

3 OCTOBER 2024

ASE 367K: FLIGHT DYNAMICS

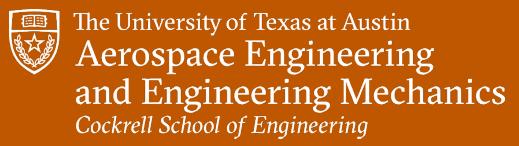
TTH 09:30-11:00
CMA 2.306

JOHN-PAUL CLARKE

Ernest Cockrell, Jr. Memorial Chair in Engineering, The University of Texas at Austin

Topics for Today

- Topic(s):
 - Translational Position
 - Rotational Orientation
 - Angular Momentum
 - Rate of Change of Angular Momentum



TRANSLATIONAL POSITION

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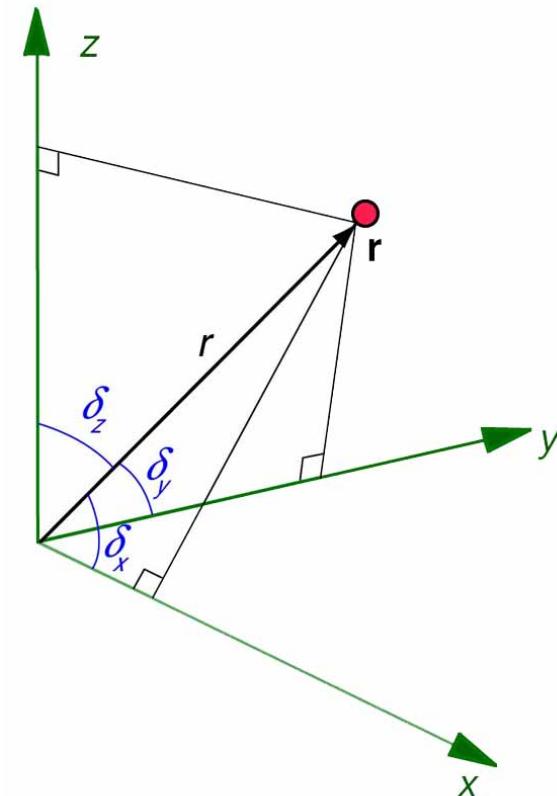
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Position of a Particle

Projections of vector magnitude on three axes

$$\mathbf{r} = \begin{bmatrix} x \\ y \\ z \end{bmatrix} = r \begin{bmatrix} \cos \delta_x \\ \cos \delta_y \\ \cos \delta_z \end{bmatrix}$$

$$\begin{bmatrix} \cos \delta_x \\ \cos \delta_y \\ \cos \delta_z \end{bmatrix} = \text{Direction cosines}$$



1. what set your course
 2. what reference frames
 do you note it in?
 3. what is its notation
 4. what is its reference frame

Measurement of Position in Alternative Frames - 1

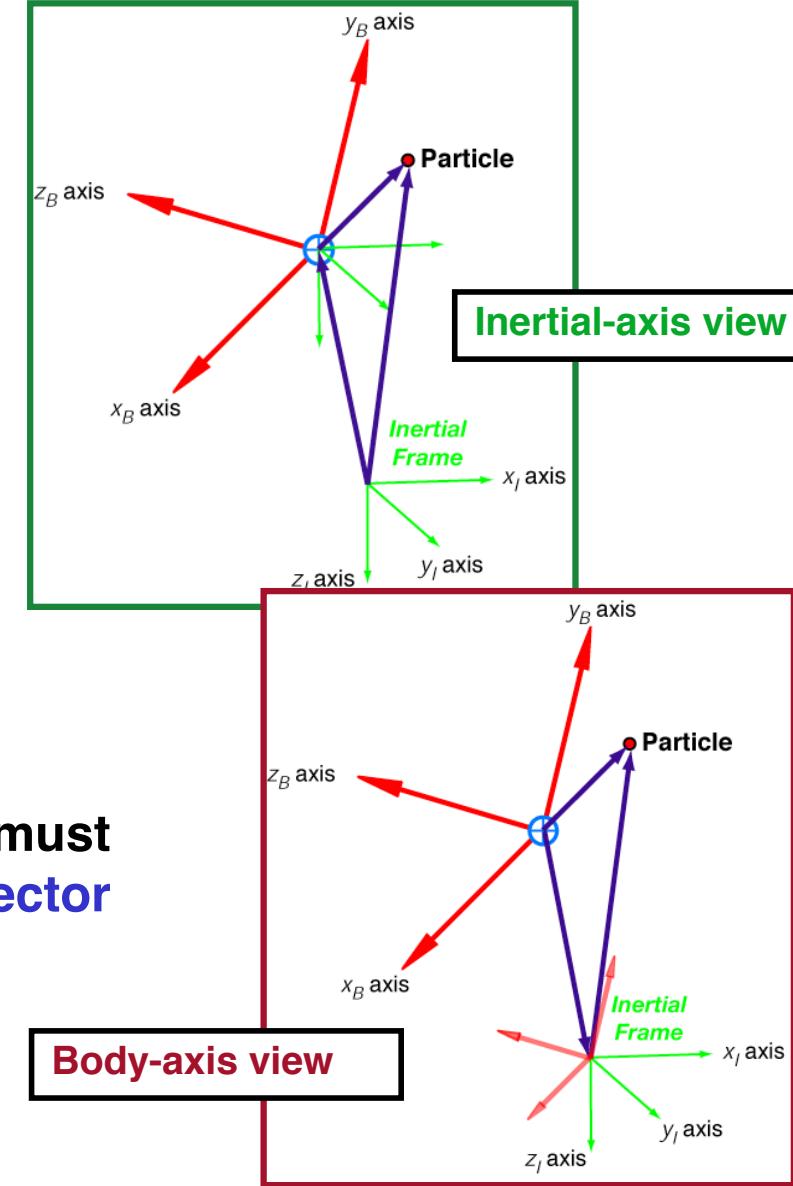
- Two reference frames of interest
 - **I:** Inertial frame (fixed to inertial space)
 - **B:** Body frame (fixed to body)

$$\mathbf{r} = \begin{bmatrix} x \\ y \\ z \end{bmatrix}$$

$$\mathbf{r}_{\text{particle}} = \mathbf{r}_{\text{origin}} + \Delta\mathbf{r}_{w.r.t.\text{ origin}}$$

Body

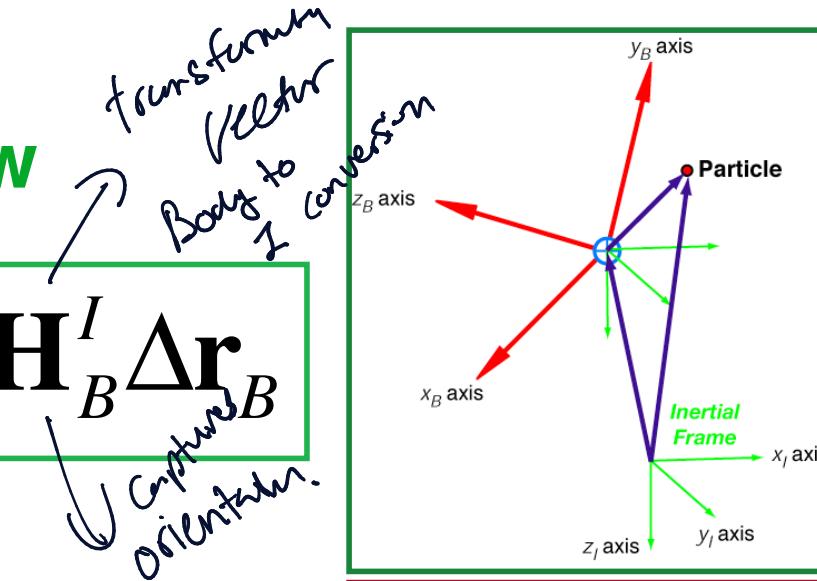
- Differences in frame orientations must be taken into account in adding vector components



Measurement of Position in Alternative Frames - 2

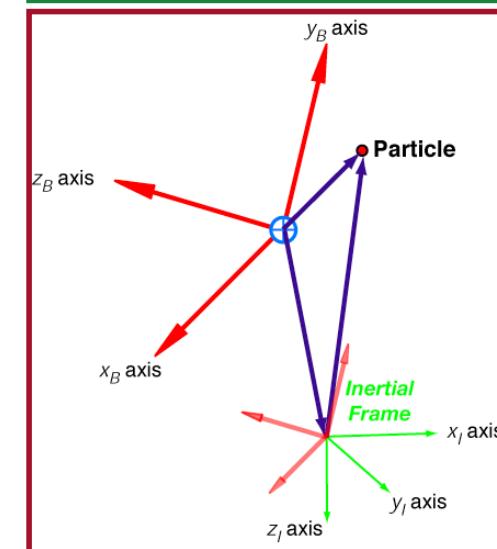
Inertial-axis view

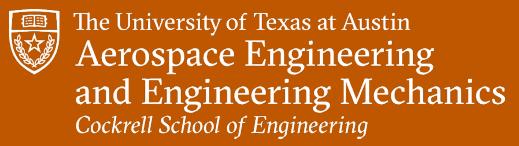
$$\mathbf{r}_{particle_I} = \mathbf{r}_{origin-B_I} + \mathbf{H}_B^I \Delta \mathbf{r}_B$$



Body-axis view

$$\mathbf{r}_{particle_B} = \mathbf{r}_{origin-I_B} + \mathbf{H}_I^B \Delta \mathbf{r}_I$$





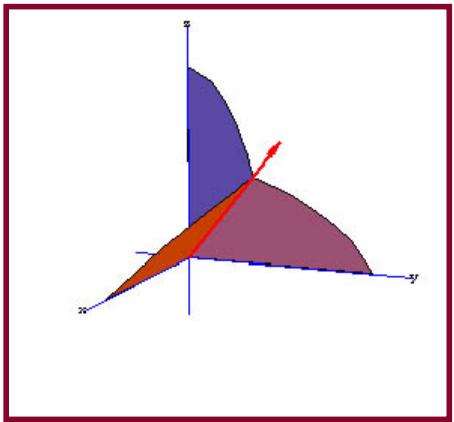
ROTATIONAL ORIENTATION

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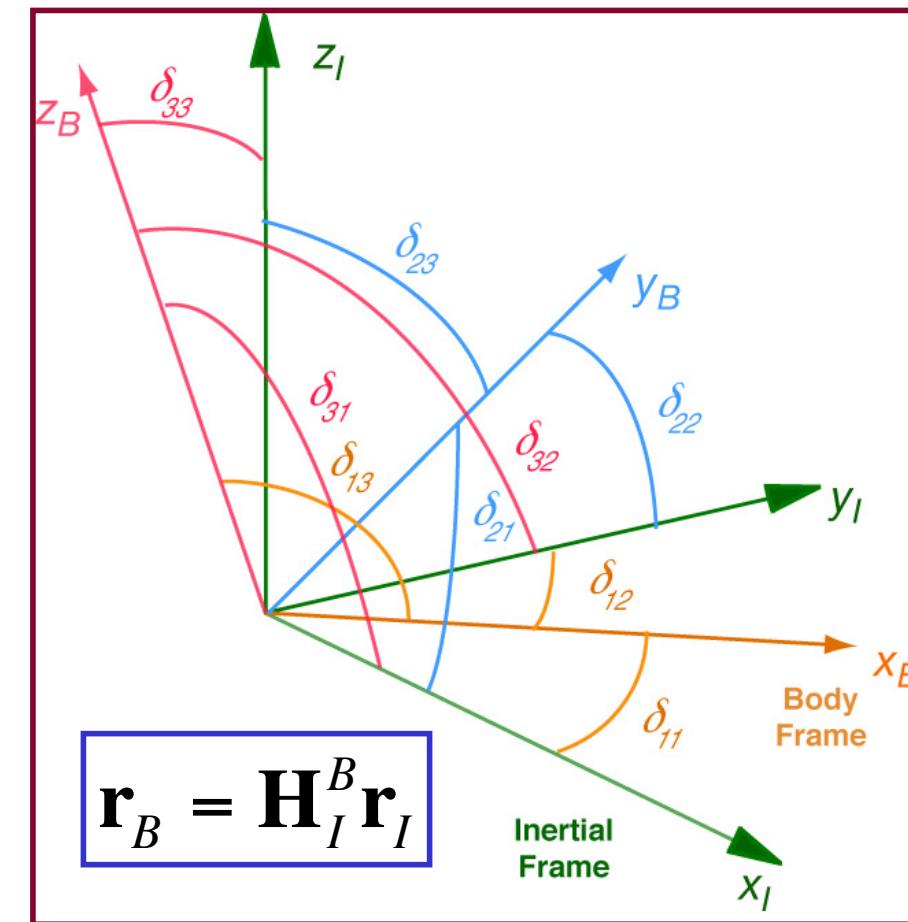
δ_{31}
 z-axis
 looks like on inertial
 Body

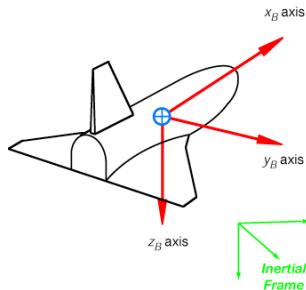


$$\mathbf{H}_I^B = \begin{bmatrix} \cos \delta_{11} & \cos \delta_{21} & \cos \delta_{31} \\ \cos \delta_{12} & \cos \delta_{22} & \cos \delta_{32} \\ \cos \delta_{13} & \cos \delta_{23} & \cos \delta_{33} \end{bmatrix}$$

- Projections of unit vector components of one reference frame on another
- Rotational orientation of one reference frame with respect to another
- Cosines of angles between each I axis and each B axis

Direction Cosine Matrix



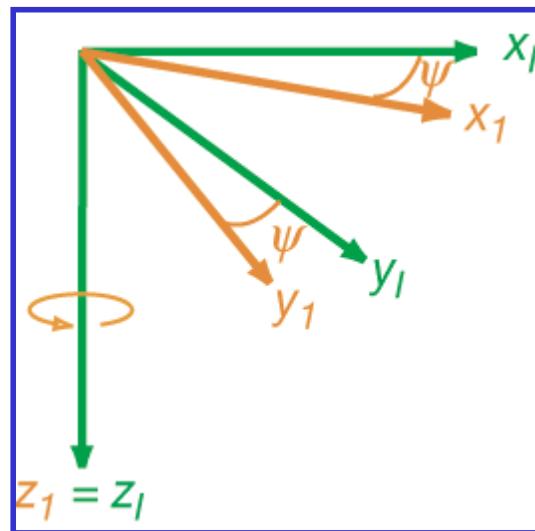


Euler Angles Measure the Orientation of One Frame with Respect to the Other

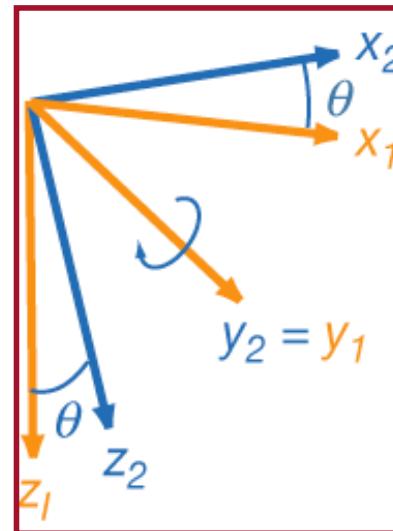
NED

- Conventional sequence of rotations from inertial to body frame
 - Each rotation is about a single axis
 - Right-hand rule
 - **Yaw, then pitch, then roll**
 - These are called **Euler Angles**

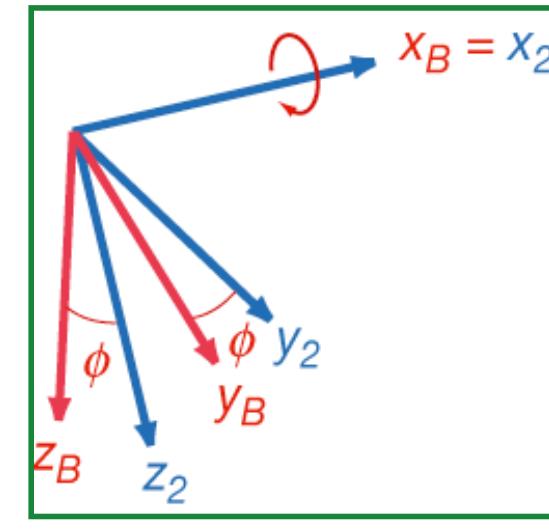
In principle
3 - 1 - 2
x - y - z



Yaw rotation (ψ) about z_I



Pitch rotation (θ) about y_1

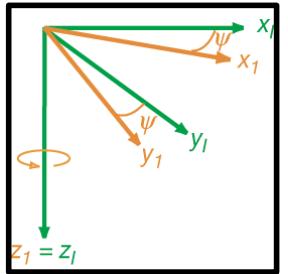


Roll rotation (ϕ) about x_2

Other sequences of 3 rotations can be chosen; however, once sequence is chosen, it must be retained

Reference Frame Rotation from Inertial to Body: Aircraft Convention (3-2-1)

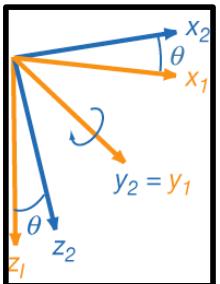
Yaw rotation (ψ) about z_I axis



$$\begin{bmatrix} x \\ y \\ z \end{bmatrix}_I = \begin{bmatrix} \cos\psi & \sin\psi & 0 \\ -\sin\psi & \cos\psi & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} x \\ y \\ z \end{bmatrix}_I = \begin{bmatrix} x_I \cos\psi + y_I \sin\psi \\ -x_I \sin\psi + y_I \cos\psi \\ z_I \end{bmatrix}$$

$$\mathbf{r}_I = \mathbf{H}_I^1 \mathbf{r}_I$$

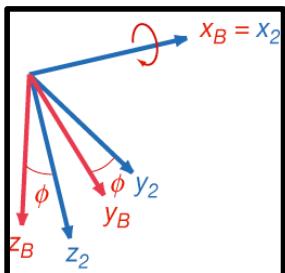
Pitch rotation (θ) about y_1 axis



$$\begin{bmatrix} x \\ y \\ z \end{bmatrix}_2 = \begin{bmatrix} \cos\theta & 0 & -\sin\theta \\ 0 & 1 & 0 \\ \sin\theta & 0 & \cos\theta \end{bmatrix} \begin{bmatrix} x \\ y \\ z \end{bmatrix}_I$$

$$\mathbf{r}_2 = \mathbf{H}_1^2 \mathbf{r}_I = [\mathbf{H}_1^2 \mathbf{H}_I^1] \mathbf{r}_I = \mathbf{H}_I^2 \mathbf{r}_I$$

Roll rotation (ϕ) about x_2 axis



$$\begin{bmatrix} x \\ y \\ z \end{bmatrix}_B = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos\phi & \sin\phi \\ 0 & -\sin\phi & \cos\phi \end{bmatrix} \begin{bmatrix} x \\ y \\ z \end{bmatrix}_I$$

$$\mathbf{r}_B = \mathbf{H}_2^B \mathbf{r}_I = [\mathbf{H}_2^B \mathbf{H}_1^2 \mathbf{H}_I^1] \mathbf{r}_I = \mathbf{H}_I^B \mathbf{r}_I$$

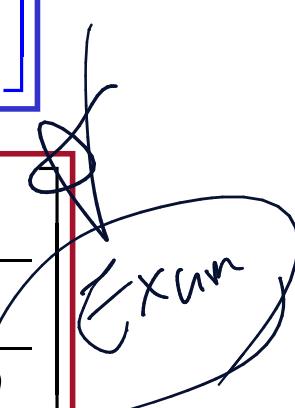
The Rotation Matrix

The three-angle rotation matrix is the product of 3 single-angle rotation matrices:

$$\mathbf{H}_I^B(\phi, \theta, \psi) = \mathbf{H}_2^B(\phi) \mathbf{H}_1^2(\theta) \mathbf{H}_I^1(\psi)$$

$$= \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos\phi & \sin\phi \\ 0 & -\sin\phi & \cos\phi \end{bmatrix} \begin{bmatrix} \cos\theta & 0 & -\sin\theta \\ 0 & 1 & 0 \\ \sin\theta & 0 & \cos\theta \end{bmatrix} \begin{bmatrix} \cos\psi & \sin\psi & 0 \\ -\sin\psi & \cos\psi & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$= \begin{bmatrix} \cos\theta\cos\psi & \cos\theta\sin\psi & -\sin\theta \\ -\cos\phi\sin\psi + \sin\phi\sin\theta\cos\psi & \cos\phi\cos\psi + \sin\phi\sin\theta\sin\psi & \sin\phi\cos\theta \\ \sin\phi\sin\psi + \cos\phi\sin\theta\cos\psi & -\sin\phi\cos\psi + \cos\phi\sin\theta\sin\psi & \cos\phi\cos\theta \end{bmatrix}$$



an expression of the **Direction Cosine Matrix**

Rotation Matrix Inverse

Inverse relationship: interchange sub- and superscripts

$$\mathbf{r}_B = \mathbf{H}_I^B \mathbf{r}_I$$

$$\mathbf{r}_I = (\mathbf{H}_I^B)^{-1} \mathbf{r}_B = \mathbf{H}_B^I \mathbf{r}_B$$

Because transformation is orthonormal

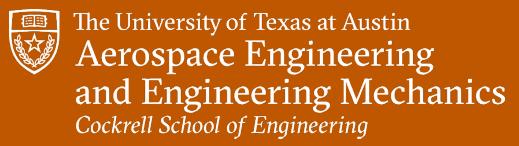
Inverse = transpose

Rotation matrix is always non-singular

$$[\mathbf{H}_I^B(\phi, \theta, \psi)]^{-1} = [\mathbf{H}_I^B(\phi, \theta, \psi)]^T = \mathbf{H}_B^I(\psi, \theta, \phi)$$

$$\mathbf{H}_B^I = (\mathbf{H}_I^B)^{-1} = (\mathbf{H}_I^B)^T = \mathbf{H}_1^I \mathbf{H}_2^1 \mathbf{H}_B^2$$

$$\mathbf{H}_B^I \mathbf{H}_I^B = \mathbf{H}_I^B \mathbf{H}_B^I = \mathbf{I}$$



ANGULAR MOMENTUM

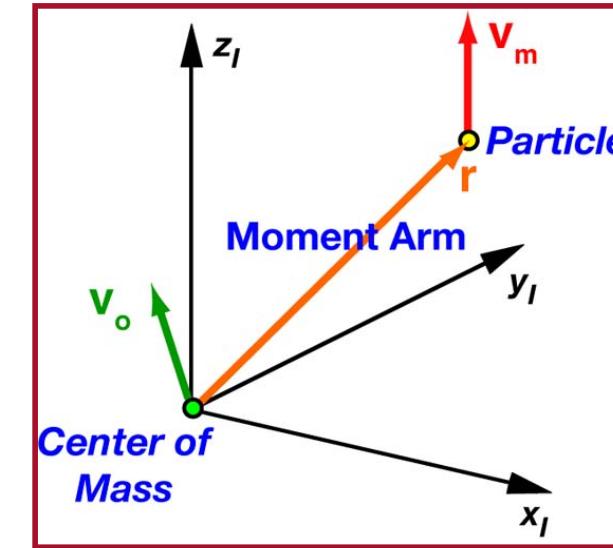
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Angular Momentum of a Particle

- Moment of linear momentum of differential particles that make up the body
 - (Differential masses) \times components of the velocity that are perpendicular to the moment arms



Angular momentum

$$d\mathbf{h} = (\mathbf{r} \times dm \mathbf{v}) = (\mathbf{r} \times \mathbf{v}_m) dm$$

$$= [\mathbf{r} \times (\mathbf{v}_o + \boldsymbol{\omega} \times \mathbf{r})] dm$$

Any object

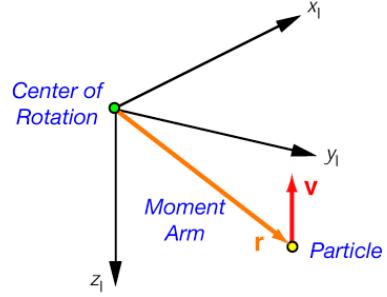
Angular momentum

$$\boldsymbol{\omega} = \begin{bmatrix} \omega_x \\ \omega_y \\ \omega_z \end{bmatrix}$$

can be decomposed into two motion

- Cross Product: Evaluation of a determinant with unit vectors (i, j, k) along axes, (x, y, z) and (v_x, v_y, v_z) projections on to axes

$$\mathbf{r} \times \mathbf{v} = \begin{vmatrix} i & j & k \\ x & y & z \\ v_x & v_y & v_z \end{vmatrix} = (yv_z - zv_y)\mathbf{i} + (zv_x - xv_z)\mathbf{j} + (xv_y - yv_x)\mathbf{k}$$



Cross-Product-Equivalent Matrix

$$\begin{aligned}
 \mathbf{r} \times \mathbf{v} &= \begin{vmatrix} i & j & k \\ x & y & z \\ v_x & v_y & v_z \end{vmatrix} = (yv_z - zv_y)\mathbf{i} + (zv_x - xv_z)\mathbf{j} + (xv_y - yv_x)\mathbf{k} \\
 &= \begin{bmatrix} (yv_z - zv_y) \\ (zv_x - xv_z) \\ (xv_y - yv_x) \end{bmatrix} = \tilde{\mathbf{r}}\mathbf{v} = \begin{bmatrix} 0 & -z & y \\ z & 0 & -x \\ -y & x & 0 \end{bmatrix} \begin{bmatrix} v_x \\ v_y \\ v_z \end{bmatrix}
 \end{aligned}$$

Cross-product-equivalent matrix

$$\tilde{\mathbf{r}} = \begin{bmatrix} 0 & -z & y \\ z & 0 & -x \\ -y & x & 0 \end{bmatrix}$$

Angular Momentum of the Aircraft

- Integrate moment of linear momentum of differential particles over the body

$$\mathbf{h} = \int_{Body} \left[\mathbf{r} \times (\mathbf{v}_o + \boldsymbol{\omega} \times \mathbf{r}) \right] dm = \int_{x_{min}}^{x_{max}} \int_{y_{min}}^{y_{max}} \int_{z_{min}}^{z_{max}} (\mathbf{r} \times \mathbf{v}) \rho(x, y, z) dx dy dz = \begin{bmatrix} h_x \\ h_y \\ h_z \end{bmatrix}$$

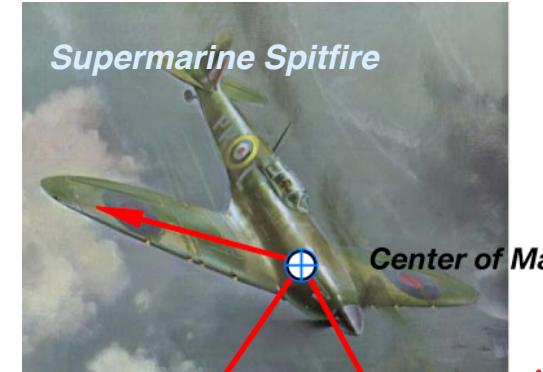
$\rho(x, y, z)$ = Density of the body

- Choose the center of mass as the rotational center

$\gamma = 0$, since
 v_o happens at
Body center

$$\begin{aligned} \mathbf{h} &= \int_{Body} (\mathbf{r} \times \mathbf{v}_o) dm + \int_{Body} [\mathbf{r} \times (\boldsymbol{\omega} \times \mathbf{r})] dm \\ &= 0 - \int_{Body} [\mathbf{r} \times (\mathbf{r} \times \boldsymbol{\omega})] dm \quad \text{It's b/c } (\mathbf{r} \times \mathbf{r}) \\ &= - \int_{Body} (\mathbf{r} \times \mathbf{r}) dm \times \boldsymbol{\omega} \equiv - \int_{Body} (\tilde{\mathbf{r}} \tilde{\mathbf{r}}) dm \boldsymbol{\omega} \end{aligned}$$

Not zero
because we're integrating



$$\boldsymbol{\omega} \times \boldsymbol{\gamma} = -(\mathbf{r} \times \boldsymbol{\omega})$$

$= \frac{r \times (\omega \times r)}{wcr \cdot r} - r(r \cdot \omega)$

$$h = r \times V_{\text{total}} \\ = r \times (\omega \times r)$$

The Inertia Matrix

$$\mathbf{h} = - \int_{\text{Body}} \tilde{\mathbf{r}} \tilde{\mathbf{r}} \boldsymbol{\omega} dm = - \int_{\text{Body}} \tilde{\mathbf{r}} \tilde{\mathbf{r}} dm \boldsymbol{\omega} = \mathbb{I}\boldsymbol{\omega}$$

$$\boldsymbol{\omega} = \begin{bmatrix} \omega_x \\ \omega_y \\ \omega_z \end{bmatrix}$$

where

$$\mathbb{I} = - \int_{\text{Body}} \tilde{\mathbf{r}} \tilde{\mathbf{r}} dm = - \int_{\text{Body}} \begin{bmatrix} 0 & -z & y \\ z & 0 & -x \\ -y & x & 0 \end{bmatrix} \begin{bmatrix} 0 & -z & y \\ z & 0 & -x \\ -y & x & 0 \end{bmatrix} dm$$

moment of inertia around X

$$= \int_{\text{Body}} \begin{bmatrix} (y^2 + z^2) & -xy & -xz \\ -xy & (x^2 + z^2) & -yz \\ -xz & -yz & (x^2 + y^2) \end{bmatrix} dm$$

around Y

around Z

Inertia matrix derives from equal effect of angular rate on all particles of the aircraft

Moments and Products of Inertia

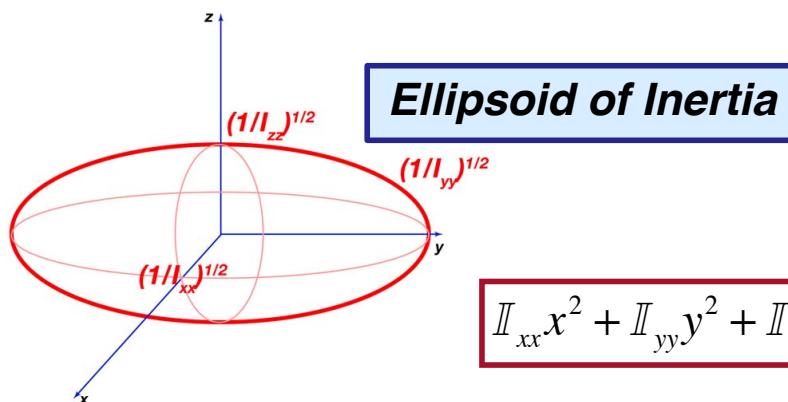
around one rotating axis affecting

$$\mathbb{I} = \int_{\text{Body}} \begin{bmatrix} (y^2 + z^2) & -xy & -xz \\ -xy & (x^2 + z^2) & -yz \\ -xz & -yz & (x^2 + y^2) \end{bmatrix} dm = \begin{bmatrix} \mathbb{I}_{xx} & -\mathbb{I}_{xy} & -\mathbb{I}_{xz} \\ -\mathbb{I}_{xy} & \mathbb{I}_{yy} & -\mathbb{I}_{yz} \\ -\mathbb{I}_{xz} & -\mathbb{I}_{yz} & \mathbb{I}_{zz} \end{bmatrix}$$

Inertia matrix

- Moments of inertia on the diagonal
- Products of inertia off the diagonal
- If products of inertia are zero, (x, y, z) are principal axes --->
- All rigid bodies have a set of principal axes

$$\begin{bmatrix} \mathbb{I}_{xx} & 0 & 0 \\ 0 & \mathbb{I}_{yy} & 0 \\ 0 & 0 & \mathbb{I}_{zz} \end{bmatrix}$$



Effect around + only affects around 2 axes

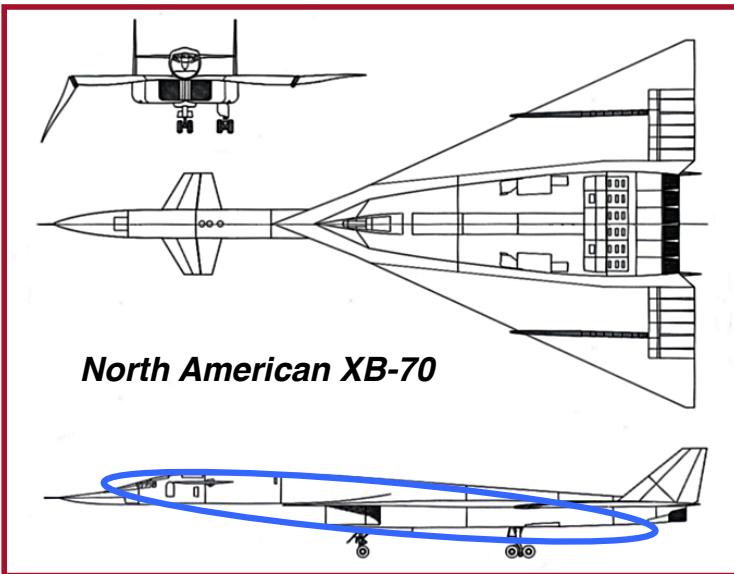
omega effect

Spanswise
Y axis for aircraft
z axis

Inertia Matrix of an Aircraft with Mirror Symmetry

$$\mathbb{I} = \int_{Body} \begin{bmatrix} (y^2 + z^2) & 0 & -xz \\ 0 & (x^2 + z^2) & 0 \\ -xz & 0 & (x^2 + y^2) \end{bmatrix} dm = \begin{bmatrix} \mathbb{I}_{xx} & 0 & -\mathbb{I}_{xz} \\ 0 & \mathbb{I}_{yy} & 0 \\ -\mathbb{I}_{xz} & 0 & \mathbb{I}_{zz} \end{bmatrix}$$

Nose high/low product
of inertia, I_{xz}



Nominal Configuration

Tips folded, 50% fuel, $W = 38,524$ lb

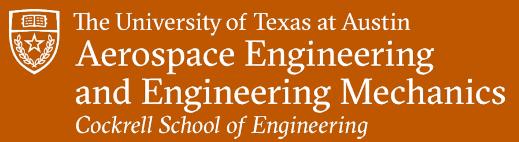
$$x_{cm} @ 0.218 \bar{c}$$

$$\mathbb{I}_{xx} = 1.8 \times 10^6 \text{ slug}\cdot\text{ft}^2$$

$$\mathbb{I}_{yy} = 19.9 \times 10^6 \text{ slug}\cdot\text{ft}^2$$

$$\mathbb{I}_{zz} = 22.1 \times 10^6 \text{ slug}\cdot\text{ft}^2$$

$$\mathbb{I}_{xz} = -0.88 \times 10^6 \text{ slug}\cdot\text{ft}^2$$



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RATE OF CHANGE OF ANGULAR MOMENTUM

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Newton's 2nd Law, Applied to Rotational Motion

In inertial frame, rate of change of angular momentum = applied moment (or torque), \mathbf{M}

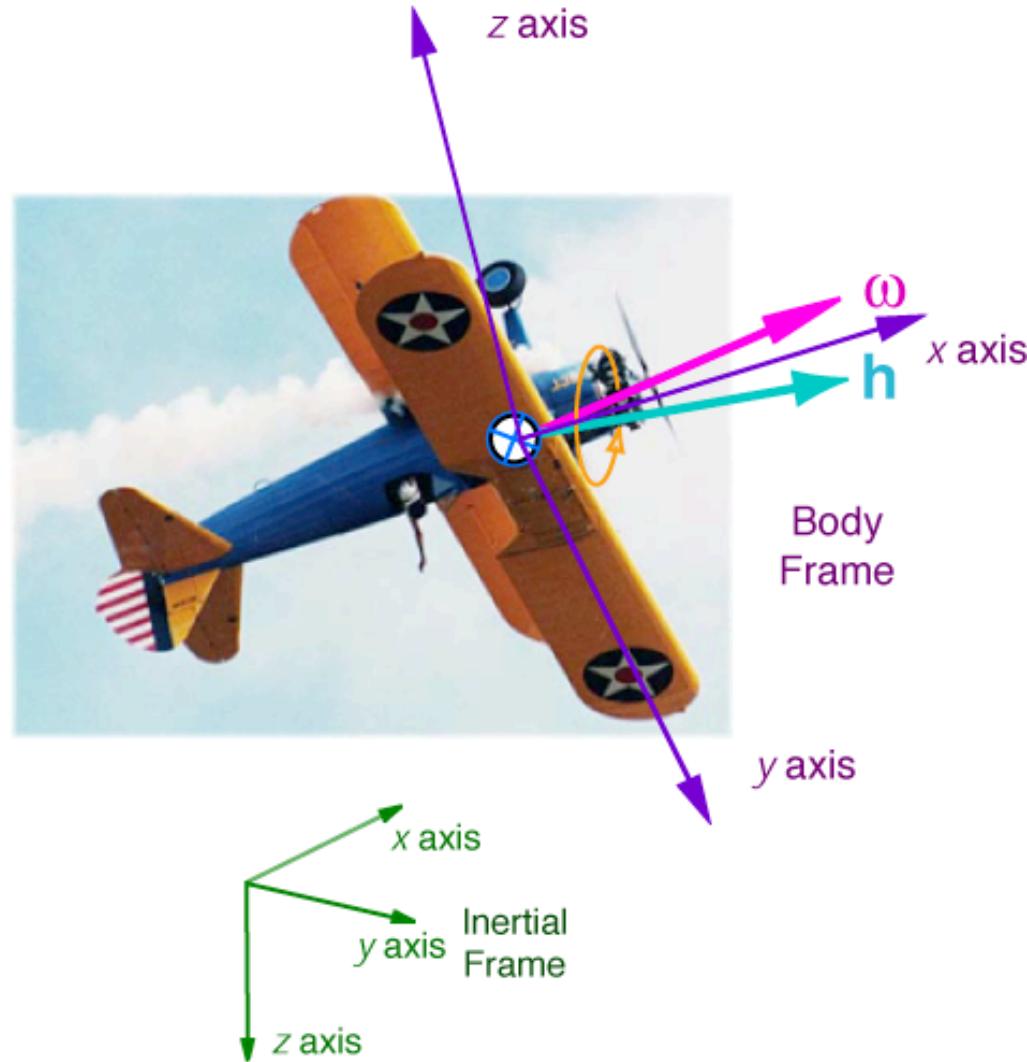
$$\frac{d\mathbf{h}}{dt} = \frac{d(\mathbb{I}\boldsymbol{\omega})}{dt} = \frac{d\mathbb{I}}{dt}\boldsymbol{\omega} + \mathbb{I}\frac{d\boldsymbol{\omega}}{dt}$$

$$= \dot{\mathbb{I}}\boldsymbol{\omega} + \mathbb{I}\dot{\boldsymbol{\omega}} = \mathbf{M} = \begin{bmatrix} m_x \\ m_y \\ m_z \end{bmatrix}$$

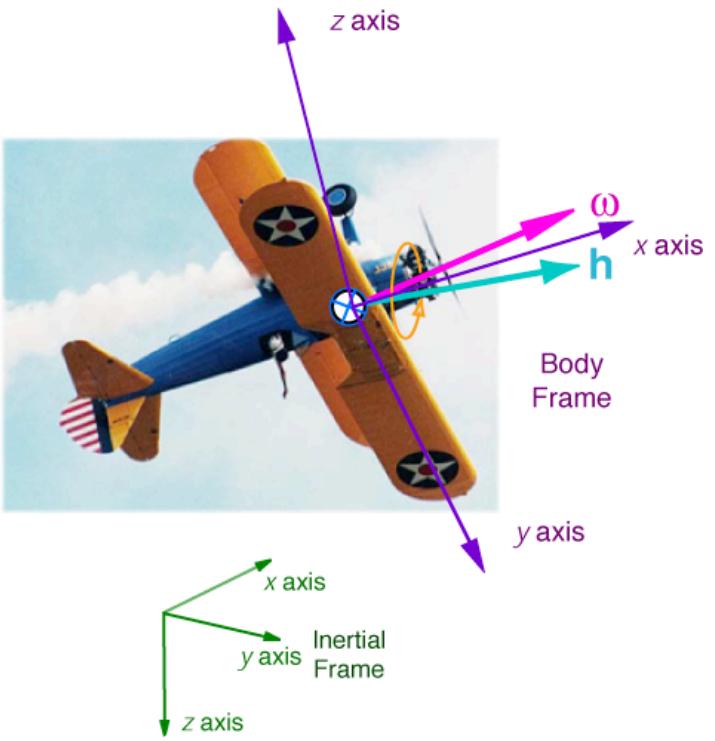
Angular Momentum and Rate

Angular momentum and rate vectors are not necessarily aligned

$$\mathbf{h} = I\boldsymbol{\omega}$$



Angular Momentum Expressed in Two Frames of Reference



- **Angular momentum and rate are vectors**
 - Expressed in either the inertial or body frame
 - Two frames related algebraically by the rotation matrix

$$\mathbf{h}_B(t) = \mathbf{H}_I^B(t)\mathbf{h}_I(t); \quad \mathbf{h}_I(t) = \mathbf{H}_B^I(t)\mathbf{h}_B(t)$$

$$\boldsymbol{\omega}_B(t) = \mathbf{H}_I^B(t)\boldsymbol{\omega}_I(t); \quad \boldsymbol{\omega}_I(t) = \mathbf{H}_B^I(t)\boldsymbol{\omega}_B(t)$$

Vector Derivative Expressed in a Rotating Frame

Chain Rule

$$\dot{\mathbf{h}}_I = \mathbf{H}_B^I \dot{\mathbf{h}}_B + \dot{\mathbf{H}}_B^I \mathbf{h}_B$$

*Effect of
body-frame rotation*

Alternatively

*Rate of change
expressed in body frame*

$$\dot{\mathbf{h}}_I = \mathbf{H}_B^I \dot{\mathbf{h}}_B + \boldsymbol{\omega}_I \times \mathbf{h}_I = \mathbf{H}_B^I \dot{\mathbf{h}}_B + \tilde{\boldsymbol{\omega}}_I \mathbf{h}_I$$

Consequently, the 2nd term is

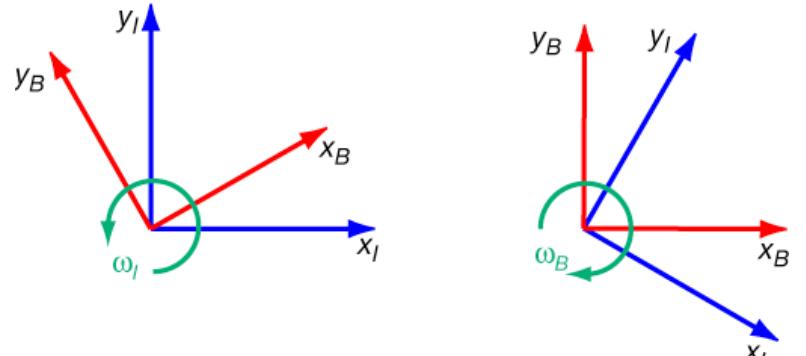
$$\dot{\mathbf{H}}_B^I \mathbf{h}_B = \tilde{\boldsymbol{\omega}}_I \mathbf{h}_I = \tilde{\boldsymbol{\omega}}_I \mathbf{H}_B^I \mathbf{h}_B$$

... where the cross-product
equivalent matrix of angular rate is

$$\tilde{\boldsymbol{\omega}} = \begin{bmatrix} 0 & -\omega_z & \omega_y \\ \omega_z & 0 & -\omega_x \\ -\omega_y & \omega_x & 0 \end{bmatrix}$$

External Moment Causes Change in Angular Rate

Positive rotation of Frame B w.r.t.
Frame A is a negative rotation of
Frame A w.r.t. Frame B

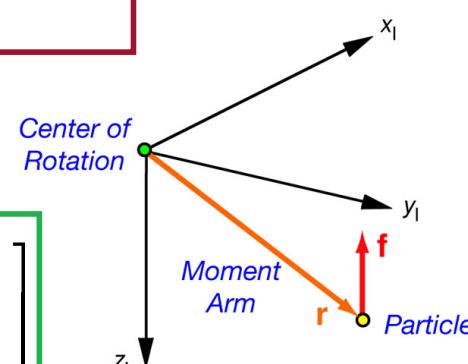


In the body frame of reference, the angular momentum change is

$$\begin{aligned}\dot{\mathbf{h}}_B &= \mathbf{H}_I^B \dot{\mathbf{h}}_I + \dot{\mathbf{H}}_I^B \mathbf{h}_I = \mathbf{H}_I^B \dot{\mathbf{h}}_I - \boldsymbol{\omega}_B \times \mathbf{h}_B = \mathbf{H}_I^B \dot{\mathbf{h}}_I - \tilde{\boldsymbol{\omega}}_B \mathbf{h}_B \\ &= \mathbf{H}_I^B \mathbf{M}_I - \tilde{\boldsymbol{\omega}}_B \mathbb{I}_B \boldsymbol{\omega}_B = \mathbf{M}_B - \tilde{\boldsymbol{\omega}}_B \mathbb{I}_B \boldsymbol{\omega}_B\end{aligned}$$

Moment = torque = force x moment arm

$$\mathbf{M}_I = \begin{bmatrix} m_x \\ m_y \\ m_z \end{bmatrix}_I ; \quad \mathbf{M}_B = \mathbf{H}_I^B \mathbf{M}_I = \begin{bmatrix} m_x \\ m_y \\ m_z \end{bmatrix}_B = \begin{bmatrix} L \\ M \\ N \end{bmatrix}$$



Rate of Change of Body-Referenced Angular Rate due to External Moment

In the body frame of reference, the angular momentum change is

$$\begin{aligned}\dot{\mathbf{h}}_B &= \mathbf{H}_I^B \dot{\mathbf{h}}_I + \dot{\mathbf{H}}_I^B \mathbf{h}_I = \mathbf{H}_I^B \dot{\mathbf{h}}_I - \boldsymbol{\omega}_B \times \mathbf{h}_B \\ &= \mathbf{H}_I^B \dot{\mathbf{h}}_I - \tilde{\boldsymbol{\omega}}_B \mathbf{h}_B = \mathbf{H}_I^B \mathbf{M}_I - \tilde{\boldsymbol{\omega}}_B \mathbb{I}_B \boldsymbol{\omega}_B \\ &= \mathbf{M}_B - \tilde{\boldsymbol{\omega}}_B \mathbb{I}_B \boldsymbol{\omega}_B\end{aligned}$$

For constant body-axis inertia matrix

$$\dot{\mathbf{h}}_B = \mathbb{I}_B \dot{\boldsymbol{\omega}}_B = \mathbf{M}_B - \tilde{\boldsymbol{\omega}}_B \mathbb{I}_B \boldsymbol{\omega}_B$$

Consequently, the differential equation for angular rate of change is

$$\dot{\boldsymbol{\omega}}_B = \mathbb{I}_B^{-1} (\mathbf{M}_B - \tilde{\boldsymbol{\omega}}_B \mathbb{I}_B \boldsymbol{\omega}_B)$$



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