## Programming Project 5 Exercises

1. Let  $\overline{\mathbf{S}}$ ,  $\overline{\mathbf{T}}$ , and  $\overline{\mathbf{V}}$  be one-electron integrals with respect to the spatial AO basis  $\{\chi_{\mu}\}$ .

$$\overline{S}_{\mu\nu} = \langle \chi_{\mu} | \chi_{\nu} \rangle \qquad \overline{T}_{\mu\nu} = -\frac{1}{2} \langle \chi_{\mu} | \nabla_{1}^{2} | \chi_{\nu} \rangle \qquad \overline{V}_{\mu\nu} = \sum_{A} \langle \chi_{\mu} | \frac{Z_{A}}{|\mathbf{r}_{1} - \mathbf{R}_{A}|} | \chi_{\nu} \rangle \qquad (1)$$

Explain why their spin-orbital counterparts S, T, V have the following block-diagonal structure

$$\mathbf{S} = \begin{bmatrix} \overline{\mathbf{S}} & 0 \\ 0 & \overline{\mathbf{S}} \end{bmatrix} \qquad \mathbf{V} = \begin{bmatrix} \overline{\mathbf{V}} & 0 \\ 0 & \overline{\mathbf{V}} \end{bmatrix}$$
 (2)

in the spin-AO basis  $\{\chi_{\mu}\alpha\} \cup \{\chi_{\mu}\beta\}$ .

2. The two-electron integrals  $(\mu\nu|\rho\sigma)$  make a 4-index tensor, which can be viewed as a matrix  $\overline{\mathbf{G}} = [\overline{\mathbf{G}}_{\mu\nu}]$  of matrices  $\overline{\mathbf{G}}_{\mu\nu} = [(\mu\nu|\rho\sigma)]$ . Explain why the spatial AO basis integral tensor  $\overline{\mathbf{G}}$  is related to its spin-AO counterpart  $\mathbf{G}$  via

$$\mathbf{G} = \begin{bmatrix} \overline{\mathbf{G}} & 0 \\ 0 & \overline{\mathbf{G}} \end{bmatrix} \qquad (\mathbf{G})_{\mu\nu} = \begin{bmatrix} (\overline{\mathbf{G}})_{\mu\nu} & 0 \\ 0 & (\overline{\mathbf{G}})_{\mu\nu} \end{bmatrix}$$
(3)

i.e. it has a block-diagonal structure with respect to each pair of indices. Note that this is one case where the index-ordering of chemist's notation has some advantages over physicist's notation.

3. The spin-orbital Fock operator is given by

$$\hat{f} = \hat{h} + \sum_{i=1}^{n} (\hat{J}_i - \hat{K}_i) \tag{4}$$

where  $\hat{h}$  is the one-electron ("core") Hamiltonian and  $\hat{J}_i$  and  $\hat{K}_i$  are Coulomb and exchange operators. Show that its matrix elements with respect to the spin-orbital basis  $\{\psi_p\}$  are given by

$$f_{pq} = h_{pq} + \sum_{i=1}^{n} \langle pi || qi \rangle . \tag{5}$$

4. Show that, if  $C_{\mu p}$  are the expansion coefficients of  $\psi_p$  in the spin-AO basis  $\{\xi_{\mu}\}=\{\chi_{\mu}\alpha\}\cup\{\chi_{\mu}\beta\}$ 

$$\psi_p = \sum_{\mu} \xi_{\mu} C_{\mu p} \tag{6}$$

then the spin-AO basis Fock matrix elements are given by

$$f_{\mu\nu} = h_{\mu\nu} + \sum_{\rho\sigma} \langle \mu\rho || \nu\sigma \rangle D_{\rho\sigma} \qquad \qquad D_{\mu\nu} = \sum_{i=1}^{n} C_{\mu i}^{*} C_{\nu i} \qquad (7)$$

where  $D_{\mu\nu}$  is called the Hartree-Fock density matrix.<sup>1</sup>

<sup>&</sup>lt;sup>1</sup>This is the spin-AO-basis representation of the one-particle reduced density matrix of a Hartree-Fock determinant.

5. Show that the electronic part of the Hartree-Fock energy

$$\langle \Phi | \hat{H}_e | \Phi \rangle = \sum_{i=1}^{n} h_{ii} + \sum_{i < j}^{n} \langle ij | | ij \rangle$$
 (8)

can be expanded in terms of spin-AO basis integrals as follows.

$$\langle \Phi | \hat{H}_e | \Phi \rangle = \sum_{\mu\nu} h_{\mu\nu} D_{\mu\nu} + \frac{1}{2} \sum_{\mu\nu\rho\sigma} \langle \mu\rho | |\nu\sigma\rangle D_{\mu\nu} D_{\rho\sigma}$$
(9)