High Performance
Computing for Weather
and Climate (HPC4WC)

Content: Introduction Lecturer: Oliver Fuhrer

Block course 701-1270-00L

Summer 2025

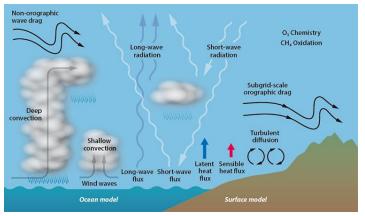


Learning goals

- Understand why weather and climate simulation requires high-performance computing (HPC)
- Understand why stencil computations are a key algorithmic motif of weather and climate models
- Understand a simple stencil program:
 higher-order monotonic diffusion (Xue 2000, MWR)

Modeling the Earth system

Physical understanding



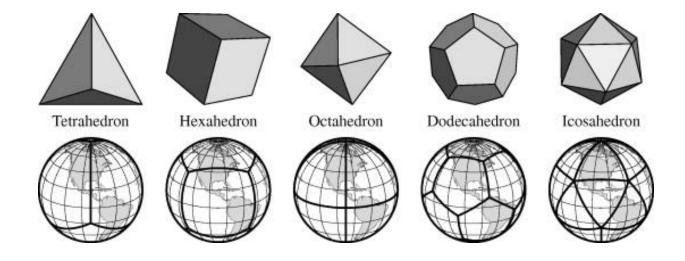


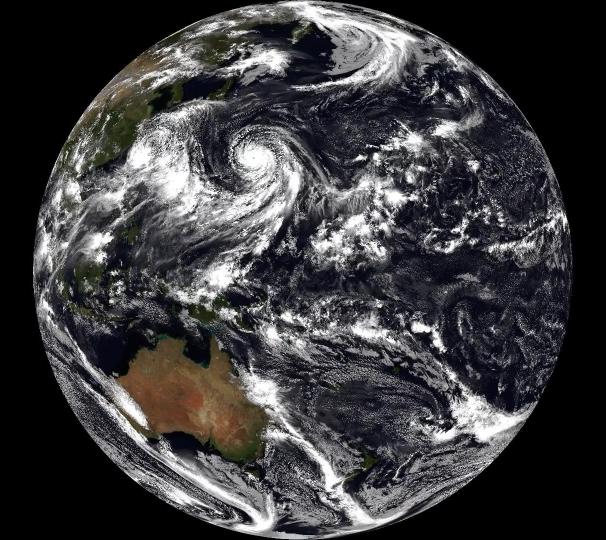
Governing equations (e.g. atmosphere)

$$\begin{split} \frac{d}{dt}\mathbf{v} &= -2\mathbf{\Omega} \times \mathbf{v} - \frac{1}{\rho}\nabla_3 p + \mathbf{g} + \mathbf{F} & \text{Conservation of momentum} \\ C_v \frac{d}{dt}\left(\rho q\right) + p \frac{d}{dt}\left(\frac{1}{\rho}\right) &= J & \text{Conservation of energy} \\ \frac{\partial}{\partial t}\left(\rho\right) &= -\nabla_3 \cdot \left(\rho \mathbf{v}\right) & \text{Conservation of air mass} \\ \frac{\partial}{\partial t} &= -\nabla_3 \cdot \left(\rho \mathbf{v}q\right) + \rho\left(E - C\right) & \text{Continuity of water vapor mass} \\ p &= \rho RT & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial t} &= \log T & \text{Equation of state} \\ \frac{\partial}{\partial$$

Variables: $\{\mathbf{v}, p, T, \rho, q\}$

Grids on a sphere



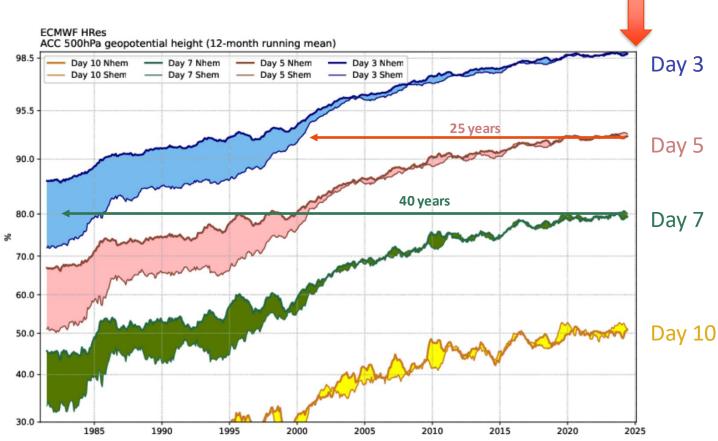




6 x 3072² x 80

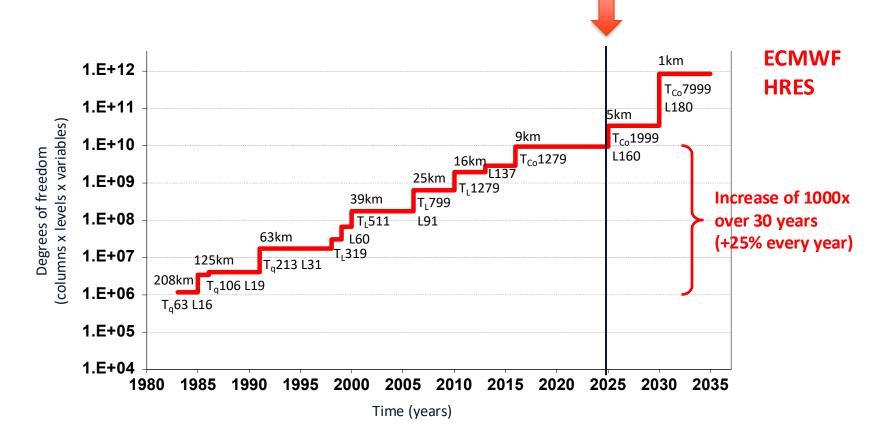
Forecast skill over time



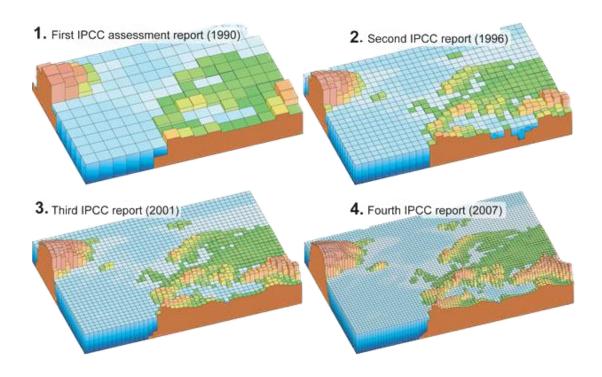


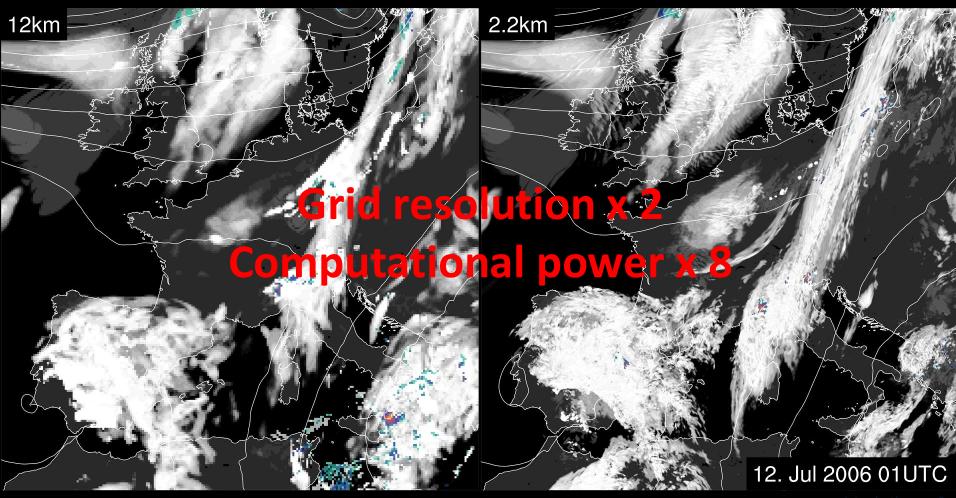
Grid resolution drives skill





Grid resolution in the IPCC





Example applications

 Swiss national weather forecast (ICON-CH1-EPS)

1.15 x 10⁸ x 80 gridpoints

 $\Delta x = 1.05 \text{ km}$

 $\Delta t = 10 s$

11 members

80 x faster than real time (0.2 SYPD)

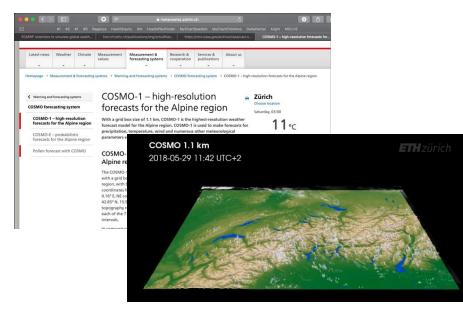
 Global coupled climate model with 1.25 km resolution (ICON Sapphire)

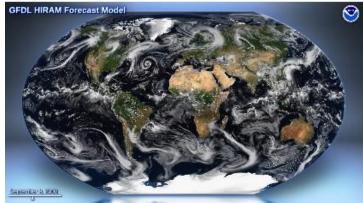
 $3.26 \times 10^8 \times 80$ gridpoints

 $\Delta x = 1.25 \text{ km}$

 $\Lambda t = 10 s$

1000 x faster than real time (3 SYPD)





Why HPC?

 Swiss national weather forecast (ICON-CH1-EPS)

 $1.15 \times 10^8 \times 80 \text{ grid points}$

 $\Delta x = 1.05 \text{ km}$

 $\Delta t = 10 s$

11 members

80 x faster than real time

 Global coupled climate model with 1.25 km resolution (ICON Sapphire)

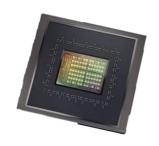
 $3.26 \times 10^8 \times 80$ gridpoints

 $\Delta x = 1.25 \text{ km}$

 $\Delta t = 10 s$

1000 x faster than real time

Currently only 100 x feasible!



NVIDIA Grace CPU @3.44 Ghz, 250 W

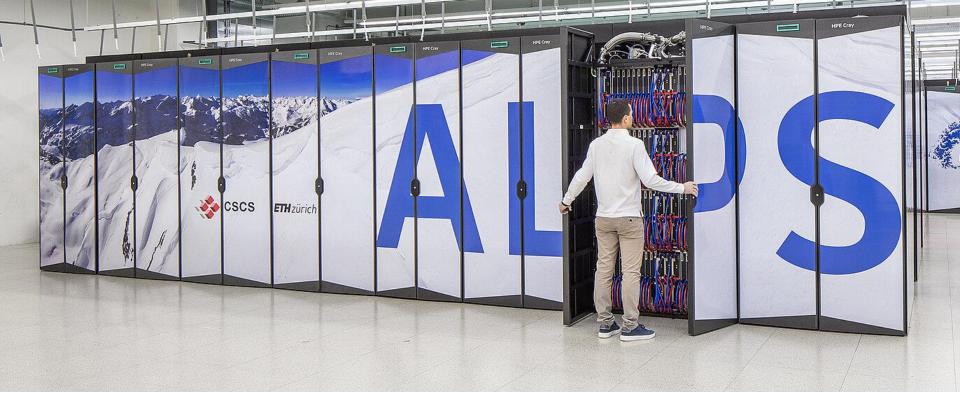
~ 1'000 CPUs

~ **250 kW** (= 320 households)

~ 500'000 CPUs

~ **125 MW** (approx. 10% of Gösgen)





14'000 CPUs

(~ 15 x more dense!)

Supercomputer Architecture

72 per socket

(Numbers are for Alps GH200 nodes and vary from system to system) Day 3 **System Cabinet** Day 2 24 per system **Blade** Day 1 56 per cabinet Node 2 per blade **Socket** 4 per node Core

High Performance Computing is the practice of aggregating computing power in a way that delivers much higher performance than one could get out of a typical desktop computer or workstation in order to solve large problems in science, engineering, or business.

Lab Exercises

01-roofline-model.ipynb

- Learn about performance metrics and how to compute theoretical peak values.
- Learn about arithmetic intensity and performance limiters.

02-stencil-program.ipynb

- Determine arithmetic intensity of a stencil program.
- Apply a performance profiling tool to gain insight into performance.
- Show limitations of the von Neumann model for understanding performance.

03-caches-data-locality.ipynb

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 higher-order monotonic diffusion (Xue 2000, MWR)

Higher-order diffusion (Xue 2000, MWR)

- Atmospheric and ocean models often need some form of numerical filtering to control small-scale noise.
- Xue 2000 (Monthly Weather Review) published a class of higher-order monotonic filters that are frequently used.
- We will focus on the 4th-order non-monotonic diffusion for this course.

2853 XUE

High-Order Monotonic Numerical Diffusion and Smoothing

MING XUE

Center for Analysis and Prediction of Storms, University of Oklahoma, Norman, Oklahoma

(Manuscript received 28 May 1999, in final form 17 September 1999)

ABSTRACT

High-order numerical diffusion is commonly used in numerical models to provide scale selective control over small-scale noise. Conventional high-order schemes have undestrable side effects, however: they can introduce noise themselves. Two types of monotonic high-order diffusion schemes are proposed. One is based on flux correction/limiting on the corrective fluxes, which is the difference between a high-order (fourth order and above) diffusion scheme and a lower-order (typically second order) one. Overshooting and undershooting found in the solutions of higher-order diffusions near sharp gradients are prevented, while the highly selective property of

The second simpler (flux limited) scheme simply ensures that the diffusive fluxes are always downgradient; the second suppose (max manners) secrete shappy shared use the suppose and a solution in 1D cases as and otherwise, the fluxes are set to zero. This much simpler scheme yields as good a solution in 1D cases as and better solutions in 2D than the one using the first more elaborate flux limiter. The scheme also preserves

monotonicity in the solutions and is computational much more efficient. The simple flux-limited fourth- and sixth-order diffusion schemes are also applied to thermal bubble convection. It is shown that overshooting and undershooting are consistently smaller when the flux-limited version of the high-order diffusion is used, no matter whether the advection scheme is monotonic or not. This conclusion applies to both scalar and momentum fields. Higher-order monotonic diffusion works better and even more so when used together with monotonic advection.

1. Introduction

AUGUST 2000

Most numerical models employ numerical diffusion or computational mixing to control small-scale (near two grid intervals in wavelength) noise that can arise from numerical dispersion, nonlinear instability, discontinuous physical processes, and external forcing. Such diffusion cannot always be substituted for by phys-

resents other processes or sources. The last term on the right-hand side (rhs) is the added diffusion term and n(=0, 2, 4, 6, ...) denotes the order. Here α_x is called diffusion coefficient. Diffusion or smoothing can also be introduced by periodically applying a spatial filter or smoother to the predicted field (e.g., Shapiro 1970, 1975). Its effect is often equivalent to applying diffusion

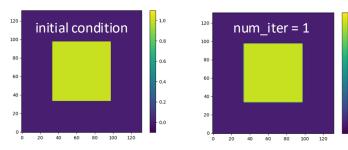
in the prognostic equation for a period of time. Second, and fourth-order formulations are most com-

Governing equations

High-order monotonic diffusion (see paper)

Simplifications (see paper)

- ignore other processes (S=0)
- 4th-order (n=4)
- without limiter



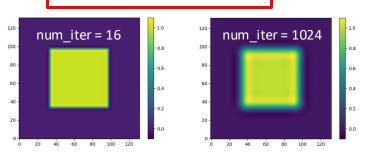
$$\frac{\partial \phi}{\partial t} = S + (-1)^{n/2+1} \alpha_n \nabla^n \phi, \tag{1}$$

$$S = 0$$
 $n = 4$ horizontal

$$\frac{\partial \phi}{\partial t} = -\alpha_4 \nabla_h^4 \phi$$

$$= -\alpha_4 \Delta_h^2 \phi \qquad \Delta_h = \nabla_h^2$$

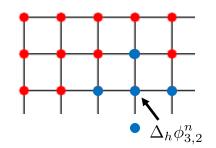
$$\frac{\partial \phi}{\partial t} = -\alpha_4 \Delta_h \left(\Delta_h \phi \right)$$



Discretization: Finite difference method (FDM)

Computational grid

Store values at gridpoints (i,j)



$$x_{i,j} = i \Delta x \quad y_{i,j} = j \Delta y$$
$$t^n = n \Delta t$$
$$\phi_{i,j}^n = \phi(x_{i,j}, y_{i,j}, t^n)$$

Spatial disretization

2nd-order centered

$$\Delta_h \phi_{i,j}^n \approx \left(-4 \phi_{i,j}^n + \phi_{i-1,j}^n + \phi_{i+1,j}^n + \phi_{i,j-1}^n + \phi_{i,j+1}^n \right) / \Delta x^2$$

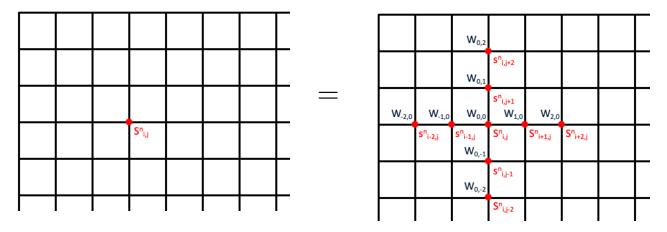
Time discretization

1st-order forward (Euler)

$$\partial_t \phi_{i,j}^n \approx \left(\phi_{i,j}^{n+1} - \phi_{i,j}^n\right) / \Delta t$$

Stencil computation (5x5, sparse)

$$s_{i,j}^{n+1} = \sum_{i_{\mathrm{rel}},j_{\mathrm{rel}}} w_{i_{\mathrm{rel}},j_{\mathrm{rel}}} s_{i+i_{\mathrm{rel}},j+j_{\mathrm{rel}}}^n$$



weighted sum of grid values compact neighborhood same pattern ∀ gridpoints A stencil computation is an algorithmic motif where the value at a certain grid point is computed from a compact neighborhood of gridpoints in it's vicinity on the computational grid using the same pattern for every gridpoint.

Stencil program $\frac{\partial \phi}{\partial t} = -\alpha_4 \, \Delta_h \, (\Delta_h \phi)$ Discretization

$$\frac{\partial \phi}{\partial t} = -\alpha_4 \, \Delta_h \, (\Delta_h \phi)$$



$$\Delta_h \phi_{i,j}^n \approx \left(-4 \phi_{i,j}^n + \phi_{i-1,j}^n + \phi_{i+1,j}^n + \phi_{i,j-1}^n + \phi_{i,j+1}^n \right) / \Delta x^2$$

$$\partial_t \phi_{i,j}^n \approx \left(\phi_{i,j}^{n+1} - \phi_{i,j}^n \right) / \Delta t$$

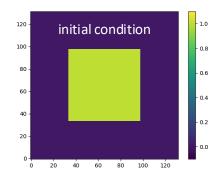
$$\alpha = \alpha_4 \frac{\Delta t}{\Delta x^2}$$
Impler

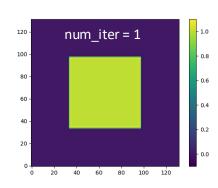


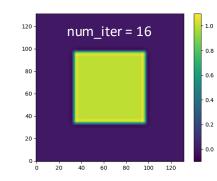
```
do num iter:
  tmp = lap(in)
  out = lap(tmp)
  out = in - \alpha out
  swap(in, out)
```

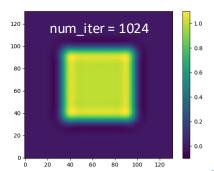
Structure of stencil_2d.F90 / stencil_2d.py

$$\begin{split} \mathsf{tmp}_{i,j,k} &= lap(\mathtt{in}_{i,j,k}) &= -4\,\mathtt{in}_{i,j,k} + \mathtt{in}_{i-1,j,k} + \mathtt{in}_{i+1,j,k} + \mathtt{in}_{i,j-1,k} + \mathtt{in}_{i,j+1,k} \\ \mathsf{out}_{i,j,k} &= lap(\mathtt{tmp}_{i,j,k}) &= -4\,\mathtt{tmp}_{i,j,k} + \mathtt{tmp}_{i-1,j,k} + \mathtt{tmp}_{i+1,j,k} + \mathtt{tmp}_{i,j-1,k} + \mathtt{tmp}_{i,j+1,k} \\ \mathsf{out}_{i,j,k} &= \mathtt{in}_{i,j,k} - \alpha\,\mathtt{out}_{i,j,k} \\ \\ \mathsf{if} \quad \mathtt{iter} < \mathtt{num_iter} \quad : \quad \mathtt{swap}\,(\mathtt{in},\mathtt{out}) \end{split}$$





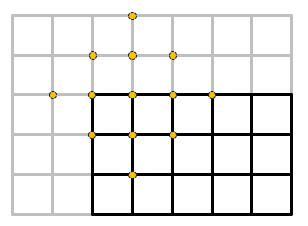


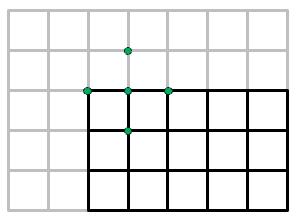


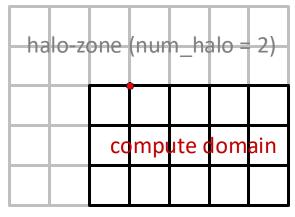
Halo points

halo_update(in)

$$\begin{array}{lcl} \underline{\operatorname{tmp}_{i,j,k}} = lap(\operatorname{in}_{i,j,k}) & = & -4 \underline{\operatorname{in}_{i,j,k}} + \underline{\operatorname{in}_{i-1,j,k}} + \underline{\operatorname{in}_{i+1,j,k}} + \underline{\operatorname{in}_{i,j-1,k}} + \underline{\operatorname{in}_{i,j+1,k}} \\ \\ \operatorname{out}_{i,j,k} = lap(\operatorname{tmp}_{i,j,k}) & = & -4 \underline{\operatorname{tmp}_{i,j,k}} + \underline{\operatorname{tmp}_{i-1,j,k}} + \underline{\operatorname{tmp}_{i+1,j,k}} + \underline{\operatorname{tmp}_{i,j-1,k}} + \underline{\operatorname{tmp}_{i,j+1,k}} \\ \\ \operatorname{out}_{i,j,k} & = & \operatorname{in}_{i,j,k} - \alpha \operatorname{out}_{i,j,k} \end{array}$$



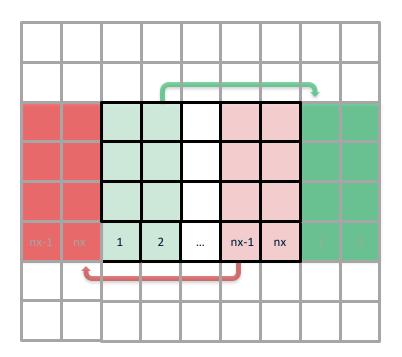


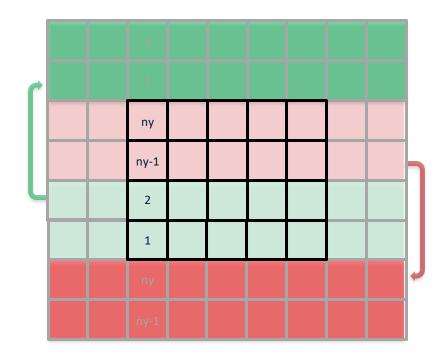


tmp(i,j,1)

out(i,j,1)

Halo update (periodic)





Not stencils

- (Some) implicit methods (e.g. vertical advection)
- Spectral methods (e.g. IFS dynamical core)
- Other algorithmic motifs
 - Reductions
 - Searches
 - -

Summary

- Spatial resolution drives required computational power $1/\Delta x^3$, this is why we need supercomputers for modeling weather and climate.
- While models use a myriad of different numerical techniques to solve the governing PDEs, the main algorithmic motif resulting from these methods are stencil computations.

Our course will **focus exclusively on stencil computations**, but the concepts are more general and transferable.

Lab Exercises

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