

# paper reading notes of

## *nonclassical photon sources*

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# Basic Concept

## 1 Jaynes-Cummings model

Jaynes-Cummings(JC) model is used to describe the coupling between atom and light within the cavity. It use quantum theory to describe both light and atom. The simplest JC model's Hamiltonian can be write as:

$$\hat{H} = \hbar\omega\hat{a}^\dagger\hat{a} + \frac{1}{2}\hbar\omega_A\hat{\sigma}_z - \hbar g(\hat{a}\hat{\sigma}^\dagger + \hat{a}^\dagger\hat{\sigma}) \quad (1.1)$$

It is obviously that it couples atom state  $|g\rangle$  and  $|e\rangle$  with number state of light  $|n\rangle$ . We can easily calculate martix elements like:

$$\langle n+1, g|\hat{H}|n+1, g\rangle = (n+1)\omega - \frac{1}{2}\omega_A \quad (1.2)$$

Therefore we can rewrite JC Hamiltonian into  $2 \times 2$  matrix as:

$$\hat{H} \begin{pmatrix} |n+1, g\rangle \\ |n, e\rangle \end{pmatrix} = \hbar \begin{pmatrix} (n+1)\omega - \frac{1}{2}\omega_A & -g\sqrt{n+1} \\ -g\sqrt{n+1} & n\omega + \frac{1}{2}\omega_A \end{pmatrix} \begin{pmatrix} |n+1, g\rangle \\ |n, e\rangle \end{pmatrix} \quad (1.3)$$

**Proof.** Now we start to solve eigenvalue problem!

$$\left[ (n+1)\omega - \frac{1}{2}\omega_A - \lambda \right] \left[ n\omega + \frac{1}{2}\omega_A - \lambda \right] - g^2(n+1) = 0 \quad (1.4)$$

$$\lambda^2 - (2n+1)\omega\lambda + (n+1)^2\omega^2 - \frac{\delta^2}{4} - \frac{\Omega_n^2}{4} = 0 \quad (1.5)$$

Easily to derive that:

$$\Delta = \delta^2 + \Omega_n^2 \quad (1.6)$$

$$\lambda = \frac{(2n+1)\omega \pm \sqrt{\Delta}}{2} \quad (1.7)$$

$$= (n + \frac{1}{2})\omega \pm \frac{1}{2}\sqrt{\delta^2 + \Omega_n^2} \quad (1.8)$$

$$= (n + \frac{1}{2})\omega \pm \frac{1}{2}\delta^2 + \Omega_n^2 \quad (1.9)$$

where

$$\delta = \omega - \omega_A \quad \Omega_n = 2g\sqrt{n+1} \quad (1.10)$$

□

This mean the state  $|n+1, g\rangle$  and  $|n, e\rangle$  are couple and struct two new state, whose eigenvalues is  $\lambda_{\pm} = (n + \frac{1}{2})\omega \pm \frac{1}{2}\Delta_n$  due to the interactions between light field and atom. And now we can findout the eigenstate of this coupling system.

**Proof.**

$$\hat{H}_n|n, \pm\rangle = \lambda_{\pm}|n, \pm\rangle \quad (1.11)$$

we can let  $|n, +\rangle$  write as:

$$|n, +\rangle = \begin{pmatrix} a \\ b \end{pmatrix} \quad (1.12)$$

then we need to solve the equation to find a and b :

$$\left[ (n+1)\omega - \frac{1}{2}\omega_A \right] a - (g\sqrt{n+1})b = \left[ (n + \frac{1}{2})\omega + \frac{1}{2}\sqrt{\delta^2 + \Omega_n^2} \right] a \quad (1.13)$$

$$(\omega - \omega_A + \sqrt{\delta^2 + \Omega_n^2})a = (2g\sqrt{n+1})b \quad (1.14)$$

$$(\delta - \Delta_n)a = \Omega_n b \quad (1.15)$$

where

$$\Delta_n = \sqrt{\delta^2 + \Omega_n^2} \quad (1.16)$$

the possible solution of a,b is :

$$a = \frac{\Omega_n}{\sqrt{(\delta - \Delta_n)^2 + \Omega_n^2}} \quad (1.17)$$

$$b = \frac{\delta - \Delta_n}{\sqrt{(\delta - \Delta_n)^2 + \Omega_n^2}} \quad (1.18)$$

Therefore

$$|n, +\rangle = \frac{\Omega_n}{\sqrt{(\delta - \Delta_n)^2 + \Omega_n^2}} |n+1, g\rangle + \frac{\delta - \Delta_n}{\sqrt{(\delta - \Delta_n)^2 + \Omega_n^2}} |n, e\rangle \quad (1.19)$$

In the same way, we can attain :

$$|n, -\rangle = \frac{\Delta_n - \delta}{\sqrt{(\delta - \Delta_n)^2 + \Omega_n^2}} |n+1, g\rangle + \frac{\Omega_n}{\sqrt{(\delta - \Delta_n)^2 + \Omega_n^2}} |n, e\rangle \quad (1.20)$$

for convenience, we better define :

$$\cos \Theta = \frac{\Omega_n}{\sqrt{(\delta - \Delta_n)^2 + \Omega_n^2}}, \quad \sin \Theta = \frac{\Delta_n - \delta}{\sqrt{(\delta - \Delta_n)^2 + \Omega_n^2}} \quad (1.21)$$

then we can simplify eigenstate as<sup>1</sup> :

$$|n, +\rangle = \cos \Theta |n+1, g\rangle - \sin \Theta |n, e\rangle \quad (1.22)$$

$$|n, -\rangle = \sin \Theta |n+1, g\rangle + \sin \Theta |n, e\rangle \quad (1.23)$$

□

<sup>1</sup> This solution are different from the textbook of JC model. It may has some wrong with that textbook.

## 2 Correlation function

## 3 Master equation

# Research Background

## 4 Original motivation of this field

In the applications of quantum information science such as: quantum computing, cryptography and metrology, generate and manipulate single photon is a necessary technique. The biggest motivation to implement high quality single photon source is that the birth of quantum computer will bring an extraordinary impact on all over the world.

## 5 Important achievements

Up to now, several systems including atom-cavity coupled with micro-cavity system, single quantum dot integrated with photonic crystal cavity, optical fibers and surface plasmons had been demonstrated that it can observe and manipulate single photon.

## 6 Applications potential of this field

# Study of scheme

## 7 Atom-Cavity coupled

Deterministic single-photon sources have been realized with neutral atoms, embedded molecules, trapped ions, quantum dots and defect centres. But different applications have different requirements. For example, quantum computing or quantum networking require that photons must also be indistinguishable and high efficiency. Unfortunately, high efficiency is hard to obtain in free space where atoms only contain between two lenses. But strongly coupling the radiating object to an optical microcavity can raise up the efficiency.

Compare to ions, atoms are largely immune to perturbations like electric patch fields close to dielectric mirrors.

Markus *et al.* report a scheme of Rb atom which can save Rb atoms in the cavity up to 30s for producing up to 300,000 photons per atom.<sup>[1]</sup>

## 8 Quantum dot - Cavity coupled

Due to their corresponding atomic-like discrete energy levels, semiconductor quantum dots are often given the moniker artificial atoms. And thanks to its property of designable, quantum dots based single photon source has the band from IR at  $1.55\mu\text{m}$  to the deep UV region at  $< 280\text{nm}$ . More than this, single photon emission from semiconductor QDs has also been realized at temperatures up to room temperature and beyond.

PART

II

PART

III

## 8.1 In the Original region 1300nm

Miyazawa *et al.* use InAs QDs obtain lowest  $g^2(0) = (4.4 \pm 0.2) \times 10^{-4}$ . [2] About 200K QKD require  $g^2(0)$  down to the value  $10^{-4}$  and the parameters  $\langle n \rangle$  and the radiative lifetime must be improved to 0.175 and 0.12 ns, respectively. Miyazawa *et al.* already achieve one of the requirements.

## 8.2 The visible

Quantum dots formed from InGaN, InP, and various combinations of type II/VI semiconductors such as CdSe/ZnSe can be used to generate single photons with wavelengths in the visible region of the spectrum. Moreover, in contrast to the long wavelength emitters discussed above, single photon emission in the visible and ultraviolet has regularly been reported from QDs at ambient temperatures of 300 K and above.

Fedorych *et al.* realized room temperature single photon emission from CdSe/ZnSSe quantum dots. They reduce the correlation function  $g^{(2)}(0)$  down to  $0.16 \pm 0.15$  at  $T = 300K$  [3].

**Definition 8.1** (Advantage). Can be obtain under room temperature, which is an key point to achieve widely applications. [Drawback] Due to high temperature, the efficiency is relatively low.

## 8.3 The UV region

Tamariz *et al.* demonstrated GaN QD single photon emitter can be operated at 300K with  $g^{(2)}(0) = 0.17 \pm 0.08$ . In addition, photon emission rates can raise up to  $6 \times 10^6 s^{-1}$  while  $g^{(2)}(0) \leq 0.5$  is maintained [4].

# 9 Optical fibers

# 10 Surface plasmons

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