* gluons).

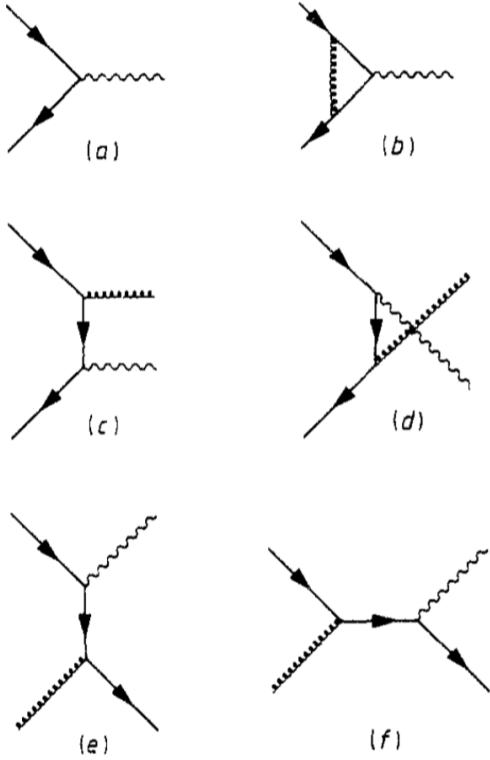


Figure 1.8: Diagrams of (a) LO Drell-Yan, (b) vertex correction, (c-d) gluon emission, (e-f) “Compton” scattering [70].

lower energy proton bombarding deuterium and tungsten at the Fermilab Main Injector (FMI, 120 GeV proton beam) [48] which is used at SeaQuest. Therefore, QCD corrections for nuclear targets (compared to free nucleons) can be considered negligible. As such, when taking ratios of cross sections of different nuclear materials, the K -factors are considered to cancel each other out.

1.4 Drell-Yan and Sea Quark Phenomenology

Interestingly enough, the nucleon is not simply a particle composed of three quarks held together by intermediating gluons. Instead, due to the coupling strength of the strong force and the ability of gluons to couple to gluons, the proton is an impressively complex structure with evanescent virtual quark pairs (and even virtual mesons). Unfortunately, there is difficulty in isolating the specific features of this *nucleon sea* using

the tried and true DIS that had so thoroughly investigated nucleon structure. In this regard, the Drell-Yan process has proven to be crucial in the exploration of the nucleon sea, despite its small cross section. In this section, I discuss the discovery of sea quark asymmetry the meson cloud model (one of many) that attempts to describe it.

1.4.1 \bar{d}/\bar{u} Asymmetry

Even though the quark sea flavor asymmetry is not the main focus of this paper, it is the flagship measurement of the SeaQuest experiment. As such, much of the work that will be explained in later chapters is common to the \bar{d}/\bar{u} analysis, and so it is worth touching on its importance here.

The Gottfried Sum Rule (Eq. 1.37) can be seen as a measure of flavor symmetry (or asymmetry) in the nucleon. If flavor symmetry were to be upheld, then the sum would strictly be $S_G = \frac{1}{3}$. The New Muon Collaboration (NMC) at CERN sought to test this by measuring the cross section ratio for DIS of muons on hydrogen to muons on deuterium [8, 11]. This cross section ratio allowed for the measurement of F_2^n/F_2^p over a Bjorken- x range of $0.004 < x < 0.8$. Incident muon energies of 90 GeV and 280 GeV were used to measure the F_2 ratio at Q^2 of 4 GeV 2 . The full results of the NMC experiment can be seen to be in clear disagreement with flavor symmetry in Figure 1.9. The extrapolated result of the NMC measurements over all x yielded

$$S_G = \int_0^1 dx \left[\frac{F_2^p - F_2^n}{x} \right] = 0.235 \pm 0.026. \quad (1.40)$$

To reconcile this result with the QPM expectation, certain assumptions need to be reevaluated. The first assumption is that the extrapolation of the NMC result to low x is correct. The second is that charge symmetry is valid. The final assumption is the flavor symmetry of the quark sea, or, $\int dx[\bar{d}(x)] = \int dx[\bar{u}(x)]$. The E-665 experiment at FNAL sought to extend the NMC low- x coverage ($10^{-6} < x < 0.3$) despite the difficulties of DIS in the extremely low- x range. The results were consistent with NMC's extrapolation [3], and so the faulty assumption must lie elsewhere. Charge symmetry violation would lead to such a discrepancy, but to date, charge symmetry has been upheld in the relevant energy regime [109]. This leaves only $\int dx[\bar{d}(x)] = \int dx[\bar{u}(x)]$ being violated to account for the difference, and if it were to be the root cause, the contribution to the violation would be

$$\int_0^1 dx[\bar{d}(x) - \bar{u}(x)] = 0.148 \pm 0.039. \quad (1.41)$$

As such, this conclusion was the first indicator pointing to the existence of more anti-down than anti-up quarks in the quark sea. This prompted further investigation of nucleon sea anti-quark asymmetry.

Following the NMC results, it was suggested [40] to exploit features of the Drell-Yan process as a direct

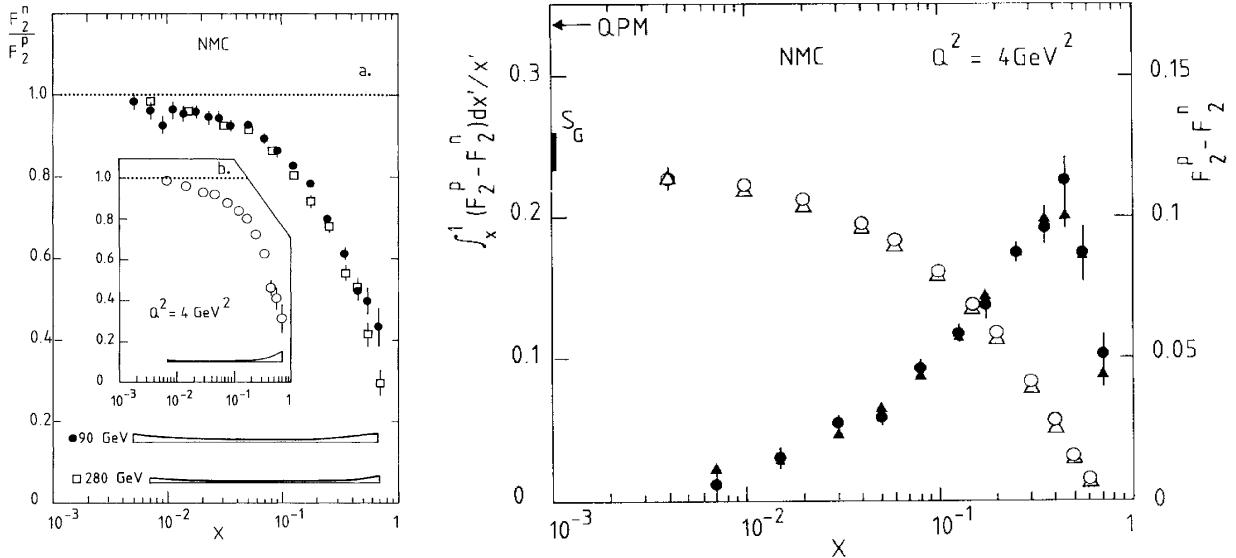


Figure 1.9: The NMC experimental measurements of (left) F_2^n/F_2^p and (right) the Gottfried Sum (S_G) with $\frac{1}{3}$ denoted by “QPM”.

probe of nucleon sea distributions. The most direct method of doing so is to measure the Drell-Yan cross-sections for both hydrogen and deuterium and take the ratio of the two. The proton-deuterium cross section can be seen as a convolution of cross sections of proton-proton and a proton-neutron

$$\sigma^{pd} \approx \sigma^{pp} + \sigma^{pn} \quad (1.42)$$

ignoring any small nuclear effects or phenomenon that may exist in the deuteron. In this picture, the deuterium-to-hydrogen DY ratio yields can be used to determine \bar{d}/\bar{u} .

The NA51 experiment at CERN [15] used a 450 GeV/c primary proton beam ($\sqrt{s} = 29$ GeV) from CERN-SPS on deuterium and hydrogen targets to measure ~ 6000 Drell-Yan dimuon events, cutting on $M_{\gamma^*} > 4.3$ GeV. The reason for the mass cut is to investigate the Drell-Yan continuum as depicted in Figure 1.2 and exclude any possible signal contamination from heavy quarkonia. The NA51 experiment looked to measure what it called the “p–n cross section asymmetry” of the Drell-Yan process:

$$A_{DY} = \frac{\sigma_{pd} - \sigma_{pp}}{\sigma_{pd} + \sigma_{pp}} = 2 \frac{\sigma_{pp}}{\sigma_{pd}} - 1 \quad (1.43)$$

$$A_{DY} = -0.09 \pm 0.02(\text{stat}) \pm 0.025(\text{sys}) \quad (1.44)$$

which was used to obtain a value for \bar{d}/\bar{u} . However, due to the NA51 spectrometer being designed to select rapidity about $y = 0$ along with having low statistics, NA51 only provided a single value for this measurement

centered at $x_F = 0$ and $x = 0.18$.

$$\left. \frac{\bar{d}}{\bar{u}} \right|_{\langle x \rangle = 0.18} = 1.96 \pm 0.15(\text{stat}) \pm 0.19(\text{sys}) \quad (1.45)$$

This provided strong evidence in confirmation that there was indeed a \bar{d} excess in the nucleon. Unfortunately, without more statistics, no x -dependent behavior of this asymmetry could be determined. This NA51 result coupled with that of NMC prompted several groups to perform global fits from DIS and DY data with the consideration that $\bar{d} \neq \bar{u}$, though no x -dependent constraints could be imposed on the $\bar{d}(x)/\bar{u}(x)$.

An improved experimental design of a “forward” muon spectrometer ($x_F > 0$) was devised and implemented by the E-866/NuSea experiment to make a more effective measurement of \bar{d}/\bar{u} . In this kinematic phase space, one can factor in what we know about the PDFs and the DY cross section to consider the measurement as being dominated by the beam quark annihilating with the target’s sea anti-quark. If we consider only the regime of $x_1 \gg x_2$,

$$\sigma^{pp} \propto \frac{4}{9} u(x_1) \bar{u}(x_2) + \frac{1}{9} d(x_1) \bar{d}(x_2) \quad (1.46)$$

$$\sigma^{pn} \propto \frac{4}{9} u(x_1) \bar{d}(x_2) + \frac{1}{9} d(x_1) \bar{u}(x_2) \quad (1.47)$$

$$\frac{\sigma^{pd}}{2\sigma^{pp}} \approx \frac{1}{2} \frac{\left[1 + \frac{1}{4} \frac{d(x_1)}{u(x_1)} \right]}{\left[1 + \frac{1}{4} \frac{d(x_1)}{u(x_1)} \frac{\bar{d}(x_2)}{\bar{u}(x_2)} \right]} \left[1 + \frac{\bar{d}(x_2)}{\bar{u}(x_2)} \right] \quad (1.48)$$

One can approximate even further that since $d(x) \ll 4u(x)$, then the relation simplifies to

$$\left. \frac{\sigma^{pd}}{2\sigma^{pp}} \right|_{x_1 \gg x_2} \approx \frac{1}{2} \left[1 + \frac{\bar{d}(x_2)}{\bar{u}(x_2)} \right] \quad (1.49)$$

The results of the E-866/NuSea measurement [60], along with the NA51 data, can be seen on Figure 1.10. A clear x -dependence can be seen as the \bar{d}/\bar{u} ratio significantly exceeds unity up to $x_2 = 0.25$, after which the ratio drops below unity, though not at the same statistical significance. This complex x -dependent behavior was not expected by global fits, and many models have been proposed to address it: meson-cloud model, chiral-quark model, Pauli-blocking model, instanton model, chiral-quark soliton model, and more. These are not discussed here, but an in-depth look into each of these can be found in some review articles by Kumano [76] and Garvey [47].

In the SeaQuest experiment, $\bar{d}(x)/\bar{u}(x)$ is to be extracted for the range of approximately $0.1 < x_2 < 0.45$. In the run time since commissioning in 2012, enough data has been analyzed to assemble a preliminary $\bar{d}(x)/\bar{u}(x)$ result, which can be seen in Figure 1.10. This measurement and its analysis are being led by

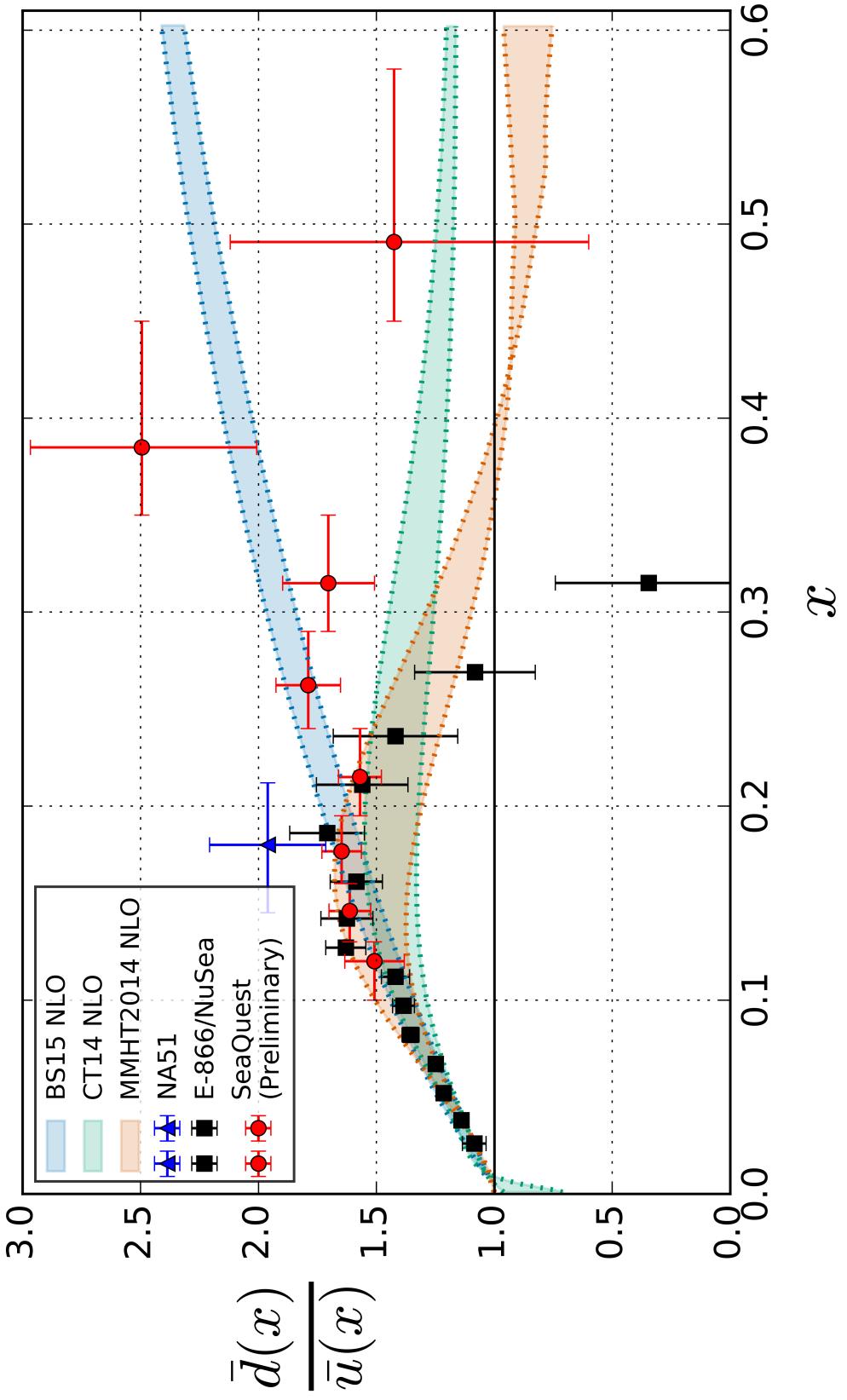


Figure 1.10: The measurements and models for the \bar{d}/\bar{u} measurement, including the NA51 [15], E-866/NuSea [60], and (preliminary) SeaQuest experiments, along with the BS15, CT14, and MMHT models.

Bryan Kerns from the University of Illinois at Urbana-Champaign. Further details regarding this topic can be found in his dissertation and forthcoming publication(s).

1.4.2 Meson Cloud Model

This asymmetry in the nucleon sea was unexpected, though there exists no physical symmetry that does require that $\bar{d} = \bar{u}$. The assumption stands that if one were to consider only the primary process of gluons splitting and producing $q\bar{q}$ pairs, then there would be no reason for $d\bar{d}$ pairs to be preferred to $u\bar{u}$ pairs. The masses of the d and the u quark are similar enough and are small when compared to the nucleon confinement scale, so they should be produced in equal numbers. Even with perturbative considerations, there does not seem to be an asymmetry in the up/down sea in excess of 1% [93]. There is an excellent discussion by Peng *et al.* [86] regarding the possibility of the origin of the \bar{d}/\bar{u} asymmetry in the nucleon sea being attributed to virtual mesons.

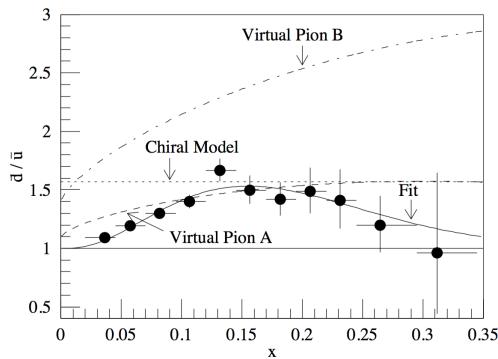


Figure 1.11: The E-866 \bar{d}/\bar{u} results compared to three different models [86]: a virtual pion model with two different $\pi NN, \pi N\Delta$ dipole form factors [75, 73] and the chiral model [102].

can be broken down into $\pi^-\Delta^{++}, \pi^0\Delta^+$, and $\pi^+\Delta^0$ states, with their coefficients being 1/2, 1/3, and 1/6, respectively.

These can account for the flavor asymmetry when considering the following: in the π^+n state there exists a \bar{d} in the “sea”, in the π^0p state there exist equal number of \bar{d} and \bar{u} ($\pi^0 = \sqrt{u\bar{u} + d\bar{d}}/2$), and in the $\pi^-\Delta^{++}$ case, there exists a \bar{u} in the “sea”. The key here is that the Δ baryon is 294 MeV heavier than the n and p , making it highly suppressed as compared to the π^+n and π^0p states. This suppression is the basic concept that can account for the \bar{d}/\bar{u} asymmetry, as the states in which the \bar{d} is produced are more likely than the

As early as 1983, before a hint of a sea quark flavor asymmetry at NA51, Thomas [103] posited that the non-perturbative interpretation of the measured proton state in terms of its principle meson-baryon Fock components results in a preference for $\bar{d} > \bar{u}$. Here, the degrees of freedom being several possible meson-baryon states:

$$|p\rangle_{\text{dressed}} = |p\rangle_{\text{bare}} + a|\pi N\rangle + b|\pi\Delta\rangle + \dots \quad (1.50)$$

Here, πN can be broken into isospin component states π^+n and π^0p , with the Clebsch-Gordon coefficients of the states preferring the π^+n state by a factor of two. The $\pi\Delta$ state

ones in which a \bar{u} is produced. Formally, one can show [101] that

$$\int_0^1 dx[\bar{d}(x) - \bar{u}(x)] = (2a - b)/3 \quad (1.51)$$

and calculations show that the relative probabilities of these configurations $a \approx 2b$ [73]. In Figure 1.11, the E-866 data for \bar{d}/\bar{u} are shown along with three models: two for the virtual pion model with different dipole form factor settings [75, 73] and one for the chiral model [102]. The increased precision and extended x -range of \bar{d}/\bar{u} as measured by SeaQuest is expected to contribute to the evaluation of these models.

1.5 The Nuclear EMC Effect

Approximately ten years before the NMC experiment's discovery of the $\bar{d}(x) \neq \bar{u}(x)$ discovery, the European Muon Collaboration (EMC) at CERN, in 1983, measured the ratios of the structure functions of iron and deuterium over a large kinematic range using the DIS of muons [14]. The result has been repeatedly scrutinized, studied, and confirmed. It was discovered that the structure function of a nucleon bound in a nucleus is fundamentally modified from that of a free nucleon. The nucleus was naively understood as a simple convolution of many nucleons, and that the partonic structure functions of bound and free nucleons should be identical. This assumption was found to not hold against actual measurements, and the observed behavior holds true across all heavy nuclei. This effect can be interpreted as a *modification of quark and gluon distributions in bound nucleons by the nuclear environment*. In this section, the history and collective behavior known as “The EMC Effect” will be reviewed, a quantitative A-dependent measure of the effect will be defined, and models that attempt to explain the phenomenon will be discussed. Lastly, the contributions of SeaQuest and previous Drell-Yan experiments to the study of the EMC Effect will be explored.

1.5.1 Historical Overview

Unpolarized inclusive lepton scattering is a process that uses an energized lepton to investigate nucleon structure, and it specifically depends on two independent variables: Q^2 and Bjorken- x (sometimes referred to as x_B , but here left as x).

$$Q^2 = -q^2 \quad ; \quad x = \frac{Q^2}{2m_p\omega} \quad (1.52)$$

where q is the transferred four-momentum, x_B is the same Bjorken- x scaling variable as is in the PDFs in Section 1.3.1 above, m_p is the proton mass, and ω is the transferred energy in the proton rest frame. At energies of $Q^2 > 2(\text{GeV}/c)^2$, the parton substructure is revealed and the inelastic structure function, F_2

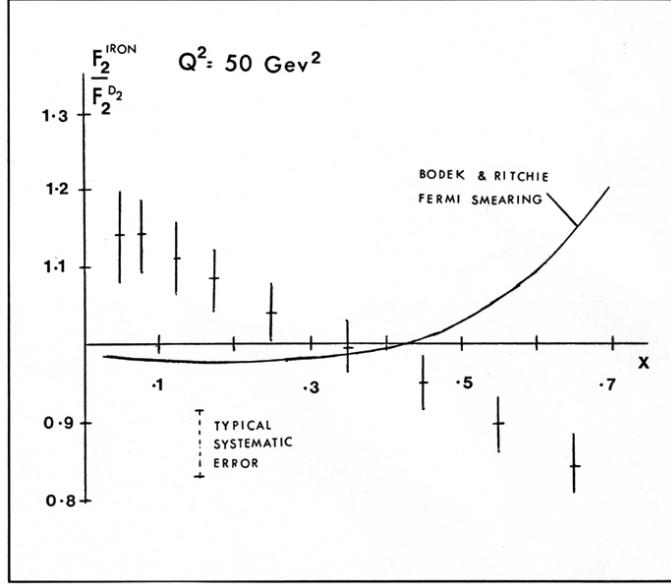


Figure 1.12: The original EMC ratio measurement by the EMC collaboration [14]. Overlaid is the naive, expected ratio taking into account only the Fermi motion of the nucleons in the heavy nucleus.

is able to be measured. The EMC group extracted the ratio $F_2^{Fe}(x, Q^2)/F_2^d(x, Q^2)$, where these are the average bound nucleon structure functions in iron and deuterium. This was done by measuring the DIS *per nucleon cross section ratio*, which can be expressed as [58]:

$$\frac{\sigma^{A_1}/A_1}{\sigma^{A_2}/A_2} = \frac{F_2^{A_1}(x, Q^2)}{F_2^{A_2}(x, Q^2)} \cdot \frac{\left[1 + 2\frac{1+\omega^2/Q^2}{R_{A_1}-1} \tan^2 \frac{\theta}{2}\right]}{\left[1 + 2\frac{1+\omega^2/Q^2}{R_{A_2}-1} \tan^2 \frac{\theta}{2}\right]} \approx \frac{F_2^{A_1}(x, Q^2)}{F_2^{A_2}(x, Q^2)} \quad (1.53)$$

where θ is the lepton scattering angle and $R_A = \sigma_L^A/\sigma_T^A$ is the ratio of the longitudinal to transverse cross section for nucleus A. The A-dependence of R_A will be discussed later, but for now, let us assume that R does not exhibit a nuclear dependence, and the per nucleon cross section ratio therefore approximates nicely to the ratio of the structure functions. For $0.4 \gtrsim x \lesssim 0.6$ where nucleon Fermi motion effects can be neglected, the expectation was to measure unity, which would indicate that the structure function of a deeply bound nucleon and loosely bound nucleons are identical. With this goal in mind, they chose ^{56}Fe and deuterium, as iron is the most tightly bound nucleus at $E_{bind}^{Fe} = 8.8$ MeV/nucleon and deuterium is as close to being a free nucleon as possible while still being a symmetric nuclei ($N = Z = A/2$) with a binding energy of only $E_{bind}^d = 2.2$ MeV/nucleon. The reason to perform this experiment is that, if this measurement showed unity in the regions not affected by Fermi motion, then heavy nuclei could be used to gain a luminosity boost when studying the free nucleon structure functions. Instead, they observed the measured cross-section ratio being 1 at $x \approx 0.3$ and then steadily declining to a value of about 0.8 at $x \approx 0.7$ (Figure 1.12)

At the time of the measurement was presented, the result astounded the physics community. They could hardly believe that high momentum transfers that were many times larger than the nuclear binding energies that the quark distributions should be altered by the nuclear mean field. This result was confirmed only three months later by the SLAC-MIT-Rochester group that re-examined its data for DIS on steel [21] and found the same behavior, as can be seen in Figure 1.13. This opened the gates to an enormous amount of activity in both experiment and theory to fully characterize the phenomenon, resulting at present more than 1100 citations of the original EMC publication [14].

The compendium of experimental data and theory papers have been discussed in great detail in the excellent review papers by Geesaman [49] in 1995 and Norton [84] in 2003. As such, I will only summarize the main theoretical ideas and the most prominent interpretation of the effect. A visual summary of EMC data can be found in Figure 1.14 where data from EMC [14], SLAC [13, 50], BCDMS [18], NMC [8], and HERMES [2] is plotted for nine different targets.

1.5.2 Universal x -dependence

The SLAC E139 experiment provides very precise information for the high- x range ($x > 0.2$) for many nuclei. The high-energy electron DIS data from SLAC E61 [99] and HERMES [2] experiments along with the muon DIS data from NMC [8] supplements this with the very low- x range. Finally, the more recent data from JLAB-E03103 [95] that uses a relatively low energy beam (5.8 GeV) provides data for the high- x Fermi motion area ($0.85 < x < 0.95$). These together help to paint a full characteristic picture of the x -dependence of this σ^A/σ^d ratio, which has been found to be universal for all $A > 2$. A visualization of some of this data for carbon and nitrogen the measurement up into four qualitative regions of this ratio can be seen in Figure 1.15.

- $0 < x \lesssim 0.06$: The “shadowing” region where the ratio is less than unity and decreasing as $x \rightarrow 0$.

The dominant contribution to the cross section in this kinematic regime is from sea quarks.

- $0.06 \lesssim x \lesssim 0.3$: The (perhaps poorly named) “anti-shadowing” region exhibits a very slight rise of the ratio above unity. In general, this range of values represents only a transition region between its two bookending effects.

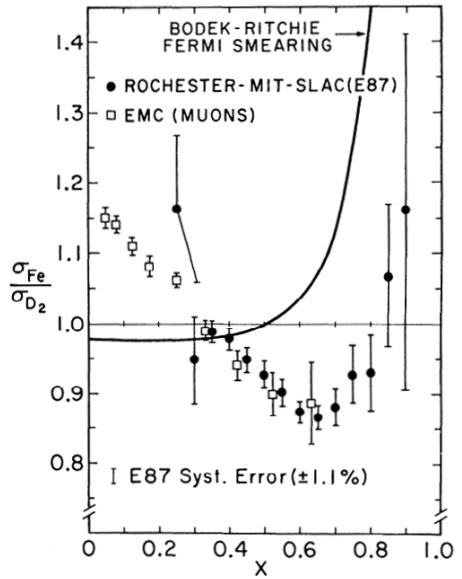


Figure 1.13: SLAC E-87 was quick to report its own finding months later confirming EMC’s observation [21].

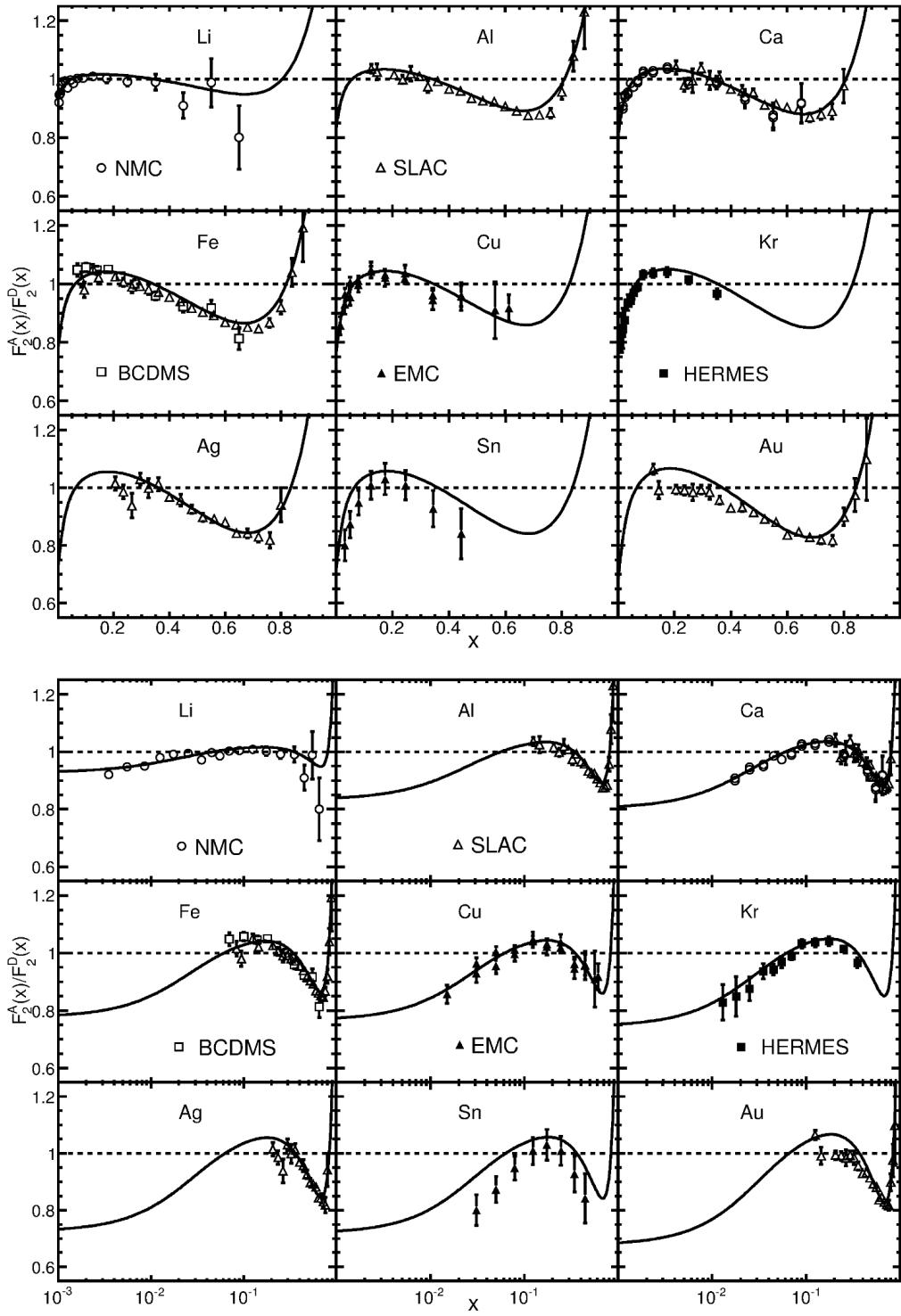


Figure 1.14: A compilation of data regarding the measurement of F_2^A/F_2^D for nine different heavy targets with an overlaid fit by Chen et al. [26].

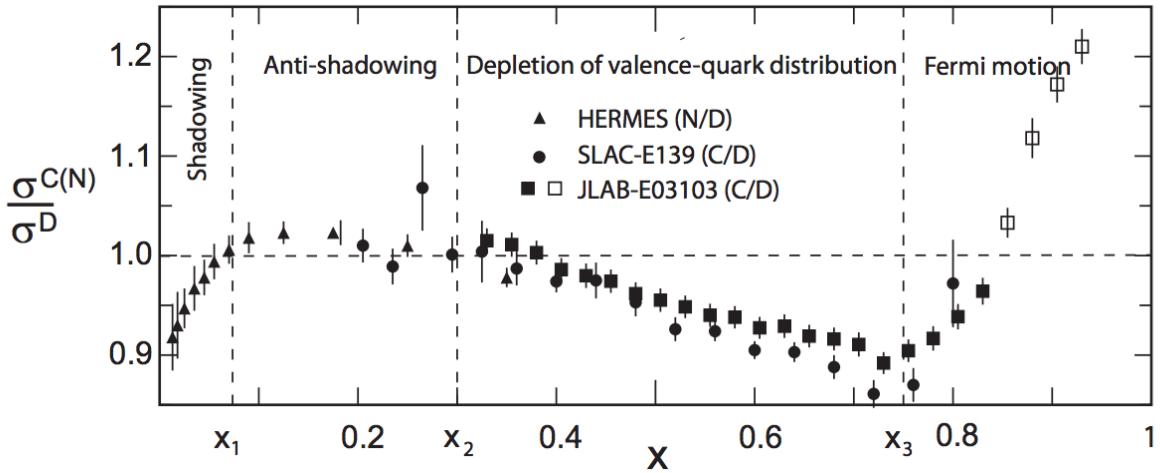


Figure 1.15: The ratio of $\sigma^{C(N)}/\sigma^d$ covering nearly the entire $x \in [0, 1]$ interval [92]. The four qualitative regions are denoted by roughly chosen boundaries x_1, x_2, x_3 (not to be confused with the beam and target x_1 and x_2 used in the rest of this paper).

- $0.3 \lesssim x \lesssim 0.8$: The ratio becomes significantly smaller than unity, decreasing with increasing x . This is of key focus for this paper, and this region is what is generally referred to as the “EMC Effect” [49] as it is the only part of the phenomenon that is not fully understood. As far as DIS goes, the sea quark contribution is essentially negligible, given the current understanding of quark PDF’s, and the behavior is dominated by valence quark distributions.
- $0.8 \lesssim x < 1$: The ratio increases rapidly with increasing x . This is dominated by a kinematic effect of the free nucleon cross section vanishing as $x \rightarrow 1$. It is also due to the effect of Fermi motion of the bound nucleons in the nucleus.

1.5.3 Nuclear Dependence of R

It was stated in section 1.5.1 that the per nucleon cross section ratio only reduces to the ratio of structure functions in the case that $R = \sigma_L/\sigma_T$ is A-independent. This particular measurement was the focus of experiments SLAC-E140 [31, 30], NMC [12], and HERMES [2]. A re-analysis of all SLAC data with improved corrections [30] showed that the difference in R between iron and deuterium is consistent with zero ($R_{Fe} - R_d = 0.001 \pm 0.018(\text{stat}) \pm 0.016(\text{sys.})$), and that there are no significantly different spin-0 constituents or higher twist effects in nuclei as compared to free nucleons. Later on, results from HERMES corroborated these findings. The ratio of R_A/R_d were derived from the NMC and SLAC results for $\Delta R = R_A - R_d$ and reported an averaged value of $R_A/R_d = 0.99 \pm 0.03$, providing further agreement on the A-independence of

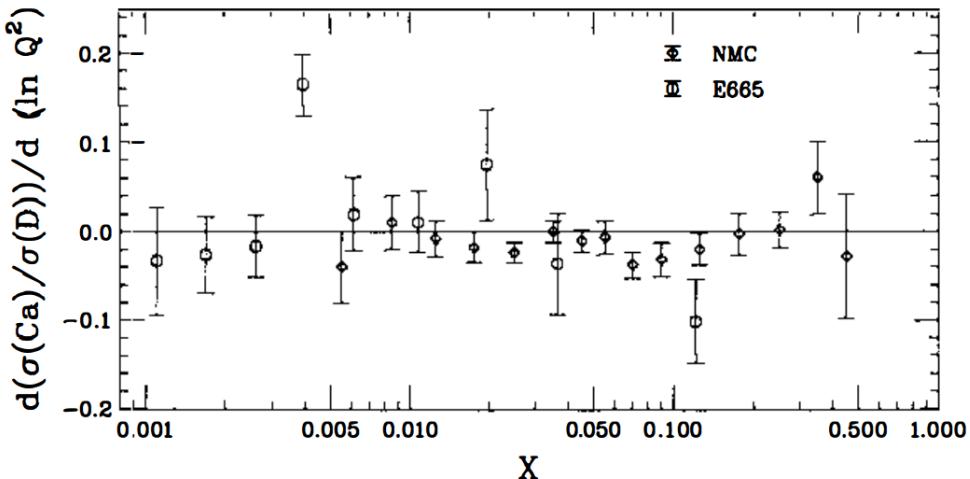
R .

A recent study by Guzey et al. [57] looked to data from SLAC E139, BCDMS, and EMC to examine the impact of R on the extraction of F_2^A/F_2^d and F_1^A/F_1^d from σ^A/σ^d data. They demonstrate the presence of a small but non-zero difference between R for nuclei and free nucleons. While this does appear to effect the extraction of the nuclear enhancement in the ratio of the F_1^A/F_1^d , the extraction of F_2^A/F_2^d appears to remain unaffected. As such, this thesis assumes that it is safe to assume that $\sigma^A/\sigma^d \approx F_2^A/F_2^d$.

1.5.4 Q^2 Dependence of EMC Effect

It is important when investigating structure functions and quark/gluon distributions to pay attention to the Q^2 sensitivity of the measurements. Nuclear structure functions can be directly interpreted as parton distributions only if Q^2 is large enough to precisely define the light-cone direction and reduce higher-twist contributions. The first condition places a restriction on x^2/Q^2 and is usually well satisfied, even at relatively low values of x and Q^2 . However, the value of the higher-twist terms is under little theoretical control, and these terms can have specific nuclear dependences.

Experimentally, the verdict was very clear from the agreement between muon and electron DIS data that any Q^2 dependence must be very small since, at the same value of x , their average Q^2 differs by more than an order of magnitude. The lack of Q^2 dependence can be seen in Figure 1.16 which shows the derivative of DIS cross section with respect to $\ln Q^2$ over a broad range of x , combining NMC and E665 data for calcium [49]. *This absence of any observable dependence of the ratio on Q^2 justifies the usual practice of combining all the Q^2 data at a fixed x -value to display the x -dependence of the ratio.* The only measurement



that seems to refute this conclusion is a measurement of σ^{Sn}/σ^C by NMC which observes a statistically significant, yet small Q^2 dependence at small values of x ($x \lesssim 0.06$) [12]. Due to a large amount of evidence to the contrary along with the fact that SeaQuest is only sensitive to $x \gtrsim 0.08$, this paper will move forward with the understanding that there is no Q^2 dependence in this ratio measurement. As such, this paper will frequently abbreviate such expressions:

$$\frac{F_2^A(x, Q^2)}{F_2^d(x, Q^2)} \rightarrow \frac{F_2^A(x)}{F_2^d(x)} \quad ; \quad \frac{2}{A} \frac{\sigma^A(x, Q^2)}{\sigma^d(x, Q^2)} \rightarrow \frac{2}{A} \frac{\sigma^A(x)}{\sigma^d(x)} \quad (1.54)$$

when in reality it is only the ratio that is Q^2 -independent.

1.5.5 Shadowing

Even though the “shadowing” region of the EMC effect ($x \lesssim 0.06$) is outside of the kinematic acceptance of the SeaQuest experiment ($x \gtrsim 0.08$), it is certainly worth briefly discussing this interpretation. This term “shadowing” has been introduced to explain the reduction of the nuclear cross sections in photoproduction, but in this particular case, has also been used in discussing very low- x modifications to the nuclear cross section in DIS.

The generally accepted explanation for this modification is that the parton distributions in bound nucleons actually *remain unchanged* compared to those in free nucleons. In the rest frame of the *nucleus*, the nuclear effect is attributed to the modification of the interaction of the virtual photon with the atomic nucleus by fluctuations of the virtual photon into quark-antiquark pairs. This pair then interacts with the nucleus via the strong interaction instead of electromagnetically. As a result of this increase in strength, the interaction no longer proceeds incoherently with all the rest of the nucleons, but *preferentially with those at the front surface*. The other nucleons in the “shadow” of the nucleons in front are much less likely to interact, and thereby do not (or much less) contribute to the cross section. As a result, when calculating the per nucleon cross section ratio $\frac{\sigma^A/A}{\sigma^d/2}$ and the deuteron essentially has the same cross section as the heavy nucleus (in this specific regard), then the ratio value will be suppressed.

Quantitatively, this behavior will occur when the $q\bar{q}$ pair fluctuation exists for longer than the mean free path of the pair in a nucleus (distance between nucleons). The fluctuation length behaves as $\Delta d \approx 1/Mx$ where M , the effective mass of the $q\bar{q}$ pair is approximately 1.1 GeV [74], yielding a fluctuation distance of $\Delta d \approx 0.2 \text{ fm}/x$. This becomes larger than the mean free path $L \simeq 2.5 \text{ fm}$ when $x \lesssim 0.08$.

1.5.6 Fermi Motion

One of the more striking features of the $F_2^A(x)/F_2^d(x)$ is the sharp rise in the ratio value at high- x ($0.8 \lesssim x \lesssim 1$). This feature arises as a result of the quantum mechanical principle of Fermi motion, which arises from the confinement of nucleons within the nucleus. If one considers the uncertainty principle inherent in all wave-like systems, $\sigma_x \sigma_p \geq \frac{\hbar}{2}$, there is a direct effect on the momentum distribution of nucleons when one considers the spatial localization of bound nucleons in a nucleus. For example, the nucleons in a tightly bound nucleus (like iron) must have a more precise spatial localization (smaller σ_x) than a nucleon in a more loosely bound nucleus (such as the deuteron). As a result, the magnitude of momentum variations will be greater in the case of iron than in the deuteron (an effect sometimes called “Fermi smearing”). The magnitude of the Fermi momenta of several nuclei can be found in Figure 1.17.

In the extreme case of a free nucleon (which deuterium is chosen to approximate as best as possible), there exists no Fermi motion, as the nucleon is unbound. As a result, the F_2 structure function must vanish as $x \rightarrow 1$ as it becomes very unlikely that a single quark will possess a large fraction of the momentum of the free nucleon without the help of quantum mechanical fluctuations and Fermi smearing. In the case of the deuteron, though, there is still some Fermi motion, as it still consists of two bound nucleons. Bodek and Ritchie were among the first to lay out a model for comparing this effect and how it would differ between different nuclei [22]. The model of how this would affect the structure function ratio is seen on the original EMC plot in Figure 1.12, given no other nuclear effects (such as the EMC Effect).

As is the case with shadowing, this regime is far beyond the kinematic acceptance of SeaQuest to measure. There is still good reason to believe that, whatever the cause of the EMC effect, it likely overlays the effect of Fermi motion, which, depending on the model, can affect the ratio measurement above $x \approx 0.5$. Even with this estimate, there are very few SeaQuest statistics at this x -range.

1.5.7 A-dependence of EMC Effect Slope

Precise measurements of $R_{EMC} = (F_2^A/A)/(F_2^d/2)$ for various nuclei, from 3He to ${}^{197}Au$ have been analyzed, and the intermediate-to-high x region ($0.35 < x < 0.7$) shows a declining behavior that can be fit with a line. The slope of this line has become the focus of study when attempting to interpret the EMC effect, as this slope (dR_{EMC}/dx) exhibits a significant A-dependence. The magnitudes of the slopes appear to increase

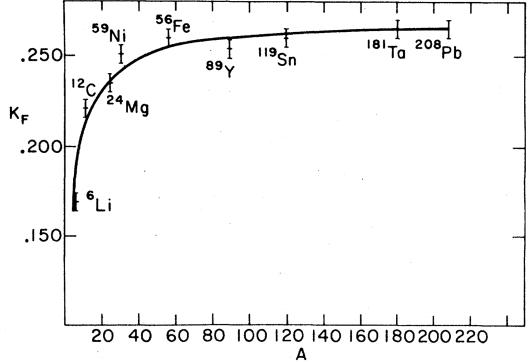


Figure 1.17: The Fermi momenta K_F (in GeV) for various nuclei of atomic weight A [22].

for nuclei with higher atomic weights, such as iron and lead.

Nucleon	dR_{EMC}/dx [95]	dR_{EMC}/dx [50]	dR_{EMC}/dx (comb.)
2H			-0.10±0.05 [53]
3He	-0.070±0.029		-0.070±0.029
4He	-0.199±0.029	-0.191±0.061	-0.191±0.026
9Be	-0.271±0.029	-0.207±0.037	-0.243±0.023
^{12}C	-0.280±0.029	-0.318±0.040	-0.292±0.023
^{27}Al		-0.325±0.034	-0.325±0.034
^{40}Ca		-0.350±0.047	-0.350±0.047
^{56}Fe		-0.388±0.032	-0.388±0.032
^{108}Ag		-0.496±0.051	-0.496±0.051
^{197}Au		-0.409±0.039	-0.409±0.039

Table 1.2: Measured EMC slopes dR_{EMC}/dx for $0.35 \leq x \leq 0.7$ [88].

with increasing A . The table of values of dR_{EMC}/dx for various nuclei can be found in Table 1.2.

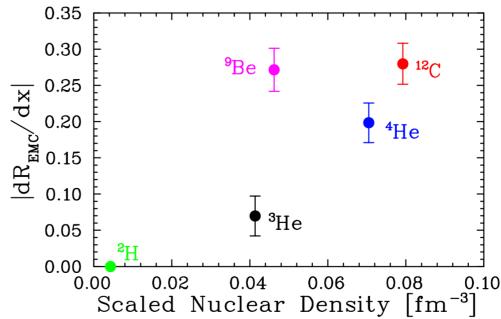


Figure 1.18: The slopes of the EMC ratio for $0.35 < x < 0.7$ vs the average nuclear density scaled by $(A - 1)/A$ [95].

The initial attempt at understanding this was to interpret it as some effect of nuclear density. Since the effect of nuclear density on a given nucleus is due to the presence of the other $(A - 1)$ nucleons, the slope $|dR_{EMC}/dx|$ was plotted against the average nuclear density scaled down by a factor of $(A - 1)/A$ in order to remove the struck nucleons contribution to the average nuclear density. The result for light nuclei was shown by Seely *et al.* to have an interesting result, as is seen in Figure 1.18. The most notable result was a definite overall trend with 9Be representing a

significant outlier. It has been theorized that 9Be is an outlier due to its nuclear structure resembling two α -particles (4He nuclei) with an extra neutron. This would mean that perhaps it is not the average nuclear density, but the *local* nuclear density or density about the struck nucleon. There is still a missing piece needed to represent this density.

1.5.8 $x > 1$ Plateaus and Short Range Correlations

What has recently spurred renewed interest in the EMC effect and its possible true origin is a series of measurements by JLAB of electron DIS cross section ratios where $x > 1$. In this kinematic space, the cross section from a free nucleon vanishes, but persists for the deuteron. It may not at first make sense that measurements can be made in this high x , but it is accessible. Such a behavior is interpreted as

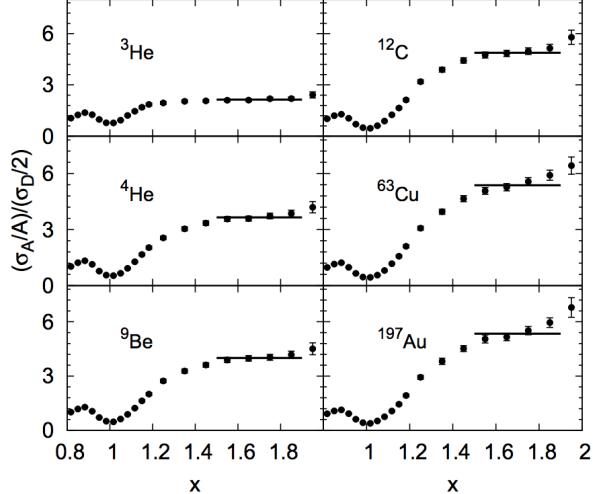


Figure 1.19: Per nucleon cross section ratios vs x at $0.8 < x < 2$ as measured by JLAB-E02019 with a line fit to the “plateau” regions at $1.5 < x < 1.9$ [43].

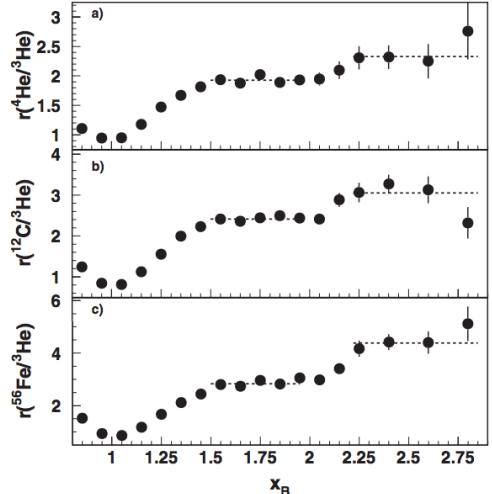


Figure 1.20: Per nucleon cross section ratios vs x at $0.75 < x < 3$ as measured by CLAS with a line fit to the “plateau” regions at $2.25 < x < 2.8$ [37].

the manifestation of short-range nucleon-nucleon correlations (three-nucleon correlations for $x > 2$). These short-range correlations happen when two (or three) nucleons happen to get very close to each other and experience a large repulsive Coulomb force, resulting in the nucleons moving apart from each other with high *individual* momentum, but low *combined* momentum.

The JLAB-E02019 experiment [43] measured cross section ratios for ^3He , ^4He , Be , C , Cu and Au to deuterium. One can see in the results shown in Figure 1.19 that the cross section ratios of $^4\text{He}/D$, C/D and Cu/D rise with x until the value plateaus at $x \approx 1.4 - 1.5$. The key observation with these plateaus is that the ratio value at this plateau increases with A . Such behavior was already observed with less accuracy by an experiment at SLAC [44], and a similar pattern can be seen in the CLAS experiment which compared nuclear cross sections to tritium (^3He) [37]. Here the measurements extend up to $x = 3$ and an indication of a second plateau is observed for $2.25 < x < 2.8$ (Fig. 1.20).

The variable $a_2(A/d)$ is introduced here as the ratio (in the plateau range) of the per-nucleon inclusive (e, e') cross sections for two nuclei. This value can be interpreted as the *ratio of probabilities to find high-momentum nucleons in those two nuclei* [108]. High-momentum nucleons were recently shown to arise primarily from nucleon-nucleon short range correlations, which is when two nucleons’ wave functions overlap briefly, causing them to have high relative momentum but low combined momentum (relative to the ~ 400 MeV Fermi motion) [7]. Thus, this $a_s(A/d)$ ratio is called the “Short Range Correlation (SRC) Scale Factor”³ The current known values for the SRC scale factors are shown in Table 1.3. The variables $a_2(A/{}^3\text{He})$

³More specifically, it is called the 2N-SRC scale factor, as there is a three-nucleon scale factor for cross section ratios with

Nucleon	$a_2(A/{}^3He)$ [38]	$a_2(A/d)$ [108]	$a_2(A/d)$ [43]
2H	0.508 ± 0.025	1	1
3He	1	1.97 ± 0.10	2.13 ± 0.04
4He	1.93 ± 0.14	3.80 ± 0.34	3.60 ± 0.10
9Be			3.91 ± 0.12
${}^{12}C$	2.41 ± 0.17	4.75 ± 0.41	4.75 ± 0.16
${}^{56}Fe$	2.83 ± 0.18	5.58 ± 0.45	
${}^{63}Cu$			5.21 ± 0.20
${}^{197}Au$			5.16 ± 0.22

Table 1.3: Measured SRC scale factors, as measured by CLAS [38] and JLAB E02-019 [108, 43].

and $a_3(A/{}^3He)$ is also defined for first and second plateau ranges, respectively, but are not emphasized here.

It was first posited by Higinbotham *et al.* [64] that the magnitude of the EMC effect in a nucleus A is related to the measured SRC scale factor of the same nucleus if both are taken relative to deuterium. In Figure 1.21, the EMC slopes from Table 1.2 are plotted against the SRC scale factors from Table 1.3 to exhibit a strong correlation indicated by the straight line. It would seem unlikely that this is coincidental, and it can therefore be asserted with some confidence that the EMC effect seen in the valence quarks is largely due to these 2N-SRC [92].

There is a future study planned at JLAB to test this 2N-SRC - EMC effect connection, experiment E1211107, which has been approved to run as part of their 12 GeV program. The experiment will study the structure function of high-momentum nucleons in deuterium by performing electron DIS from nucleons in deuterium and identifying nucleons that are knocked out *backwards* towards the direction that the beam came from [62]. Its primary goal will be to simultaneously measure the EMC effect and SRC scaling factors in a wide variety of light and heavy nuclei.

respect to tritium.

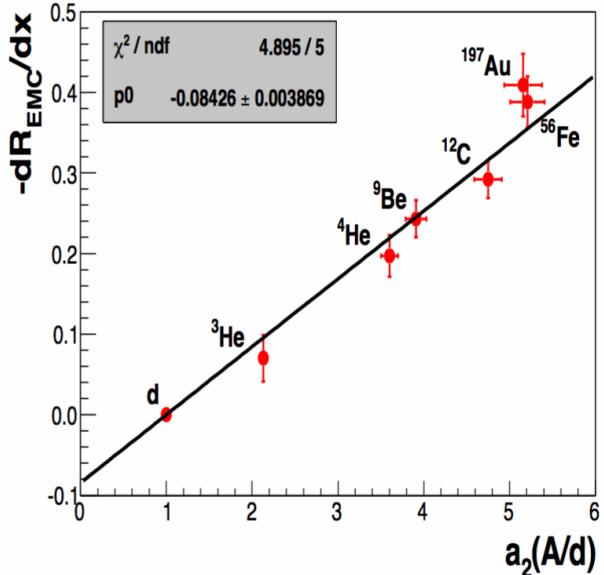


Figure 1.21: The EMC slopes plotted against the 2N-SRC scale factors [61]. The fit is constrained to go through deuterium at $(1,0)$.

1.5.9 Models For the EMC Effect

The A-dependent modification of the nucleon has been such an intriguing phenomenon that so many models have arisen attempting to reproduce the data, leading those who study it to give it the moniker “*Every Model is Cool*” [80]. Two models are discussed here that have the potential to be evaluated by the SeaQuest measurements.

Pion Cloud Model

A proton is not simply three quarks, and a nucleus is not simply a convolution of nucleons. The nucleons themselves are bound by the exchange of mesons at intermediate and long range. These mesons carry a certain fraction of the momentum of the nucleons, and a modification of the meson exchange field in a nuclear medium was predicted by Coester to lead to an enhancement in the DIS and DY EMC ratio [19]. Later in Figure 1.24, E-772’s DY nuclear cross section ratio results show strong disagreement with this virtual meson exchange model. Jung and Miller [68] have reevaluated the predictions of Berger and Coester [19], which leads to a more flat enhancement to the cross section ratios. SeaQuest will be able to provide data at higher- x where models may be able to be more easily distinguished from each other.

Rescaling

In this classification of models, the EMC effect is to be explained by a change in the scale of Q^2 used to probe the nucleus as compared to a free nucleon. Alternatively, the momentum fraction x of the probed parton can be scaled commensurately with the size of the nucleus.

In the case of Q^2 rescaling, it is explained that [92] the strength of the strong force between quarks is ascertained not just by the $1/Q$ at which they are probed, but also by the volume in which they are confined, defined by the nucleon radius r_A [27, 83]. It is therefore relevant to parametrize the strong coupling constant strength not by just Q^2 , but by $(Q \cdot r_A)^2$. If the confinement scale is changed for a nucleon inside of a nucleus for any number of reasons, then the parton distributions for nuclei A and B will be related by

$$q_A(x, Q^2) = q_B(x, \xi \cdot Q^2) \quad (1.55)$$

where ξ is the *rescaling* parameter determined by the nucleus radii (confinement scales) r_A and r_B . When considering “dynamical rescaling” models, ξ is defined as

$$\xi = \left(\frac{r_A}{r_B} \right)^{\frac{2\alpha_s(\mu_A^2)}{\alpha_s(Q^2)}} \quad (1.56)$$

where μ_A is a low-momentum cutoff for radiating gluons.

As a result of this Q^2 “dynamical rescaling,” for the case of $r_A > r_B$ then the structure function ratio $F_2^A/F_2^B > 1$ for low x and $F_2^A/F_2^B < 1$ for large x . A benefit of this model is that it only requires a change in confinement scale of the nucleon, and the mechanism by which this occurs is open ended. Some possible causes could be the “swelling” of nucleons, multi-quark “bags,” short range correlations, color flow through the nucleus [92], or one of any other model that possibly changes the nucleon confinement scale within a nucleus.

In the case of x -rescaling, the EMC effect is explained by nuclear binding and Fermi motion corrections [46, 98]. In this case, the universal x -dependent behavior is explained by x having to be scaled up to a $x_A > x$. In the case of DIS, a nucleon i with momentum p_i , the Bjorken- x variable is replaced by:

$$x = \frac{Q^2}{2p_i q} \rightarrow x_A = \frac{Q^2}{2(M + E_i)\nu - 2\vec{p}_i \cdot \vec{q}} \quad (1.57)$$

where \vec{q} is the momentum of the virtual photon and E_i is the removal energy of the nucleon, which can vary between -9 and -27 MeV, depending on the nucleus and whether or not the nucleon in question is a proton or neutron (Fig. 1.22).

As a result of the x -rescaling, for the same scattered lepton in DIS, the momentum fraction x_A is probed with regards to the structure function. Since the structure function falls off quickly with increasing x , the nuclear structure function ratio depletion is fully explained.

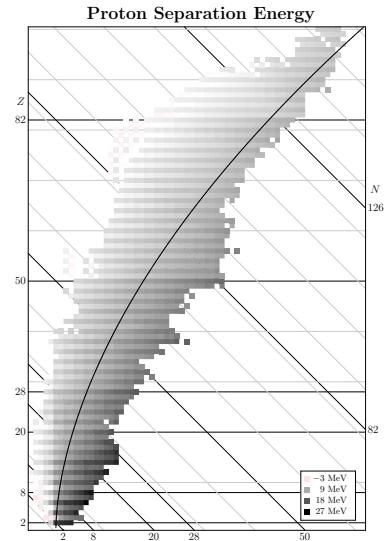


Figure 1.22: The amount of energy required to remove a proton from a nucleus as a function of Z and N [32]. This value, E_i , is used in the x -rescaling model.

1.6 The Role of Drell-Yan in Studying the Nuclear Quark Sea

There have been many models attempting to explain the A -dependence of such phenomena as the EMC effect. Each of these models provides a prediction regarding the individual sea/valence quark/antiquark distribution functions. In order to test these predictions, one can take stock of the many nuclear probes and assess the capabilities of what each can contribute.

In general, since deep inelastic scattering of electrons and muons is an electromagnetic interaction, it is not innately sensitive to the whether or not the struck quark was a valence quark or a sea quark and

can only measure the charge-weighted (E&M+weak) sum of parton distributions. Through neutrino and antineutrino DIS it is possible to discern between valence and sea distributions. It is the case that most models make predictions for valence quark distributions at mid-to-high x that are not in tension with each other. As such, measuring the valence distributions would not help much in discriminating between different models. Models do, however, differ in their predictions of the sea quark distributions at higher x ($x \gtrsim 0.2$) - but at high momentum fraction, the distributions of sea quarks diminish very quickly (Fig. 1.7). As such, antineutrino DIS has been unable to deliver precise measurements of sea quark distributions at high- x yet.

The Drell-Yan process, $pA \rightarrow l^+l^-$, uses the s -channel version of the same DIS process, but due to the nature of requiring interacting $q\bar{q}$ pairs, it provides a tool that allows one to cleanly and specifically investigate the sea quark distributions. With a forward spectrometer (such as is the case for E-772 and SeaQuest), where one accepts the $x_F > 0.3$ one is able to study the Drell-Yan interaction in the specific case of a valence quark from a free proton (the beam) annihilating with an antiquark from the quark sea in the target. As a result, a study of this process for various nuclei can provide a direct, precise measurement of the antiquark sea distributions that can assist in evaluating the plethora of EMC effect models out there, along with refining constraints on global fits of PDFs.

In this paper, the A-dependence of the Drell-Yan cross sections is studied, which directly measures the A-dependence of the antiquark distributions [19]. We define this ratio, R^{DY} as the ratio of Drell-Yan differential cross-sections, using Eq. 1.13 to get

$$R^{DY} \equiv \frac{d^2\sigma_A}{dx_1 dx_2} \Bigg/ \frac{d^2\sigma_d}{dx_1 dx_2} \quad (1.58)$$

$$\approx \frac{\sum_f e_f^2 \cdot q_f^{\text{beam}}(x_1) \bar{q}_f^A(x_2)}{\sum_f e_f^2 \cdot q_f^{\text{beam}}(x_1) \bar{q}_f^d(x_2)}. \quad (1.59)$$

Assuming that the \bar{u} distribution is roughly similar to the \bar{d} distribution in the target, this term is dominated by up-quark contributions by a factor of eight: $e_u^2 = 4e_d^2$, and there are twice as many up-quarks from the beam proton as down quarks [49]. As a result, we can approximate

$$R^{DY} \approx \frac{e_u^2 \cdot q_u^{\text{beam}}(x_1) \bar{q}_u^A(x_2)}{e_u^2 \cdot q_u^{\text{beam}}(x_1) \bar{q}_u^d(x_2)} \quad (1.60)$$

$$= \frac{\bar{q}_u^A(x_2)}{\bar{q}_u^d(x_2)} \quad (1.61)$$

meaning that by measuring R^{DY} for pA to pd collisions with a forward spectrometer, a good measure of the nuclear modification of the anti up quark sea distributions can be achieved.

1.6.1 Previous Measurements

Even though Drell-Yan events for $x_F > 0.3$ do directly provide a measure the sea quark distribution, there is a limited amount of data available. FNAL-E-772 [6] was a 800 GeV fixed target experiment which provided a measure of σ^A/σ^d for C, Ca, Fe , and W as a function of kinematic variable x_2 . Soon to follow was FNAL-E-866 [106], also a 800 GeV proton beam fixed target experiment, which contributed a somewhat similar measurement of the DY cross section ratiof for W and Fe with respect to Be . The results which can be seen in Figures 1.24 and 1.23 are able to show only a limited x -range of measurements of $0.05 \leq x_2 \leq 0.25$ for E-772 and $0.02 \leq x_2 \leq 0.1$ (again, the sea quark distribution drops off sharply with increasing x). This data does show that at low- x_2 ($x_2 \lesssim 0.07$), there is indication of shadowing occurring. This makes sense for DY as well as DIS, even though there is no need the $\gamma \rightarrow q\bar{q}$ fluctuations, as the number of sea antiquarks increases greatly at this low range of x and would thereby provide the same nucleon “shadow” as was described in Section 1.5.5.

For values above this shadowing region, the ratios appear to be consistent with unity. The range of x covered line up with the *transition region* (sometimes called “anti-shadowing”) between the shadowing and EMC effect region. In some DIS measurements, this region can be characterized by having a value in excess of unity. A measurement of unity in the Drell-Yan cross section ratio would seem to imply that whatever this effect might be, it is likely not caused by any nuclear modifications to the sea quark distribution.

Regarding the other models, there was a great expectation that the Drell-Yan cross section would sizeably increase due to virtual pions which mediate the internuclear force. It was surprising that E-866 and E-772 observed no modification of the sea quark distributions in large nuclei, which brings into question the long-standing model of mesons being exchanged between nucleons as the mechanism by which nuclei are held together [25].

To be able to truly draw any conclusions regarding various models, greater precision is needed, along with data at a higher- x range where models are more greatly differentiated. The SeaQuest experiment will attempt to do just this.

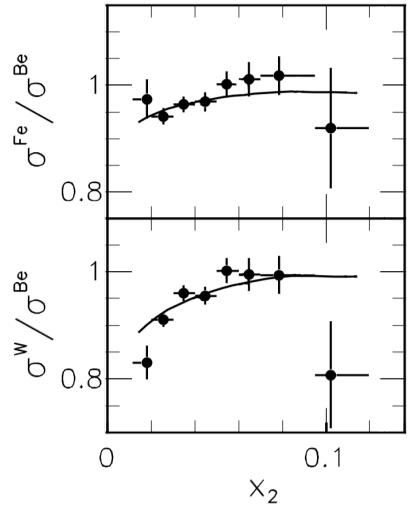


Figure 1.23: E-866 Drell-Yan σ^A/σ^{Be} for iron and tungsten as a function of x_2 [106].

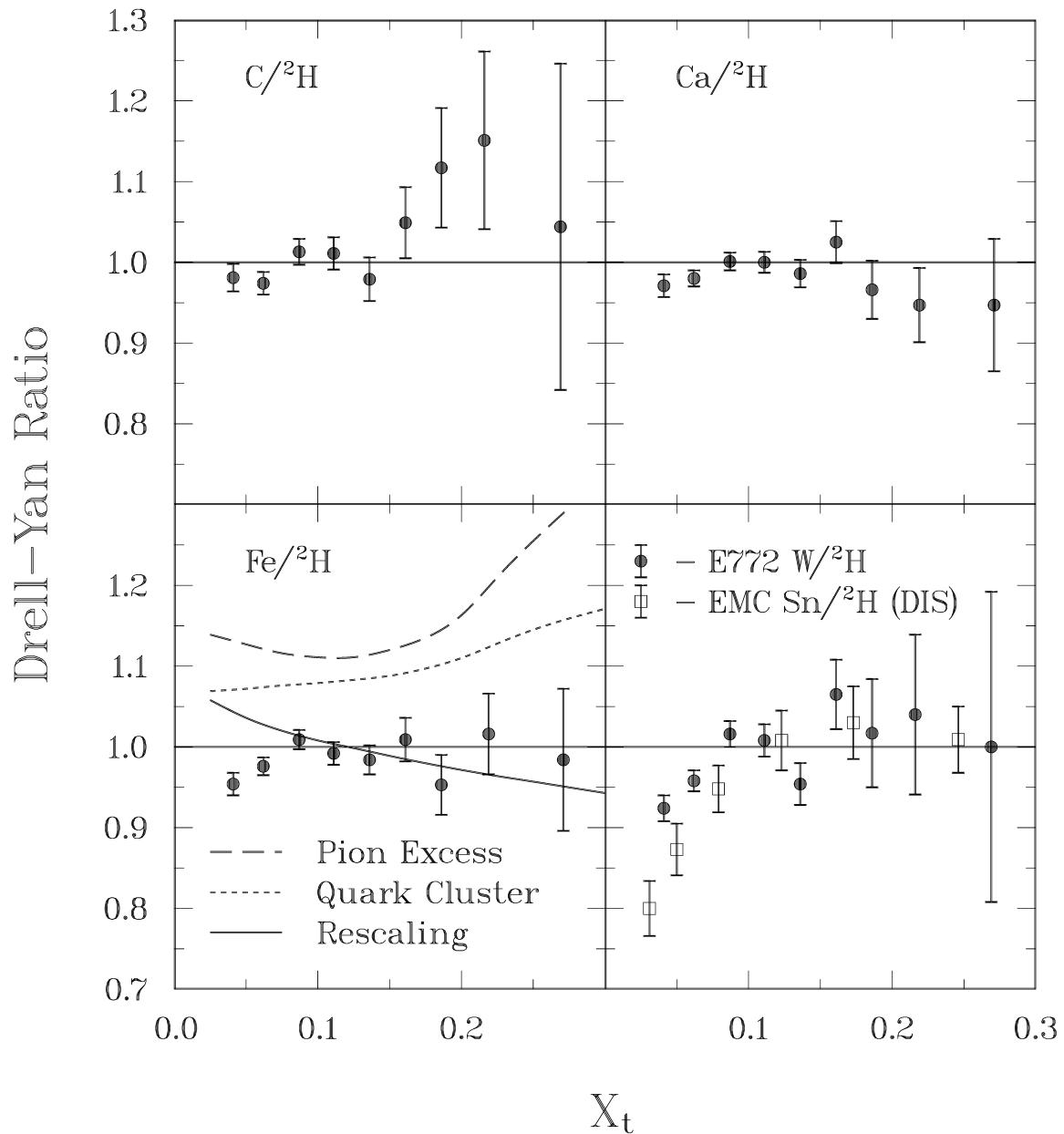


Figure 1.24: E-772 Drell-Yan $\frac{A}{{}^2H}$ ratios for several nuclear targets as a function of x_2 (x_t) [6].

1.6.2 Nuclear Dependence of Drell-Yan at SeaQuest

The SeaQuest Experiment was proposed soon after the conclusion of the E-866 experiment. The main motivation was to explore the $\bar{d}(x)/\bar{u}(x)$ value observed by E-866 to more precisely ascertain its behavior as a function of x . With the ratio value being as large as ~ 1.7 at its peak and with it having a decreasing value at higher x (though with low statistical certainty), no models are able to accommodate such behavior, and so more data is needed.

Sea quark distributions at higher- x the SeaQuest collaboration at Fermilab is designed to detect proton-induced Drell-Yan cross sections on several targets: liquid hydrogen (1H), liquid deuterium (2H), carbon (^{12}C), iron (^{56}Fe), and tungsten (^{184}W). The beam source for the experiment is the Fermilab Main Injector which provides a 120 GeV proton source, yielding a \sqrt{s} of 15.1 GeV. This relatively lower beam energy as compared to 800 GeV ($\sqrt{s} = 38.8$ GeV) for E-772/E-886 provides an opportunity to study parton distributions at larger x . This is due to the fact that, as seen in Eq. 1.13, for a fixed value of x , the Drell-Yan cross section is inversely proportional to the square of the center-of-mass energy, s . Another benefit of utilizing a lower-energy beam is that the primary background, J/Ψ production, which limited the instantaneous luminosity in E-866, scales with s , and is thereby reduced. This allows SeaQuest the usage of a more intense proton beam, which is very useful when it comes to measuring a rare nuclear process like Drell-Yan.

Even though the emphasis of SeaQuest is on the \bar{d}/\bar{u} measurement, the periodic data taking on deuterium and the heavy nuclear targets will allow for a measurement of $R^{DY} \approx \bar{u}^A(x)/\bar{u}^D(x)$ for a large range of A and x , providing a more thorough glimpse into the modification of sea quark PDFs in the presence of a nuclear medium. The normalized (x_1, x_2) and (x_F, M_{γ^*}) coverages for SeaQuest are shown in Figures 1.25 and *mass vs xF here!*, as generated by Geant Monte Carlo simulation.

At the time of this writing, SeaQuest is the only fixed target proton-nucleon Drell-Yan experiment currently taking data in the world.

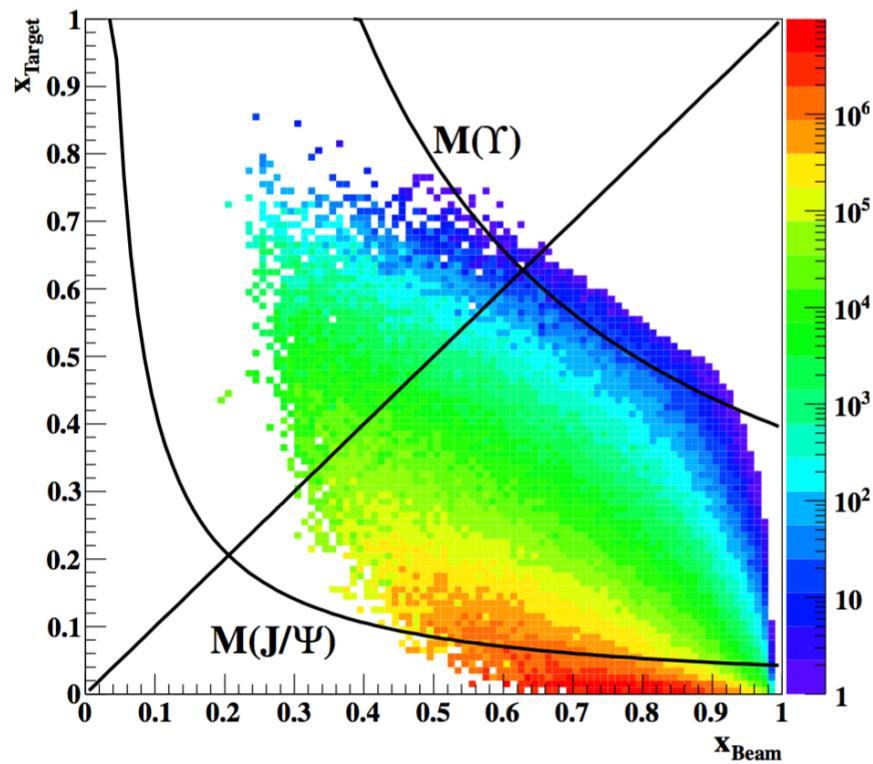


Figure 1.25: The x_1 and x_2 acceptance of the SeaQuest spectrometer, with curves showing lines of constant mass where there are mass cuts applied at $M_{J/\Psi} < M_{\gamma^*} < M_\Upsilon$ [94].

Chapter 2

Apparatus

SeaQuest is the operational name of Fermilab Experiment 906 (*E-906*) performed at its Neutrino-Muon (*NM*) experimental area. The experiment was designed to take *high-intensity beam* at relatively *low center-of-mass energy*, provide *good mass resolution*, and allow for *accurate target-to-target systematic normalization*. The apparatus consists of a moving target table, two dipole magnets, 8 hodoscope planes, 24 drift chamber planes, and 4 proportional tube planes. Upstream of the target table (towards the beam source), there is also a Čerenkov counter for beam intensity monitoring and there are several segmented wire ionization chambers (SWIC's) for beam profiling.

2.1 Apparatus Overview

SeaQuest is a fixed-target experiment. In this style of an experiment, a stationary target is placed in the path of an accelerated beam of particles, as opposed to *collider* experiments where two accelerated beams are directed against each other, in opposite directions. The proton beam interacts with the target material and produces a variety of daughter particles. These daughter particles are tracked through a forward spectrometer and selectively filtered dependent on the purpose of the study.

The tuned and monitored 120 GeV proton beam is sent from the Fermilab Main Injector (MI) where the beam protons strike one of 7 targets. The high-momentum charged particles that are produced are focused into the acceptance of the experiment's detectors with the solid iron dipole magnet, FMAG, or NM3S. This solid focusing magnet also sweeps away low-momentum particles and acts as a beam dump / absorber.

The SeaQuest Spectrometer (Fig. 2.1) consists of the aforementioned focusing magnet, several tracking chambers that record the positions of charged particles through the length of the spectrometer, and an analyzer magnet to bend the particles between tracking stations. The spectrometer measures particle momenta by recording the bend of each charged particle as it passed through the analyzer magnet, where the magnetic field is known. This is performed by reconstructing the trajectory of a particle in one-half of the spectrometer (before the analyzer magnet) and then similarly reconstructing the trajectory of particles

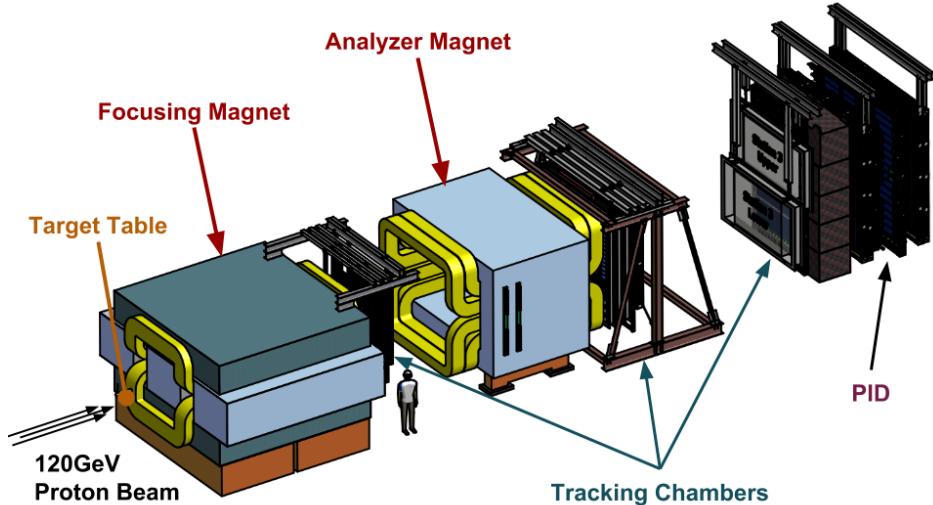


Figure 2.1: Perspective view of the SeaQuest spectrometer apparatus.

in the other half. If two trajectories can be matched up, then the particle momentum can be extracted by taking the ratio of the magnet's p_T -kick to the change in the track's direction.

The spectrometer's geometric design and event triggering selection is optimized to detect oppositely-charged pairs of muons while minimizing the sensitivity to various sources of unwanted backgrounds. Positive identification of muons is achieved by requiring signals in the hodoscopes for known muon “*roads*” along with requiring signals in the proportional tubes located at the farthest end of the experiment, past an iron wall. Electrons and any hadrons are stopped by the solid iron focusing magnet and iron wall farther down while muons will pass through them unencumbered.

The coordinate system is defined as the following: the z -axis points along the beam direction, the y -axis points upwards vertically, and the x -axis lies along the horizontal direction in such a way that a right-handed coordinate system is formed. The terms *upstream* and *downstream* are often used when referring directions or regions in the experimental hall. *Upstream* often refers to the direction towards the beam source ($-z$ direction) while *downstream* refers to the $+z$ direction of the beam trajectory. The origin of the coordinate system was chosen to be the point where the proton beam meets the *upstream*-facing surface of FMAG, the solid focusing magnet.

2.2 Main Injector Proton Beam

The Fermilab Main Injector (MI) receives protons that have been accelerated by the Radio Frequency Quadrupoles (RFQ), the Linear Accelerator (LINAC), and the Booster, and it continues to accelerate them

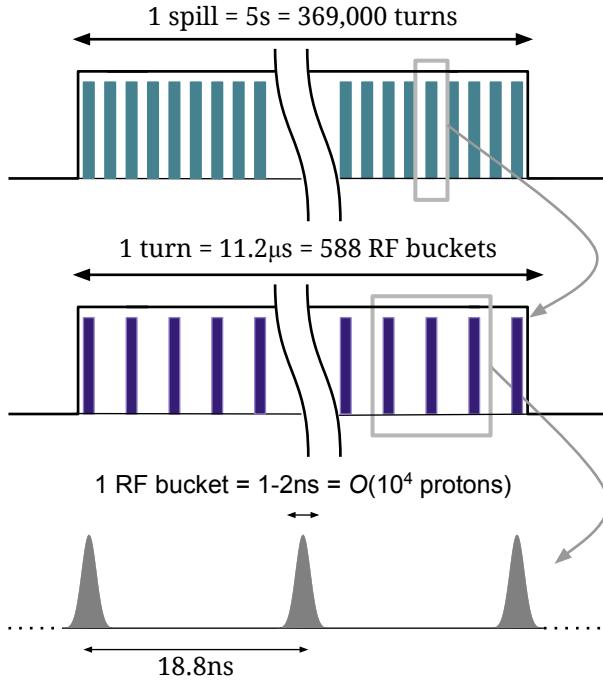


Figure 2.2: Spill structure of the beam delivered to SeaQuest.

from 8 GeV up to the nominal energy of 120 GeV. Along the way, radio-frequency cavity (RFC) accelerators in the LINAC and the MI ‘‘bunches’’ up the protons such that the beam has its characteristic 53.1 MHz structure. After the period of acceleration, the protons are then ‘‘scraped off’’ slowly with each passing *turn* of the collected proton beam and sent down the Neutrino-Muon (NM) beam delivery line for approximately five seconds of every minute, called a ‘‘slow spill’’ or just ‘‘spill’’. Beam is extracted using a resonant process, and the extracted beam retains the 53.1 MHz structure of the Main Injector RFC. Each bunch, or ‘‘bucket’’, of protons is less than 2 ns long and the time between bunches is approximately 18.8 ns. The spill structure of the beam is depicted in greater detail in Figure 2.2.

The beam sent to SeaQuest is not uniform in time throughout the spill. There are beam bunches in the MI that are intentionally left empty so that the abort kickers can ramp to a full field during a gap in the beam. There are also bunches left empty to allow the injection kickers to inject 8 GeV protons from the Booster without disturbing bunches of protons already in the Main Injector. Typically, 498 of the 588 ‘‘RF buckets’’ in the Main Injector contain protons during the SeaQuest slow spill cycle. It is the case, however, for SeaQuest, that the intensity of the bunches corresponding to these 498 full buckets varies greatly throughout the slow spill. On *average*, each bucket will have $\mathcal{O}(10^4)$ protons and the spill has an intensity of approximately 2×10^{12} protons per second and therefore about 1×10^{13} protons delivered per

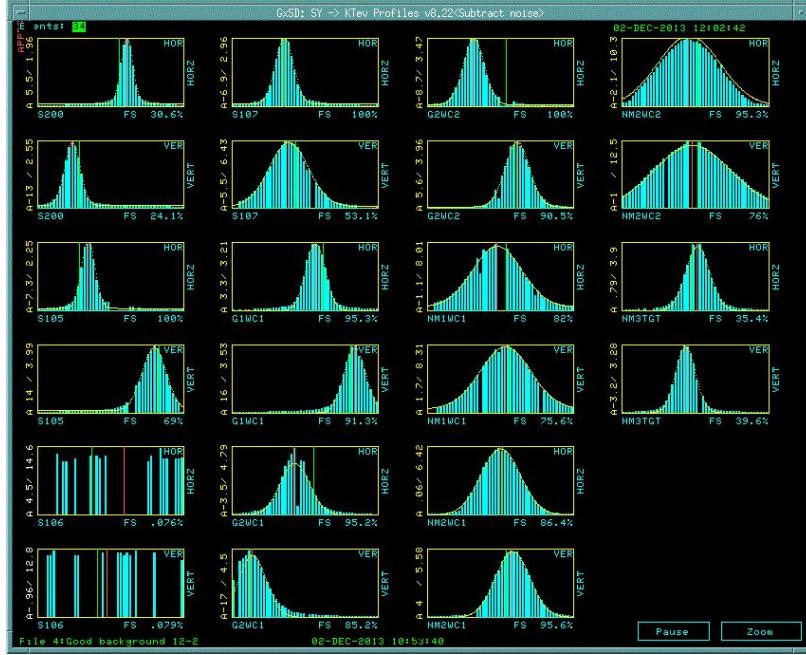


Figure 2.3: Beam profile detailed by SWIC detectors along the NM beam line.

spill.

Several guiding and focusing magnets bend and deliver beam to the NM beamline which serves both the test beam facility and SeaQuest at NM4. The beam is focused to a width of $250\ \mu\text{m}$. The profile, position, and intensity are measured along the NM beamline by several detectors. The intensity of the beam is monitored by an ion chamber (IC) and a secondary emission monitor (SEM) in the NM3 sector of the beamline. The beam profile and position are monitored by SWICs and beam position monitors (BPMs), respectively. The Accelerator Control Network (ACNET) display of the SWIC readout can be seen in Figure 2.3. The closest BPMs and SWICs to the spectrometer were located in NM2 enclosure. The beam profile does not maintain its $250\ \mu\text{m}$ shape and spreads slightly as it moves towards the spectrometer. The final beam profile is measured by inspecting the upstream-facing side of the solid targets, and it was found to be approximately 6mm wide by 1mm high.

2.3 Beam Intensity Monitor

SeaQuest's trigger system (described in detail later) mostly fires on fake dimuons caused by two low p_T muons from unrelated pion decays. The hits in downstream hodoscopes from the pions combined with hits in the upstream hodoscope from two other unrelated particles frequently add up to a false dimuon signal. Since this type of fake trigger involves four unrelated particles, the probability that a trigger will occur

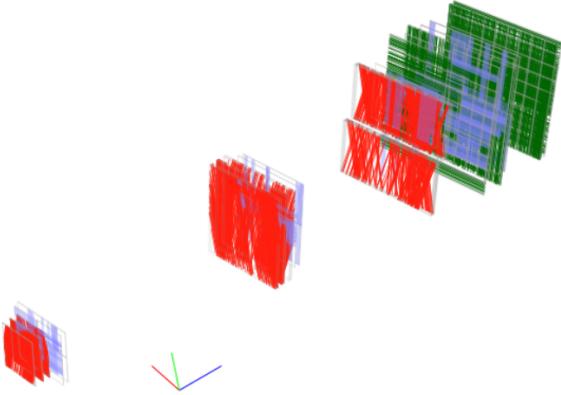


Figure 2.4: A single high-intensity event with majority of all detector elements firing off. White space within the rectangles indicates inactive elements whereas red, blue, and green represent elements which have fired during that event. Track reconstruction in these cases is impossible.

increases with I^4 , where I is the intensity of the beam bucket; the number of protons in the triggered beam bunch.

The SeaQuest data acquisition system (also described later in detail) can read out approximately 3000 events per second without significant dead time. During the commissioning run of SeaQuest, the trigger rate was very high and the trigger dead time was close to 100%. These triggers were taken at such high beam intensities that the occupancy of all SeaQuest detector elements was more than 50%, making pattern recognition essentially impossible (see Figure. 2.4). The Beam Intensity Monitor (BIM) was designed to solve this problem.

The SeaQuest Beam Intensity Monitor (BIM) senses when the beam intensity is above a programmable threshold. If an RF bucket with an intensity above this threshold is detected, the BIM sends a signal to inhibit certain triggers until the intensity once more falls below the threshold. The inhibit threshold is tuned frequently as trigger and beam conditions change, but it is typically set at $\sim 95,000$ protons per RF bunch. For reference, a full RF bucket at an intensity of 2×10^{12} protons per spill is $\approx 10,000$ protons.

The beam intensity is measured using an atmospheric pressure gas Čerenkov counter. A gas mixture of 80% Argon and 20% CO₂ is used as the Čerenkov radiator. The counter and readout electronics were designed to have $O(ns)$ time resolution, and a linear response over a large dynamic range. A diagram of the counter is shown in Figure 2.5. A 45 degree aluminized Kapton mirror directs light to a single photomultiplier tube. A *baffle* of black construction paper held parallel to the mirror ensures that the proton path length through the light-radiating gas with respect to the mirror is independent of beam position. A two-inch diameter 8-stage photomultiplier tube (PMT) is positioned close to the mirror so that all Čerenkov light

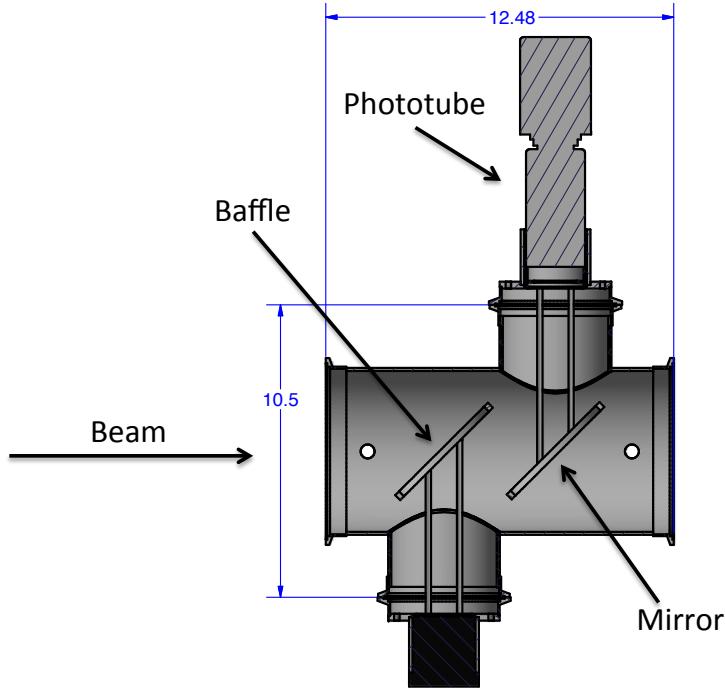


Figure 2.5: The Beam Intensity Monitor (BIM) Čerenkov counter. Measurements are in inches.

created between the baffle and the mirror falls directly on the aperture of the PMT. It was observed during the commissioning run that after exposure to $\approx 3 \times 10^{17}$ protons (~ 3 weeks of uninterrupted usage), the mirror reflectivity is significantly reduced in the beam spot, and the mirror then needs to be replaced.

The signal from the BIM is integrated and digitized using a custom charge (Q) integrator and encoder (QIE) integrated circuit board, which comes from a family of circuits used first by the KTeV experiment at Fermilab[110]. The chip is clocked with the Main Injector RF clock and provides an ADC (analog-digital conversion) every 18.8 ns clock cycle. The light incident on the photomultiplier tube is attenuated using neutral density filters (NDF's) so that the QIE least count corresponds to ~ 30 protons per beam bunch.

In addition to inhibiting triggers during high-intensity periods of beam, the BIM readout module also provides critical information used to calculate the number of protons incident on the SeaQuest targets while the experiment is ready and able to trigger. This value is needed to normalize SeaQuest cross section measurements. The BIM readout module provides the following:

- Sum of all ADC signals for the entire spill (QIESum).
- Sum while inhibit is asserted at trigger logic.
- Sum during trigger dead time.

- A snapshot of beam intensity 16 buckets before and after the triggered RF bucket

These are used to calculate a ratio of protons that were ‘live’ (the experiment can trigger) via the following:

$$\%live = \frac{QIESum - (inhibit\ sum + dead\ time\ sum)}{QIESum} \quad (2.1)$$

$$liveProton = totalProton \cdot \%live \quad (2.2)$$

where “totalProton” is the intensity value recorded from the SEM detector located just upstream of the BIM Čerenkov counter. The SEM itself is calibrated by foil activation. The snapshot of the triggered RF bucket intensity along with the 32 surrounding RF bucket intensities is used for studies and corrections of the rate-dependent effects on detector efficiencies and reconstructed measurements.

2.4 The SeaQuest Targets

A wide range of atomic weights (from 2 to 184) is required to do an A-dependence study of the Drell-Yan process. At SeaQuest, the targets used are $^1H(\ell)$, $^2H(\ell)$, C , Fe , and W . In addition to the two liquid targets and the three solid targets, two positions on the target table were used for measuring background signal rates: an empty flask, identical to the flasks used for the 1H and 2H targets, and a single empty solid target holder. Colloquially speaking: the 1H target is interchangeably referred to as the liquid hydrogen target, ℓH_2 , $LH2$, or H_2 ; 2H is likewise referred to as the liquid deuterium target, ℓD_2 , $LD2$, or D_2 ; the empty flask is referred to as the “Empty” target; the empty solid target holder is referred to as the “None” target.

These are all mounted on a laterally-moving, remotely positionable table (in the $\pm x$ direction), able to move over a range of 91.4 cm. The table’s center is located at $(0, 0, -1.25)$ meters, directly in front of the upstream face of FMAG, the solid iron focusing magnet. Because of the ~ 5.0 cm diameter of the targets and the 6x1mm dimensions of the beam, the targeting efficiency was 100%. The details of the target materials are summarized in Table 2.1, and the layout of the target table can be seen in 2.6.

The H_2 gas used is “Ultra High Purity 5.0 Grade” or 99.999% pure. The deuterium has come from two different sources. The first of these is a Fermilab-provided supply of gas left over from previous bubble chamber experiments. This gas was known to have a small hydrogen contamination and was measured by mass spectroscopy to have a composition of 85.2% D_2 , 12.7% HD , 1.2% 4He , and 0.8% H_2 by mole. As the analysis of experimental data commenced, handling the ramifications of the D_2 impurity came under focus.

Position	Material	Density [g/cm ³]	Thickness [cm]	Interaction Length	Spills/ Cycle (%spills)
1	H_2	0.07065	50.8	0.06902	10 (43%)
2	Empty	NA	NA	0.0016	2 (9%)
3	D_2	0.1617	50.8	0.1144	5 (22%)
4	None	NA	NA	0.0	2 (9%)
5	Iron	7.874	1.905	0.1135	1 (4%)
6	Carbon	1.802	3.322	0.0697	2 (9%)
7	Tungsten	19.30	0.953	0.0958	1 (4%)

Table 2.1: Characteristics of the seven SeaQuest target positions. The “Spills/Cycle” is only a typical configuration and can vary according to needs and running configurations. The non-zero interaction length of the empty flask is due to the 51 μm -thick stainless steel end-caps of the flask and the 140 μm -thick titanium windows of the vacuum vessel that contains it.

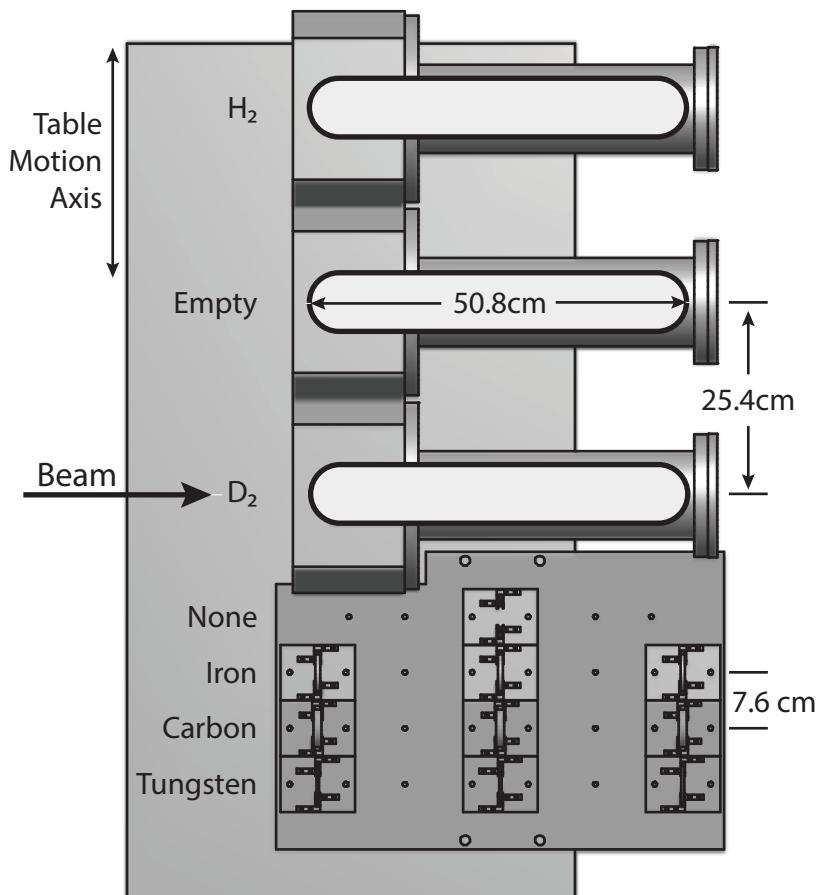


Figure 2.6: The layout of the target table and its seven target positions, as seen from above.

Unexpected bottle-to-bottle variation in contamination became evident, and the sample-taking methodology itself for spectroscopy became suspect of introducing contamination. In order to no simplify analysis and reduce the substantial complexity and cost of further gas analysis, SeaQuest switched to commercially available “Research Grade” D₂, which is better than 99.6% pure with virtually all HD to balance. The data analyzed in this paper deals with the impure D₂ target material before this switch. Further information on the D₂ composition and how it is handled in the context of analysis will be covered in Chapter 5.

Each of the three solid target positions is divided into three disks of 1/3 the total thickness provided in Table 2.1. These disks are spaced 25.4 cm apart to approximate the distribution of the liquid target, thereby minimizing target-dependent variation in spectrometer acceptance. The one exception to this is that during the Run II period the iron disks were more closely spaced (17.1 cm). The decision to place these iron disks closer together than the rest during Run II is still unclear.

The target table is able to move between two different targets in about 30 seconds. This allows a change between targets in the 55 seconds between successive spills. With this frequent target interchange, the systematic uncertainties associated with drifts in beam characteristics, monitor gains, and detector efficiencies are reduced to a minimum when investigating A-dependent ratios. How much beam time each target received is determined by interaction lengths of the targets along with the amount of statistics desired for certain targets. As the flagship measurement of SeaQuest is the \bar{d}/\bar{u} asymmetry, more emphasis was placed on the hydrogen and deuterium than the nuclear targets. The spills per cycle and beam time allocation can be found in Table 2.1.

2.5 Focusing and Analyzing Magnets

Two large dipole magnets are used in the experiment to be select forward going ($x_F > 0$) dimuons, reject low-momentum particles, and analyze their kinematic characteristics. The most upstream magnet, denoted “FMAG”, is a solid iron A-frame magnet with an aperture of 1.22m in the x -direction and 66 cm in the y -direction. It is assembled from 43.2 cm x 160 cm x 503 cm iron slabs, as shown in Fig. 2.7. The magnet has no air gap, and the iron has extremely high purity, allowing a 2000 A excitation current to generate a nearly constant, central magnetic field of 1.9 Tesla (yielding a 2.91 GeV/c total magnetic deflection). The field is generated by exciting the embedded aluminum “*bedstead*” coil to 2000 A at 25 V (50 kW). The current exciting FMAG is monitored by the Fermilab ACNET system and is broadcast to the SeaQuest slow data acquisition system every acceleration cycle. The excitation is also input to the beam-disabling safety system in order to prevent the beam from hitting the SeaQuest spectrometer when FMAG is not fully powered.

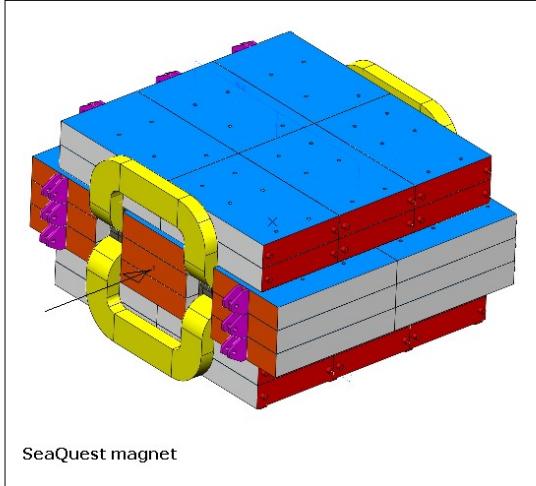


Figure 2.7: Perspective drawing of FMAG’s aluminum coils embedded in an arrangement of iron slabs.

FMAG also acts as the beam dump for the 120 GeV beam. There is a 5cm diameter by 25 cm deep bore drilled into the upstream end of FMAG (recall, this is the origin of the experiment’s coordinate system). The 120 GeV protons that do not interact in the SeaQuest targets 125 cm upstream of FMAG, interact in the central iron slab. Most of the 2.0 kW beam power is dissipated in this slab and is eventually conducted to the coils and external surfaces to be radiated away.

The downstream magnet, denoted “KMAG”, is a 300 cm long iron rectangular magnet with a 289 cm wide by 203 cm high central air gap. It was originally constructed by the KTeV collaboration[5] at Fermilab. It is excited to a central field of 0.4 T (0.402 GeV/c magnetic deflection) by 1600 A at 270 V (430 kW). The spatial distribution of the magnetic field in KMAG was measured by the KTeV group and re-verified by SeaQuest. In normal running conditions, both FMAG and KMAG bend muons horizontally in the same direction. This two-magnet configuration is often referred to as a focusing spectrometer.

The 2.91 GeV/c and 0.402 GeV/c magnetic deflections deliver a transverse-momentum (p_T) kick along the to charged particles passing through the spectrometer. The magnets bend the paths of the muon in the $\pm x$ direction, with the sign depending on the orientation of the magnetic fields and the particles’ charges. Between Run II and Run III of data taking, the current direction was reversed, thereby reversing the direction of the magnetic fields. During Run II, the magnetic fields were pointing in the $-y$ direction, and in Run III, the magnetic fields were flipped to point in the $+y$ direction. This was done for two reasons: (1) to identify any left-right asymmetries in the experiment, and (2) to limit the amount of radiation on the electronics in the experimental hall, as a large number of positive particles were being swept directly towards the electronics racks during Run II.

2.6 Beam Dump, Shields, and Absorbers

In order to prevent damage to the downstream detectors from the beam and reduce signals from incidental radiation, the spectrometer is designed with a beam dump and two hadron absorber walls. Approximately 125 cm downstream of the target table is the water-cooled beam dump whose upstream face is located at (0.0, 0.0, 0.0) m. The beam dump is one of the many solid iron 5m blocks that fill and surround the FMAG coils. The whole length of the beam dump along the beam axis is equivalent to ≈ 35 nuclear interaction lengths of iron.

Between the downstream face of FMAG and Station 1, there is a 2 cm thick wall of borated polyethylene which is put in place as a fast neutron shield. This material is 5% boron by weight, with the rest being polyethylene. The polyethylene contains high hydrogen content, making it an effective fast neutron radiation shield, slowing down the fast neutrons down to thermal speeds. The boron in the material provides attenuation of thermal neutrons, thus reducing the levels of capture-gamma radiation elsewhere in the experiment. Borated polyethylene at this thickness is a common and optimal neutron shielding material for areas of low to intermediate neutron flux where the temperature is below 82°C [96]. These conditions make the downstream side of FMAG ideal for its placement.

Farther downstream, there is another hadron absorber wall located between Station 3 and Station 4. The absorber wall consists of a stack of 98 cm thick iron blocks. This is an equivalent of ≈ 6 nuclear interaction lengths. The purpose of this wall is to identify muons at the rear of the apparatus by effectively blocking all other types of particles. The only charged particles which can penetrate this absorber wall that weren't swept away by either magnet are muons.

2.7 Tracking Detectors

The tracking detectors are the instruments used for measuring the values of the kinematic variables of the dimuon pairs. Several different types of detectors are grouped together to form a detector *Station*. The types of detectors used are hodoscopes, wire chambers, and proportional tubes. There are four stations throughout the experimental hall that provide tracking information at different points along the spectrometer, numbered from 1 to 4 in order of increasing z . Station 1 is located between FMAG and KMAG. Station 2 is located at the downstream face of KMAG. Station 3 and 4 are just upstream and downstream, respectively, of the iron absorber wall. The Station layout can be seen in Fig. 2.8

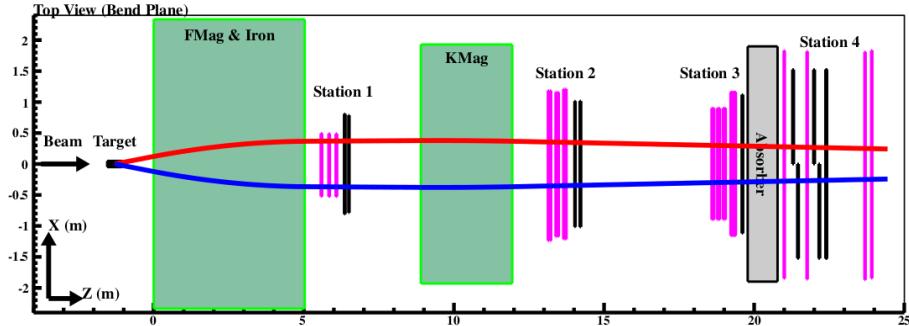


Figure 2.8: Spectrometer layout of FMAG, KMAg, and Detector Stations 1-4.

2.7.1 Triggering Hodoscopes

Hodoscope arrays are located at each of the four detector stations. These detectors' primary usage is to select events with two opposite-signed muon tracks in them. Certain 'roads' through the spectrometer are defined in the fast trigger logic, and when two desired roads are observed in a given event, the trigger system tells the data acquisition systems to record that event's data. In addition to this, the hodoscopes provide analysts with the ability to discard or ignore certain hits in adjacent chambers for which there is no corresponding nearby hodoscope hit. This is useful in decreasing the hit multiplicities in the wire chambers, which in turn decreases the combinatoric complexity of reconstruction algorithms.

Each of the eight hodoscope planes is split into two halves: top and bottom in the case of planes with vertically-oriented paddles, or left and right for planes with horizontally-oriented paddles (denoted by 'T', 'B', 'L', and 'R', respectively). In each half-plane, the hodoscopes are a set of long rectangles arranged 'picket fence'-style with a small 0.3175 cm overlap, as to prevent any particles from possibly slipping between paddles. A single hodoscope detector element is composed of plastic scintillator material connected to Philips XP 2008 photomultiplier tubes (PMT) by plexiglass light guides. Stations 1, 2, and 4 each have two hodoscope planes, with planes of both vertically- and horizontally-oriented paddles (for measuring in x and y , respectively). Station 3 consists only of a vertically-oriented plane, and thus measures in the $x-z$ direction only, which is in the experiment's $x-z$ bend plane. The hodoscope planes are named according to detector station and the direction that they measure. For example, the y -measuring hodoscope plane in Station 2 is called "H2Y". The individual half-planes are named according to detector station and which half it is of the two. For example, the top half of the x -measuring hodoscope plane in Station 1 is referred to as "H1T". As such, the "H3X" detector is composed of "H3T" and "H3B". The detailed specifications of each hodoscope plane are given in Table 2.2. A precise alignment of the hodoscopes was achieved by examining the distributions of positions of tracked muons at each hodoscope plane when a given hodoscope element in that plane was

fired.

Detector	Paddle Width [cm]	Overlap [cm]	# of paddles (per half-plane)	Width × Height (per half-plane) [cm] × [cm]	<i>z</i> -position [cm]
H1X	7.32	0.3175	23	162 × 70	666
H1Y	7.32	0.3175	20	79 × 140	653
H2X	13.00	0.3175	16	203 × 152	1421
H2Y	13.00	0.3175	19	132 × 241	1400 (L), 1406 (R)
H3X	14.59	0.3175	16	228 × 168	1959
H4X	19.65	0.3175	16	305 × 183	2234 (T), 2251 (B)
H4Y1	23.48	0.3175	16	152 × 366	2130 (L), 2146 (R)
H4Y2	23.48	0.3175	16	152 × 366	2200 (L), 2217 (R)

Table 2.2: Parameters of all hodoscope planes. *z*-positions of H2Y and H4 half-planes are offset slightly due to the half-planes themselves overlapping.

2.7.2 Drift Chambers

Each of Stations 1, 2 and 3 is equipped with a drift chamber (DC) to measure the passing *x* and *y* positions of muons at its *z* location, with each DC flat vertical to the *z* axis, and a drift cell is of the box shape. These measured positions are critical for reconstructing the trajectory of muons, and thereby their kinematics. Each DC contains six drift chamber planes, arranged in three pairs with parallel wire orientations (each pair referred to as a “view”). Wires oriented vertically (*x*-measuring) are referred to as being in the “X” view and at angles of $+14^\circ$ and -14° with respect to the *y*-axis are the “V” and “U” planes, respectively. The second plane in each view is offset from the first by one-half of a wire-to-wire distance (“cell width”) in order to resolve the left-right ambiguity of drift direction. This offset plane of each pair is referred to as the *primed* plane, and is denoted with a ‘*r*’. So, in each DC, there are X, X’, U, U’, V, and V’ planes, with primed and unprimed planes (like X and X’) constituting a view.

The individual drift chambers at Stations 1 and 2 are called “D1” and “D2”, respectively. Station 3 has two drift chambers since the desired acceptance area it has to cover is substantially larger than one DC can cover. These are split vertically to cover the top and bottom halves, and are called “D3p” and “D3m” where “p” and “m” stand for “plus” and “minus”. Table 2.3 summarizes the parameters of the DC’s. D1 and D3m have been upgraded during the data taking, as listed in Tab. 2.4. This original and upgraded versions are referred to as DN.1 and DN.2, respectively.

The acceptance size of each chamber has been adjusted with a Drell-Yan event simulation in order to be as sensitive as possible to the x_2 range of interest. Particularly, the greater the acceptance width is, the higher the reach in x_2 is. This makes the hit-rate tolerance of the chambers a key feature because the spectrometer

Chamber	Plane	Number of wires	Cell width [cm]	Width \times height [cm] \times [cm]	z -position [cm]
D1.1	X	160	0.64	102 \times 122	617
	U, V	201	0.64	101 \times 122	± 20
D1.2	X	320	0.50	153 \times 137	617
	U, V	384	0.50	153 \times 137	± 1.2
D2	X	112	2.1	233 \times 264	1347
	U, V	128	2.0	233 \times 264	± 25
D3p	X	116	2.0	232 \times 166	1931
	U, V	134	2.0	268 \times 166	± 6
D3m.1	X	176	1.0	179 \times 168	1879
	U, V	208	1.0	171 \times 163	± 19
D3m.2	X	116	2.0	232 \times 166	1895
	U, V	134	2.0	268 \times 166	± 6

Table 2.3: Parameters of all chambers. Those of primed planes are almost the same as of unprimed planes. The z -positions of U and V planes are relative to those of X planes.

Run	Period	Chamber combination
1	2012 Mar.-2012 Apr.	D1.1, D3m.1
2	2013 Nov.-2014 Aug.	D1.1, D3m.2
3	2014 Nov.-2015 May	D1.1, D3m.2
3	2015 Jun.-2015 Jul.	D1.2, D3m.2
4	2015 Sep.-	D1.2, D3m.2

Table 2.4: Combination of D1 and D3m chambers per data taking period.

is exposed to a large number of background particles, particularly near the edges where the desired x_2 events occur. It is particularly significant for the most-upstream station (i.e. Station 1) which receives the highest hit rates. Experimental data shows that the rate tolerances are 3.0 MHz/wire at D1, 1.6 MHz/wire at D2 and 0.7 MHz/wire at D3 with a beam intensity of 5×10^{12} protons/spill. The gas-amplification gain should not be degraded under these hit rates.

For Run 2 and beyond, the gas mixture used for almost all the chambers is Argon:Methane:CF4 (88%:8%:4%) with a drift velocity of about $20 \mu\text{m}/\text{ns}$. A “fast gas” mixture used for D1.2 (upgraded D1) is Argon:CF4:Isobutane:Methylal (68%:16%:13%:3%) with a drift velocity of $50 \mu\text{m}/\text{ns}$ and thereby a better hit-rate tolerance. This is “fast” in comparison to the $\approx 20 \mu\text{m}/\text{ns}$ rest of the DCs. The spatial resolution of each plane is required to be $400 \mu\text{m}$, which corresponds to a momentum resolution of $\Delta p/p = 0.03 \cdot p$ (GeV/c). The resolution of dimuon invariant mass is dominated by the multiple scattering in FMAG; the chamber momentum resolution is about 10% of the total mass resolution at maximum.

The D1.1, D2, and D3m.1 chambers have been inherited from previous Drell-Yan experiments that have been conducted at Fermilab. Chambers D2 and D3m.1 have their origin in SeaQuest [81] and D1.1 is from SeaQuest’s direct predecessor, E-866/NuSea [60, 104]. Since these chambers have not been used for decades since E-866, they had to be refurbished by restringing $\approx 30\%$ of their sense wires due to them being loose or broken. Newly supplied electronic readout boards were also mounted on these chambers.

The D3p and D3m.2 chambers were designed and constructed specifically for this experiment in order to cover the large acceptance required at Station 3. D3p was newly constructed by the TokyoTech SeaQuest collaborators and was shipped from Japan to Fermilab. The first part of data taking, Run 1, was carried out using D3m.1 while preparing for the construction of D3m.2. The newer D3m.2 is wider than D3m.1 by 25 cm at each side, allowing the high- x_2 statistics on \bar{d}/\bar{u} and $Y_A/Y_{^2H}$ to increase by $\approx 20\%$ at $x_2 \sim 0.3$ and $\approx 10\%$ at $x_2 \sim 0.4$. The operational stability also improved, as D3m.1 suffered from frequent dead/noisy wires, HV trips, and leak currents. The D1.2 chamber was also designed and constructed for this experiment by the University of Colorado Boulder. As it is wider than D1.1 by 25 cm at each side and greater hit-rate tolerance, the anticipated statistics is expected in the high- x_2 region is expected to increase still more. It was installed in the experimental hall near the end of Run 3.

2.7.3 Proportional Tubes

Downstream of Station 3 and the 1 m thick iron hadron absorber wall is Station 4 with its hodoscope planes and proportional tube detectors (prop tubes). At Station 4, the only beam-induced particles that remain that can leave tracks in an ionization detector are high energy muons. The prop tubes enable the task of

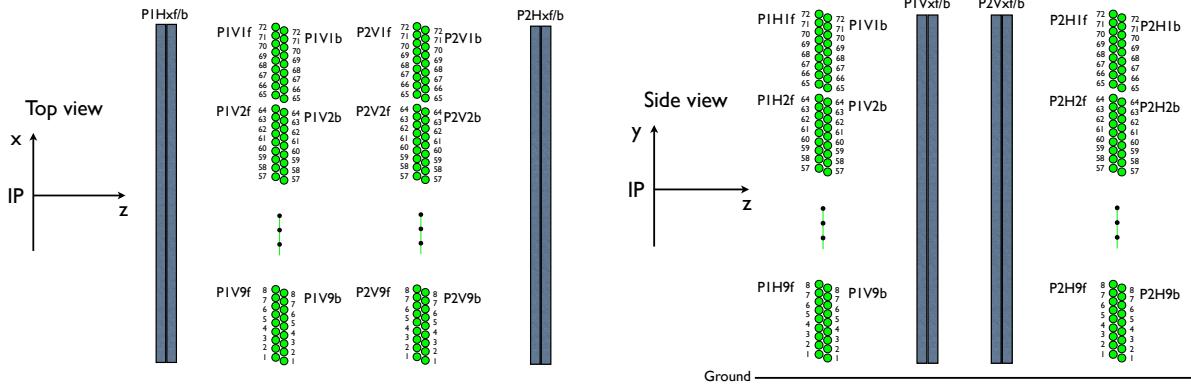


Figure 2.9: Proportional tube top ($x - z$ plane) view. Figure 2.10: Proportional tube side ($y - z$ plane) view.

muon particle identification (PID) for the experiment and consist of four planes. Each detector plane is made of 9 prop-tube modules, with each module assembled from 16 12-ft long 2" diameter prop-tubes staggered to form two sub-layers. The first and fourth planes are oriented along the horizontal direction (tubes parallel to the floor) to provide positional measurements in y , as shown in Fig. 2.9. The second and third planes are arranged vertically to measure x -position as shown in and Fig. 2.10.

Each plane of proportional tubes is composed of 9 *modules*, each having a set of 16 proportional tubes: 8 in line in a horizontal or vertical orientation and 8 more offset (“primed”) from the first 8 by 0.5 inches. This offset structure prevents any muons from going undetected between tubes and allows for left-right disambiguation for two hits in adjacent primed-unprimed sub-plane pairs in the same module. The modules are labeled 1–9 in increasing order from left to right for the x —measuring planes and from 1–9 in increasing order from top to bottom. So, the top module of the first vertical plane of prop tubes is P2H1. Additionally, in each module, the primed and unprimed sub-planes are referred to as *f* for “front” (facing upstream) and *b* for “back” (facing downstream). This substructure can be seen in detail on Fig. 2.9 and Fig. 2.10.

A single prop tube is made of 2-inch diameter aluminum tubes with a wall thickness of 1/16 inches. The central anode wire is a gold-plated 20 μm diameter tungsten wire kept at approximately 1.95 kV. Considering the staggered nature of each plane’s substructure, one can in principle achieve a spatial resolution of 0.3 mm. During Runs 2 and 3 of data taking, a resolution of 0.5 mm for high energy muons was observed, which is more than sufficient for muon identification purpose. The gas mixture for prop-tubes is P-10 (Ar:Methane = 90:10) mixed with a 10% CF₄ gas (Ar:CO₂:CF₄ = 70:20:10) which yields the maximum drift time about 400 ns. With this maximum, the prop tubes can handle a singles rate up to 2 MHz while normal operational hit rates are typically below 1 MHz.

A typical desired high-energy muon within spectrometer acceptance will traverse through two prop-tubes in each plane and induces hit signals on two anode wires. The path of a track is reconstructed from the drift time measured on the two anode output, with a custom TDC board that provides 0.44 ns timing resolution. With the hit information reconstructed from readouts of the forward and backward planes in $x - z$ and $y - z$ direction, precise reconstruction of the track trajectories can be obtained. Ideally, 8 hits from the 4 planes are used to form a track pointing back to the target. If such is the case, a candidate muon track is successfully identified.

Detector Plane	Number of modules	# tubes per module	Width (x) \times height (y) [cm] \times [cm]	z -position of front (back) sub-plane [cm]
P1H	9	16	368 \times 368	2,099 (+4)
P1V	9	16	368 \times 368	2,175 (+4)
P2V	9	16	368 \times 368	2,367 (+4)
P2H	9	16	368 \times 368	2,389 (+4)

Table 2.5: Parameters of the four proportional tube planes.

2.7.4 Mass Resolution from Chamber Resolution

The mass resolution of the spectrometer is, in part, limited by the resolution of the tracking chambers and the distances between them. For an arbitrary set of hit positions in Stations 1, 2, and 3 (r_1, r_2, r_3) which are separated from each other in z by distances z_{12} and z_{23} , with KMAG in between Stations 1 and 2 supplying a P_{kick} , one can derive $\Delta P/P$. A line, or track segment, is first reconstructed between r_2 and r_3 . The slope (and its uncertainty) of this track segment in the $z - x$ plane is:

$$s_{23} = \frac{r_3 - r_2}{z_{23}} , \quad \Delta s_{23} = \frac{1}{z_{23}} \sqrt{\Delta r_3^2 + \Delta r_2^2} \quad (2.3)$$

The values of Δr_n here are the position resolutions of the individual tracking chambers at each Station. One can use this slope to project a trajectory through KMAG and onto Station 1. The momentum is calculated by seeing where the actual hit position is in Station 1 and looking to the distance (d) between where the particle hit the station and where it would have hit had KMAG's magnetic field not existed. This is called a *sagitta analysis*, as the magnetic field moves the particle's trajectory as if along a circle, and we are looking to the particle's path along the sagitta of that circle. Ultimately, the momentum uncertainty is related to this distance as:

$$\frac{\Delta P}{P} = \frac{P}{P_{kick}} \Delta d \quad (2.4)$$

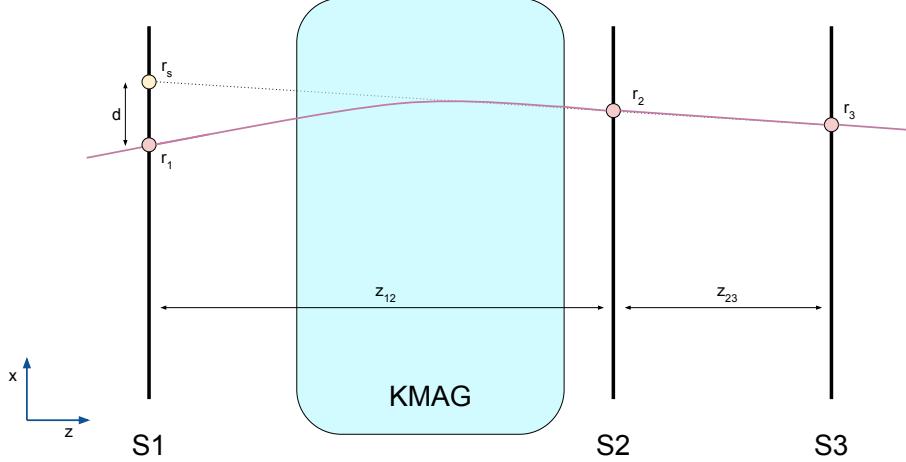


Figure 2.11: A simplistic depiction of a track passing through from Station 1, through KMAG where its path is bent by the magnetic field, and then straight on through Stations 2 and 3.

The expected sagitta point (r_s) is located at:

$$r_s = s_{23}z_{12} + r_2 \quad (2.5)$$

$$\Delta r_s = \sqrt{\Delta s_{23}^2 z_{12}^2 + \Delta r_2^2} = \sqrt{\Delta r_3^2 \left(\frac{z_{12}}{z_{23}}\right)^2 + (1 + \left(\frac{z_{12}}{z_{23}}\right)^2) \Delta r_2^2} \quad (2.6)$$

And the distance, d follows as:

$$d = r_s - r_1 \quad (2.7)$$

$$\Delta d = \sqrt{\Delta r_s^2 + \Delta r_1^2} = \sqrt{\Delta r_3^2 \left(\frac{z_{12}}{z_{23}}\right)^2 + (1 + \left(\frac{z_{12}}{z_{23}}\right)^2) \Delta r_2^2 + \Delta r_1^2} \quad (2.8)$$

The momentum resolution due to chamber resolution can therefore be described as:

$$\frac{\Delta P}{P} = \frac{P}{P_{kick}} \sqrt{\Delta r_3^2 \left(\frac{z_{12}}{z_{23}}\right)^2 + \left(1 + \left(\frac{z_{12}}{z_{23}}\right)^2\right) \Delta r_2^2 + \Delta r_1^2} \quad (2.9)$$

The dependence of the resolution on P/P_{kick} works such that the larger the P_{kick} is (w.r.t. particle momentum P), the more the track is bent, and the more precisely the original momentum is defined. Besides chamber resolution, other contributions to this momentum resolution are from multiple scattering through FMAG's iron beam dump and from the spectrometer's angular resolution.

The mass resolution is linked to momentum resolution via the following relation

$$\frac{\Delta M}{M} = \frac{\Delta P}{2P} \quad (2.10)$$

The factor of two arises from the fact that the Drell-Yan dilepton mass comes from two independent muons with the same momentum resolution. This mass resolution is important because it, in turn, contributes to the $\Delta x_2/x_2$ resolution, which is the key dependent variable for SeaQuest measurements.

$$\frac{\Delta x_2}{x_2} \approx 0.57\Delta x_F + 0.012M^2\frac{\Delta M}{M} \quad (2.11)$$

2.8 Trigger

The SeaQuest Trigger System is designed to quickly select candidate dimuon events from the high-rate, high-background environment using discriminated hodoscope signals.

2.8.1 Design Requirements

The trigger must select events of interest to the main physics goals with good enough efficiency (maximum signal-to-background ratio) to facilitate high-statistics analyses. In general, it is optimized to accept high-mass (4-10 GeV) dimuons originating from the targets and beam dump. The event selection is also designed to intentionally reduce acceptance of dimuons from other, higher-rate sources, such as J/Ψ decays and non-Drell-Yan dimuons originating from the beam-dump, though enough J/Ψ events are triggered on to allow high-statistics analysis of J/Ψ physics. Other goals of the trigger design are to keep the triggering rate low enough to maintain an acceptable DAQ livetime and to keep the trigger's internal throughput deadtime-free. By design, there should be no bottleneck at this stage, and the trigger should be capable of firing on any and all RF-buckets while the DAQ is live.

The hardware, firmware, and design should be sufficiently flexible to quickly accommodate changes in the spectrometer, beam conditions, and physics goals. Any changes to the geometric acceptance of the spectrometer, such as new/moved detectors or changes in the magnetic fields must be immediately reflected in the trigger selection in order to maintain high signal efficiency and good background rejection power. Similarly, a change in the beam duty factor or intensity should be accompanied by a change in the event selection, ensuring trigger rate optimization and zero internal deadtime. Finally, the trigger should be capable of specific modifications to facilitate special runs for other physics goals. For these reasons, the design of the trigger system underscores a significant need for flexibility in its hardware and design.

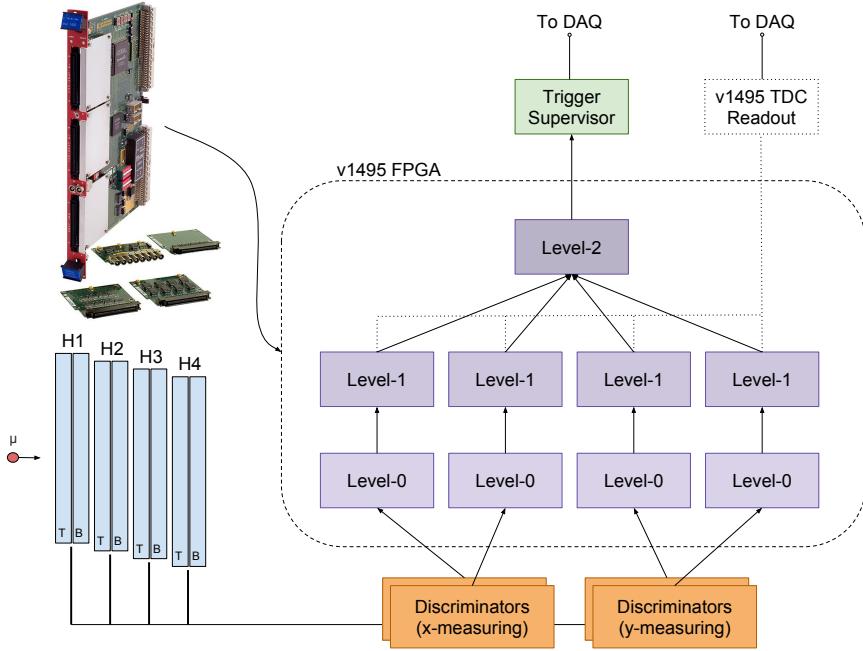


Figure 2.12: The trigger system at SeaQuest is composed of 9 CAEN v1495 FPGA modules. These modules output a trigger signal to the Trigger Supervisor, which tells the DAQ when to record data.

Lastly, the design of the trigger system should include self-diagnostic capabilities, allowing for constant monitoring of the trigger system's performance. Internal pulser-testing is employed to test the function of each compiled firmware every time the trigger logic changes. For added transparency, data from the internal TDC's is used by online and offline software to monitor the self-consistency of the trigger for each recorded physics event.

2.8.2 Trigger Hardware

The triggering system, which is depicted in Fig. 2.12, begins with the Stations 1-4 hodoscope arrays. Hodoscopes are used for this task as the time resolution and high speed is limited only by the speed of light through the scintillator paddles and secondary emission of electrons (through the stages of the PMTs. This occurs over the span of nanoseconds, which allows for fast triggering on RF buckets that occur every 19 ns during spills. The hodoscope planes - both x- and y-measuring planes - are separated into top and bottom sections as far as triggering is concerned. The signals from the hodoscope arrays are passed through a set of discriminator modules, which will only pass a signal on to the rest of the trigger if there is enough of a signal from the PMT's past a preset threshold.

The output of the discriminators is passed along to one of the nine CAEN v1495 FPGA (Field Pro-

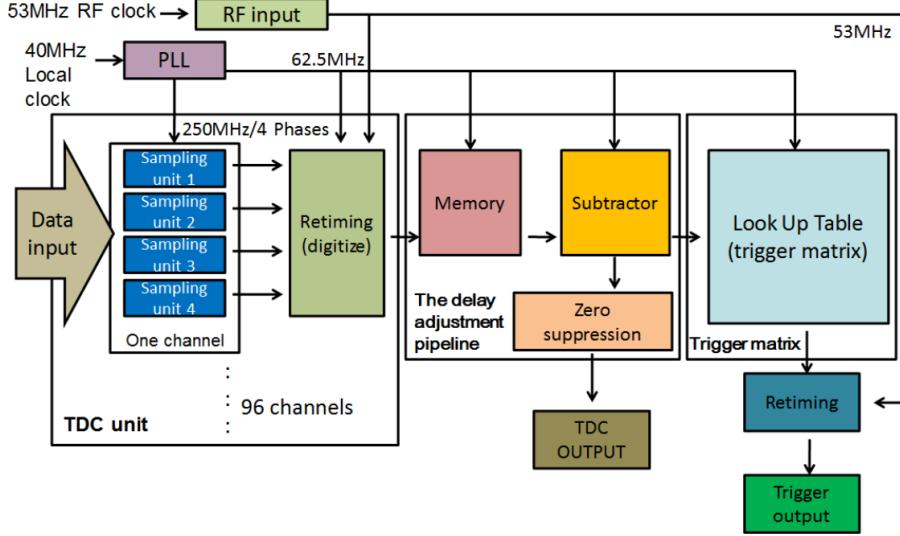


Figure 2.13: A block diagram of the major functions of the v1495 FPGA [97].

grammable Gate Arrays) VMEbus (Versa Module Europa) modules [24]. This hardware trigger consists of a single decision stage with a three-step parallel pipeline. These steps are referred to as the *level-0*, *level-1*, and *level-2* stage triggers. The detailed internal flowchart of the CAEN v1495 electronics can be seen in Fig. 2.13.

At the first step, level-0, there is nothing actively performed during nominal data taking. The signal is passed on through, unaltered, to level-1. As it stands, level-0 is not used for any of the triggering modes. This stage is, however, extensively used for the pulser testing and self-diagnostic purposes. The level-0 stage is capable of being programmed to pulse arbitrary patterns on to the level-1 stage, which is very useful for observing how level-1 (and subsequently, level-2) will behave in a controlled setting. Whenever a new firmware is installed on the v1495, a pulser test is performed to ensure that all design requirements are met before uploading the firmware for production data taking.

The next level-1 trigger records the hit signals from the x- and y-measuring hodoscope hits split into top and bottom partitions with respect to their location in the SeaQuest spectrometer. The hit patterns are tested via a lookup table against a preselected set of hodoscope hit patterns, or “*trigger roads*”. Each trigger road corresponds to a set of four hodoscopes: one from each station, with all being in the top or bottom half of the spectrometer. In this lookup table, there is a charge and approximate p_T value for each preselected road. These are known via studies of Monte Carlo studies of muon paths from di- or single-muon events originating from the targets or beam dump. The roads were selected such that as many good dimuons are selected while reducing roads that are dominated by background signal that might cause trigger and/or DAQ deadtime.

Trigger	Condition	Sign	# μ	Prescale
FPGA-1	$T \wedge B$	Opposite	2	1
FPGA-2	$(T \wedge T) \vee (B \wedge B)$	Opposite	2	1000
FPGA-3	$T \wedge B$	Same	2	123
FPGA-4	$T \vee B$, all p_T	+ \vee -	1	30000
FPGA-5	$T \vee B$, $p_T > 3 \text{ GeV}$	+ \vee -	1	2000
NIM-1	Y-coincidence	NA	NA	30000
NIM-2	X-coincidence	NA	NA	1000
NIM-3	Random RF	NA	NA	1000

Table 2.6: Trigger settings for Run 3 and beyond. Prescale figures shown are typical values and were adjusted as experimental needs were tuned. NIM-1 and NIM-2 triggers’ exact conditions could be changed and reconfigured from the counting room.

During data taking, only the x-measuring hodoscopes were used for the v1495 trigger. As such, the two level-1 stages that look to the y -measuring hodoscope signals were unused. The level-1 trigger logic identifies four-out-of-four X1-X2-X3-X4 coincidences, which are characteristic of high p_T single muons produced from the targets or beam dump. Each time a candidate coincidence that satisfies a preselected road is found, the level-1 step outputs certain logical bits indicating the track’s charge, detector half (top or bottom), and p_T bin.

The third step of the triggering logic is the level-2 trigger, or the “*Track Correlator*” stage. This step takes the outputs of level-1 (charge, detector half, p_T bin) and finds if any combinations satisfy any of the preprogrammed triggering modes. If one of these modes was satisfied, a signal is sent to the Trigger Supervisor (TS) module that communicates with the DAQ system to record an event at a specific synchronized RF bucket.

2.8.3 Triggering Modes

The v1495 FPGA had five configurable sets of physics triggers were used in the trigger system for the majority of production-level data taking. These five are referred to interchangeably as “Matrix” and “FPGA” triggers, followed by the number 1-5 indicating which of the five. These modes are described in detail in Table 2.6.

Different trigger settings were used for capturing data for different purposes. The primary triggering mode is FPGA-1, which looked for a combination of two roads of opposite sign that reside in opposite halves of the spectrometer. The FPGA-2 looked for good muon pairs that occurred in the same half. This, however, turned out to have a combinatoric problem that had two adverse effects: track reconstruction was difficult for these types of events, and the trigger fired off too often, causing a high trigger and DAQ deadtime. As such, FPGA-2 was disabled for much of data taking. FPGA-3 trigger looked for the same kind of signal as FPGA-1, but with same-sign muon pairs. This type of event is useful for analyzing the

experiment's combinatoric background. FPGA-4 and -5 are “singles” triggers that record events that are useful for estimating backgrounds and detector efficiencies. Each of these has a “*prescale factor*” that limits how many of each mode actually fires the trigger. For example, FPGA-3 has a prescale factor of 123, which means that, within a single spill, there must be 123 FPGA-3 events before it tells the DAQ to record one. This keeps certain high-frequency trigger modes from dominating the readout and causing undue deadtimes.

In addition to the v1495 FPGA triggering mechanism, the outputs of the hodoscope discriminators are also sent up to the control room to a set of NIM logic modules. These signals, however, do not describe each hodoscope paddle but instead are *ANDs* of the tops and bottoms of each plane. For example, if one or more paddles in H1T fires, then a positive signal is sent out indicating that H1T fired. These signals can be used to form a rudimentary trigger on situations such as 4- or 3-out-of-4 coincidences in the X- or Y-direction hodoscope planes. In general, NIM-1 is used for Y-coincidences, and NIM-2 is used for X-coincidences.

The third NIM trigger mode is of particular importance, as it is used for estimating a great deal regarding backgrounds and rate dependence issues. The NIM-3 trigger is called the “Random RF” trigger, as it is controlled by two different clocks beating against each other. When the two clocks are in coincidence, they will send a signal to the Trigger Supervisor to record the event corresponding to the next immediate RF bucket. This trigger allows the experiment to get a random sample of the events that are occurring at the spectrometer, notably without a bias that a selective trigger can incur.

2.9 Data Acquisition Systems

Create flowchart!

The whole of the SeaQuest experiment spans two sectors of the NM beamline (NM3 and NM4) and records not only detector data but also accelerator and atmospheric conditions. Designing a data acquisition (DAQ) for this experiment, therefore, requires a bandwidth and timing which cannot be provided by a single central computer system. SeaQuest has therefore employed four separate DAQ subsystems called “Main DAQ”, “Scaler DAQ”, “Beam DAQ” and “Slow Control”. The Main DAQ records the main detector information and the trigger timing. Scaler DAQ records various scaler readouts once per spill to reduce bias from the dead time of Main DAQ. Beam DAQ records information from a beamline Čerenkov detector read out by its QIE board. The Slow Control is a catch-all for everything else, from magnet current to radiation monitors that are read out once per spill.

2.9.1 Main DAQ

The MainDAQ is powered by a Jefferson Lab developed software package named CODA (CEBAF¹ On-line Data Acquisition). The Trigger Supervisor could receive up to 12 different kinds of triggers. The first four triggers (FPGA 1-4) can be prescaled up to 24 bits. The second four triggers (FPGA 5, NIM 1-3) can be prescaled up to 16 bits. The rest of triggers (NIM4, flush trigger, Beginning of Spill (BOS), End of Spill (EOS)) are not prescalable. The Main DAQ can configure which triggers are enabled at the Trigger Supervisor level and what the prescale factors are for each triggering mode. Once the TS receives a trigger, it will count the prescale factor and, based on how many triggers of that type have fired so far, it will decide if the Main DAQ will accept the trigger or not.

Once the TS accepts the trigger, it sends an “accepted trigger” signal to the other front end electronics, such as the detectors’ TDC readouts. Facilitating intercommunications, the Main DAQ system is connected to the TS and the CPU’s in the VME crates via a local network. These CPU’s are known colloquially at SeaQuest as readout controllers, or “ROC’s”. The TS connects with each VME crate, and when the TS sends an “accepted trigger”, an interrupt signal is sent to each of the ROC’s. The ROC’s will then start reading the various front-end electronics (FEE).

The TS also sends a signal to the QIE of the BIM, and the QIE retains information regarding the ± 16 RF buckets around that triggered RF bucket to assist in investigating the beam quality around each of our event triggers. The output of the QIE is encoded into a scaler latch card format. If the beam intensity in RF buckets neighboring the trigger is higher than the user-select threshold, then the board will issue a veto signal to the TS to ignore this trigger.

The Main DAQ’s deadtime is considerable, as it has to communicate with the TDC’s, which have the longest copy-in-progress time of $32\ \mu s$. The average TDC readout time is approximately $300\ ns$ per 32bits (one hit), and as such, the slowest ROC which contains 7 TDCs has, on average, $150\ \mu s$ deadtime. This deadtime is accounted for under the umbrella term of “*busy*” time and factored into the calculation of the “live protons” received by the experiment.

2.9.2 Scaler DAQ

The Scaler DAQ is used to ensure that the detector and trigger systems are working properly. The readout controller CPU is an MVME5500 [82] with four scalers. One of these scalers is triggered by the coincidence of a $7.5\ kHz$ gate generator and the beam spill signal. This records the $7.5\ kHz$ response of the hodoscope arrays. The other three scalers are triggered by the BOS or EOS signals and record spill-level rates. Data

¹The old acronym inside of an acronym bit. CEBAF: Continuous Electron Beam Accelerator Facility

collected by these spill-level scalers are the number of times each Main DAQ trigger is satisfied (both raw and accepted), the intensity of the beam, and the rates of the hodoscope arrays. The readout of this VMEbus DAQ is done using CODA very similarly to the Main DAQ but on a completely different system. An independent program analyzes the data in real-time in order to monitor the performance of the detectors and triggers, as well as the quality of the beam.

2.9.3 Beam DAQ

The Beam DAQ is composed of a Čerenkov detector in the proton beam (the Beam Intensity Monitor discussed earlier in this chapter), a QIE board, and a custom C++ program to control their operation and read out. The Beam DAQ is responsible for recording the 53 MHz structure of the beam, i.e. the intensity of each RF bucket. Its calculation of the 53 MHz duty factor $DF = \frac{\langle I \rangle^2}{\langle I^2 \rangle}$ is the primary measure of beam quality that accelerator operators use for tuning. In short, the duty factor is a measure of the *uniformity* of the intensity of the beam. If we were to ignore the buckets intentionally left empty, and if it were to be the case that every RF bucket had the same number of protons, the 53 MHz duty factor would be 100%.

There four types of data recorded by the QIE have already been covered in the Beam Intensity Monitor section: QIESum, beam inhibit, busy time, and the ± 16 RF bucket intensities. The Beam DAQ commences read out of each of these blocks of data when the EOS signal is seen. The block of QIE data for all buckets is about 300 MB. To read this much data in time to analyze it and be ready for the next spill, the DAQ program utilizes multithreaded processes. Three threads are used to read the data from the QIE board's three ethernet chips, and up to eight threads are used to analyze the data. Analyzed data is displayed on a public web page so that shifters and accelerator operators can monitor the quality of the beam (Fig. 2.14). The fully processed data is also written out to tab-delimited ASCII files, which are archived and also uploaded to our online MySQL server.

2.9.4 Slow Control Readout

The Slow Control readout system is an aggregation of many different readouts from various parts of the experiment. This is also the umbrella term for other critical monitoring and bookkeeping tasks that must be performed. The one common thread to them all is that they only *need* to be read out / performed once per spill in order to give perspective on the conditions of the experiment for each given spill. The term “*Slow*” in Slow Control is adapted from the term “Slow Spill”, which is used to characterize the method of beam delivery to SeaQuest. The Slow Control is run by an array of Python, Perl, and Bash scripts that, for the most part, run independently. These scripts primarily interact with and read out from EPICS, ACNET,

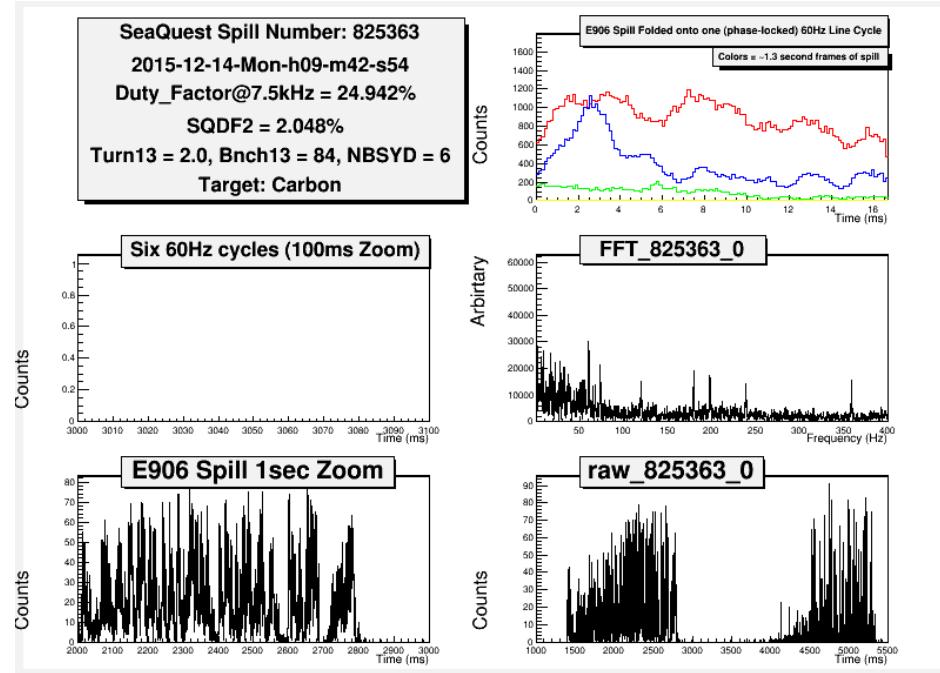


Figure 2.14: A standard Beam DAQ display showing beam characteristics and spill data. The bottom two plots show the output of the QIE on a ms scale.

and the local Spill Counter.

EPICS, ACNET, and Kiethley

The EPICS software package (Experimental Physics and Industrial Control System) is used for communicating any number of variables across the experiment's local network. EPICS is used for feeding readouts from many different sources and formats, and making them available to any other system that may have need of certain information, but does not natively communicate with a different system's format. Bridging this gap, EPICS is able to gather information regarding the target, beam, and environmental data into one place.

The EPICS server is installed on SeaQuest's target control computer, where the target position, pressures, temperatures, and proximity sensors are monitored and controlled. These values are critical to track and monitor, as they affect what target is in position at any given spill, along with information that may infer the cryogenic liquid targets' densities. Additionally, there is information regarding the operational mode of the target table, such as whether or not the target table is being manually controlled or if it is automatically repositioning itself. In Slow Control parlance, information from the target computer that is sent to the EPICS server is labeled as "Target" type data.

Several values are read into the EPICS server from the Accelerator Network (ACNET), which monitors

many critical systems. Beam intensities (S:G2SEM) and FMAG / KMAG currents are two of the major factors that must be factored in for any analysis. There are additional readouts with esoteric names that can record anything from beam pipe vacuum to upstream radiation monitors. It is important to note that the FMAG and KMAG variables from ACNET are read out to the EPICS server every 3 s, particularly because it is unsafe to deliver beam to the sensitive detectors without FMAG being fully powered. As such, the FMAG current is read into the beam interlock system. In the Slow Control feed, this type of information is labeled as “Beam” data.

Finally, there is a Keithley multimeter in SeaQuest Hall near the electronics racks by Station 3. This device has the responsibility of relaying information regarding ambient temperature, air pressure, and humidity up to the control room. This information on its own is not typically used for any analyses, but in the case that, for example, a rise in humidity results in any electrical device failures, the monitoring data exists to diagnose it as such. This data is gathered by a ‘kscan’ program that is run on the target machine, and its operation typically takes 11s to return a result. Data from the Keithley multimeter that is put into the Slow Control feed is labeled as “Environment” data.

Spill Counter

Perhaps one of the most critical components of the data acquisition is the coordination between the many sources of data in the DAQ. The Main DAQ, Scaler DAQ, Beam DAQ, and Slow Control do not share a single clock or network, so some measures need to be taken to ensure that data from different sources regarding a single spill’s worth of data actually all correspond to the same spill. The concept of a locally maintained Spill Counter arose as a solution to this problem, and it is handled in the form of a simple Python script and a flat text file containing a single integer.

The enacted solution begins with placing an initialized spill counter ($spillID$) value (currently $O(10^5)$) into an otherwise empty text file at a specified location. Then, when the ACNET readout receives a BOS/EOS signal from AD, a simple python script performs the following:

- Read the current $spillID$ from the file
- Broadcast this value over the EPICS server
- Insert a single ASCII-encoded event into the Main DAQ and Scaler DAQ CODA streams containing only that $spillID$ number
- Waits for 20 s while other scripts may want to use the current $spillID$ value to connect its value to the spill that just ended

- Updates the file to contain spillID + 1

This simple approach has been very effective in its usage so far, and the spillID incremented via this method has been shown to be accurate upon scrutiny of data contents across DAQs. This can be easily done by looking at a series of full and empty spills, as it is obvious which spills had a nominal beam spill and which did not by investigating the output of each DAQ. The weakness lies in the fact that this system relies on a single file, which is susceptible to any number of file access and timing issues. At some point in the future, it would be advisable to have all spillID-using systems to read the current value from the broadcasted EPICS server value instead of having several applications attempting to access and read the single file (even if just read-only).

Python Slow Control Script

With all of this going on to aggregate data from several sources, there is one final script that takes this aggregated data and makes it readily available for the rest of the DAQ systems. A `read_slowcontrol.py` script, once every 60 s supercycle, puts the following together into a single ASCII text file:

- Clock time (“vxticks”) of the Trigger Supervisor
- Spill count data
- Beam data
- Environmental data
- Target data
- Timestamp

Once all of this is written to a file, this is sent to the Main DAQ and Scaler DAQ to be integrated into the full CODA file as a plain text event with a simple, tab-delimited format.

With data in all these several formats, some coordinated effort has to be put into processing and storing it. Further, with these disparate sources, there must be some coordinating factor that unifies it so that users can associate any one piece of information with the rest. The following chapter discusses the solutions implemented by SeaQuest in the methods and choices of processing and storing all of this data.

Chapter 3

Photomultiplier Tube Base Upgrade

During Run I of SeaQuest, observations of hodoscope wire maps (as in Fig. 3.1) suggested an apparent drop in expected performance in the y -measuring hodoscopes. While this performance was most obviously seen in the y -measuring hodoscope planes, the x -measuring planes were likely also affected. This effect was assumed to be due to high-intensity RF buckets that caused very high multiplicity in all of the detectors in the spectrometer for that event. The result of these intense events seemed to push the PMTs and/or their PMT base electronics past their operational capacity.

The understood cause of this “*sag*” in performance, as it came to be called, was due to a destabilization in the voltage divider in the PMT base. This critical component holds each dynode stage at a specific voltage, and when this destabilizes and is unable to maintain an appropriate voltage difference between dynode stages, inefficient performance of the PMT results.

During the Fall of 2012, prototyping and testing were performed with the goal in mind being to assemble a new base for the Philips XP-2008 PMTs [87] and compare its performance to the original PMT base and to some modern, high-performance Hamamatsu PMTs. Once a base design tested well, the new bases would be manufactured and installed in the existing frames of the original PMT bases.

3.1 PMT Basic Construction and Operation

Figure 3.2 shows a schematic design of a typical photomultiplier tube and base setup. It consists of a photocathode that is followed by an electron multiplier section (or dynode string) then an anode from which a final signal is delivered. During operation, a high voltage is applied to the photocathode, dynodes, and anode in such a way that there’s a potential “ladder” going from stage to stage. When an incident photon from the hodoscope scintillator paddle hits the photocathode, an electron is emitted via the photoelectric effect [39]. The voltage difference between the cathode and dynode stages draws the emitted electron to the dynodes, and each time an electron hits a dynode, some of that electron’s energy is transferred to other electrons in the dynode. These electrons then are emitted and become accelerated towards the next dynode

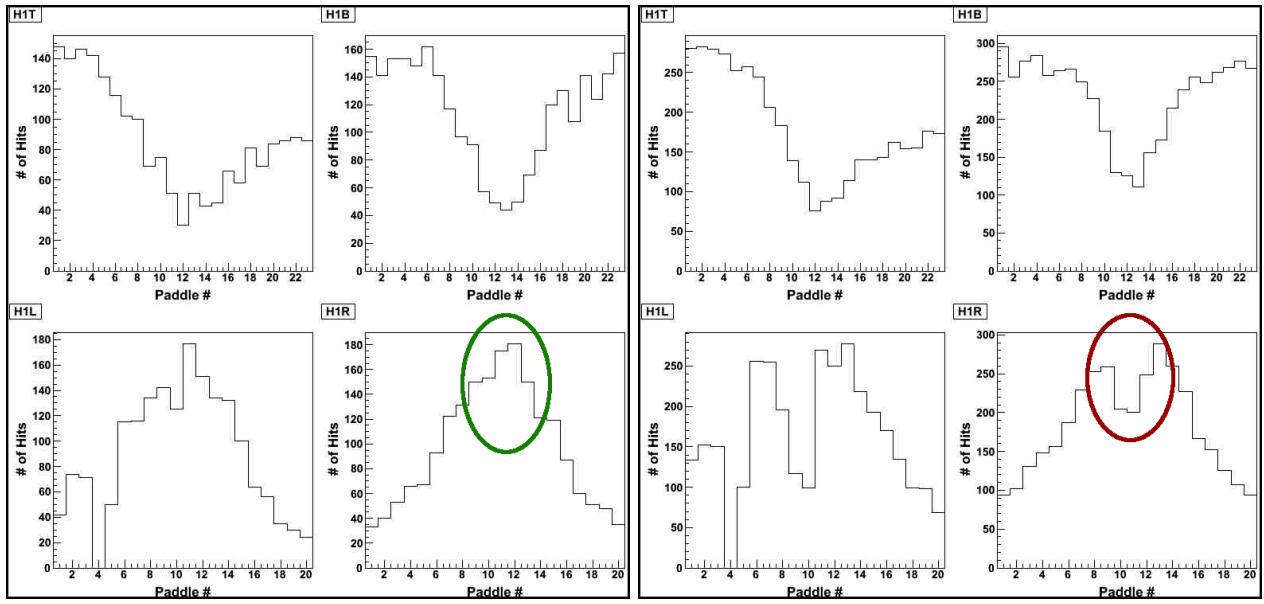


Figure 3.1: (Left) Histogram of hodoscope ‘hits’ in a typical event; (Right) Histogram of high-intensity event, with marked sagging most noticeably in the middle of the y-measuring hodoscopes

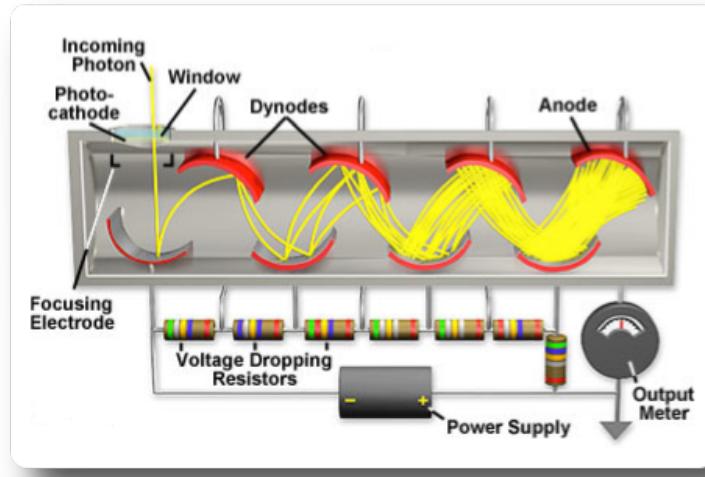


Figure 3.2: A diagram of typical PMT operation. The circuit controlling the voltage-dropping resistors is the part that was upgraded in this chapter.

stage. This process is called secondary emission, and by the time the process is repeated, there is a cascade or avalanche of electrons that land on the anode, resulting in a signal that can be amplified and analyzed.

It is the case that the voltage divider ultimately supplies the electrons that are emitted in this signal cascade. If too many photons and resultant electron cascades occur, the dynode stages' voltage divider will destabilize as they attempt to resupply the dynode stages with electrons. The problem that was experienced at SeaQuest was that these high-intensity events were flooding the PMTs with photons, causing this “saturation” which caused this destabilization and the inefficient performance that was observed. The goal specifically was to test out modern base designs that provided for added stability to the performance of the voltage divider, even under high rates.

In general, each base divides around a -1500 V potential total over the photocathode (K), ten dynode stages (D1-D10), and the anode (A). There are two currents that are referred to here:

- Signal Current: This is the signal that passes over the anode, which is the end-result of the cascading secondary emission electrons from each dynode stage.
- Bleeder Current: This is the current through the voltage divider. It is termed the “bleeder” current since the compounding electrons in the signal current must be “bled” from the current through the voltage divider.

Throughout these voltage base designs, capacitors are commonly implemented in the latter dynode stages where the most electrons are emitted. These capacitors, when charged, are able to replenish the lost charge on its corresponding dynode stage in the event that an intense light pulse induces a large signal current. As the capacitor is able to hold its own charge, this resupply can occur without requiring the charges to be drawn from the bleeder current, thereby keeping the voltage across the dynode stages more stable.

3.2 PMT Base Design Iterations

There were several iterations of base design to determine which was best to approach for full base production and installation at SeaQuest. The core addition was the inclusion of transistors between dynode stages, according to the improvements suggested by C.R. Kerns in his paper regarding high-rate PMT bases [71]. Common solutions to destabilization in PMT bases have been to have (1) very large capacitor banks with charges $> 10^3$ times greater than the time-averaged dynode current and/or (2) miniature on-board, separately-powered Cockcroft-Walton power supplies for the final dynode stages. A Kerns-style transistorized base allows for a light-weight, small size, and simple base that does not require extra power supplies or voluminous energy storage capacitors.

In general, there are three important features that were tuned in this set of prototypes that affected the performance of the phototubes:

- Lower resistance
- Transistors (with protective diodes)
- Higher capacitance between dynode stages
- Distribution of voltage division

The lower overall resistance of the voltage divider increases the bleeder current. This means that the base will be more capable of handling high-intensity, as it will be better able to replenish the charges on each dynode stage in the case of a large signal. Typically, the larger the bleeder current, the larger the signal current can be without destabilizing the voltage divider. Higher rates usually put higher demand on the signal current, so by reducing the overall resistance, one can easily increase the rate capability. The shortfall here is that with voltage constant and resistance decreased, according to Ohm's Law ($V = IR$), the current will increase. As a result, the power dissipated by the circuit ($P = I^2R$) will go as I^2 , and the PMT base may heat up to critical temperatures faster than it can dissipate the heat as there is no significant ventilation in the PMT base enclosure. The Philips XP-2008 manual quotes that for continuous usage and storage the ambient temperature should not exceed $50^\circ C$ ($122^\circ F$). In addition to heat concerns, there's typically a power rating for the class of small on-board resistors that were planned to be used. Approaching or exceeding that power limit would run the risk of burning out a resistor and rendering the base inoperable.

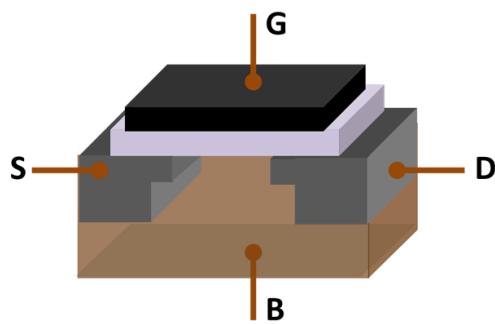


Figure 3.3: MOSFET showing gate (G), body (B), source (S) and drain (D) terminals. The gate is separated from the body by an insulating layer (white) [85]

the dynodes. Wherever transistors are used, diodes are also implemented to prevent the unlikely case of a

The use of a specific type or micro-transistor, a metal-oxidesemiconductor field-effect transistor (MOSFET), is introduced here to maintain the proper voltage division. In general, MOSFET transistors have an *source*, a *drain*, and a *gate*, where current flows freely through from the source to the drain, gate permitting (Fig. 3.3). If at any point a certain voltage across the gate of the transistor is not supplied (here, the voltage across dynode stages), then the source-to-drain current through the transistor is stopped until the proper gate voltage is restored. This helps greatly to "intelligently" regulate the voltage across

Recommended voltage divider

Type A for maximum gain

K	D1	D2	D3	D4	D5	D6	D7	D8	D9	D10	A	(total :12)
2	1	1	1	1	1	1	1	1	1	1	1	

Type B for bets timing / linearity compromise

K	D1	D2	D3	D4	D5	D6	D7	D8	D9	D10	A	(total :18.25)
2	1	1	1	1.25	1.25	1.5	2.25	2.25	2.5	2.25		

K: photocathode Dn: dynode A: anode

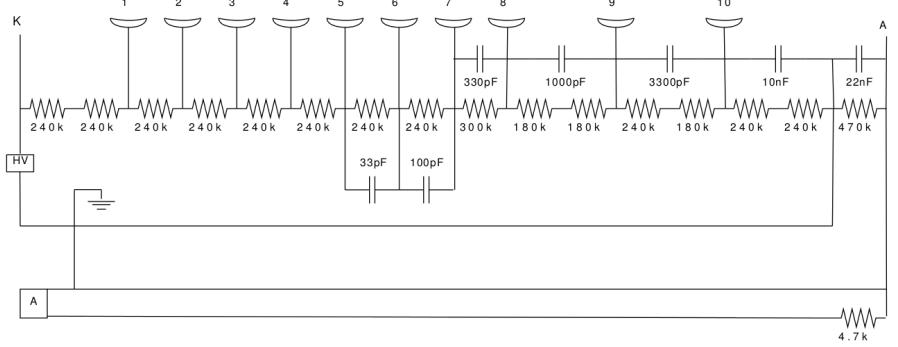
Figure 3.4: Suggested voltage division schemes for gain vs. timing/linearity compromise [87].

current moving across the transistors' gate in the wrong direction. This protects the transistors from being damaged particularly when powering the circuit on and off.

Having capacitors along the higher dynode stages allows for a quick resupply of charges to the dynodes, thus maintaining proper voltage division. This, however, is only a stop-gap measure and is only effective in cases of high instantaneous currents. It is the case that, should there be a constant too-high rate of operation, the capacitors will not be able to recharge themselves. Typical recharge time for the capacitors discussed in this section can range from 0.1-1 ms.

Finally, the specific division of voltage across each stage, from D1 to A, has an influence on the behavior of the PMT operation. As we see from the operations manual of the phototube in Fig. 3.4, in the case of a progressively increasing voltage division, there is a good compromise between timing and linearity. With respect to phototube operation, "linearity" is the quality that the amount of charge deposited on the anode is linearly proportional to the energy of the incident photon. "Timing", on the other hand, is the quality that the time it takes for a high-energy photon signal and a low-energy photon signal to progress through the stages should be the same. For the purpose of optimization at SeaQuest, we wish to optimize the amount of signal (i.e. amplification or *gain*) that the phototube can accommodate. For this, the recommended voltage partitioning is flat from D1-A [87].

It should be noted that the prototype iterations of PMT base design changes were not intended to cover the entire phase space of these tunable parameters. The purpose of these tests was to get a sense of how changing one or more of these parameters could affect the PMT high-rate capabilities. The decision for a final base design had to be decided on with relative speed to get them manufactured, built, and installed before Run II of the experiment. As such, true optimization of all parameters could not be achieved within the scope of these tests.



NOTES: Ground of the HV supply to the ground of the anode output separated, where they were connected with the 22nF capacitor and 470k resistor in parallel, in order to

- I) keep the noise on the HV ground away from the anode
- II) to avoid non-linear effects coming from the final 10th dynode

On some bases, the ground of the HV supply and anode output are shorted together with a braid.

Precision: +/- 22% (S), 33pF only +/- 10% (K)

Figure 3.5: The original PMT base inherited from the ARGUS and HERMES experiments.

3.2.1 Original Base

The base that came attached to the PMTs were manufactured specifically for use by the ARGUS experiment, which was a relatively (by SeaQuest standards) lower-rate collider experiment that used e^+e^- annihilation at the *DORIS II* ring at DESY. After their tenure at ARGUS, they were handed down to the HERMES experiment located at the *HERA* polarized electron accelerator at DESY.

Though no actual circuit diagram was documented for the original PMT base, it was dissected and each component was measured. The results can be found in Fig. 3.5, and its voltage division seen in Fig. 3.6. It features a simple string of resistors with capacitors of increasing capacitance along the last six stages. The voltage division can be recognized to be similar to the timing-linearity compromise scheme described in Fig. 3.4. With a total resistance of approximately $3.95\ M\Omega$ and the operational voltage of -1500 V, the expected standing (bleeder) current even when sitting in the dark is expected to be 0.38 mA. With this in mind, we would expect the voltage divider to destabilize when the signal current approaches this value.

3.2.2 Prototype Base v1

Once the task was set to update the PMT base design, the Fermilab Particle Physics Division was consulted on the matter. In 2010 a base design with similar goals for the exact same PMT model was designed by Sten Hansen [59]. The circuit diagram for the new base can be found in Fig. 3.7.

Here, the resistance was significantly reduced by a factor of about 2.9, allowing for much more bleeder current, without exceeding or closely approaching the on-board resistors' power rating. Also, the voltage division was designed to be relatively "flat" (Fig. 3.8) across stages from D1 to A, which is stated to be recommended for optimal gain. With a total resistance of approximately $1365\ k\Omega$ and the operational voltage

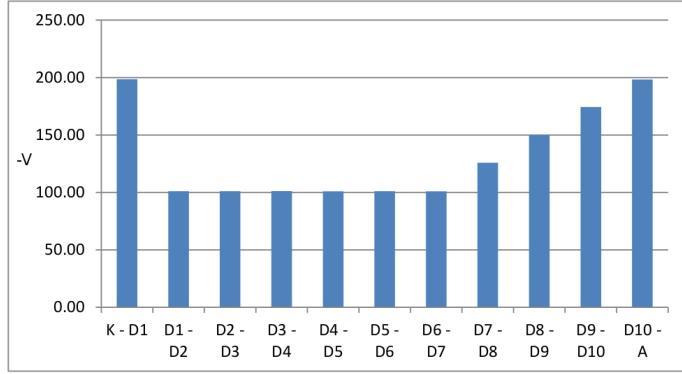


Figure 3.6: The voltage division between subsequent stages for the original PMT base design when supplied with -1500 V.

of -1500 V, the expected standing (bleeder) current even when sitting in the dark is expected to be 1.1 mA. Already here, we see that this design parameter alone suggests its ability to withstand $\sim 3x$ more signal current as compared to the original base.

The introduction of MOSFET transistors is seen between each stage from D7 to D10. Sending the current in parallel over a $1\text{ M}\Omega$ resistors allows the gate of the transistor to measure the voltage without drawing much current. The Zener dynodes are there in place before each transistor gate to ensure that current only goes one way across the sensitive gate channels.

It is also notable that there are banks of capacitors in parallel across the higher dynode stages. Since there is higher current through this circuit than the original base under the same voltage, there will be a greater demand placed on the capacitors to resupply the dynode stages with spent charge. The two 10 nF capacitors in parallel across each stage (which amount to 20 nF total) is significantly greater than the capacitance across the stages of the original base.

3.2.3 Prototype Base v2

The first modification made to the prototype board was to keep everything identical except for the total resistance of the circuit. This was accomplished by halving the resistance of each of the first six stages (R6-R13 on Fig. 3.7) from to increase the bleeder current. The resulting current of the base at -1500 V is at around 2.2 mA. The voltage division retained the same values as described in Fig. 3.8.

3.2.4 Prototype Base v3

In the case that destabilization was occurring prior to the dynode stages with the added capacitors and transistors, the third prototype was decided to take the Prototype v1 design and add more “transistorized”

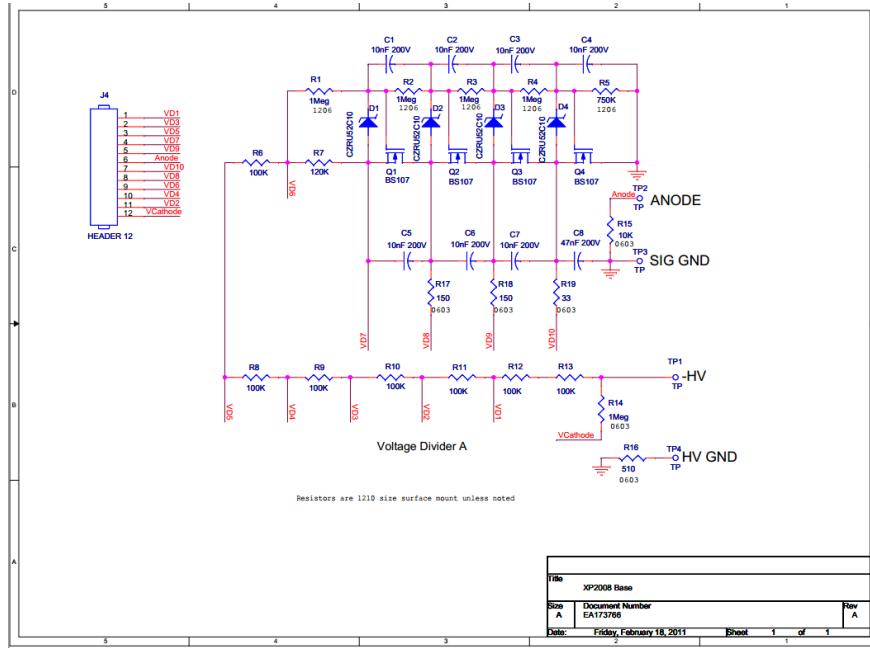


Figure 3.7: The Prototype v1 board circuit diagram received from Fermilab Particle Physics Division [59]. The parts are denoted as R: resistor, C: capacitor, Q: MOSFET transistor, D: Zener diode.

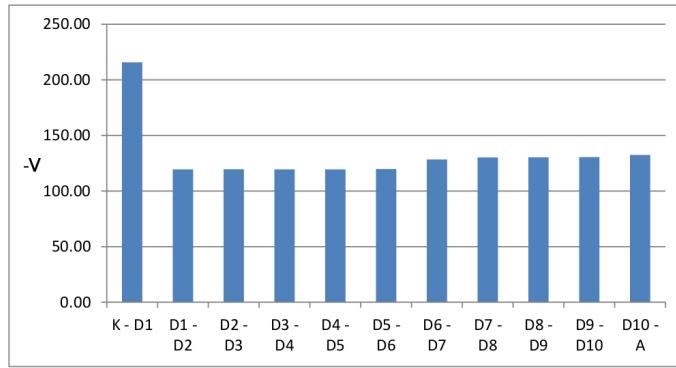


Figure 3.8: The voltage devision between subsequent stages for the Prototypes v1, v2, and v3 PMT base designs when supplied with -1500 V.

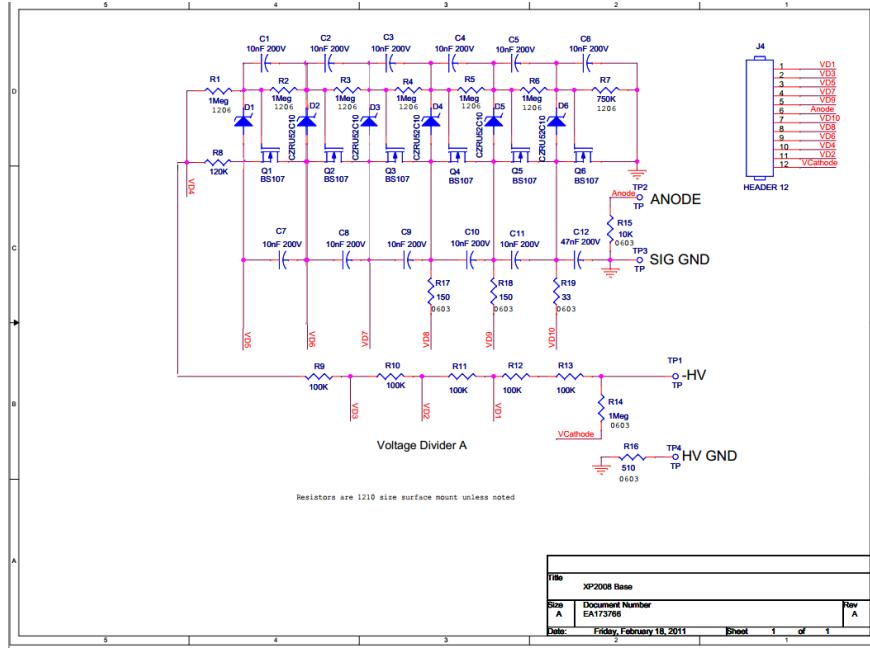


Figure 3.9: The Prototype v3 board: 3 more transistorized stages than the Prototype v1 design.

stages earlier on. This entailed extending the parallel configuration of capacitors, transistor, diode, and $1 M\Omega$ to the D5-D6 and D6-D7 stages. This prototype configuration can be seen in Fig. 3.9. The voltage division was not significantly altered by this change and remained relatively the same as what is described in Fig. 3.8.

3.2.5 Prototype Base v4

The final modification arose from a suggestion from a Fermilab collaborator. It was suggested that it may significantly extend the dynamic range of the tube/base by increasing the voltage drop in, specifically, the last stage relative to the other stages. This reduced to simply replacing R5 (of Fig. 3.7), a $1 M\Omega$ resistor, with a $1.5 M\Omega$. All of the rest would remain unchanged from Prototype v1. The premise for this modification was in the case that the final batch of electrons needed help being “swept” to the anode with a higher voltage difference. The change applied resulted in the voltage distribution according to Fig. 3.10.

3.3 PMT Base Comparisons

There is a specific difficulty with the objective to increase the rate capability of our PMT's. This difficulty is that there was not a known instantaneous intensity or target rate capability to attain. The intensity that caused the original PMT performance to sag is unknown, and if it was known, it would be difficult to

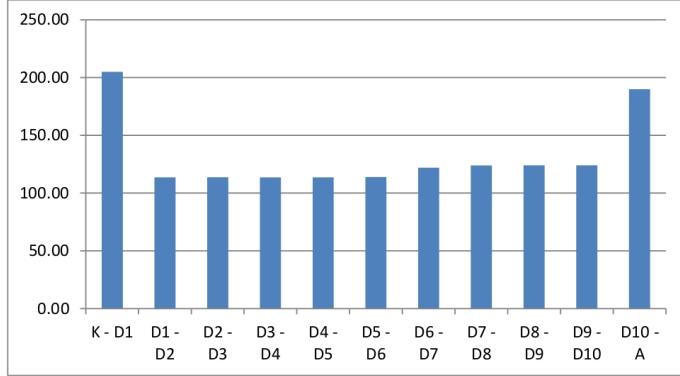


Figure 3.10: The (negative) voltage between subsequent stages for the Prototype v4 PMT base.

match the intensity with an experimental setup. As a result, the objective of these tests was to *compare* the performance of the same PMT under controlled conditions using the original and various prototype bases.

Due to the effects of using different PMTs and due to temperature and humidity fluctuations, the PMT behavior can be somewhat variable from test to test. For this reason, one can only reasonably compare results within each base comparison test, and not across different comparison tests. Each was performed on different days, and sometimes with different PMTs.

3.3.1 Testing Apparatus, Measurements, and Procedure

In this experiment, a PMT attached to a PMT base was placed into a light-tight box facing a fast-pulsing 470nm wavelength LED, which provides the driving photonic signal.

PMT Test Setup

The LED light source assembly consisted of a machined aluminum block that housed the LED, which was a fast-pulsing, narrow beam, water-clear blue LED, peaking in the $\lambda = 450$ nm region. This block had a sliding aluminum insert which had a 1/4" depression used for inserting a 1" diameter NDF of arbitrary optical density [9]. The LED, according to specifications, has a N ns rise time and a N ns fall time, which is important when considering pulsing frequencies on the order of 30 MHz.

The LED itself was driven by an Agilent 33520 Function / Arbitrary Waveform Generator [4] capable of generating signals up to 30MHz. The LED's unattenuated intensity was far too much light to perform any useful test with such photosensitive hardware. Its intensity was attenuated by use of a neutral density filter (NDF), with a rating $D = 3.0$, where the NDF allows 1 in 10^D photons through (1 in 1000 for $D = 3.0$). When testing began, a $D = 4.0$ NDF was used, but the amount of light reaching the PMT was so small that no decrease in PMT performance for any of the bases was observed all the way up to the highest LED

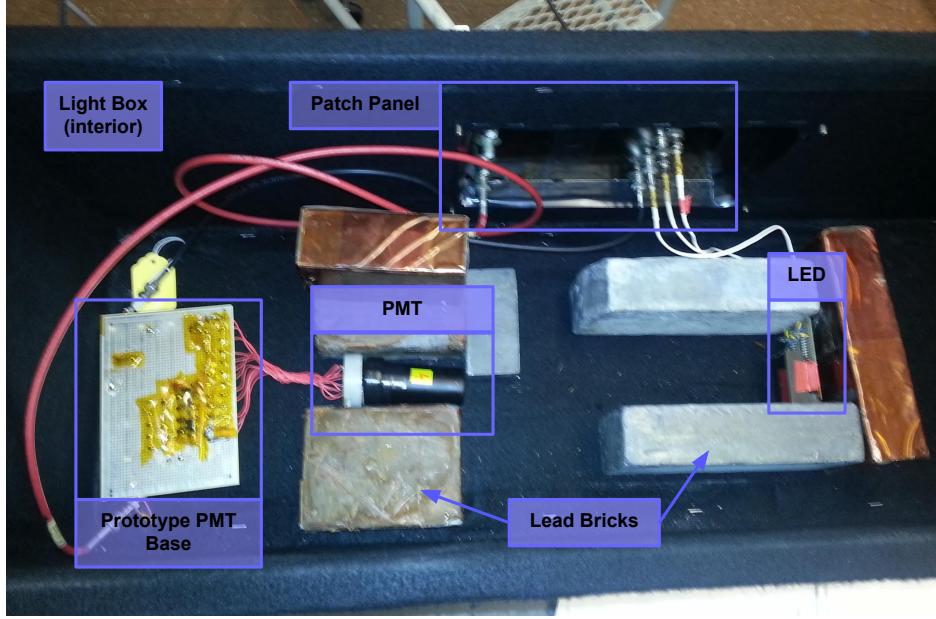


Figure 3.11: Inside of the light box, with the prototype board (left) wired up to a Philips XP-2008 PMT (middle), facing a fast-pulsing LED source (right).

frequencies. As such, the $D = 3.0$ NDF was chosen for this study.

These were all kept within the light-tight box¹ loaned to the SeaQuest collaboration from the Daya Bay group. This light box was equipped with a patch panel that had both BNC and HVBNC connectors by which to power the PMT base and read out its signal. The light box and all of the installed components can be seen in Fig. 3.11.

A simple data acquisition was assembled from an amplifier, discriminator, and scaler in order to observe that the PMT was functioning properly and firing off at the rate that the LED was set to pulse at.

Measured and Calculated Quantities

The PMT base was powered by a high voltage supply, and an ammeter was connected between the two in order to measure the amount of current drawn, or bleeder current, from the HV power supply. The PMT signal was processed by an oscilloscope, averaging the pulse over 300 pulses. The primary measurement was measuring the area of the averaged pulse ($V \cdot s$).

We calculated the signal current from the anode as:

$$Q_{pulse} = \frac{\int V dt}{R} \quad (3.1)$$

$$I_{signal} = f Q_{pulse} \quad (3.2)$$

¹These light-tight boxes are paradoxically referred to as both “light boxes” and “dark boxes”, as they’re used for testing *light-sensitive* equipment and they’re made to be very *dark* inside.

where f here is the driving frequency of the pulsing LED, R is the termination resistance of the signal ($50\ \Omega$), and $\int Vdt$ is the integrated area of the averaged PMT pulse. The amplitude of the pulses was also measured, as it was important in determining if it were feasible to remove the typically noisy hodoscope amplifiers from the SeaQuest stations 1 and 2 DAQ setup.

The measured and calculated quantities of interest are the bleeder current, the averaged signal amplitude, averaged signal area, and signal current over the anode.

Due to the limitations of the instruments used in these tests, there is some systematic uncertainty to consider. The ammeter used to measure the bleeder current between the HV supply and the PMT base used an analog gauge that had closely spaced ticks at 0.2 mA intervals. As such, the systematic uncertainty for the bleeder current is considered to be $\pm 0.05\text{ mA}$. Also, due to pulse-to-pulse variations which cause a mild amount of smearing in the average, the measurement of the amplitudes and signal areas have an uncertainty of $\pm 5\text{ mV}$ and $\pm 0.1\text{ nVs}$, respectively. The uncertainty in the signal area translates to an uncertainty in the calculated signal current of around $\pm 0.04\text{ mA}$. These fluctuations are considered negligible with respect to the quantities measured.

Test Procedure

Most of the prototyping was performed over the span of two weeks. As new suggestions were made at base improvements and the iterations v1-v4 progressed, more comparison tests were performed. The procedure for the each test was the same and is as follows:

1. Connect the base to the PMT and to the HVBNC and BNC connectors.
2. Close the lid to the light box.
3. Gradually power the PMT (-250 V every 5 s) until it is as -1500 V.
4. Watch the scaler counter on the DAQ to see if it is rapidly counting. No (or very few) counts on the scaler. means the box is light-tight.
5. Let the PMT and base warm up to near-operational temperatures with its standing current for a minute or two.
6. Power on the function generator and LED, setting the function generator to create a 16 ns square wave pulse to the LED circuit (the minimum achievable pulse width [4]) at a frequency of 10 Hz.
7. For each of several frequencies from 10 Hz to 30 MHz, perform the following:
 - (a) Observe the scaler to ensure that the count is increasing at the same rate as is set by the function generator. This confirms that both the PMT, base, function generator, and LED are generally working as planned.

- (b) Trigger on the signal from the PMT output.
 - (c) Average the signal over 300 triggered pulses and then freeze the frame.
 - (d) Record the current draw from the HV supply (bleeder current).
 - (e) Measure and record the pulse amplitude with the oscilloscope cursor.
 - (f) Measure the pulse area using the signal integration feature, setting vertical cursors at the zero-intercept on each side of the pulse.
 - (g) Return to triggering mode and increase the function generator frequency as warranted.
8. Power off the function generator and LED circuit.
 9. Step the PMT voltage down slowly to 0 V.
 10. Open the light box and replace the base for subsequent tests as necessary, repeating the steps above.

3.3.2 Original vs. Prototype v1

The first test was to compare the original base with the first iteration of the prototype base. It was generally expected that this new design with more modern technology would exceed the performance of the original base, but the question was twofold: to what extent does it perform better and what is the baseline by which to compare future improvements?

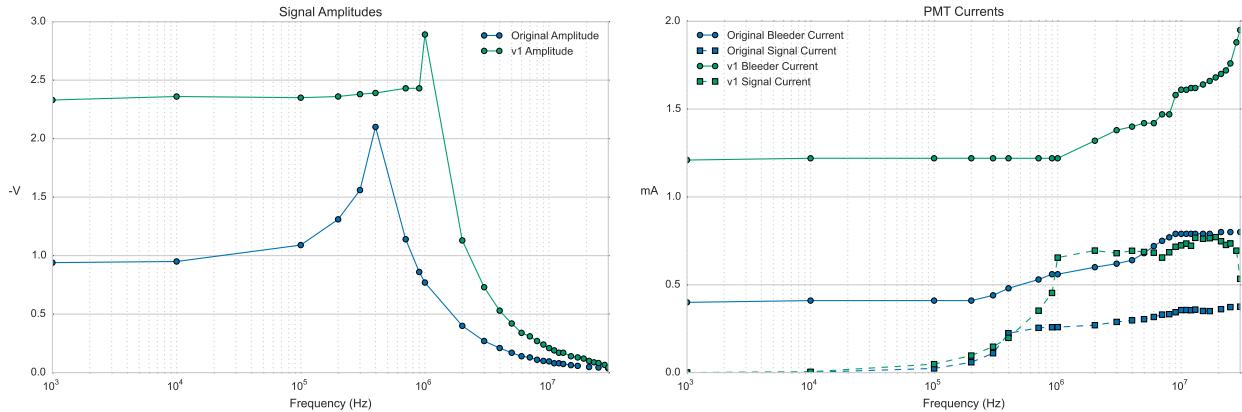


Figure 3.12: Measurements of the (a) signal amplitudes and (b) bleeder and signal currents in the original and prototype v1 PMT bases. The HV was inadvertently set to -1600V V for this test instead of the intended -1500 V.

The first important observation is the baseline amplitudes of the signals from the two bases, as seen in Fig. 3.12(a). While the absolute voltage itself is not as useful since the LED doesn't necessarily reflect actual operating conditions, we do see that the amplitude from the new base design is increased by a factor of more than two. This in itself is significant in when considering the limits of the discriminators that were used in the hodoscope arrays for Stations 1 and 2. With even a modest boost to gain and signal amplitude,

the PMT signals should be large enough that an amplifier is not required in order to get a discriminator to trigger on a PMT pulse. Seeing as the amplifiers used had often caused troubles with the amount of electronic noise they tended to imbue, this was an encouraging observation.

Next, we look at the behavior of the amplitudes as the frequency progresses. It is safe to assume that the unchanging behavior in the lower frequencies indicate that the voltage divider is stable and that there is no voltage breakdown yet. At higher frequencies, there is a qualitative peak in amplitude for both bases, followed by a steep decline. This rise and then sharp fall would seem to indicate the end of stable PMT base operation, with the fall indicating the onset of degradation of performance.

To better identify the conditions that cause this onset, we look to currents applicable to both bases (Fig. 3.12(b)). In this particular test, we see the original base peak in amplitude at ~ 4 MHz and prototype base v1 peak at ~ 10 MHz. Looking to the currents, we see that these particular points mark a certain point for both where the signal current reaches almost exactly 50% of the bleeder current. The baseline improvement with this base over the original one was estimated to extend the operational range by a factor of three – the point where the amplitude drops below the low-frequency baseline.

Around the time of these tests, Sten Hansen conducted a SPICE (Simulation Program with Integrated Circuit Emphasis) simulation of the prototype circuit and was able to reproduce what was observed. In particular, it was observed that the gain began to change dramatically when the anode (signal) current reaches a significant fraction of the total current drawn from the HV supply; a current that is much less than 100%. It was concluded that 50% was consistent with his findings.

3.3.3 Original vs. Prototype v1 vs. Prototype v2

Several suggestions were made in altering the prototype base in order to maximize this improvement, including lowering the overall resistance of the circuit. If the signal current grows at the same pace, then by increasing the bleeder current, there is a chance that a higher frequency signal will be required to bring the signal current up to the point where the divider breaks down, which we estimated above to be $\sim 50\%$ of the increased bleeder current.

Looking to Figure 3.13(a), we see that this basic premise is technically sound. The prototype v2 base exceeds the dynamic range of both the original base and the prototype v1 base by factors of approximately a factor of four and two, respectively. This is a significant improvement, and this, at first, seems to be an avenue to pursue. However, an issue arose after increasing the LED frequency from 13 MHz to 14 MHz when a resistor burned out. The test continued with the rest of the bases, but the results paint a clear picture of what happened.

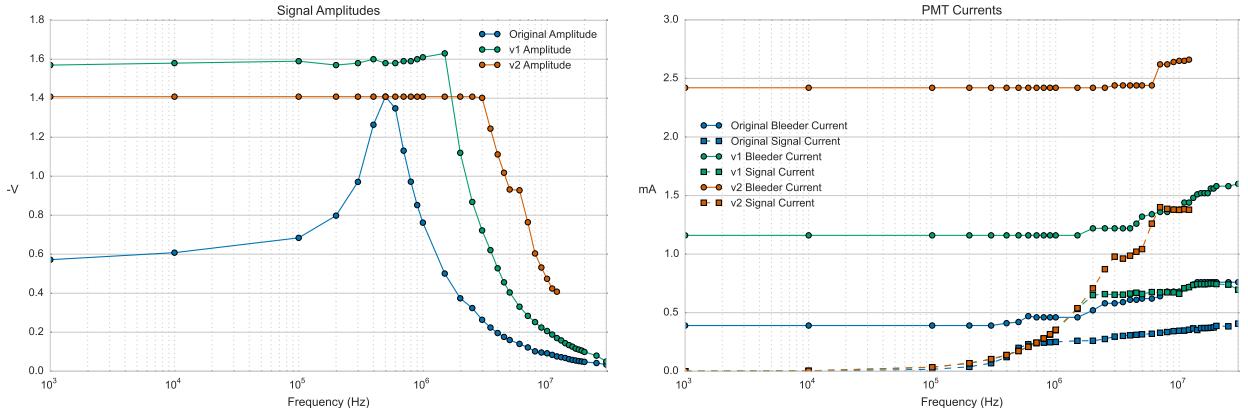


Figure 3.13: Measurements of the (a) signal amplitudes and (b) bleeder and signal currents in the original, prototype v1, and prototype v2 PMT bases.

By halving the resistance, the bleeder current doubled (Figure 3.13(b)). This caused the power being dissipated by the resistors to exceed their power rating. The prototype bases were built with 0.25 W resistors, and as the bleeder current rose well above 2.5 mA, one of the resistors failed. This did not bode well for how a smaller on-board resistor might fare with the same current.

To further exacerbate matters, it was found that the LeCroy 1440 HV supplies that were to be used could only supply a maximum of 2.5 mA per channel [77]. Seeing as these power supplies had proven themselves sensitive and prone to communication and performance failures, it was deemed inadvisable to demand anywhere near the maximum current from these HV supplies.

For these reasons, it was decided that the new PMT base design of prototype v1 not be modified in this manner which might suffer from overheating, component failure, and a too-high current demand.

3.3.4 Prototype v1 vs. Prototype v3

Perhaps it was the case that the “transistorized” steps were doing their job well and the destabilization was occurring at prior stages. Since the only downside of adding additional steps was added cost, a base with more transistorized stages was created as depicted in Fig. 3.9.

As is quite evident in Figure 3.14, there was no significant difference in the performance of the two prototype bases. As such, this modification was not adopted into the final design.

3.3.5 Prototype v1 vs. Prototype v4

The final test modification of increasing the resistance and thereby the voltage drop across the final D10-A stage was tested against the prototype v1 base.

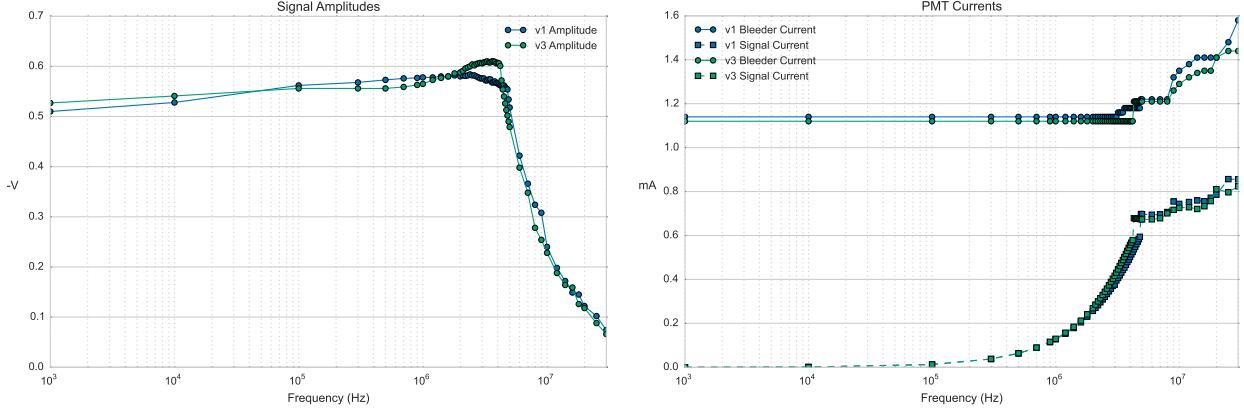


Figure 3.14: Measurements of the (a) signal amplitudes and (b) bleeder and signal currents in the prototype v1 and prototype v3 PMT bases.

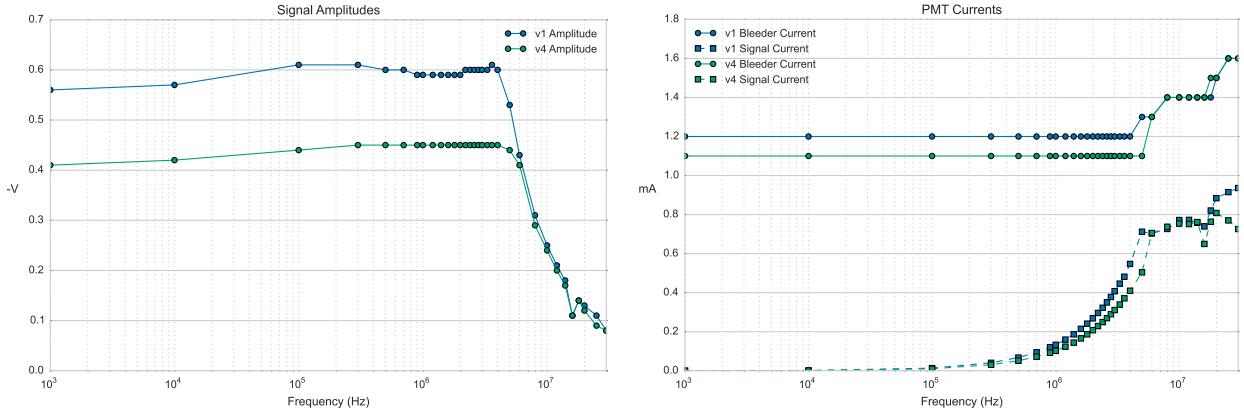


Figure 3.15: Measurements of the (a) signal amplitudes and (b) bleeder and signal currents in the prototype v1 and prototype v4 PMT bases.

This modification turns the voltage divider slightly away from the “optimum gain” configuration and more towards the “optimal timing / linearity compromise” configuration. This is seen to have a direct effect on how a loss in gain can be observed between the v1 and v4 base in Fig. 3.15(a). While the rate capability does seem to slightly increase from 4 MHz to 5 MHz, the value of having a higher gain was deemed more desirable. Higher gain, again, means that the then-used noisy amplifiers would no longer be needed, and the marginal loss in rate capability was considered a reasonable tradeoff.

3.3.6 Base Comparison Conclusions

After these tests were conducted and the information regarding each of the bases’ performances was evaluated, it was decided that the prototype v1 base design was the optimal base configuration to move forward with. This base design yielded a clear improvement in both signal amplitude ($\sim 2x$) and rate capability

($\sim 3x$) at the expense of more complex circuitry and an increased supply current draw. It then became the task to fit this circuit onto small printed circuit boards (PCBs), and then manufacture, install, and test them.

3.4 Base Manufacturing and Installation

It was decided that the original base structure was well-packaged and well-shielded with its μ -metal canister. As such, a two-component board was designed by Sten Hansen that consisted of a 2.5" diameter board and a small rectangular "daughter board". By June of 2013, the mother and daughter boards (attached) had been manufactured and delivered to UIUC for assembly into the rest of the base.

Since the sockets that coupled the board to the PMT pins was not commercially available, they had to be salvaged by deconstructing the original bases. This consisted of cutting the components between the two boards of the original base and then using a solder vacuum to separate the board from the socket pins. The socket and the other structural elements of the PMT base were separated and prepared for construction with the new base PCB (Fig. 3.16).



Figure 3.16: (Left) The components of the disassembled original base and the manufactured new base PCB with attached daughter board. (Right) The original and new bases fully assembled, side-by-side.

The deconstruction and re-soldering of the bases took place over the span of about three weeks. Each base was tested to provide the same voltage division across stages and was tested in operation inside the light box apparatus by observing its signal characteristics and the rate of the scaler counting (which should match up with the LED frequency).

After all bases were constructed and tested, the final task was to find the optimal operating voltage for a PMT assembly that would be suitable for detecting minimum ionizing particles passing through a

scintillator paddle similar to the ones used at Stations 1 & 2 at SeaQuest. We also want the minimum possible voltage, as higher voltages correspond to increased noise. This was performed by assembling three stacked scintillator paddles, each with a PMT attached to their light guide. Their outputs were passed through a discriminator and then to a coincidence unit, which in turn had its output sent to a scaler counter. The top and bottom PMTs were powered to -1500 V, and their coincidences were counted on one channel of the scaler counter. Simultaneously, the middle paddle, set to a variable voltage setting, had its coincidences with the top and bottom paddles counted on the other channel of the scaler counter. The ratio of these two counts ($\epsilon = (\text{triple coincidence}) / (\text{top} - \text{bottom coincidence})$) described the efficiency of the middle hodoscope. The voltage of the middle PMT is varied until a plateau in the efficiency is identified and recorded to be the ‘nominal voltage’ of this reference PMT.

All of the constructed PMT bases and the “reference” PMT and base were then sent up from UIUC to FNAL to be installed into Stations 1 & 2 and used for “gain matching”. The gain matching is done by exposing a scintillator paddle to a γ -ray source while connected to a PMT and adjusting its voltage until it produces the same signal amplification, or *gain*, as the rest of the PMTs. The original PMT that the rest are matched to is the reference PMT from the previous efficiency study.

Matching the gain is not a simple task, and the typical approach is to take advantage of a physical phenomenon called the Compton Effect occurs in scintillator material when a γ -ray scatters off of it. Often times, the γ -ray of energy E will not deposit the whole of its energy, but only a portion of it. The spectrum of the photons that scatter off will have energy

$$E' = \frac{E}{1 + \frac{(1-\cos\theta)E}{m_e c^2}} \quad (3.3)$$

where θ is the angle of deflection and m_e is the electron mass. This, in turn, deposits transferred energy $E_T = E - E'$ into the scintillator which has a maximum value of

$$(E_T)_{max} = E \left(1 - \frac{1}{1 - \frac{2E}{m_e c^2}} \right) \quad (3.4)$$

This energy is known as the Compton Edge, as it is characterized by a very sharp falloff in the energy spectrum of energy deposited by a source of a specific energy. The particular energy of this edge is not important for the purpose of gain matching. Instead, it’s used as a qualitatively sharp feature that can be clearly identified when examining the energy spectrum of photons in scintillator material, as seen in Figure 3.17. As such, it can be used as a landmark to use for adjusting the voltages of each PMT in a way such that the gain of each is the same.

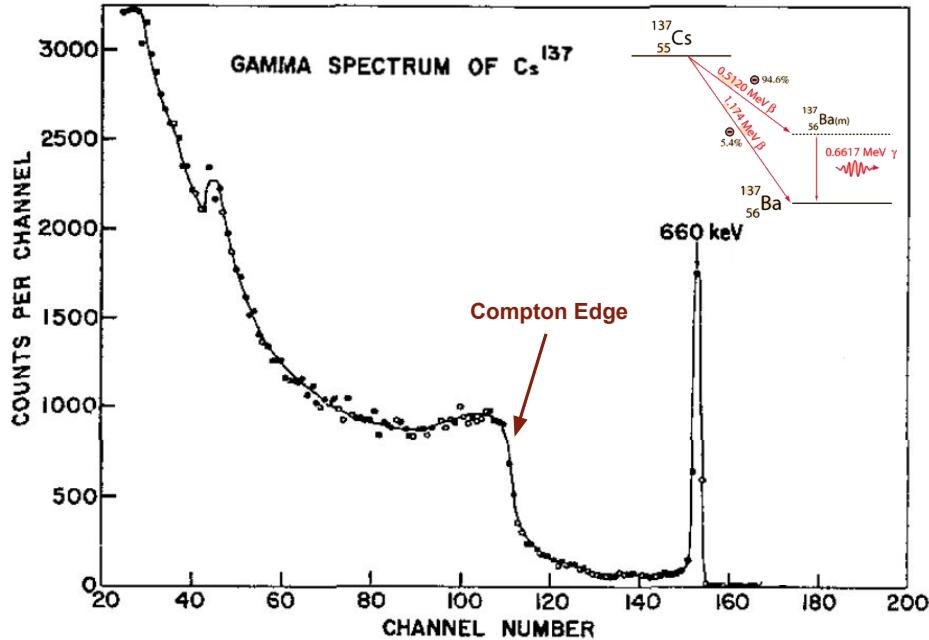


Figure 3.17: The γ -decay energy spectrum of Cs-137 as seen through a scintillator-PMT-qVt setup [107]. Here, the Compton edge can be seen at channel 110.

Once installed onto a scintillator paddle, the reference PMT was powered to nominal voltage and a radioactive Cs-137 source was pointed at the paddle. Since the PMT signal has some *linearity* to its signal, the amount of energy of incoming photons will translate approximately linearly to the amount of signal out of the PMT. The signal from the PMT (with the Cesium source pointing at the scintillator paddle) is passed through a qVt (charge-voltage-time) module in “Q-mode” to histogram the total integrated charges of the PMT signals. The qVt’s channel location of the Compton edge for this reference PMT at this nominal voltage is recorded. Then, every other PMT is powered and measured through the qVt with the Cesium source, and its voltage is adjusted until the Compton edge falls on the same channel. This procedure is generally known as “gain matching”.

3.5 New Base Performance

All of this work is all for naught if the new bases do not hold up to even the highest of beam intensities at SeaQuest. After all of the aforementioned testing, installation, and gain matching, data taking resumed. In Figure 3.18, the hit per event distributions are shown for x - and y -measuring hodoscopes as for different intensity ranges. It is difficult to compare these plots to those in Fig. 3.1, as there was no measure of per-event intensity at the time, and the beam back when that measurement was made was less intense (though

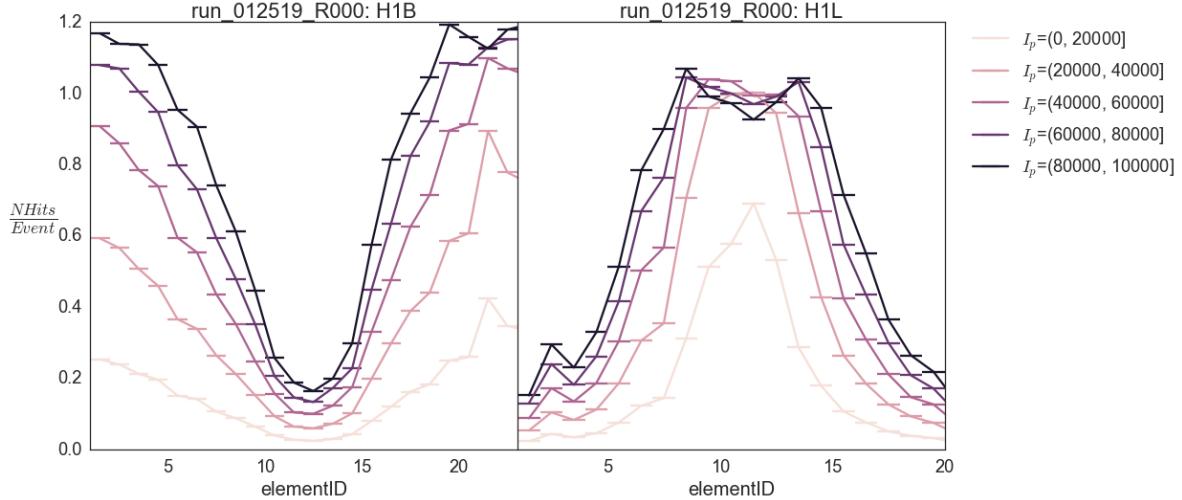


Figure 3.18: The number of hits per event for the Station 1 hodoscopes for several intensity ranges. Shown is (left) the top x -measuring array and (right) the left y -measuring array. I_p is in terms of number of incident protons for the event.

just as non-uniform).

Regardless, the new readouts can be interpreted positively. One can see in the y -measuring array that the shape does flatten out at the peak for higher intensities. It can be noted that this flattening is much less severe than the “sag” that was initially observed. More importantly, when observing the absolute shape of the distribution of the x -measuring array, the shape remains intact and uniformly increasing all the way up to the highest intensities. Seeing as the x -hodoscopes are the only arrays that are currently being used for triggering and analysis, this is the most important behavior to observe.

Along with this confirmation, the amplitude output by the PMT signals was above the minimum threshold of the discriminators being used. As a result, the noisy amplifier modules were able to be removed from the signal pipeline and thereby reduced the number of false hits due to electronic noise.

Chapter 4

Data Productions

Data analysis is a highly dynamic activity - particularly for a nascent, large-scale experiment. As a result, having a flexible and scalable data management and analysis framework is of particular importance. Different data components may be needed at different points of study, new readouts may need to be added in the middle of data taking, and analysis styles themselves may differ from user to user – all of which are either facilitated or impeded depending on the flexibility of the framework implemented.

In an attempt to provide such a framework to its users, SeaQuest makes use of the capabilities of an array of MySQL servers, leveraging relational database technologies. While MySQL and database systems (RDBMS) are nothing new, even to physics experiments, they are not widely adopted among the physics community. Often, in the cases that they are, it is to ineffective or inefficient result.

During my tenure at SeaQuest, I was primarily in charge of the effort to decode the raw, serialized data from the data acquisition systems and store them in a manner in which it could be readily accessible by collaborators worldwide. Beyond that, emphasis was placed on curating the tracked physics data and preparing it in such a way that facilitates quick and simple analysis from many different analytical software packages.

In this chapter, the details regarding the processes and formats that define the online and offline data productions will be described in some technical detail. The technologies used to perform this were a combination of compiled C/C++ codes, MySQL, ROOT, and the Fermilab Computing Division’s primary computing grid. The end product provided to all collaborators was in the form of unified “schemas” (or, groups of tables) possessing all of the combined *run*, *spill*, *event*, *slow control*, and *tracked* information, along with specific data quality flags in place to assist in the selection of good, analyzable data. Finally, there is a retrospective discussion on the scalability and best practices, followed by speculation on how future data management technologies might be used at SeaQuest and other experiments.

4.1 Raw Data Processing

The three raw outputs of the data acquisition systems, as described above (and in Chapter 2) are Main DAQ CODA files, Scaler DAQ CODA files, and Beam DAQ ASCII files. Each raw data file corresponds to the data taken from certain subsystems over approximately one to two hours of running time. These three types require varying degrees of de-serialization, parsing, processing, and storage – a process as a whole defined as *decoding*.

All raw data files are backed up to long-term tape storage (managed by FNAL Computing Division), and the decoded and processed data gets stored on one of four MySQL servers to be used for analysis by the collaboration. Data is also output to ROOT files for ease of use by independent tracking programs.

4.1.1 CODA Event Format

A single CODA file can be described as being a chain of events. These can generally be divided into “CODA Events” for those related to the CODA file itself, and “Standard Physics Events” when containing experimental data read out by CODA and its subsystems. Each event can be represented by an array of unsigned 32-bit integers. When using the CODA EVIO (event input/output), one whole event is read out into such an array, which can vary in size from 6 to $\sim 50,000$ integers long, depending on the type of content.

CODA Events primarily correspond to the actions taken by a shift-taker using the *rcgui* application on the data acquisition computer. The types of events that fall under this category are: Sync, Prestart, Go, Pause, and End. It should be noted that the Sync and Pause events do not generally get used at SeaQuest. These provide markers in the CODA file of what commands were given to CODA through the course of taking data.

Before data taking can begin The readout controllers and the trigger supervisor need to be loaded with a designated firmware in order to operate properly. Before a CODA file is started and created, this “Download” is performed, and then CODA is ready to transition from the downloaded state to the prestarted state. The CODA file is then initialized with a prestart event, whose format can be seen in Figure 4.1. The most important information stored in the prestart event is the run number, which is used universally through the experiment.

Once in a prestarted state, a “Go” action is initiated to transition it into data taking mode. This action generates a Go event. If ever paused, a Go event would be created once data taking was resumed. Go events contain the current time and the number of events in the run so far, also seen in Figure 4.1. It should be noted that the “time” stored in these events is not trusted, as the data acquisition computers are not synced to an official external source.

event_length = 4		
event_type = 17	1	0xCC
time		
run number		
run type		

event_length = 4		
event_type = 18	1	0xCC
time		
(reserved)		
number of events in run thus far		

Figure 4.1: (Left) The Prestart event, with event type 17; (Right) The Go event, with event type 18[100].

event_length		
event_type	0x10	0xCC
data bank		
data bank		
...		

fragment_length			
(code & status)	ROC#	1	counter
data words			

Figure 4.2: (Left) A standard physics event, with various event types, filled with a series of data banks; (Right) The ROC data bank format[100].

The last CODA Event of importance is the End Event. This contains the number of events in the run, but primarily serves as a marker to any programs reading the CODA file that the EOF (end of file) has been reached, and to commence any post-run processes.

Most of the events in the CODA file, however, will be “Standard Physics Events”. The contents of these can be varied, but can be characterized generally as having a header containing the total length of the event (in number of 32-bit integer “data words”), followed by several “data banks”, which are typically filled with a series of ROC outputs (Fig. 4.2). The contents and formats of the ROC data banks are completely controlled by the ROC programming, and whatever set of 32-bit words the ROC has to read out, it will store in its “data words” block.

4.1.2 ROC Data Bank Formats

4.1.3 Decoding Raw Data

The CODA file decoding process is nearly identical in the cases of the Main DAQ and Scaler DAQ outputs, and only differ by content; the Scaler DAQ reads out specific scaler data, and the MainDAQ contains TDC readouts of the detectors. For each one to two hour *Run*, the CODA files can be well-described as the following sequence of events (and the data they contain):

1. Prestart Event (Run data)
 2. Begin Spill Event (Spill data, Scaler readout)
 3. Many Physics Events (Event data, TDC readout)
 4. End Spill Event (Spill data, Scaler readout)
 5. SlowControl Event (Slow control readout, Spill ID readout)
 6. Spill Counter Event (Spill ID readout)
- ...(Repeat 2-6 for each *Spill*)

The decoding programs use C and C++ in conjunction with Jefferson Lab's CODA I/O library[100] to read these events and parse them according to their individual formats. Data from these CODA events are decoded and placed into hierarchical categories.

Run Level Data

Run-level data contains data and metadata pertaining to the entirety of the run that is recorded. At the time of the Prestart Event, the date and time of the run are stored, along with a readout of the specific settings of all non-trigger TDC boards. Settings regarding which triggers are enabled and what each of their prescale factors is also stored in the run-level data.

After the End Event is encountered, metadata is aggregated and stored regarding such items as the number of chamber hits, the types of triggers that were fired, the target positions used, average magnet currents, and other useful metrics. All of these fields of settings and aggregated values can be used later on to quickly determine if a run is analyzable, or if it has characteristics that are desirable for specific niche analyses.

Spill Level Data

The *Beginning of Spill* (BOS) and *End of Spill* (EOS) events bookend the set of physics events for a given spill. At each BOS and EOS events, the 140 MHz VME scalers are read out. At the beginning of the spill, all scalers are zeroed out, and then read out again after the spill has ended.

Slow Control events are read out between spills, which contain data regarding the current spill identifier number, target systems, beam and radiation monitors, and environmental readings.

The spill identifier (*spillID*) is what is used to synchronize the data together across various data acquisition systems. As such, the *spillID* is read out redundantly in both Slow Control and Spill Counter events (which contain only the *spillID* value) to ensure that the data is appropriately labeled.

When the End Event (for the whole CODA file) is reached, the independently-recorded Beam DAQ data (recorded in an ASCII file) is read and stored with the rest of the Spill-level data.

Event Level Data

For each spill, $\sim 3k$ events are triggered to be recorded, though this number can vary greatly with the particular beam settings of a run. With each event, three types of information is stored: the trigger which fired the event, a measure of the beam intensity per RF bucket, and the full detector readouts. The detector readouts require the most processing of all the rest of the data. The CODA files contain the hardware addresses of each detector *hit*, along with a *TDC time*. The following steps briefly summarize the processing steps:

1. Mapping: Map the hardware address to a detector name and detector element number
2. Timing: Classify hits as in-time or not and calculate *drift time* from TDC time
3. R-T (time-to-space): Translate *drift time* to *drift distance*
4. After-Pulse Elimination: Remove hits that result from signal reflection and other electronic artefacts
5. Trigger Road Reconstruction: Use *v1495* TDC hits to reconstruct possible trigger roads that may have fired
6. Hodoscope Masking: Remove drift chamber hits that have no adjacent hodoscope hit
7. Trigger Road Masking: Same as hodoscope masking, but only using hodoscopes from reconstructed trigger roads

This fully processed data is then stored into one the experiment's MySQL databases and/or a ROOT file.

4.2 Online and Offline Processing

There are two running modes of raw data processing: on-line and off-line. For on-line mode of productions, all Run- and Spill-level data is decoded, but only 1-in-*n* Physics Events are processed, where *n* is typically between 3 and 15, depending on the event data taking rates. This “*sampling mode*” is used in order for the decoding to reliably keep up with even high-intensity beam data. This data pipeline is important for on-line

monitoring of the experiment by shift crews. With this up-to-the-minute stream of data, any abnormalities in duty factor, wire maps, et al. can be observed directly and acted upon quickly.

For off-line, or “batch”, productions, a large group of categorically similar runs is defined, and the chain of production processing is initiated. The steps of this process is generally decoding, tracking, archiving, and merging. The logical blocks of runs are typically defined by “roadset”, or, the version of the set of trigger roads used in the L1 trigger matrix. Differences in these trigger matrices can affect the distributions to a degree commensurate with the difference in the set of roads that are triggered on by the FPGA-1 trigger. As such, each block of data using a given roadset should be analyzed individually and then have its results combined with those of other roadsets via a weighted average.

The decoding and tracking is performed on Fermilab Computing Service’s FermiGrid, which provides the computing resources necessary to process hundreds of runs simultaneously. The tracking has also expanded its capabilities to run on the Open Science Grid, which is a grid computing network that pools together computational resources from national laboratories and universities throughout the world[29]. This allows the intensive reconstruction algorithms to perform its tracking on thousands of grid nodes at once.

A single decoding job submission will output the processed data to one of the four available MySQL servers and also to a ROOT file. The processed data in ROOT form at this stage is called the “digit” data. Whether or not the decoding sends Hit and TriggerHit information to the MySQL schemas is an option that can be set at the time of batch processing. The current operating mode is to perform batch processing, sending Hit/TriggerHit data to MySQL for only 1-in-50 runs. This allows users to capture a glimpse of full hit distributions for runs within a roadset while preventing a single pass of the experiment’s data from filling up all available data storage space.

After this step, jobs will be submitted to run one or both of the two tracking programs based on the ROOT file and/or the MySQL data. Once the tracking is completed, the ROOT file is archived on the Fermilab BlueArc NAS backup system for storage and for tracking.

Upon the completion of decoding and tracking of a specified range of runs, all of their Run-, Spill-, and Event-level data, along with its tracked data, is combined into a single *merged* schema. These *merged* schemas are mirrored across all four of the MySQL servers for optimal redundancy and availability.

4.3 RDBMS Data Structure

Add pros/cons. Discuss motivation.

The processed data is primarily stored in MySQL Server 5.1 databases. MySQL is an open-source

Relational Database Management System (RDBMS) developed by Oracle that is well-suited for the storage and responsive querying of hierarchical data.

Each run is decoded into its own schema, and contains its own instances of all tables of a specified design. The tables are all *join*-able to each other by sharing *foreign keys* with each other in the form of the *runID*'s, *spillID*'s, and *eventID*'s. The contents of the tables are *indexed* in such a way that *joins* and queries gain a speed performance boost, but this comes at the cost of disk space.

The data on the server is world-wide accessible and can be queried using the standard querying language. The queried data can be directed to any analysis code in any programming language due to the large array of MySQL API's available.

4.3.1 Data Schemas

4.3.2 Schema Table Organization

4.4 Data Quality

Talk about dataQuality bits, the levels they're used at, and how to use them in selecting data.

4.5 Discussion

4.5.1 Retrospective on Scalability

4.5.2 Key Practices for Success

4.5.3 Future Technologies

Chapter 5

Analysis

CHAPTER STATUS: IN PROGRESS

5.1 Reconstruction of SeaQuest Dimuons

Once the raw hit and timing information has been gathered and processed by the production level software, the reconstruction software is run on it in order to find tracks that pass through the spectrometer. The ultimate purpose of the detectors in the spectrometer is to identify the 4-momenta of particles that occur as a result of high energy collisions and particle decays. Individual particle tracks and their 3-momenta are identified by tracking their paths through the known magnetic fields. These tracks are then matched up to find possible dimuon vertices. Reconstructed tracks, dimuons, and their constituent hits along with their full defining set of kinematics are provided as output for analysts to use in their studies. This tracking software, known as “kTracker”, was developed and is maintained by Kun Liu of Los Alamos National Laboratory. In this section, an overview of the tracking algorithm will be given in how the kTracker steps through the mass quantities of raw data in order to provide the reconstructed set of data used for SeaQuest’s analyses. Details of data reduction, track finding, track fitting, and vertex fitting will be discussed broadly.

5.1.1 Data Reduction

As the number of hits in the tracking chambers increases, it becomes more and more difficult to successfully perform pattern recognition in identifying real tracks. As a result, a litany of methods are performed in order to reduce the hits registered down to hits that are likely to have come from physical tracks and not unwanted noise and radiation. The sum of these data reduction measures thereby serve to make events that were once un-trackable more likely to be trackable. For events that would already be able to be tracked, these cuts still serve to reduce the run time of the tracking algorithm and curtail the number of erroneous tracks that could be reconstructed purely due to combinatorics and coincidences. These measures can be grouped into two groups: hit removal and event multiplicity cuts.

Timing and Masking

The first of the hit removal measures is the **in-time cut**. When an event is triggered by the trigger supervisor, each TDC issues a “stop time” by which each hit’s timing is measured against to derive its TDC time. Numerous “trigger timing” studies of the detectors’ and the TDCs’ behaviors have identified a certain range of values known as an “in-time window.” Any hits that occur outside of this timing window can be eliminated from tracking, as there is no possibility that the hit could have come from the span of time when the detectors were triggered. It should be noted, as it will later be relevant, that the drift chambers have in time windows as large as ~ 600 ns, which is equivalent to ± 15 beam pulses (RF buckets). This means that this alone will not do a very good job of paring the hits down to the relevant hit data.

Due to the behavior of the electronics of the drift chambers and their communications with the ROCs, there is an *echoing* effect that causes a signal to bounce back and forth along the signal line. The real signals that result from this echoing are called “afterpulses,” and they should not be used for reconstruction. Measures have been taken to reduce this effect in the form of adding ferrite cores at the ends of the cable ribbons connecting the electronic readout boards to the TDC cards, but afterpulses are still not absolutely avoided. Building upon the “in-time” classification, an the **“afterpulse removal”** criteria is defined such that only the first in-time hit for each element of each detector is accepted. One minor flaw with this method of data reduction is that It is possible that the first in-time hit for an element is actually the afterpulse of a hit outside of the in-time window. In these cases, the actual desired in-time hit will be removed.

The hodoscopes and their relatively high-precision timing can be used to eliminate spurious hits with great efficacy by a method called **“hodoscope masking.”** The hodoscopes have two qualities: the speed at which they operate, yielding a precise in-time window, and the fact that they are required to have fired in order for the trigger to have fired. The conclusion that one can arrive at is that if a chamber wire fires off and there is no adjacent hodoscope paddle fired off within the same event, then these chamber hits can be removed. From the detector specifications and the geometry of the experiment, a lookup table can be made associating a hodoscope paddle with a swath of chamber plane elements (and vice versa). This many-to-many relationship is used to eliminate all drift chamber hits that don’t have at lease one associated in-time hodoscope hit.

Taking this even further, there is also **“trigger road hodoscope masking”** that is applied as a superlative of standard hodoscope masking. In this data reduction method, the set of in-time hodoscope hits is matched against the set of trigger roads (H1-H2-H3-H4 paddle combinations) used to trigger on the data was triggered with. First, the finite set of trigger roads that could have fired the trigger are determined. Then, hodoscope masking is performed once more, but the masking considers *only* the hodoscopes that are

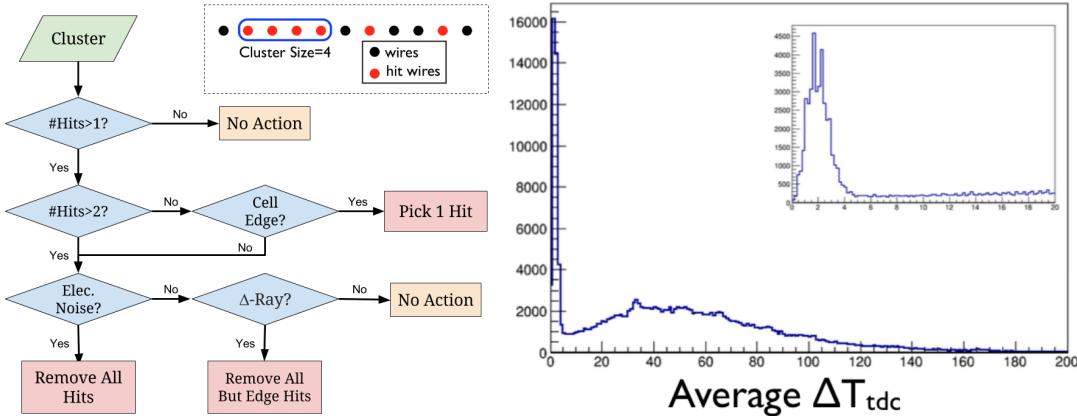


Figure 5.1: (Left) A flow chart mapping out the decisions in how to treat clusters. (Right) A histogram of time differences of hits within a cluster. The peak below 10 ns is from electronic noise, and is the deciding factor in electronic noise classification.

a part of those reconstructed trigger roads.

Cluster Removal

The final measure taken to reduce the number of background hits is the **cluster removal**. A cluster is defined as any grouping of two or more adjacent hits within a drift chamber plane. The focus here is to either determine which wire is the best to use of a cluster or determine if the whole cluster should be removed. To decide on this, the three known sources of clusters are evaluated: electronic noise, cell-edge hits, and delta rays. Specific criteria are imposed in order to determine the cluster type and the action to take, as described by the flowchart in Figure 5.1.

Attached to the drift chambers, there are electronics boards (ASDQ) that relay the signals from the sensory wires to the TDCs. Each of these boards are connected to eight or sixteen wires. Similar to the afterpulse source, and it's possible that several adjacent wires can fire off due to noise in a single board. In these cases, none of these hits are likely due to any actual physical hits and should be completely remove where found. These types of clusters of hits are classified by the criterion of having very similar timing ($\Delta t \leq 10$ ns), which is logically to be expected if noise from the electronics board was at fault for the signals. This cut is applied only to clusters of > 3 hits for stations 1, 2, and 4 detectors. At station 3, there is a marked proclivity for electronic noise, and so the threshold for the number of hits is lowered to > 2 hits while narrowing the timing difference to $\Delta t \leq 8$ ns.

The next type of clusters to classify is “cell-edge” clusters, which is defined as cases where there are two adjacent hits in a single plane where both of them have a drift distance that is greater than 90% than the half of the cell width. In these cases, it is very likely the case that an energetic muon passed right between

two sense wires, and the hit was registered by both wires. In these cases, it is advantageous to reduce the number of hits and thereby reduce the combinatoric complexity for the tracker to simply reduce the number of hits here from two down to one. The hit that is retained is the one with the smaller drift distance.

The final known type of clusters that can occur are from *delta-rays*, which are (relatively) low-energy knock-on electrons that scatter away along the $x - y$ -plane of the detector. While these electrons are lower energy, they can still ionize the gas and incur hits in the detector. These clusters differ qualitatively from the electronic noise clusters in that, due to the speed of the traversing delta rays and the drift speed of the chamber gas, their timing difference will be $>10\text{ ns}$ and will be missed by the electronic noise cluster removal. If a such a cluster is found, it is desired to keep the original muon hit and remove the rest of the hits incurred by the delta ray. In order to do so, the two hits on the edges of the cluster are kept while the rest are discarded.

Event Removal

After imposing all of these data reduction methods, there is still a chance that the event's contents would make it prohibitive for tracking. The primary criteria for such classification is the *multiplicity cut*. If there is an extremely large number of hits within one of the detector groups, the tracking algorithm will end up spending a very long time working through all possible combinations of hits in order to find good tracks – too long to be considered reasonable. As such, certain multiplicity limits (shown in Table 5.1) are imposed on the data to abort tracking on certain events. If the occupancy of any of the detector groups exceeds the limit imposed on it, the event tracking will be aborted regardless of the occupancy of the rest of the detector groups.

Finally, there is an additional multiplicity cut on the number of reconstructed possible *trigger roads* for an event. If the in-time hodoscope hits can match five or more trigger roads of a single sign (positive or negative muon roads denoted with positive or negative roadID, respectively), then event reconstruction is aborted. In these cases, most, if not all, chamber hits in that detector half will be masked by the road masking, and it's likely the combinatorics will be too much for the tracking to reasonably handle.

Det. Group	Hit Limit
D1	250
D2	200
D3p	150
D3m	120
Prop Tubes	250
H1T + H1B	10
H2T + H2B	10
H3T + H3B	10
H4T + H4B	10

Table 5.1: The multiplicity limits imposed on detector groups by the tracking.

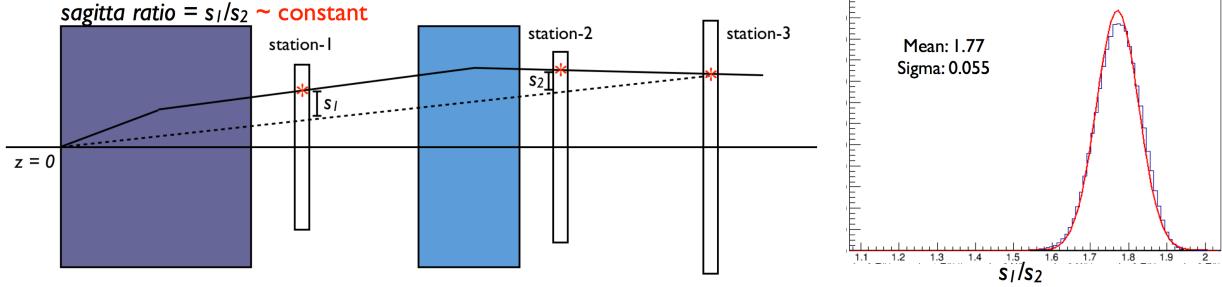


Figure 5.2: (Left) The definition of the sagittas s_1 and s_2 , with the red asterisk denoting the actual hit positions in the stations. (Right) The ratio of the *true* ratios of s_1/s_2 across muons of all kinds of kinematics (within acceptance) shows a single mean value with a defined variance.

5.1.2 Track Finding

Once the set of hits for a single event have been narrowed down to only the most plausibly useful hits, the next step is to begin the track finding. This consists of testing out combinations of hits between stations in order to arrive at a set of likely track candidates. The procedure for doing so can be broken down into four phases: local triplet reconstruction, St. 2-3 tracklet reconstruction, St. 1 projection of tracklets, and global track reconstruction.

The local triplet reconstruction can be thought of as the process of turning the hits of several wires in a drift chamber into a single set of $x - y$ position points within the chamber. Even further than that, a “stub” of a track is formed since a rudimentary slope can be estimated from the information garnered from the hits. The process begins with the combination of hits from the primed-unprimed planes. Having a pair of adjacent primed and unprimed hits allows for the elimination of the wire chamber left-right ambiguity, and better constrains the estimated position and slope. Then, all hits in the X-view are paired with all overlapping U-view hits. Those overlapping points from the X-U combinations are matched with the possible V-view hits. Each of these X-U-V combinations is then fit to an $x - y - z$ -position and a slope. In forming these triplets, a single hit may be used in multiple triplets.

Once this has been done for all Station 2 and 3 hits, the set of their track stubs are matched up to form straight “tracklets” between the two stations. Tracklets are then kept according a certain set of criteria: if their constituent stubs sufficiently point towards each other, if the tracks vaguely point back towards the target, if the tracklet points at hodoscopes (H2, H3, and H4) that fired, and if they point to proportional tubes (3 out of 4 planes) that fired. The absorber wall between Stations 3 and 4 provide a good requirement with looking to the corresponding prop tube hits, as only muons should have made it past the iron wall and fired the prop tubes.

With a set of reasonable tracklets, they are each projected through KMAG and onto Station 1 to find

possible matching hits there. The method by which this is done makes use of an analysis of the *sagitta ratio* if known DY muon tracks. The magnetic field of KMAG applies a p_T -kick of 0.404 GeV on the muon, and the influence on its deflection is based purely on the p_z of the muon at hand (assuming charge is known). As such, a large set of DY muons within the spectrometer acceptance are generated via Monte Carlo in order to measure two quantities in relation to St.3-St.1 and St.2-St.1. In Figure 5.2, the track path of a single muon is described with the influence of the magnets being treated as a single-impulse p_T kick at its center. With a line drawn between the tracklet position at Station 3 and the target ($z = x = y = 0$), the quantities s_1 and s_2 are defined, which are the distances between this line and the actual hits at Station 1 and Station 2, respectively. These studies show that the ratios of these two values across various muon characteristics is distributed as a gaussian with a defined mean and variance ($\mu = 1.77, \sigma = 0.055$). Since this ratio and s_2 is known for a given tracklet, s_1 can be approximated and used to constrain where the hit would actually be on Station 1 within a range of ± 5 cm. With this window established, triplets are matched up at Station 1 with the tracklet in order to find candidate *global tracks*.

Finally, with this set of candidate global tracks is narrowed down based on certain characteristics. This is useful as the track fitting phase which comes next is relatively intensive, so if there is any way to eliminate unlikely global tracks it should be done. The first step towards this end is the removal of hits with large *residuals*, where hit residuals are defined as the difference between where the hit triplet is and where the track is projected to pass through at that z position. Tracks with residuals greater than three times the resolution of the chambers are removed one at a time, with the track being refitted after each removal. Once this is done and all hits are within this acceptable residual range, then quality cuts are applied to the tracks. These cuts are (1) the hodoscope paddles along the track path must have fired, (2) the track momentum must be $5 < p_{tot} < 100$ GeV, (3) the track must have at least four hits in each station, (4) there must be at least one hit in each view (prime-unprime plane pairs) at each station, and (5) the multiple scattering in the iron absorber wall must not result in a scattering angle of more than 30 mrad. With these quality cuts in place, the remaining tracks are passed on to the Kalman filter track fitting algorithm.

5.1.3 Kalman Filter Track Fitting

For a given track, each positional measurement throughout the spectrometer has a well-defined uncertainty to them. When these spatial information is used to reconstruct the kinematics of a muon, the sum of these *fuzzy* measurements can be combined to render a better overall fit than any single measurement could make on its own. The tracking procedure has adopted an established algorithm to do just this: the *Kalman Filter*[69]. The Kalman filter was developed as a signals processing algorithm to clean up a true signal out

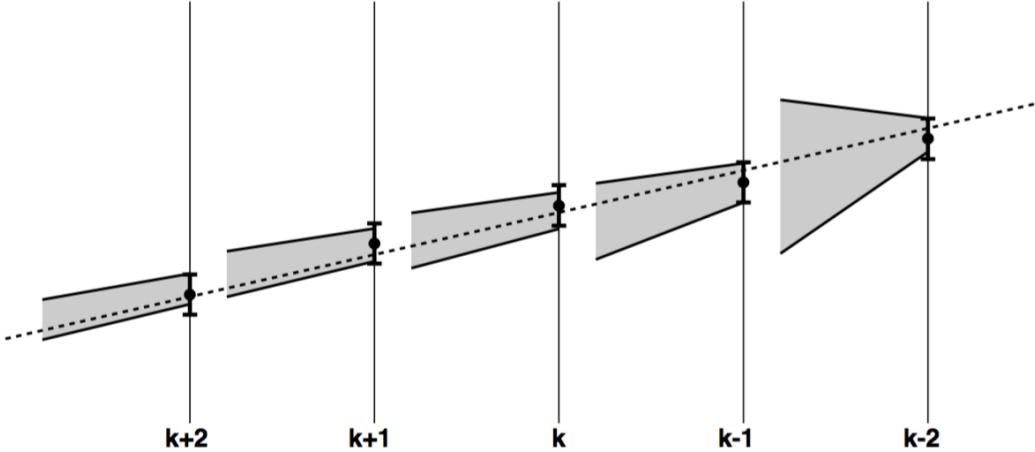


Figure 5.3: Shown is the application of the Kalman filter technique to track fitting [105]. The vertical lines correspond to the detector planes with indicated on it the measurement points and their errors. The cones represent the uncertainty in the reconstructed track parameters. As can be seen the knowledge of the track parameters is step wise updated with each measurement (Kalman filter technique applied from right to left).

of some noisy measurements, and it has since been adapted to many different uses, including particle physics track fitting.

Before the algorithm must be applied, one must first set up a physical model that will propagate one physical state to another; in the case here, behavior of linear kinematic motion that takes a particle from Stations 3 to 2 to 1 are defined, including the process of passing through KMAG and its fringe field. Once this is set up, the model is used to propagate the state of the muon from one station to the other, evaluating at each step the kinematics, along with their associated uncertainties. This estimation of the measured parameters is combined with the state vector of the muon at that point to create a weighted average used to scale the measurements by their expected accuracy, which is largely dominated by the resolutions of the drift chambers. After the weighted average is calculated, the model propagates on to the next state. The algorithm repeats, ending up with a final state with a defined covariance.

Compared to the generalized Kalman Filter procedure where the propagation step is done analytically, the tracking algorithm uses Geant4 simulation to perform a dynamic state propagation from one station to another. This method is preferable for our purposes since Geant4 is well-tested and robust in its ability to handle relativistic interactions with matter and properly simulate the effects of fringe magnetic fields. Also, due to the procedure by which SeaQuest reconstructs tracks and the backwards-forwards time invariance of particle kinematics, the Kalman Filter is used to step *backwards* in time. It is primarily done in this way since the downstream end of the experiment is generally “cleaner” than the upstream end, and the straight path between Stations 3 and 4 are generally well-defined. The Kalman algorithm is iterative, and starting

out with a less uncertain starting state leads it to converge to a stable final track state faster (fewer epochs).

In the fitting of tracks, a state vector consists of the momentum and the position (p_x, p_y, p_z, x, y, z). The tracklets, which allow for a good starting state, are then propagated upstream towards Station 1, with the measurements being updated as more steps are taken. Figure 5.3 provides a depiction of stepping (from right to left) through the stations with the uncertainty in the observables updating at each step. The uncertainties of the measurements shown stem from known chamber resolutions. The propagation between Stations 2 and 1 and between Stations 1 and the target area are performed using measured magnetic field maps and Geant4 simulations of particle interactions through matter. At the projected target area, the final observables are estimated. At this point, the process is reversed, and the particle is propagated forward in time, finely adjusting the parameters to minimize the χ^2 of the track. If the χ^2 of the global track converges, then it is passed on to the vertex finding/fitting stage of the tracking algorithm.

5.1.4 Vertex Fitting

With the fitted global tracks, they are matched up with each other to create a two-particle vertex with all its kinematics (combined from constituent tracks) estimated, along with a goodness of fit of the tracks to the vertex itself. The vertex-finding and -fitting algorithm used in the kTracker is adapted from the approach taken by Gorbunov and Kisiel at the ALICE collaboration[51]. This approach, again, makes use of the Kalman Filter algorithm.

Before tracks can be combined to vertices, they are propagated upstream from Station 1 through FMAG via a method called a magnet ‘swim.’ In this process, the solid iron magnet is considered as many small slices in z , and a commensurate amount of energy loss and p_T -kick is applied at each z -slice. Once it emerges from the upstream face of FMAG, the track is then projected as a straight line to the target region, and it is these straight segments which are used as inputs to the vertex fitting process.

In the vertex fitting, a state is defined as being the position of the (presumably) dimuon vertex (x_d, y_d, z_d). This is first initialized to $(0, 0, \bar{z})$, where \bar{z} is the average of all tracks’ points of closest approach to the z -axis. *talk with Evan and Jason about how track pairs are picked and vertex finding is done, then finish this subsection.*

Express the calculation of $M, dpx, dpy, dpz, x1, x2, xF, costheta, phi, phi-gamma$ from two muon momenta.

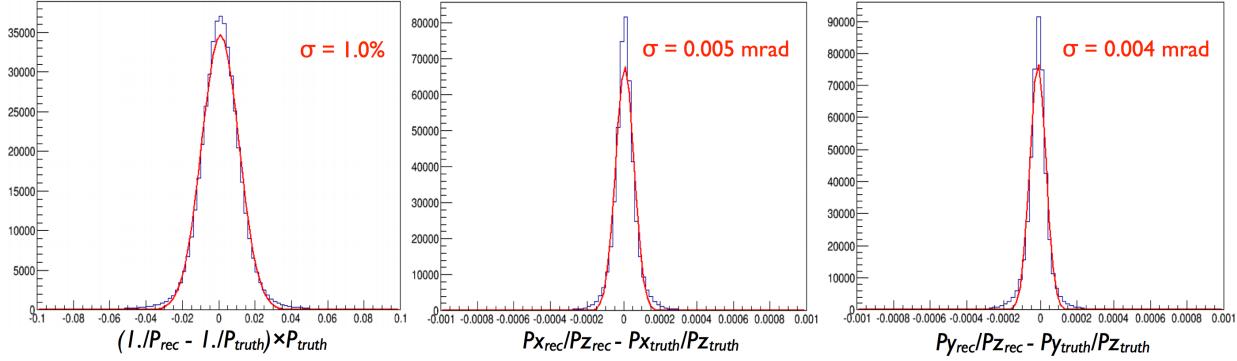


Figure 5.4: The residuals of reconstructed track kinematics against MC truth values.

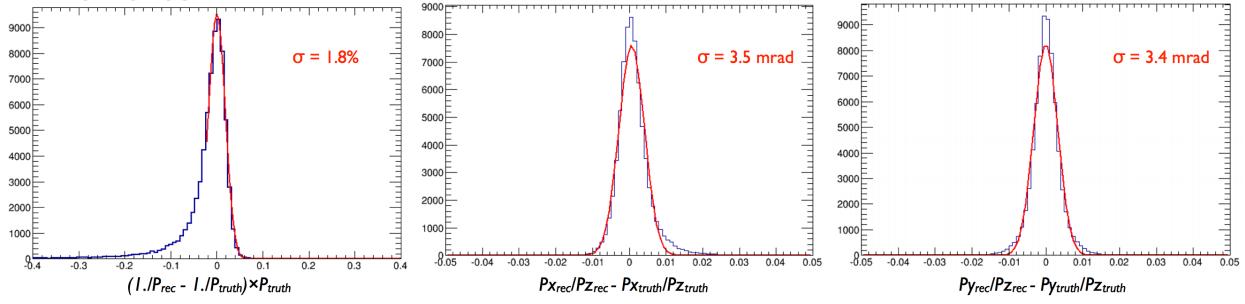


Figure 5.5: The residuals of fitted track vertex kinematics against MC truth values.

5.1.5 Monte Carlo Validation

With the tracking algorithm established, the next logical step is to measure its accuracy and efficiency, given a sample of perfectly clean Monte Carlo generated events. These events use a geometry definition that represents the latest survey numbers describing the hall, and they also simulate the ‘digitization’ of the hits (e.g. spatial coordinates in detectors into wire hits). Tracking is then run on these clean events to see if the developed algorithm can extract what is generated. In Figures 5.4 and 5.5, the residuals are shown, describing the reconstruction resolution of the track and vertex momenta. Overall, the tracking can reconstruct tracks and dimuon p_{tot} within a $\sigma = 1.8\%$ uncertainty. The residuals on track p_T/p_z is very good ($\sigma = 0.005$ mrad), but the effects of multiple scattering through FMAG and the vertex fitting smears out this to an residual uncertainty of $\sigma = 3.5$ mrad.

Regarding efficiencies on clean MC, the tracking was able to find and fit two muon tracks (i.e. χ^2 converged in fitting for both) at an efficiency of 92.2% of thrown events. Once the effects of FMAG and vertex fitting is factored in, the rate at which the tracking is able to find one muon pair vertex is 69.2%. This tracking efficiency and it’s target-, kinematic- and intensity-dependence will be investigated in great depth later on in the Rate Dependence section.

5.2 Analysis Cuts

The tracking software is concerned only with reconstructing possible tracks of relativistic charged particles through the SeaQuest spectrometer. It is, however, completely agnostic as to whether or not the charged particle was a muon that originated from a Drell-Yan interaction between the beam and the target. Though the tracker is agnostic as such, it should be noted that certain measures are in place in the design of the experiment that serve to increase the chances that a reconstructed track is a true DY muon track. Such examples are the choice of trigger matrix and the construction of an absorber wall before Station 4. Despite these measures, there are many cases in which we would wish to discard some signals, events, or whole spills in order to keep the data from becoming contaminated by anything but true $p - A$ Drell-Yan target dimuons.

As such, a set of cuts must be imposed to the tracked set of dimuon yields

- Drell-Yan Selection Cuts: Exclude kinematic regions where the signal is dominated by background
- Target Selection Cuts: Dimuons can be created in our beam dump. These cuts seek to remove the signals that likely originate from the dump.
- Trigger Selection Cuts: Many triggers were used to take data, and we only wish to analyze data triggered by a certain trigger configuration.
- Intensity Cuts: Data taken at unusually high intensities where components of the experiment can perform abnormally should be excluded.
- Quality Cuts: Most cuts fall under this umbrella. These include, but are not limited to: hardware settings, DAQ quirks, dimuon/track goodness-of-fits, etc.

In this section, these cuts are discussed in the order of tiers of data taking: Run, Spill, Event, Dimuon, Track. The intent with all of these cuts is to be as conservative as possible to ensure that all possible contamination is removed from the data despite the inevitable reduction in statistics.

5.2.1 Data Selection

The first cut on analyzable data is the choice of dataset to use. This document presents and discusses the EMC ratio results for SeaQuest roadsets #57, #62, and #67. Each of these roadset settings were in place for the span of months and contain varying amounts of statistics. In total, many millions of tracks and dimuons of varying quality are found through the track reconstruction and vertex fitting procedure discussed in the previous section. Several cuts regarding quality of track/dimuon fit and experimental conditions are discussed below and applied to these datasets in order to isolate signal dimuons from the *target*.

Table 5.2: The experiment-specific data sources and their details.

Schema	Date Range	# Dimuons	Roadset Notes
merged_roadset57_R005_V001	dummy	dummy	dummydummydummy
merged_roadset62_R005_V001	dummy	dummy	dummydummydummy
merged_roadset67_R005_V001	dummy	dummy	dummydummydummy

The names of the SeaQuest MySQL merged productions along with their date range, magnet polarity, respective statistics, and notes regarding roadset differences are presented in Table 5.2.

5.2.2 Run Level Cuts

It is an outspoken goal of the SeaQuest collaboration to absolutely keep as much data as possible. This means, in general, that the largest chunk of data that we want to allow to be thrown away is on a spill-level. That notwithstanding, there are a certain set of factors that could lead an entire run’s (~ 1 hr’s) worth of data to be excluded.

One factor is the amperage settings of the focusing and spectrometer magnets. If these are not set nominally (2000 A and 1600 A, respectively), the tracking will improperly reconstruct tracks, as it depends on a specific p_T -kick value from each magnet in order to discern sensible tracks from the data. For this reason, if the averaged magnet currents are below 50% of their nominal values, the run is excluded from analysis.

Another crucial characteristic for a run is if it contains Spill values. This can be attributable to a malfunction in the Beginning of Spill / End of Spill triggers or a malfunction in the Spill Counter function. Regardless of cause, without this spill label, it becomes impossible to relate data taken to any other critical values such as target position or beam intensity.

The rest of the run-level cuts can be summarized by the umbrella term of “DAQ issues”. If the Slow Control failed to be recorded for any reason and there are zero Beam, Target, or Environment table values in a run, it is excluded. The same goes for the BeamDAQ and Scaler readouts. Perhaps most prevalent of the DAQ issues is the case of a run that failed upon starting. These resulted in empty 32 kB files with no usable data.

Beyond these issues, there are also ranges of runs for which there was some larger issue at hand. The best example of this is the period of data taking during Roadset 62 where the automatic target control was malfunctioning. The shift crew controlled the target position manually, but the actual target that was in the path of the beam during the spill could not be verified. As such, these runs were excluded. There is very little indication from the data itself that would incriminate itself as being untrustworthy, and so these

are specifically excluded by analyzers. In practice, the runs are removed from the data set via an excluded range of spillIDs.

5.2.3 Spill Level Cuts

Seeing as most of the experimental readouts pertaining to the condition of the experiment are read out only once per spill, it is a natural result that a large variety of quality cuts that are applied are evaluated at the spill level. A summar of all such cuts are summarized in Table 5.3, but I will discuss a few key conditions here briefly.

One feature of a spill is the “dutyfactor53MHz”, which is the aggregated duty factor ($\frac{\langle I \rangle^2}{\langle I^2 \rangle}$) measurement over the whole spill as calculated by the Beam Intensity Monitor discussed in an earlier chapter. This an important measure that directly describes the quality of the beam provided to the experiment. A value that is too low corresponds to a spill that had the majority of its protons in just a small number of RF buckets. A duty factor value that is too high exceeds the capability of the main injector, and must therefore be erroneous. A visualization of good and bad duty factor values over a certain range of spills can be seen in Figure

add scatterplot of dutyfactor with good values in black and bad values in red to show how much they're outliers.

The ratio of “acceptedMatrix1” to “afterInhMatrix1” is another quality factor that may not be intuitively obvious. The acceptedMatrix1 value is the number of accepted triggers recorded by the DAQ, and the afterInhMatrix1 value is the total number of FPGA1-triggered events that were not inhibited by the BIM. Putting bounds on the ratio of these two allows the exclusion of two things: (1) spills with high DAQ dead time and (2) spills where, for whatever reason, only part of the spill was recorded. The latter of these could be the case if a spill was being partially recorded after data taking for a run had just begun *how??*.

add line of 3 scatterplots of accepted, afterInh, and ratio with good values in black and bad values in red to show how much they're outliers.

Other bounds such as the “S:G2SEM” and “QIESum” are measures of total delivered beam in units of protons and QIE, respectively. By putting bounds on these values, we avoid spills that are empty (no beam delivered) and spills that have far too much beam delivered. The lower bounds are non-zero due to a pedestal in each measurement.

add both on same plot, but with a scale to bring them near each other. g2sem: good=blue, bad=violet, qiesum: good=orange, bad=reddish

Other sweeping criteria for qualifying a spill as good are that spillID $\neq 0$ and that there must be no

Table	Value	Roadsets	Good range
Spill	dataQuality	all	0
	targetPos	all	$\in \{1, 2, 3, 4, 5, 6, 7\}$
Target	TARGPOS_CONTROL	all	= Spill table ‘targetPos’ value
Scaler	TSgo	57, 59 61 62, 67, 70	[1e3, 8e3] [1e3, 1.2e4] [1e2, 6e3]
	acceptedMatrix1	57, 59 61 62, 67, 70	[1e3, 8e3] [1e3, 1.2e4] [1e2, 6e3]
	afterInhMatrix1	57, 59 61 62, 67, 70	[1e3, 3e4] [1e3, 1e5] [1e2, 1e4]
	<u>acceptedMatrix1</u> <u>afterInhMatrix1</u>	57, 59 61 62, 67, 70	[0.2, 0.9] [0.0, 0.9] [0.2, 1.05]
	S:G2SEM	all	[2e12, 1e13]
	F:NM3S (FMAG Current)	57, 59, 62, 67, 70 61	> 1000 [200, 500]
	F:NM4AN (KMAG Current)	all	> 1000
BeamDAQ	QIESum	all	[4e10, 1e12]
	trigger_sum_no_inhibit	all	[4e9, 1e11]
	inhibit_block_sum	57, 59, 61 62, 67, 70	[4e9, 1e11] [4e9, 2e11]
	dutyfactor53MHz	57, 59, 61 62, 67, 70	[15, 60] [10, 60]
	#Tracks/Spill	all	> 0

Table 5.3: The experiment-specific data sources and their details.

Min SpillID	Max SpillID	Details
371870	376533	Trigger timing shift
378366	379333	Trigger timing shift
394287	409540	Run 3 commissioning period
394287	414555	LD2 flask filled with LH2
416207	424180	Manual target control
482574	484924	Magnet flipped before roadset changed
526201	526364	Improper QIE inhibit setting
581369	582460	KMag off
675814	684041	Drift Chamber 1 gas flow problem
684663	689443	Incorrect D3p in-time window

Table 5.4: Specific spill ranges of excluded data for all roadsets.

duplicates of any of the values listed in Table 5.3. It is necessary due to the relational nature of the data that there must be a real spillID by which one can determine all other spill-relevant values; cases where spillID=0 are cases in which no Spill Counter event was available, and these should be excluded. Regarding duplicate values, if such exist for one spillID, it is almost impossible which is the correct one, so the spill should not be considered.

Finally, certain periods of data taking are known by the collaboration to be problematic and not suitable for analysis. There exist periods of time where the trigger timing was off by a certain margin and degraded the efficacy of the trigger’s ability to trigger on dimuon events. There’s also a range of spills for which the magnets were flipped before a new trigger roadset could be devised for such a configuration. These are defined as a series of “bad spill ranges” and can be found summarized in Table 5.4.

As briefly described in the chapter on database design, the addition of a “dataQuality” field was touched on. Here, this field is used as a bitmask in order to identify whether or not a spill passes or fails a given cut by flagging a bit as 0 (pass) or 1 (fail). This allows for analyzers to easily select only spills that pass all criteria by simply requiring *dataQuality* = 0.

5.2.4 Event Level Cuts

The bare minimum requirement for the Event level data is that it must have an entry in all event-level tables: Event, kEvent (the kTracker’s event meta-data), and QIE. If any of these is missing, then the event is discarded. More specific criteria can be imposed to consider an entry to be “missing,” such as when a row in the QIE table reports all “-1” values for the BIM RF readouts. In these instances, there was a malfunction in the BIM, and the “-1” acts as a *sentinel value* to indicate that the information does not exist for that event.

In addition to missing data, certain criteria can be placed on the Event data in order to exclude poor-quality or unwanted events. The simplest cut that is imposed is the requirement that the event that was triggered must have satisfied the FPGA-1 trigger. The different trigger configurations are designed on optimized for different purposes; the FPGA-1 trigger is designed to select high-signal, low-background Drell-Yan (opposite sign muon pair) events, and by removing other types of triggers such as the NIM-3 (random RF) and FPGA-4 (singles) that exist in the data, it is ensured that there will be no effect due to different triggering systems when performing analysis.

The cut on beam intensity is another that ensures that we exclude data that would introduce undue volatility and uncertainty into the measurements. Due to the varying beam intensity at SeaQuest, there are inefficiencies which are corrected via weighting of events (this process is described in great detail later). Cutting on very high intensities removes the region where uncertainties in the rate dependence fits can cause large uncertainties in correction weights.

The cuts described here and more are summarized in Table 5.5.

5.2.5 Dimuon Cuts

Dimuons are selected based on many criteria: the known acceptance of the spectrometer, physical constraints, goodness of vertex fits, and more. The precise description can be summarized by the following: we are looking for two tracks of particles with opposite signs that can be *well-fit* to originating from the target table. Stepping through the cuts and their justifications:

- **dx , dy , dz :** These define the fitted point of origin (“dimuon (x , y , z)” for the dimuon pair. If the vertex was outside the (dx , dy) region about the beam path, then this is probably not a real dimuon. If it is outside the desired dz range, then it probably did not come from the target.
- **dpx , dpy , dpz :** These define the total momentum of the dimuon (“dimuon momentum (p_x, p_y, p_z)”). If $|p_{x/y}| > 2 \text{ GeV}$, then one or both muons would pass outside of the acceptance of the experiment. Likewise, if $|p_z| < 30 \text{ GeV}$, one or both of the muons would have such low momentum as to end up bent out of acceptance of the spectrometer by the magnets. Finally, since a 120 GeV proton beam is used for SeaQuest $p_z > 120 \text{ GeV}$ must only come from unrelated muons that are reconstructed to this high momentum value.
- **$\text{sgn}(px1)$, $\text{sgn}(px2)$:** These are terms for the momentum of the positive and negative track, respectively. Requiring a specific sign for the x -momentum for each track acts as a safety check that the tracking identified the charge of the particle correctly.

- **sgn(roadID1), sgn(roadID2)**: This is another check that ensures that the muons came from opposite sides of the experiment (top-bottom). Such should be the case when triggered by a FPGA-1 event, but this enforces it.
- **trackSeparation, chisq.dimuon**: These are measures of the goodness of fit that two tracks came together to form at a single vertex. The track separation is the distance between the points of closest approach from each track to the beam axis (z-axis). The χ^2_{dimuon} is a measure of the quality of tracks combined with how well the two tracks converge to a single point.
- **mass**: There is a di-lepton range of invariant masses that are book-ended by the J/Ψ and the Υ in which Drell-Yan should be the only source. A cut on this range excludes any background that would be introduced by real J/Ψ s and Υ s.

5.2.6 Track Cuts

Going even deeper to remove signal contamination, if one of the two tracks for even a *good* dimuon is found to fail certain checks, the entire event is excluded from analysis. Again, stepping through each cut:

- **numHits**: A track passes through 18 drift chamber planes. Even if the chambers are somewhat inefficient, it is unlikely that the number of hits corresponding to a single track would be less than 15.
- **roadID \neq 0**: It is possible that a reconstructed track passes from the top half of the experiment to the bottom (and perhaps bends back up to the top). By the roadID naming convention, these are assigned a roadID of 0, and since we are looking to only analyze cases where there is one “top” track and one “bottom” track, these should be excluded. By imposing this requirement, such tracks are removed.
- **chisq/(numHits-5)**: Each track is described as a muon with five physical parameters ($x_1, x_2, p_T, \theta, \phi$). This expression described the χ^2 per degree of freedom, and can be cut on to select only good quality tracks that fit the hit pattern well.
- **pz1**: This is the momentum of the track at Station 1. In studies of signal background sources, one of the primary contributors is combinatorial background, which can be analyzed by looking to like-sign triggered events and studying their kinematics. As can be seen in Figure 5.6, depending on the number of hits, applying a cut on pz1 can remove the strong peak of combinatoric background from unrelated muons.
- **chisq_dump-chisq_target**: These are the χ^2 s of a track when forced to go through the beam dump and the target, respectively. By requiring this quantity to be above a certain value means, roughly

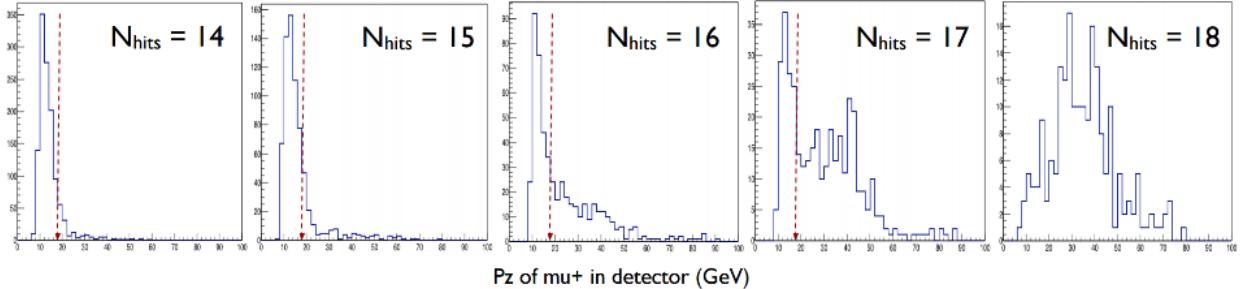


Figure 5.6: The p_z distribution of muons that form like-sign dimuon pairs, which is a source of background for DY signal events. In cases where $14 \leq \text{numHits} < 18$, imposing a cut on $p_z > 18$ GeV removes a majority of this background.

speaking, that the track is very likely to have come from the target and not the beam dump.

5.3 Normalization of Relative Yields

The base measurement for the studies presented in this paper is, fundamentally, the number of dimuon events on a given target. This is known as the *raw yield* (Y). The ratio of these yields for a given target is used to study the kinematic dependence of the *per nucleon cross section ratio* that defines these phenomenon, but there must be a target-to-target normalization for such measurements.

$$R \equiv \frac{F_2^A}{F_2^D} \approx \frac{\sigma_{\mu\mu}^A}{\sigma_{\mu\mu}^D} \frac{2}{A} = N_{A/D} \frac{Y^A}{Y^D} \quad (5.1)$$

Here, A is the mass number of the target nucleus, and $N_{A/D}$ is the normalization factor, which can be expressed as:

$$N_{A/D} = \left(\frac{N_0^D \cdot \text{lt}^D \cdot T^D(\xi)}{N_0^A \cdot \text{lt}^A \cdot T^A(\xi)} \right) \cdot \left(\frac{2 \cdot n_D \cdot L^D}{A \cdot n_A \cdot L^A} \right) \cdot \frac{\varepsilon^D}{\varepsilon^A} \cdot \frac{\bar{\Omega}^D}{\bar{\Omega}^A} \cdot \frac{\epsilon^D}{\epsilon^A} \quad (5.2)$$

where N_0 is the number of incident protons, lt^A is the % live time on the target, T^A is the target transmission (or beam attenuation) factor, n_A is the number of nuclei per unit volume, L^A is the target thickness, ε^A is the detector efficiency, $\bar{\Omega}^A$ is the spectrometer acceptance, and ϵ^A is the tracking reconstruction efficiency. One of the key benefits of performing a ratio measurement is that, to good approximation, the terms regarding chamber efficiency, tracking efficiency, and acceptance should be equal across targets and should therefore cancel each other out. How well this assumption holds up will be assessed and evaluated.

With the normalization applied, this ratio R can be accurately studied by being binned in one or more kinematics in order to observe interesting behaviors and compare them against previous measurements and various models. Once The cuts described in the previous section are applied, the surviving dimuons are all

Table	Value	Good range	Type
kTrack	numHits	> 14	Quality
	roadID	$\neq 0$	Quality
	z0	(-400, 200)	Quality
	chisq/(numHits - 5)	< 5	Quality
	pz1 (Where numHits \neq 18)	> 18	Quality
	chisq_dump - chisq_target	> 10	Target
kDimuon	dx	(-2, 2)	Quality
	dy	(-2, 2)	Quality
	dz	(-300, 200)	Quality
	dpx	(-3, 3)	Quality
	dpy	(-3, 3)	Quality
	dpz	(30, 120)	Quality
	xB	(0, 1)	Quality
	xT	(0, 1)	Quality
	xF	(-1, 1)	Quality
	trackSeparation	(-250, 250)	Quality
	chisq_dimuon	< 15	Quality
	px1	> 0	Quality
	px2	< 0	Quality
	dz	(-300, -60)	Target
Event	mass	> 4.2	Drell-Yan
	MATRIX1	1	Trigger
kEvent	status	0	Quality
QIE	$\forall(RF+08, RF+07, \dots RF-08)$	< Beam.Inh_thres	Quality
	$\forall(RF \pm XY)$	> 0	Quality
	triggerCount	≥ 0	Quality
	intensity_p	< 60000	Quality

Table 5.5: A list of all quality and kinematic cuts for selecting good quality dimuons from the target.

considered to be *good quality target Drell-Yan dimuons*, and all analyses and corrections described from here on refer to them only.

This section will discuss the several components of the target-to-target relative normalization from the determination of beam intensity, target thickness, live time, chamber efficiency, and acceptance. Beam intensity and live time information was obtained from the BIM and scaler data. Acceptance was calculated with Monte Carlo simulation. Finally, chamber efficiencies and their ratios per target were studied with NIM3 random trigger data embedded onto clean MC. The tracking efficiencies are highly rate-dependent, and will be discussed at length in another section.

5.3.1 Live Proton Count and Beam Attenuation

The measurements of the number of incident protons, live time, and beam attenuation combine to give a figure for the total number of protons on target for which data was recorded. The important sources to consider here are the upstream SEM detectors that record the number of incident protons, the QIE recorded sums of protons (in QIE units), the number of QIE units for which there was either a trigger inhibit or DAQ deadtime. In practice, SeaQuest calculates a value called “live protons” *for each spill*:

$$N_0^A = \text{G2SEM} \quad ; \quad \text{lt}_A = \frac{\text{QIESum} - \text{inhibit} - \text{deadtime}}{\text{QIESum}} \quad ; \quad \text{liveProton} = N_0^A \cdot \text{lt}_A \quad (5.3)$$

this *liveProton* value is integrated, good spill by good spill, for every target in each data set.

What are average livetimes for the experiment? Deviations?

It is also important to factor in the effects of inelastic scattering of the beam within the target, causing the effective amount of beam-on-target to decrease with every step dz through the target. This is different for each target based on target nuclear densities and cross sections. With this, the liveProton value can be corrected. Consider incident proton number N_0 , where dN is the number of protons attenuated when passing through target slice dz , and L is the total length of the target. The function for the number of protons remaining at a given position z is

$$N(z) = N_0 e^{-\frac{\rho_A}{\text{NIL}_A} z} = N_0 e^{-z/\lambda} \quad (5.4)$$

where ρ_A is the density of the material, NIL_A is the nuclear interaction length of the target, and $\lambda = \frac{\text{NIL}_A}{\rho_A}$.

When integrating over the length of the target, one arrives at the total [protons · cm] exposed to the target. Assuming that each proton interacts independently, one can divide this integral by the length of the target

Target	Target Position	Length[cm]	n_A [mol/cm ³]	ρ [g/cm ³]	NIL [g/cm ²]	λ [cm]	ξ	$T(\xi)$
D_2	3	50.80	0.081	0.163	71.8	439.41	0.1156	0.944
Fe	5	1.91	0.141	7.874	132.1	16.78	0.1136	0.945
C	6	3.32	0.150	1.802	85.8	47.61	0.0698	0.966
W	7	0.95	0.105	19.300	191.9	9.94	0.0958	0.954

Table 5.6: Details pertaining to the measurables and physical properties individual nuclear targets.

to arrive at *effective* number of protons incident on the target

$$N_{\text{eff}} = \frac{1}{L} \left[N_0 \lambda \left(1 - e^{-\frac{L}{\lambda}} \right) \right] = N_0 \left[\frac{1 - e^{-\xi}}{\xi} \right] = N_0 T(\xi) ; \quad \xi = L/\lambda. \quad (5.5)$$

The normalization of the transmission or beam attenuation factor T^D/T^A is therefore only a function of the length of the target, its density, and its nuclear interaction length. In total, the product of $N_0 \cdot \text{lt} \cdot T(\xi)$ represents the effective protons on target which was recorded by the data acquisition systems. Values of ξ and $T(\xi)$ can be found in Table 5.6.

5.3.2 Target Density Ratios

The number of nucleons that the effective protons will interact with should be considered when performing the target-to-target normalization. Aside from the background “targets” and liquid hydrogen, there are four nuclear targets in this A-dependence study. The nucleons in liquid deuterium are viewed as essentially free particles, and therefore proved an isoscalar reference by which the heavier targets are compared to. In all, the nuclear targets used are D_2 , C, Fe, and W, and their attributes pertinent to target density ratios can be found in Table 5.6. The yield ratios are normalized to the number of nucleons in each target by the normalization component $(2 \cdot n_D \cdot L^D)/(A \cdot n_A \cdot L^A)$.

Uncertainties in densities, target lengths? Find solid target with largest uncertainty, estimate LD2 uncertainty the best we can, get % systematic from the ratio of them.

5.3.3 Detection Efficiency

The key cause of detector inefficiency is high occupancies from high intensity events. Also, the total number of tracks that fired off elements of the many detectors is dominated by events generated from the 503 cm iron beam dump. The difference in interaction lengths of the targets should have only a small effect on the absolute chamber and hodoscope efficiencies. Further, any longer term time-dependent drifts in detector efficiencies will largely cancel out due to the frequent minute-by-minute interchange of targets. It is therefore

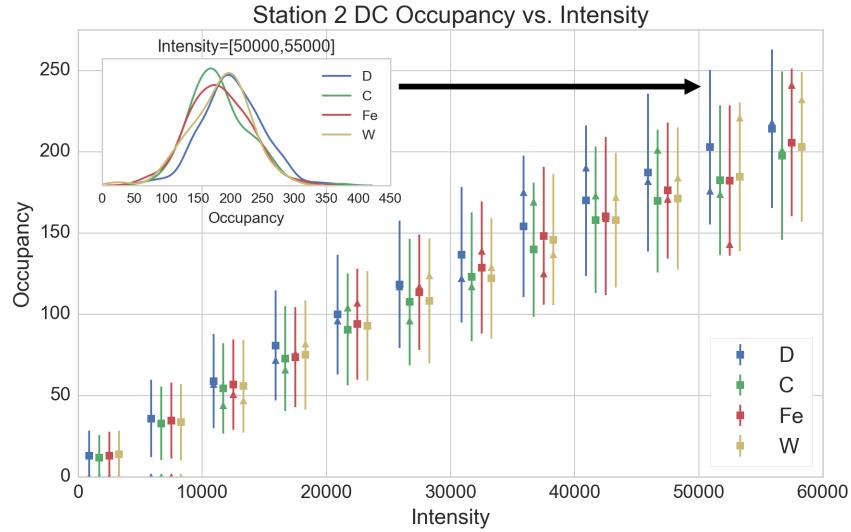


Figure 5.7: The Station 2 drift chamber occupancy for the four targets as a function of intensity. The error bars show a systematic deviation of occupancies from the mean for a given intensity bin. The inset plot shows a KDE plot of distributions for the intensity bin in which the values differ the most. The triangle markers represent the modes of the distributions.

justifiable to assume that there is no target-to-target chamber efficiency dependence to normalize to. The overall ratio of detector efficiencies between targets is therefore taken to be $\varepsilon^D/\varepsilon^A = 1$.

It is however possible that pions created in the targets have more of a chance to decay into muons that proceed to add to the background hit rates of the experiment; pions created in the beam dump are much more likely to get re-absorbed before they decay. To investigate the possible magnitude of this contribution to occupancy, detector occupancies were studied for the different targets in unbiased NIM3 random events as a function of chamber intensity. We see in Figure 5.7 that they are not significantly different from each other, though there is a slight difference in occupancies between the deuterium target and the rest. However, since the occupancy distributions at a given intensity bin do not change much for a given intensity, it is safe to assume that the efficiencies likewise do not change by much from target-to-target.

To override any concern of chamber efficiencies having an effect on any absolute or ratio measurements, a study was performed to measure the effect of chamber efficiencies on the track reconstruction algorithm. This was done by embedding clean Monte Carlo-generated Drell-Yan dimuons into real NIM3 random event hit data. Then, a constant 94% chamber efficiency was applied by removing a random 6% of the chamber hits, and tracking was performed on this sample. The 94% figure reflects the best estimate of overall chamber efficiencies at SeaQuest. Following that, a rate-dependent chamber efficiency was imposed of the form

$$\varepsilon(I) = 0.95 - a \cdot I \quad (5.6)$$

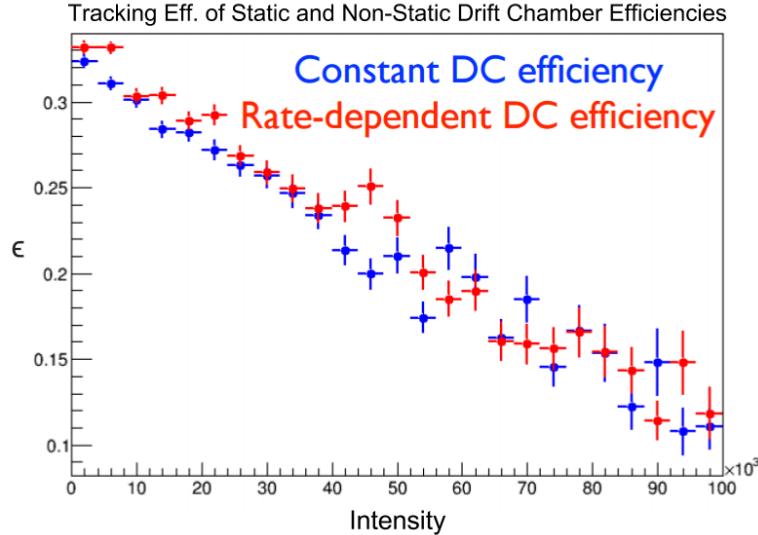


Figure 5.8: The ratio of successful reconstruction to all thrown by MC for two chamber efficiency models. The linear decreases in efficiency are not significantly different from each other.

where a was set to $2\text{e-}6 \text{ protons}^{-1}$. This resulted in a linear drop in efficiency with it dropping to as low as 75% at the highest intensity regions considered (100000 protons). The tracking algorithm was then run on this data sample.

For even the constant 94% efficiency case, the tracking is expected to decrease in *tracking efficiency* (a very different topic to tackle, as we will approach later in this chapter) as intensity increases and pattern recognition becomes more difficult as overall occupancies increase. The important observation to make is how this intensity-dependent tracking efficiency compares to the case of decreasing efficiency. For both cases, the ratio of successful reconstructions to all thrown dimuons (before even the trigger acceptance is imposed) can be seen in Figure 5.8. It is perhaps surprising to see that even at high intensities, the linear drop in tracking efficiencies does not change between the two cases. The reason for this is due to the fact that the tracking algorithm is very *robust* in the way it is able to handle missing hits, as it only requires a few hits in each station (3 out of 6) to successfully find a track.

Putting this together:

- Differences in chamber efficiency between targets could only come from different resultant chamber occupancies as a result of that target being in the beam's path.
- Chamber occupancies are only slightly different for deuterium when compared to the other nuclear targets, and so if efficiencies are different between targets, the modification will likely be very small.
- Even in the case of vastly different efficiencies, the tracking efficiency is unchanged.

- As a result, effect of chamber efficiency on the yields will be the same for all targets.
- $\varepsilon^D/\varepsilon^A$, a measure of how the chamber efficiency effects the deuterium target yields as compared to a nuclear targets, can therefore be considered to be 1.

5.3.4 Relative Acceptance Correction

The *true* acceptance as a function of any set of kinematic variables \bar{x} can be expressed as:

$$A(x) = \lim_{N \rightarrow \infty} \frac{Y_{accepted}(\bar{x})}{Y_{thrown}(\bar{x})} ; \quad N \equiv \int Y_{thrown}(\bar{x}) dx \quad (5.7)$$

where *thrown* denotes Drell-Yan simulated events from Monte Carlo which generates every kinematic value by sampling over the distributions of that variable with a random number generator. The simulation includes such processes as multiple scattering and energy loss that occur due to particle passage through materials (such as our targets and beam dumps). Survey numbers of the experimental hall are used to build an accurate representation of the apparatus and its sensitive detector regions. Muon pairs over much of the kinematic phase space and was weighted according to the known Drell-Yan cross section. This does not factor in any higher order nuclear effects or anything like \bar{d}/\bar{u} asymmetries.

The pairs that qualify as *accepted* are those that satisfy the following two criteria: (1) were able to trace their way through the spectrometer without leaving the sensitive range of the detectors at all stations, and (2) they followed paths that would have satisfied the un-like-sign dimuon trigger (FPGA1). Such data was written to file in a format that matches real experimental data and is then ready for analysis. These events are weighted according to the kinematics and DY cross-section as mentioned above, but for analysis, these weights when analyzed were divided by N , the number of thrown events.

Approximately 800,000 dimuons were thrown for each target to have the acceptances analyzed. For an absolute acceptance, a study of a full 4π (or 2π) simulation is required, where 4π is the solid angle of an entire (semi-)sphere. In such simulations dimuons are thrown in every physically possible direction, and then one can derive the absolute acceptance of the target and/or spectrometer when taking into account this source of *thrown* events. However, for the purposes of these studies, we are only concerned about $\bar{\Omega}^D/\bar{\Omega}^A$, or the relative acceptance correction. As we see in Fig. 2.6, we see that though the spatial distribution of the liquid deuterium flask and the solid targets are similar, they are fundamentally differently distributed. The density of the deuterium in the flask is constant throughout the length of the flask while the disks are discrete and separated. As a result, it perhaps may be the case that due to beam attenuation, there is a spatial difference between where most dimuons are generated between the different targets. The cause of

any acceptance difference, however, is moot; the goal here is to estimate the relative acceptance correction factor between deuterium and the nuclear targets.

The target dependence of the acceptance was studied via yield ratios of heavy targets to that of deuterium for the valid kinematic regions of the experiment. These regions are $4.2 \text{ GeV} < M_{\gamma^*} < 10 \text{ GeV}$, $0.1 < x_2 < 0.5$, $0.35 < x_1 < 0.9$, $0.0 < x_F < 0.9$, and $0.0 \text{ GeV} < p_T < 3.0 \text{ GeV}$. Since no nuclear effects are simulated by the Monte Carlo generator, this ratio should lie at 1; any deviation of this must be purely from the acceptance effects due to the slight differences in the geometries and densities of the targets. The results can be seen in FigureThe acceptances of solid targets were calculated by MC to be Z.ZZZ larger than that of the liquid deuterium target. As such, a correction factor of $\bar{\Omega}^A/\bar{\Omega}^D = Z.ZZZ$ is included in the target-to-target normalization.

5.4 Dimuon Weights and Corrections

Weighting of events is a standard procedure for mapping one distribution onto another. The most common example is with Monte Carlo (MC) simulation weighting. A simple physics event MC meant to simulate a well-defined process will ‘throw’ an event with certain kinematics randomly drawn from known distributions via *inverse transform sampling*. By doing this, the simulation is already close to being physical, but according to the cross section of a given process, a certain combination of kinematics will be decidedly more or less probable. A weighting subroutine calculates this likelihood for each event and assigns a weight with respect to how likely that event is based on the kinematics thrown. When binning this weighted data, the adjusted number of events in a bin is given by

$$N_{bin} = \sum_{i \in bin} w_i ; \quad \sigma_{bin} = \sqrt{\sum_{i \in bin} w_i^2} \quad (5.8)$$

which reduces to the Poissonian statistical picture when all weights are 1.

Weighting in this manner can be used in any case where it is desirable to map one distribution to another, as in the case of applying corrections in the case of efficiencies and background subtractions. The most intuitive connection between the use of weights and corrections is in the case of efficiency corrections. Let us say that a dimuon with a certain set of kinematics, for whatever combination of effects, has a reconstruction efficiency of 80%. This means that four out of five actual dimuons with these kinematics will be reconstructed. If you know this efficiency ϵ_{recon} , you can apply a weight as $w = 1/\epsilon_{recon}$ to each dimuon, as each dimuon represents, in part, a larger population that is not fully represented. In the case of our example, a single dimuon represents $1/0.80 = 1.25$ dimuons in order to make up for the ones that are

missed due to that specific inefficiency.

At SeaQuest, there is an issue of rate dependence, and as the sources of this effect is identified, we apply a correction for each source in the form of an applied weight. Here, we discuss the reconstruction efficiency and empty/none target background subtraction corrections.

5.4.1 “kEfficiency” Correction Using Messy/Clean Data

It is a matter of combinatorics at SeaQuest that, as the intensity increases, the number of background hits on all of the detectors increases. As the number of hits in all the detectors increases, the harder it becomes for the tracking to identify the actual dimuon(s) from the whole event. This is regarded to be an intensity-dependent tracking efficiency. Seeing as the primary tracking algorithm for SeaQuest is kTracker, this is known as the “kEfficiency”, and its intensity, target, and kinematic dependency makes it a factor to model and correct for.

The procedure developed by Evan McClellan of UIUC for calculating this kEfficiency is to take compare the performance of the tracking on a clean MC sample of dimuons to an analogous sample that’s mixed in with real intensity-dependent background. To do this, a *single* sample of MC-generated dimuons and all of their hits is generated; this sample is denoted as the “*clean*” sample. Then, this sample is mixed with all of the hits from a NIM3 event from the experimental data. The NIM3 trigger, as is discussed in Chapter 2, is the randoms trigger, and by using these events we avoid any effects due to trigger selection bias. This mix of MC dimuons and NIM3 triggered data is denoted as the “*messy*” sample. These NIM3 events all have an intensity, and by embedding a clean MC dimuon into a situation typical of an event at a given intensity, we get a basis by which to judge how efficiently the tracking can reconstruct a dimuon in the clean sample versus the messy sample as a function of intensity (and other kinematic variable).

It is important to note that the DY cross section is so small that there is a negligible chance that a given randomly-triggered NIM3 event will contain an actual dimuon, nonetheless a dimuon matching the characteristics of the MC dimuon that is embedded. As such, it is assumed that the dimuon embedded in the NIM3 event is the only dimuon to find in the event. It is also justified to assume that the kinematics of a dimuon is in no way correlated to the intensity, and the intensity of an event relates only to the amount of background hits observed. The relation to the intensity and the number of background hits can be observed in Figure 5.9.

Once the messy and clean samples are prepared, the tracking software is run on them and the two outputs are compared. The first step is to bin the resulting reconstructed dimuons into intensity bins. This binning uses the weights of the dimuons that is assigned by the GMC that created them. The weights are used

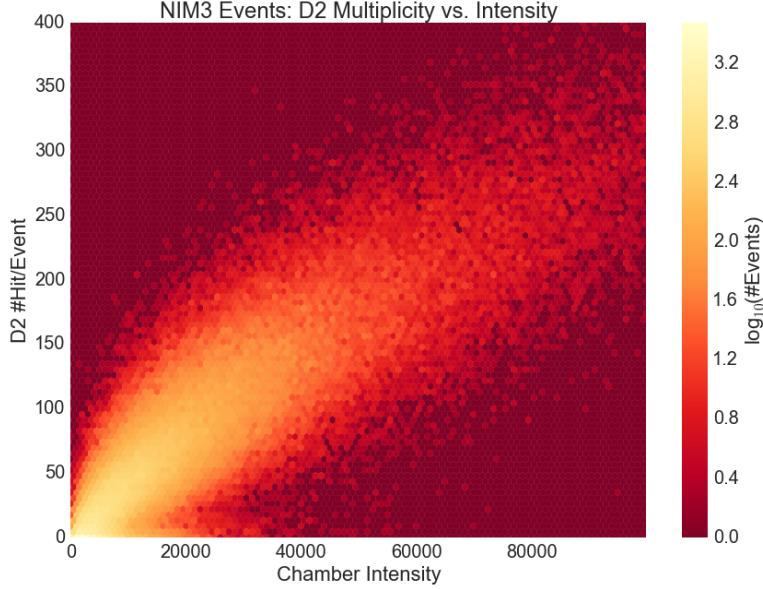


Figure 5.9: The linear tendency of detector multiplicity to increase with intensity is observed by looking to unbiased NIM3 events. Here, the occupancy per event of the drift chambers at station 2 are shown as a function of intensity.

instead to avoid undue influence on the efficiency from dimuons that are highly unlikely to occur. The value of the bin and its uncertainty are calculated according to Eq. 5.8.

For each matching bin, the ratio of $\frac{\#_{messy}}{\#_{clean}}$ values is used to calculate the efficiency of the tracking for that bin. The uncertainty of that efficiency is more complicated, as the relation between the messy and clean samples is of a *binomial* nature. Each dimuon successfully reconstructed in the messy sample represents a *positive* outcome from the number of possible outcomes which is represented by the clean sample. With this kind of relationship, binomial uncertainty calculations must apply. For large enough N number of trials and efficiency $\epsilon \notin \{0, 1\}$,

$$\epsilon = \frac{N_+}{N} ; \quad \delta\epsilon = \sqrt{\frac{N_+ N_-}{N^3}} = \sqrt{\frac{\epsilon(1-\epsilon)}{N}} \quad (5.9)$$

but when weighted trials are involved, things become more complicated in calculating the statistical error.

The calculation of this statistical error[78] is as follows:

$$\epsilon = \frac{\sum_{i \in +} w_i}{\sum_i w_i} \quad (5.10)$$

$$\delta\epsilon = \frac{\sqrt{\sum_{i \in +} w_i^2 \left(\sum_{i \in -} w_i\right)^2 + \sum_{i \in -} w_i^2 \left(\sum_{i \in +} w_i\right)^2}}{\left(\sum_i w_i\right)^2} \quad (5.11)$$

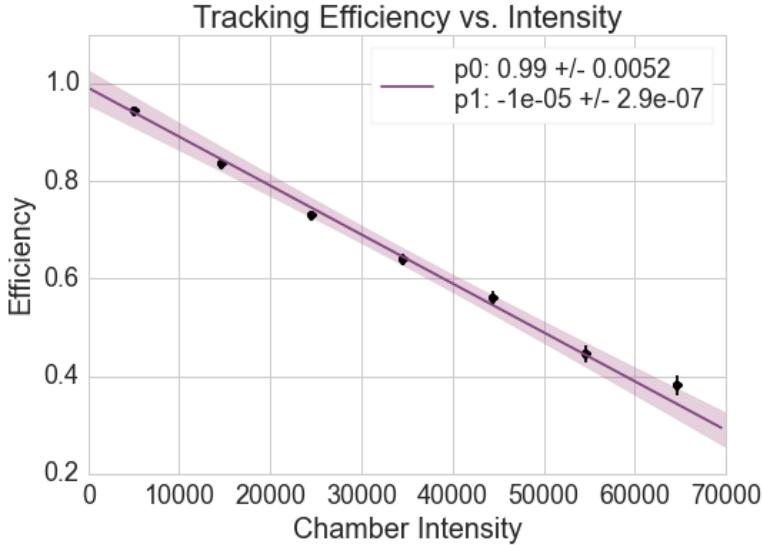


Figure 5.10: The ratio of “messy” to “clean” dimuon samples from the Deuterium target in Roadset 67 data. This shows the linear relationship of the tracking efficiency to the intensity of the events. The shaded band indicates the 95% confidence band of the fit.

where

$$\sum_{i \in -} w_i = \sum_{i \in +} w_i - \sum_{i \in +} w_i \quad (5.12)$$

$$\sum_{i \in -} w_i^2 = \sum_{i \in +} w_i^2 - \sum_{i \in +} w_i^2 \quad (5.13)$$

For a set of efficiencies, a linear function is fit to the efficiencies. This fit function would then be used to calculate the efficiency for any given dimuon that occurs at any given intensity. The linear parameters are fitted to the data (which is weighted by its statistical uncertainties) by a χ^2 -minimization procedure to render a line of best fit along with a confidence interval (Fig. 5.10)

Further, these fits have a statistically significant kinematic dependence, so it becomes important to unfold the kinematic space and perform this procedure for several kinematic bins in one or more dimensions to assure an accurate efficiency calculation. Of the six defining kinematics of the dimuon, we examine the dependence of the fit on five of them: $(x_1, x_2, p_T, \theta, \phi)$, neglecting the azimuthal production angle ϕ_{γ^*} . Each of these five kinematics is divided into three bins (low, medium, high) values to observe on which kinematics the linear fit depends. The results are shown in Figure 5.11. The curves produced seem to indicate that there only exists a kinematic dependence on the x_1 and x_2 kinematics. If this is the case, then the clean and messy samples can be split two-dimensionally in these two kinematics, and a fit can be made to both variables. When analyzing dimuons, its kinematics can indicate which fit to use to calculate the tracking

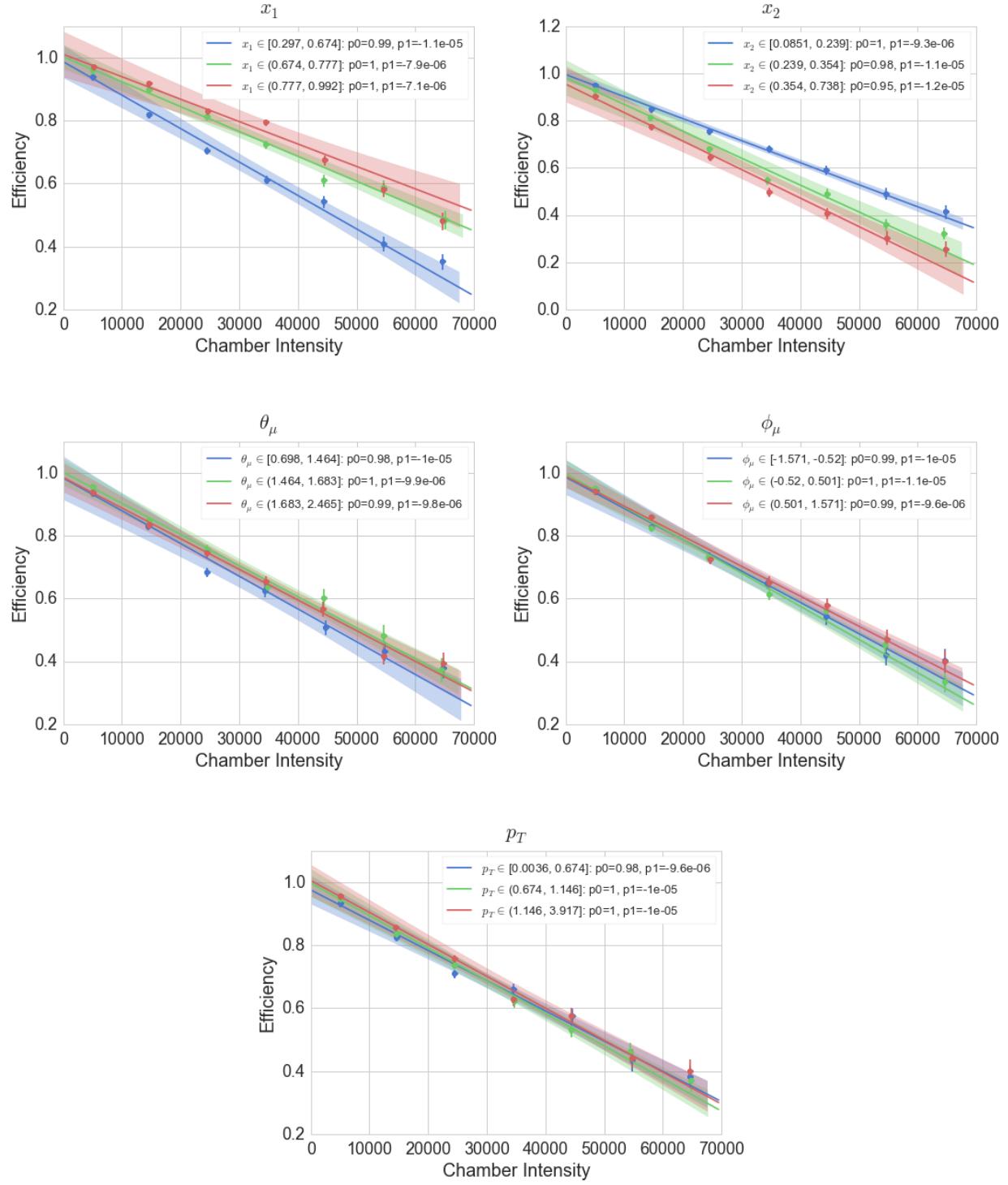


Figure 5.11: Tracking efficiency with the data broken into three statistically equivalent bins in the five primary kinematics, $(x_1, x_2, \theta_\mu, \phi_\mu, p_T)$. There appears to be a significant kinematic dependence on x_1 and x_2 .

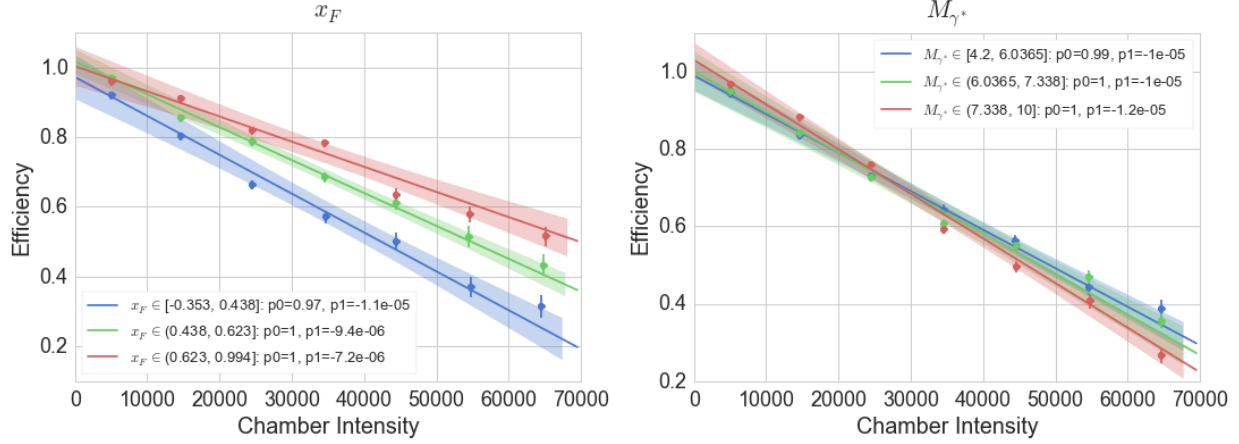


Figure 5.12: The kinematics x_F and M_{γ^*} are investigated. A clear kinematic dependence exists in x_F while M_{γ^*} is largely consistent.

efficiency based on the intensity of the event, and a weight can be calculated.

But before we come to any clear conclusions, let us investigate two more kinematic phase spaces. We know from the discussion in Chapter 1 that $\{x_1, x_2\}$ can be used almost interchangeably with $\{x_F, M_{\gamma^*}\}$, and each of x_1 and x_2 depend on both x_F and M_{γ^*} . As such, if only one and not the other shows to influence the tracking efficiency curves, then only that one would be needed for the tracking efficiency correction. We see the behavior of the efficiency curves as a function of x_F and M_{γ^*} in Figure 5.12. It can be concluded that since there is no substantial mass dependence observed, the tracking efficiency has a kinematic dependence solely on x_F .

Another consideration is with respect to whether or not there is a target dependence to factor into the correction. For all of the above “kEfficiency” plots, only deuterium data is used. The kEfficiency curves for deuterium, hydrogen, carbon, iron, and tungsten can be found on Figure 5.13. It can be concluded that there is enough of a difference between deuterium, iron, and the rest to justify calculating and applying this correction on a target-by-target basis.

Also, as a sanity check, there should be no time-dependence of this function. A quick check comparing different sets of data from two different roadsets quickly confirm that there is no such time dependence (Fig. 5.13).

As a result, the final calculation of an tracking efficiency for a given dimuon event should depend on (1) target, (2) x_F , and (3) intensity. In order to apply this as a weight, we first define $kEff_{t,\bar{x}}(I)$, a set of linear fit functions: one for each target (t) and for each kinematic bin (\bar{x}), which may be binned in any number of kinematic dimensions. In the case of this analysis, the kinematic binning will only be in x_F . The procedure in weighting each dimuon will then be to take each dimuon event i and its chamber intensity I_i , originating

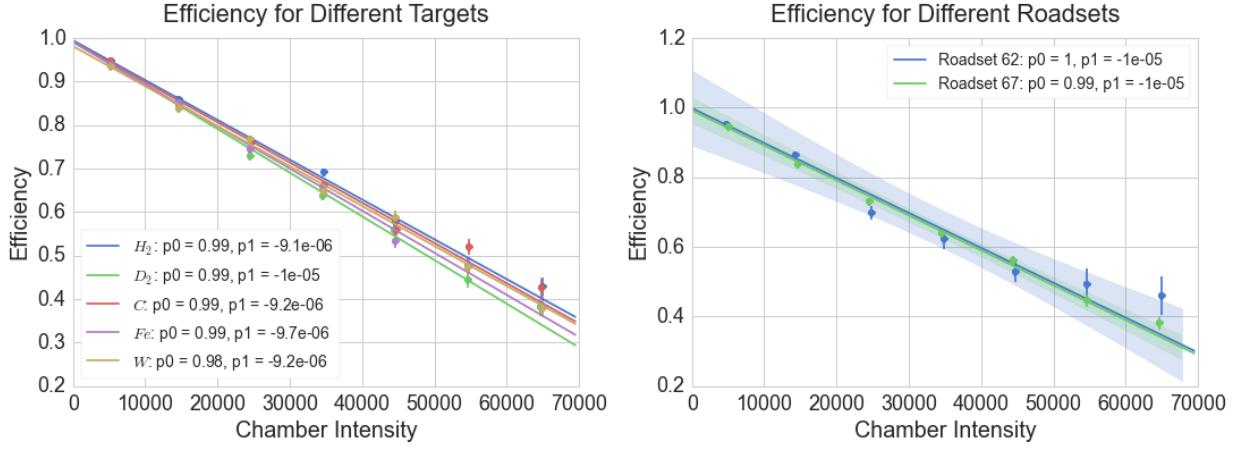


Figure 5.13: Comparing the tracking efficiency curves across (left) the five targets and (right) two temporally different subsets (different roadsets) of data. Confidence bands removed from the target plot for the sake of clarity.

from target t_i in x_F bin x_{F_i} , create a weight as

$$w_i = \frac{1}{kEff_{t_i, x_{F_i}}(I_i)} \quad (5.14)$$

5.4.2 The Naive Empty/None Correction

First a few values must be defined. An overall efficiency based on ‘contributions’ from the Empty/None background can be defined as

$$\epsilon_b = 1 - \frac{Y_b/p_b}{Y_t/p_t} = 1 - \left(\frac{p_t}{p_b} \right) \left(\frac{Y_b}{Y_t} \right) \quad (5.15)$$

where Y is the dimuon *yield* from a given target and p is the integrated “live proton” count for the same given target. Here, $b \in \{\text{Empty, None}\}$ and $t \in \{H, D, C, Fe, W\}$. It is yet to be conclusively decided by the collaboration whether or not the Empty and None targets are to be combined for these calculations.

The next step is to adjust the yields with the $kEff_{t, \bar{x}}(I)$:

$$Y_a \rightarrow Y'_a = \sum_{i: i_t = a} \frac{1}{kEff_{t_i, x_i}(I_i)} \quad (5.16)$$

The yield efficiency when factoring in background then becomes

$$\epsilon'_b = 1 - \left(\frac{p_t}{p_b} \right) \left(\frac{Y'_b}{Y'_t} \right) \quad (5.17)$$

5.4.3 Rate Dependent Empty Correction

Upon inspection, the Empty/None target data has shown to exhibit a rate dependence, with many more events per proton being found at higher intensities. This can be corrected for, but the events should first be weighted by their tracking efficiencies. Unfortunately since there are no *messy/clean* embedded data sets for the empty flask, the “kEfficiency” fit must be extrapolated from the fits for the deuterium and hydrogen targets. If we are to assume that the difference in fits is a function of the beam’s interaction with the target, then we can arrive at the *kEff* exponential fit constants for the empty/none targets (A_E) for each kinematic bin from the exponential fit constants from the hydrogen and deuterium targets ($A_{D/H}$). Here, we only additionally need the percentage of beam that interacts with the target (X_t), along with a few other values characteristic to the targets found in Table 5.6.

$$X_t = 1 - e^{-\frac{L_t}{\lambda_t}} \quad (5.18)$$

$$A_E = A_H - X_H \cdot \frac{\Delta A}{\Delta X} \quad (5.19)$$

$$= A_H - X_H \cdot \frac{A_D - A_H}{X_D - X_H} \quad (5.20)$$

With this fit constant in hand and the kEfficiency correction applied, one can plot the yields of empty flask and none events as a function of chamber intensity. This plot and its shape can be found in Figure *MAKE FIGURE FOR THIS*. Due to the low-statistics nature of the empty and none events, they are combined across all roadsets and kinematics instead of being broken down. As such, this correction is applied solely based on chamber intensity with no other consideration for roadset or kinematics. The fit that is applied is of the form: $f(I) = p_0 \cdot (1 + p_1 \cdot I_c)$. This form is chosen because the scale of the final function will be a calculated from the data used, and as such, p_1 is the only parameter that is needed. The value calculated was:

$$p_1 = (6.576 \pm 1.566) \times 10^{-5} \quad (5.21)$$

The fit function used, $C_{norm}(I_c) = B \cdot C(I_c) = B \cdot (1 + p_1 \cdot I_c)$ can be interpreted as the probability that a given target dimuon actually arose from something other than a target interaction (beam dump, upstream interactions, flask container). Here, B is the normalization factor, where the condition should stand that the sum of $C_{norm}(I_c)$ over all target dimuons should be equal to the number of events from the empty and none targets. Both sums need to be corrected for kEfficiency and scaled by the number of live protons.

$$\sum_T \frac{B \cdot C(I_c)}{P_T \cdot K_T(I_c)} = \sum_E \frac{1}{P_E \cdot K_E(I_c)} \quad (5.22)$$

$$B = \frac{P_T}{P_E} \cdot \frac{\sum_E \frac{1}{K_E(I_c)}}{\sum_T \frac{C(I_c)}{K_T(I_c)}} \quad (5.23)$$

$$\delta B = \frac{P_T}{P_E} \cdot \frac{\sqrt{\sum_E \frac{1}{K_E(I_c)^2}}}{\sum_T \frac{C(I_c)}{K_T(I_c)}} \quad (5.24)$$

This normalization constant is calculated for each target and each roadset. With this in hand, a weight for correcting for the empty target correction can be applied as $(1 - C_{norm}(I_c))$. One can then combine this with the weighting for the kEfficiency correction (as the weighting should combine multiplicatively):

$$\epsilon_{tot} = \epsilon_{kEff} \cdot \epsilon_{empty} = \frac{1 - C_{norm}(I_c)}{K_T(I_c)} \quad (5.25)$$

It is worth noting that this model does not take into account the attenuation of beam throughout the different target materials, which could alter the scale of the correction slightly. However, without having a figure as to how much of the empty target signal comes from the beam dump and how much of it comes from upstream of the target area, it is currently infeasible to appropriately account for the degree to which a beam attenuation factor should be worked in.

Uncertainties that arise as a result of the uncertainties in the $C(I_x)$ fit and the calculation of B are difficult to propagate analytically. What is performed instead is a an independent variation of p_1 and B by -1σ , 0, and $+1\sigma$ and observing the magnitude of which the per nucleon cross section ratio changes as a result. This systematic uncertainty due to this correction is calculated in the Systematic Uncertainties section of this chapter.

5.5 Combinatorial Background

Discuss analysis investigating results for the like-sign reconstruction to estimate the combinatorial background. Hopefully it will be intensity-dependent.

5.6 ld_2 Contamination

Cite Paul's summary of the contamination correction During Run III of the SeaQuest experiment, analysis of the deuterium source used indicated that it was impure to a significant degree. This calls for a significant correction since most measurements, yields, ratios, and analyses assume that the targets are pure. An impure deuterium target changes many factors including the averaged mass number of the fluid and its

Species:	Test:	Runs: 3230-11143		11144-11477		11478-14652	
		1	2	1	2	1	2
H_2		1.23	1.225	0.25	0.25	4.603	4.591
HD		16.187	16.171	8.393	8.391	6.983	6.982
D_2		79.78	79.729	90.871	90.808	80.504	80.443
He		2.20	2.188	0.018	0.014	<0.01	<0.01
CH_4		<0.009	<0.009	<0.007	<0.007	<0.008	<0.008
H_2O		<0.009	0.073	0.038	0.099	0.033	0.061
HDO		0.015	0.026	0.024	0.033	0.021	0.04
D_2O		<0.007	<0.009	<0.007	<0.008	<0.009	<0.012
Ne		0.007	0.065	<0.027	<0.027	<0.03	<0.03
N_2		0.428	0.419	0.323	0.323	6.395	6.382
C_2H_6		<0.006	<0.006	<0.005	<0.005	<0.006	<0.006
O_2		0.085	0.098	0.068	0.068	1.389	1.416
Ar		0.005	0.006	0.004	0.003	0.063	0.064
CO_2		<0.002	<0.002	0.011	0.009	0.01	0.01
C_3H_8		<0.006	<0.006	<0.005	<0.005	<0.006	<0.006

Table 5.7: Results of the isotropic analyses of the liquid deuterium samples. Two tests are applied to each sample, and all numbers are in percentages.

overall density. In this section, the contamination measurements and the correction procedure is discussed.

5.6.1 Contamination Measurements

The SeaQuest deuterium target was emptied and filled several times during data taking. At each of these times, a sample was taken from either the target or from the deuterium gas cylinder used to fill the target. These samples were sent to Los Alamos for isotopic analysis where each sample was analyzed twice. The results of these analyses are shown in Table 5.7. In each case of parallel analyses on a single gas sample, the two results were in agreement. Upon examination of the deuterium compositions, some problems are immediately clear. Specifically, the ratios of H_2 , HD , and D_2 are different between the three samples, but the bottles were supposed to have been filled from the same source many years ago. Second, the some of samples had a significant contribution of other gases, including N_2 and O_2 . This indicates an additional contamination of air. A sample of the first H_2 target was analyzed and found to be $\sim 98\%$ pure with negligible contributions from other species.

The source of the deuterium used at SeaQuest is from a bubble chamber experiment at Fermilab. At some undocumented point in the past, the contents of the bubble chamber were transferred into a tube trailer for storage. At a later time, the D_2 was transferred into cylinders for storage. Discussions with gas storage experts indicate that there was likely no special treatment of the cylinders other than being evacuated when they were filled with D_2 , with the desired special treatment being to “bake” the cylinders, or heat them up

	Run	3230-	11144-	11478
	Range:	11143	11477	14652
Species:				
H_2		1.26	0.25	4.99
HD		16.65	8.44	7.59
D_2		82.08	91.31	87.42

Table 5.8: Renormalized compositions (in percentages) when considering only the hydrogen isotopes.

	Run	3230-	11144-	11478
	Range:	11143	11477	14652
Atom:				
H		9.6	4.5	8.8
D		90.4^{+5}_{-0}	95.5^{+0}_{-5}	91.2^{+4}_{-1}

Table 5.9: Atomic percentages of the deuterium targets with asymmetric uncertainties in the d percentage.

to release residual hydrogen that tends to diffuse into the metal surface of their inner walls. In addition to that, the containers used to take samples from the target or cylinders were also not “baked,” so existing estimates of the contamination may *only* be from the sampling containers – or the contamination could be doubly worsened by contamination in the sampling containers. As such, effects of this contamination are treated as a systematic error with asymmetric uncertainties. After run 14652, a commercial deuterium gas has been used which has an analyzed purity of $> 99.9999\%$, and no systematic uncertainty will need be applied.

Contamination from air in the sample gases is unlikely to have come from the targets. Prior to filling, the targets undergo several cycles of evacuation and refilling with deuterium before the final fill is done. This cycle should flush out all residual air from the target flask. Further, if there was air supply being fed into the target, all but the lowest temperature gases should have been removed by the liquid nitrogen “cold trap.” Aside from hydrogen isotopes, helium is the only other component found in the samples. Since the target temperature were approximately 23 K, any helium would have been in a gas state and not in the bulk liquid in the flask. Thus, the gases from the analysis shown in Table 5.7 can be renormalized to consider only the hydrogen isotopes, which can be seen in Table 5.8.

Until the true compositions can be ascertained, it has been recommend that analyses use the atomic percentages provided by the Los Alamos analysis as shown in Table 5.9, which represents the best understanding of the deuterium at this point. To assign a systematic uncertainty to these numbers, the recommended approach is to use asymmetric uncertainties, assuming that one the extreme values for H and D. Due to the fact that there is a complete inverse correlation between hydrogen and deuterium an uncertainty is assigned to the deuteron percentage only.

5.6.2 Contamination Correction

With the assumption that the heavier atoms are caught in the cold trap and the gaseous helium stays at the top of the flask and away from the path of the proton beam, a contamination correction is implemented based on the atomic percentage composition detailed in Table 5.9. The nuclear interaction length, density, and atomic mass of the liquid deuterium target is corrected based on the contamination. With A as the atomic mass, λ as the nuclear interaction length, ρ as the density, and $F_{H/D}$ as the proportion of H or D in the target, the following are assumed to be approximately true:

$$A = A_H \cdot F_H + A_D \cdot F_D \quad (5.26)$$

$$\rho = \rho_H \cdot F_H + \rho_D \cdot F_D \quad (5.27)$$

$$\frac{1}{\lambda} = \frac{F_H}{\lambda_H} + \frac{F_D}{\lambda_D} \quad (5.28)$$

$$\lambda = \frac{\lambda_H \cdot \lambda_D}{F_H \cdot \lambda_D + F_D \cdot \lambda_H} \quad (5.29)$$

Due to the contamination, some of the signal events with the liquid deuterium target in place will actually come from hydrogen instead of a deuteron. With $\hat{\sigma}_{pH}$ and $\hat{\sigma}_{pD}$ representing the *measured* target cross section, and σ_{pp} and σ_{pn} representing the proton-proton and proton-neutron cross section,

$$\hat{\sigma}_{pH} = a \cdot \sigma_{pp} + b \cdot \sigma_{pn} \quad (5.30)$$

$$\hat{\sigma}_{pD} = c \cdot \sigma_{pp} + d \cdot \sigma_{pn} \quad (5.31)$$

$$\sigma_{pp} = \frac{b \cdot \hat{\sigma}_{pD} - d \cdot \hat{\sigma}_{pH}}{b \cdot c - a \cdot d} \quad (5.32)$$

$$\sigma_{pn} = \frac{a \cdot \hat{\sigma}_{pD} - c \cdot \hat{\sigma}_{pH}}{a \cdot d - b \cdot c} \quad (5.33)$$

Here, the coefficients a, b, c , and d denote the number of protons and neutrons per atom for the different targets. We know that $a = c = 1$, and since the hydrogen target is considered to be pure, $b = 0$. The remaining value, d , is the fraction of deuteron atoms in the sample (F_D above). With this, we can arrive at an expression for $\sigma_{pD}^{\text{pure}}$.

$$\sigma_{pH}^{\text{pure}} = p = \hat{\sigma}_{pH} \quad (5.34)$$

$$n = \frac{\hat{\sigma}_{pD} - \hat{\sigma}_{pH}}{d} \quad (5.35)$$

$$\sigma_{pD}^{\text{pure}} = p + n \quad (5.36)$$

$$\sigma_{pD}^{\text{pure}} = \frac{\hat{\sigma}_{pD} - \hat{\sigma}_{pH} \cdot (1 - d)}{d} \quad (5.37)$$

In practice, this correction is applied bin by bin, utilizing the weighted dimuon yield for each target in place of $\hat{\sigma}$. The d used depends on the run ranges applicable to the roadset being analyzed, and the values are derived from those in Table 5.9.

5.7 x_A Nuclear Scaling Variable

When discussing heavy nuclei and, a certain shift in understanding of scaling variables should be considered. The use of the Bjorken- x scaling variable up until now should be noted to be the parton momentum fraction of the proton $x_p = Q^2/(2m_p\omega)$. It has been recently discussed by Frankfurt & Strikman[45] and implemented by Hen *et al.*[63] that the *structure functions of nucleons bound in nuclei should be extracted in the reference frame of the nucleus*. This is done by considering to the x_A scaling variable, defined as:

$$\frac{x_A}{A} = \frac{Q^2}{2q \cdot P_A} ; \quad x_A = \frac{Q^2}{2q \cdot P_A/A} = \frac{AQ^2}{2\omega m_A} = x_p \cdot \frac{Am_p}{m_A} \quad (5.38)$$

where q and P_A are the 4-momentum vectors of the virtual photon and target nucleus respectively, and m_A is the mass of the target nucleus. For the same values of Q^2 and ω , x_A differs from x_p by the ratio of the bound nucleon mass to the free mass. Therefore, a cross section measured at Q^2 and ω on a nucleus A will depend on the nucleon structure function evaluated at x_A rather than x_p ¹. For converting from the commonly used x_p over to x_A for the different targets, we apply a factor of $\frac{Am_p}{m_A}$, which for deuterium, carbon, iron, and tungsten are *AAAAA, BBBB, CCCCC, and DDDDD, respectively*.

As a result, this means that the standard EMC cross section ratio is actually proportional to the nucleon structure function in nucleus A evaluated at parton momentum x_A divided by the nucleon structure function in deuterium evaluated at parton momentum fraction x_d ($x_d = 2Q^2/2m_d\omega = x_p \cdot 2m_p/m_d$)[63]. For a symmetric nuclei, this can be written as:

$$\frac{2}{A} \cdot \frac{\sigma_{DY}^A(x_p, Q^2)}{\sigma_{DY}^d(x_p, Q^2)} = \frac{F_2^A(x_A, Q^2)}{F_2^d(x_d, Q^2)} \quad (5.39)$$

where $\frac{F_2^A(x_A, Q^2)}{F_2^d(x_d, Q^2)}$ is the ratio of structure functions at the same Q^2 , but at different x . Since we want to compare the structure functions at the same parton momentum fractions, we correct this by using

$$\frac{F_2^A(x_A, Q^2)}{F_2^d(x_d, Q^2)} = \frac{F_2^A(x_A, Q^2)}{F_2^d(x_A, Q^2)} \cdot \frac{F_2^d(x_A, Q^2)}{F_2^d(x_d, Q^2)} \quad (5.40)$$

$$\frac{F_2^A(x_A, Q^2)}{F_2^d(x_A, Q^2)} = \frac{2}{A} \cdot \frac{\sigma_{DY}^A(x_p, Q^2)}{\sigma_{DY}^d(x_p, Q^2)} \cdot \frac{F_2^d(x_d, Q^2)}{F_2^d(x_A, Q^2)} \quad (5.41)$$

where here we have $\frac{F_2^d(x_d, Q^2)}{F_2^d(x_A, Q^2)}$ as a correction factor to arrive at $\frac{F_2^A(x_A, Q^2)}{F_2^d(x_A, Q^2)}$, which is the ratio of structure functions in the different nuclei evaluated at the same parton momentum fraction (which is what we're

¹This parlance is in terms of DIS and not Drell-Yan, but the final correction factor still stands for Drell-Yan cross-section ratio measurements with $Q^2 \rightarrow M_{\gamma^*}^2$ and $\omega \rightarrow ???$

looking to extract). The correction factor can be evaluated using well-known parameterizations of the the structure function of the deuteron (*see L.W. Whitlow, Phys. Lett. B282 1992 and P.E. Bosted, Phys. Rev. C77 2008*). In *Figure refAAA*, the correction factor $\frac{F_2^d(x_d, Q^2)}{F_2^d(x_A, Q^2)}$ is shown as a function of x_p for different targets.

In summary, the traditional EMC per-nucleon cross section ratio measurement has yielded a ratio for $\frac{F_2^A(x_A)}{F_2^d(x_d)}$, when what is really desired is $\frac{F_2^A(x_A)}{F_2^d(x_A)}$. The addition of this $\frac{F_2^d(x_d)}{F_2^d(x_A)}$ correction factor is all that is needed to convert the measurement to the desired result. This correction factor will be applied to the SeaQuest result.

5.8 Isoscalar Corrections for ^{183}W and ^{56}Fe

Give brief correction in approximating $n==p$ conversion. Be sure to show results for both iso-corrected and not!

For asymmetric nuclei ($N \neq Z$), following Aubert *et al.* and Bodek *et al.* [14, 21], an additional isoscalar correction factor (R_{ISO}) is applied to the yield ratio, approximating the ratio to a hypothetical nucleus with equal numbers of protons and neutrons ($N = Z = A/2$):

$$R_{ISO}(x) = \frac{A}{2} \cdot \frac{F_2^p(x, Q^2) + F_2^n(x, Q^2)}{Z \cdot F_2^p(x, Q^2) + N \cdot F_2^n(x, Q^2)} = \frac{A}{2} \cdot \frac{1 + R_{np}(x, Q^2)}{Z + N \cdot R_{np}(x, Q^2)} \quad (5.42)$$

$$\frac{\sigma_{DY}^A(x_A, Q^2)_{ISO}}{\sigma_{DY}^d(x_A, Q^2)_{ISO}} = \frac{\sigma_{DY}^A(x_A, Q^2)}{\sigma_{DY}^d(x_A, Q^2)} \cdot R_{ISO}(x_A) \quad (5.43)$$

where F_2^p and F_2^n are the free nucleon structure functions and $R_{np}(x, Q^2) = F_2^n(x, Q^2)/F_2^p(x, Q^2)$ ($R_{np}(x) \approx 1 - 0.8x$). The free nucleon structure functions are extracted from world data on DIS scattering [*citation needed*] (*Bodek Phys Rev D20 1979, Arrington Phys G Nucl Part Phys 36 2009*) off of deuterium and hydrogen with corrections accounting for Fermi motion of nucleons in deuterium.

Applying this correction to the structure function ratio extracted in Equation 5.41, the same extraction for asymmetric nuclei becomes (*add vertical space between equations.*)

$$\frac{F_2^A(x_A, Q^2)}{F_2^d(x_A, Q^2)} = \frac{2}{A} \cdot \frac{\sigma_{DY}^A(x_p, Q^2)}{\sigma_{DY}^d(x_p, Q^2)} \cdot \frac{F_2^d(x_d, Q^2)}{F_2^d(x_A, Q^2)} \cdot R_{ISO}(x_A) \quad (5.44)$$

$$\frac{F_2^A(x_A, Q^2)}{F_2^d(x_A, Q^2)} = \frac{2}{A} \cdot \frac{\sigma_{DY}^A(x_p, Q^2)_{ISO}}{\sigma_{DY}^d(x_p, Q^2)_{ISO}} \cdot \frac{F_2^d(x_d, Q^2)}{F_2^d(x_A, Q^2)} \cdot \frac{R_{ISO}(x_A)}{R_{ISO}(x_p)} \quad (5.45)$$

5.9 Summary of Analysis Steps

Due to the dynamic nature of the SeaQuest analysis at the time of the writing of this document, it is important to lay out the precise procedure of the application of all of these analysis steps. This ensures an ability to reproduce results, understand differences between methods, and allows those new to analysis to have a roadmap for getting started with analysis. Pragmatically speaking, the process taken in this analysis going from the raw dimuon yields to the final extraction of $F_2^A(x)/F_2^d(x)$ can be outlined in several finite steps:

- Get raw yields after all analysis cuts (Section X).
- Get live proton count (Section X).
- Calculate the kEfficiency for targets H_2 , D_2 , C , Fe , and W (Section X).
- Extrapolate kEfficiency for Empty/None target.
- Calculate $C(I_c)$ and B_T empty target correction fit and parameters (Section X).
- Calculate weighted yields:

$$Y'_T = \sum_T \frac{1 - B_T \cdot C(I_c)}{K_T(I_c)} ; \quad \delta Y'_T = \sqrt{\sum_T \left(\frac{1 - B_T \cdot C(I_c)}{K_T(I_c)} \right)^2} \quad (5.46)$$

- **???** Adjust yields by combinatoric background correction factor.
- Adjust the deuterium yields for the hydrogen contamination (Equation X).
- Take ratios of these adjusted yields Y'_A/Y'_d
- Apply the target-to-target normalization factor to the ratio:

$$N_{A/D} = \left(\frac{N_0^D \cdot \text{lt}^D \cdot T^D(\xi)}{N_0^A \cdot \text{lt}^A \cdot T^A(\xi)} \right) \cdot \left(\frac{2 \cdot n_D \cdot L^D}{A \cdot n_A \cdot L^A} \right) \cdot \frac{\varepsilon^D}{\varepsilon^A} \cdot \frac{\bar{\Omega}^D}{\bar{\Omega}^A} \cdot \frac{\epsilon^D}{\epsilon^A} \quad (5.47)$$

- Apply the $\frac{F_2^d(x_d, Q^2)}{F_2^d(x_A, Q^2)}$ correction factor to the ratio.
- With asymmetric nuclei ($^{56}Fe, ^{138}W$), apply $\frac{R_{ISO}(x_A)}{R_{ISO}(x_p)}$ isoscalar correction factor.

It is understood that this procedure is general enough that it can be performed with any kinematic binning. Here, the x_A conversion and the isoscalar correction are applied to the calculated, binned, adjusted yield

ratios. It is also perhaps suitable to apply these to each event with precise considerations for each event's invariant mass.

5.10 Systematic Uncertainties

Chapter 6

Results

CHAPTER STATUS: NOT STARTED

6.1 Dimuon Yield Ratios

Give EMC ratios, maybe parton energy loss ratios. If there's time, maybe quarkonia ratios. Show for separate roadsets and with/without corrections. Show heavy-to-heavy ratios, too.

Chapter 7

Discussion

CHAPTER STATUS: NOT STARTED

Chapter 8

Conclusions

CHAPTER STATUS: NOT STARTED

We conclude that Bryan Dannowitz likes coffee.

Appendix A

Topic A

Appendix B

Topic B

References

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