

In this section, in anticipation of the thermal conductivity predictions in Section IV, we discuss two possible sources of error in the VC-predicted normal mode properties.

E. Discussion

For disordered systems, it is generally only possible to assign a ~~unique~~ lifetime and group velocity to vibrational modes in the low-frequency, propagating limit. ^{we believe that the} The VC-predicted group velocities, particularly for $v_g(\nu) \ll v_s$ or $v_g(\nu) \approx 0$, are an underprediction of the representative velocity scale ~~for~~ ^{used in Eq. (16) to calculate} the thermal diffusivity of high-frequency modes in the disordered lattice ~~as calculated from Eq. (16)~~ ^{statement supported}. This is demonstrated by considering the AF theory predicted mode thermal diffusivities, ^{in the} which are finite for the heavily disordered LJ alloy at a concentration of 0.5 (Fig. 7). ^{while} For phonons, the thermal diffusivity ^{from Eq. (16)} can be zero because of the VC predicted group velocities, ^{this result} which is not consistent with the AF theory predictions. ^(REF.)

Predictions from model disordered systems demonstrate the existence of a plateau of the thermal diffusivity at high frequency, which is consistent with the minimum phonon mean-free path hypothesis⁶³ used in most models of thermal transport in disordered materials.^{26,64,65} The concept of a vibrational mean free path is only valid, however, for low-frequency propagating modes in disordered systems.⁵⁰ The more fundamental property is the vibrational mode lifetime⁵⁶ or thermal diffusivity.^{7,8,25,50}

The VC-NMD and Gamma-NMD predict lifetimes ~~that~~ are generally larger than the IR limit

for LJ argon and its alloys (see Fig. 4). VC-ALD predicts essentially monotonically decreasing lifetimes with increasing frequency for the LJ argon alloys [Fig. 5(a)]. Because

VC-NMD and VC-ALD use the same values for $v_g(\nu)$, the phonon mode diffusivities will be underpredicted for VC-ALD compared to VC-NMD for the LJ argon alloys ~~with a significant number of lifetimes below the IR limit (see Fig. 5)~~.

There are thus two underpredictions to consider when interpreting the thermal conductivity predictions in Section : (i) underprediction of the thermal diffusivity that results from using the VC group velocities for VC-NMD and VC-ALD, and (ii) the underprediction of the mode lifetimes for LJ argon alloys by the VC-ALD perturbative models.

IV. THERMAL CONDUCTIVITY PREDICTIONS

The thermal conductivities of the LJ systems can now be predicted from Eq. (1) using the vibrational mode properties predicted by the VC-NMD and VC-ALD methods. Given the discussion regarding the VC-predicted mode properties in Section III E, we also predict ther-

The thermal diffusivity is the fundamental transport property

Why do you need to say both of these?

diverging, doesn't seem relevant, other than first part of first sentence, I suggest removing

Falling below the IR limit.

(needed? the section is already short!

The thermal conductivities predicted by HC-NMD and HC-ALD system size-dependent [i.e., $k = k(N_0)$]

(do you also use this procedure for GK?)

mal conductivity using the equilibrium MD-based GK method, which is a top-down ~~method~~ ^{approach} that does not make any approximations about the nature of the normal modes. Thermal conductivities predicted by the GK method naturally capture all scattering mechanisms.^{32,66,67} The main challenge in the GK method is how to specify the ~~converged value~~ of the thermal conductivity from the integral of the heat current autocorrelation function, ~~which is determined by the maximum of this integral.~~ The heat current was computed every 10 time steps from the same atomic trajectories (positions and velocities) used for the VC-NMD and Gamma-NMD methods.

so how did you do it?

→ To predict a bulk thermal conductivity, ~~finite-size extrapolation~~ ^{a linear procedure} is used, whereby

$$\frac{k(N_0)}{k_0} = 1 - \frac{c_0}{N_0}, \quad (19)$$

where c_0 is a constant, ~~k_0 is the extrapolated bulk thermal conductivity~~, and $k(N_0)$ is the size-dependent thermal conductivity.¹³ For VC-NMD and VC-ALD, the validity of the finite-size scaling requires that the low-frequency modes in the system be dominated by phonon-phonon scattering (i.e., $\tau \propto \omega^{-2}$) and follow the Debye approximation with respect to the group velocity DOS.^{13,14} For the LJ argon alloys, this requirement is satisfied for modest system sizes (for $N_0 = 6$ to 12) so that both VC-NMD and VC-ALD predictions can be extrapolated to a bulk value.

Eq. (19)

which one? both?

For SW silicon alloys, the thermal conductivity is dominated by low-frequency modes, so that large system sizes are needed to satisfy the extrapolation requirements and only GK and VC-ALD can be used ($N_0 \leq 42$ in the present study, similar to the converged system sizes in⁴¹). This highlights the efficiency of the VC-ALD method, which is necessary when computationally expensive DFT calculations are used (Section).^{4,10,11,14,68,69}

integrate w/ section V
p 27

Bulk thermal conductivity predictions are made for VC-NMD, VC-ALD, and GK are tabulated in Table I and plotted in Fig. 7. While agreement between the three methods is found for the perfect crystal, VC-NMD and VC-ALD underpredict the alloy thermal conductivities compared to GK. Also plotted ^{in Fig. 7} is the high-scatter thermal conductivity prediction,

explain how you specified k_{HS}

k_{HS} [Eq. (3)]. The underprediction is modest for VC-NMD, 20% or less for all concentrations. In Section IIID, we argued for the existence of a minimum mode thermal diffusivity, D_{HS} . As shown in Fig. 7, the ^{thermal} diffusivities of many high-frequency modes in the LJ alloys, predicted by both VC-NMD and VC-ALD, fall below this limit. Based on this observation, we propose that any thermal diffusivity below the limit be set to D_{HS} . The results of this adjustment, referred to as VC-NMD* and VC-ALD*, are plotted in Fig. 7. The adjusted

From 75
and included in Table I.

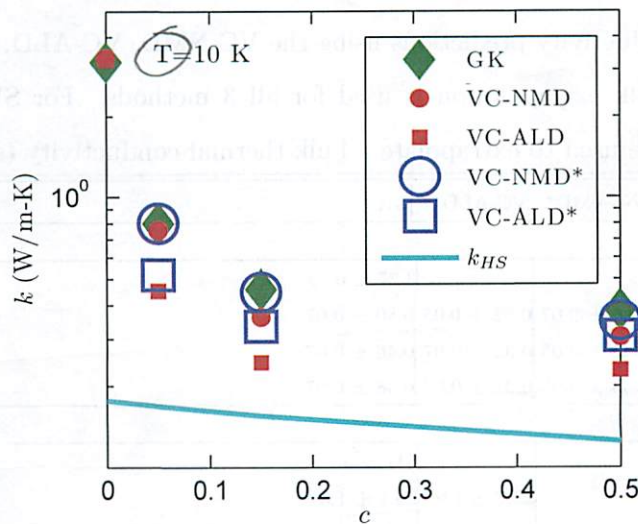


FIG. 7: Thermal conductivity predictions for LJ argon alloys at $T=10$ K using the VC-NMD, VC-ALD, and GK methods. The high-scatter thermal conductivity prediction k_{HS} (see Eq. (3)) and the high-scatter adjusted VC-NMD* and VC-ALD* are also plotted. The adjusted VC-NMD* is brought into agreement with GK method predictions, while the correction VC-ALD* is only marginally changed.

thermal conductivities predicted by VC-NMD* are within 10% of GK for all concentrations. Combined with D_{HS} , ^{we believe that} the VC-NMD predicted thermal diffusivities are good representations for the explicitly disordered modes present in the MD simulation.

The VC-ALD method underpredicts the thermal conductivity of the LJ argon alloys, where ^{the} underprediction is worst for a concentration of 0.05, ^{where} k_{VC-ALD}/k_{GK} is 0.56. By applying the ^{HS} high-scatter-limit adjustment VC-ALD*, the thermal conductivities are brought into marginally better agreement, worst for a concentration of 0.05, where k_{VC-ALD^*}/k_{GK} is 0.65. ^{As seen in Fig. 5(b),} The VC-ALD method fails to accurately predict the high-frequency mode thermal diffusivities for LJ argon alloys, which can be seen in Fig. 5(b). Since the group velocities are the same for VC-NMD and VC-ALD, the underprediction of the high-frequency thermal diffusivities, and hence thermal conductivity, is due to the underprediction of the high-frequency mode lifetimes, ^{from the ALD compared to VC-NMD, which we know are accurate values due to their} [Fig. 5(a)]. ^{agreement with Gamma-NMD.} Thus the adjustment VC-ALD* fails to bring VC-ALD into agreement with VC-NMD and GK.

↑ integrate discussion of Fig. 5(c) here (from p 19), I think it will help your argument and lead well into next section

TABLE I: Thermal conductivity predictions using the VC-NMD, VC-ALD, and GK method. For LJ argon alloys, the bulk extrapolation is used for all 3 methods. For SW silicon alloys, only VC-ALD and GK can be used to extrapolate a bulk thermal conductivity (see Section IV).

c	VC-NMD	VC-ALD	VC-NMD*	VC-ALD*	GK
LJ					
0.00	3.33 ± 0.12	3.45 ± 0.05			3.27 ± 0.09
0.05	0.76 ± 0.07	0.45 ± 0.05	0.80 ± 0.07	0.52 ± 0.05	0.80 ± 0.07
0.15	0.36 ± 0.04	0.24 ± 0.03	0.45 ± 0.05	0.33 ± 0.07	0.46 ± 0.07
0.50	0.31 ± 0.04	0.23 ± 0.03	0.35 ± 0.05	0.31 ± 0.07	0.38 ± 0.07
SW					
0.00		484 ± 21			517 ± 33
0.05		23.5 ± 1.9		23.7 ± 1.9	20.1 ± 1.5
0.15		11.9 ± 0.9		12.0 ± 0.9	9.90 ± 0.9
0.50		11.1 ± 0.9		11.2 ± 0.9	9.30 ± 0.9

don't shrink text for submission

The failure of VC-ALD to predict the thermal conductivities of the LJ alloys is due to an underprediction of the high-frequency mode lifetimes, which dominate the thermal conductivity [see Sections ---, Figs ---].

V. SW SILICON To provide a contrast,

Because of the discrepancies between the VC-NMD and VC-ALD predicted mode properties and thermal conductivities for the LJ argon alloys, we now predict the mode properties and thermal conductivity for ^{bulk and alloyed} SW silicon alloys which are "stiffer" (larger phonon group velocities and lifetimes). ^{the perfect crystal and an alloy} The lifetimes for a concentration of 0.5 predicted by VC-NMD and VC-ALD are plotted in Fig. 8(a). The VC-NMD predicted lifetimes are generally larger than the IR limit for SW silicon alloys, similar to VC-NMD predictions for the LJ argon alloys (Fig. 4). Unlike the LJ argon alloys, the VC-NMD and VC-ALD predicted lifetimes agree over most of the frequency spectrum, except at the highest frequencies, where VC-ALD underpredicts VC-NMD and falls below the IR limit. The high-frequency plateau of the VC-NMD predicted ~~mode~~ lifetimes for LJ argon (Fig. 4) is not seen for SW silicon.

As shown in Fig. 8(c), the thermal conductivity of SW silicon alloys is dominated by low-frequency modes. As seen in Fig. 8(b), VC-NMD and VC-ALD both predict a significant number of modes with $D_{ph}(\nu) < D_{HS}$ for both the LJ argon and SW silicon alloys, ^{the} similar to the findings for the LJ alloys [Fig. 5(c)].

From p24 (X) For SW silicon, we can only make bulk thermal conductivity predictions for VC-ALD and GK (see Table I and Fig. 9) because computational demands limit the system sizes accessible with VC-NMD. ~~From Table I,~~ The thermal conductivities predicted for SW silicon alloys by VC-ALD and GK are in agreement within ~~less than~~ 20%, ~~even without the adjustment~~ VC-ALD, where VC-ALD ~~actually~~ overpredicts compared to GK. The predicted thermal conductivities for SW silicon alloys at all concentrations ^{are} over an order of magnitude larger than the high-scatter prediction, k_{HS} [Eq. (3)]. ^{comment on Fig 8(c) here} Because thermal transport in SW silicon is ^{dominated by} low-frequency ^{HS} ~~dominant~~ ^{modes}, however, the high-scatter adjustments VC-NMD* and VC-ALD* vary by only one percent compared to the unadjusted VC-NMD and VC-ALD.

For the large concentrations and mass ratios considered in this work, the terms higher order terms in the Tamura perturbation theory are order 1 and larger at high frequencies ^{As discussed in Section 9.3} It is possible that higher-order interactions in the Tamura theory are ~~also~~ responsible for the discrepancy of the lifetimes predicted by VC-NMD and VC-ALD in SW silicon at the highest frequencies, but this discrepancy is unimportant to the overall thermal transport. (I would end this section here)

The LJ argon and SW silicon alloys studied in this work have different ranges of phonon

nothing new here- needed?

frequencies, lifetimes, group velocities, and total thermal conductivity. For bulk silicon, the thermal conductivity is dominated by low-frequency modes,^{5,11,30,31} which is also true for bulk and alloyed SW silicon (see Fig. 8). For SW silicon, VC-ALD predicts thermal conductivities in reasonable agreement with the explicitly disordered GK method (Fig. 9). For LJ argon, VC-ALD underpredicts the high-frequency phonon lifetimes and thermal diffusivities (see Fig. 5), leading to an underprediction of thermal conductivity when compared to VC-NMD and GK (see Section IV).

The results for the SW silicon and LJ argon alloys suggest that thermal modeling of ordered and disordered lattices can be separated into two broad groups: low-frequency dominated and full-spectrum materials. Materials dominated by low-frequency modes tend to have high thermal conductivities, which is due to their large group velocities and long lifetimes.^{30,31,70,71} These low-frequency modes follow closely the scalings predicted by the perturbative VC-ALD models, which are valid at low-frequencies.

LJ argon is a material whose thermal transport has significant contribution from high-frequency modes, even for the bulk [see Fig. 5 (c)]. This high-frequency range is where we predict that the perturbative Tamura theory will have non-negligible contributions from higher-order interactions (see Section IIIC2). While the higher-order interactions in the Tamura theory are also predicted to be non-negligible for SW silicon, this does not affect the thermal conductivity predictions significantly because high frequencies ^{modes} are unimportant to thermal transport. This is also true for the thermal conductivity spectrum of SiGe alloys from predictions^{9,11,70} Experimental measurements^{30,31,71} For example, the thermal conductivity of SiGe alloys exceeds the high-scatter limit by more than an order of magnitude at room temperature for all compositions.^{26,30,31}

Based on the thermal conductivity predictions for VC-NMD* and the well-defined peaks in the structure factors (Fig. 3), the reduction of group velocities in disordered lattices due to zone folding seems to be an underprediction of the group velocity of moderate to high frequency modes.⁴⁵

VI. SUMMARY

~~The concept of simple alloying remains at the forefront of efforts to control or minimize the thermal conductivity of semiconducting and thermoelectric materials (cite SiGe nanoporous,~~

comment on HS

all of this is better suited to the summary

Summary needs organizational work.

I suggest-

1. Quick review of everything you did

2. General discussion of low- ω vs. Full spectrum materials

3. Why VC-ALD is important

4. conditions under which VC-ALD (Pymu) will breakdown, effects on k prediction

5. Recommendations
* try not to repeat yourself- aim for one line in flow *

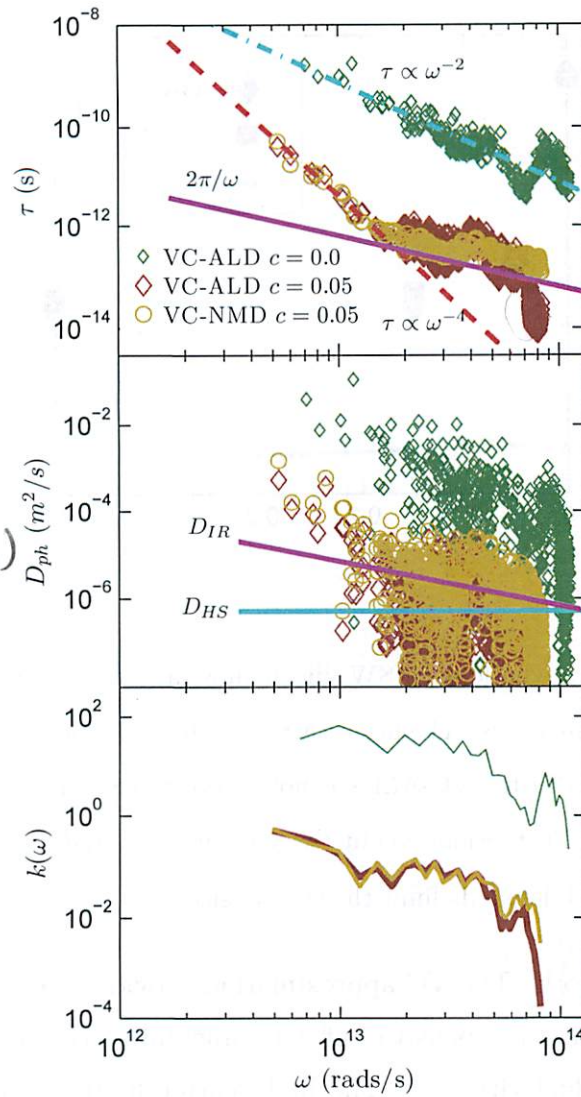


FIG. 8: (a) predicted lifetimes for VC modes using VC-NMD and VC-ALD for SW silicon. (b) predicted VC mode thermal diffusivities, compared to the AF,HS limit. (c) the thermal conductivity frequency spectrum, which is peaked at low frequency, in contrast to LJ argon (Fig. 5).

PbTe) Results in this work suggest that the lower limit for the vibrational mode thermal diffusivity in alloys is $(1/3)v_s a$.

The results in this work support the idea of a minimum thermal diffusivity for the vibrations in disordered lattices.(cite) Although this minimum thermal diffusivity is usually interpreted as a minimum mean free path, we find that concept is not necessary for inter-

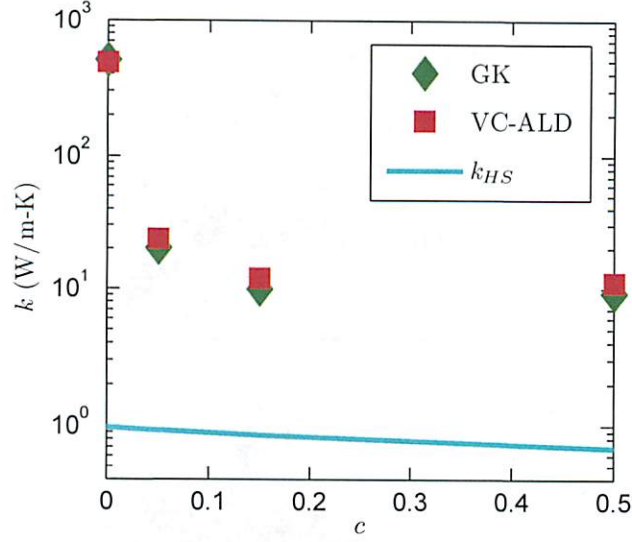


FIG. 9: Thermal conductivity predictions for SW silicon alloys at a temperature of 300 K using the VC-NMD, VC-ALD, and GK methods. The high-scatter thermal conductivity prediction k_{HS} (see Eq. (3)) is also plotted. The adjusted VC-ALD* is not shown since it differs by only one percent compared to VC-ALD. The bulk thermal conductivity cannot be predicted using the VC-NMD method because computational demands limit the system sizes.

preting the results of this work. The VC approximation provides a computationally cheap framework, which is essential for expensive but experimentally accurate *ab initio* methods for predicting thermal conductivity.(cite) The high-scatter limit of thermal diffusivity is more useful for examining the thermal transport in alloys under the framework of the VC approximation. The fundamental quantity is the mode lifetimes and the group velocity is an approximation, and expressed together as thermal diffusivity they can be interpreted in the presence of disorder.

The VC approximation results in two underpredictions for modes at high frequencies: (i) underprediction of the disordered vibrational thermal diffusivities because of the VC predicted group velocities, and (ii) underprediction of the disordered mode lifetimes by the perturbative VC-ALD method. The validity of the VC approximation has been verified for SW silicon alloys, which is a model system dominated by low-frequency modes. The underprediction of both the mode group velocities by the VC and the lifetimes by VC-ALD

occurs at the highest frequencies for SW silicon alloys, but these modes are unimportant to thermal transport. This is the plausible explanation for the success of the VC approximation phenomenologically and predictively for low-frequency dominated, high thermal conductivity materials.(cite)

For LJ argon alloys, which have significant thermal transport from high-frequency modes, VC-NMD and VC-ALD both underpredict the thermal conductivity compared to the system-level GK method. VC-NMD underpredicts only modestly, and can be brought into good agreement with GK by applying the high-scatter limit adjustment VD-NMD*. For VC-ALD, the high-scatter limit adjustment VC-ALD* still underpredicts compared to GK and VC-NMD*, suggesting that the perturbative Tamura theory is not appropriate for LJ argon alloys. This may be true for the high-frequency modes of any disordered lattice,(cite) and the high-scatter limit $D_{AF,HS}$ should be considered whenever the perturbative VC-ALD method is used.

Acknowledgments

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↑ put names in alphabetical order

Appendix A: NMD using Non-Exact Normal Modes

For a normal mode of the lattice supercell used for the MD simulations, the ~~autocorrelation of the total normal mode energy~~ ^{is an} autocorrelations ~~are damped~~ ^{function} exponentials with a decay time $\tau(\nu)$ and the kinetic energy autocorrelation is a ^{projecting onto} ~~cosinusoidal~~ exponentially-damped oscillation with frequency $2\omega(\nu)$.(cite AHRT) When using the VC normal modes to ~~map the~~ ^{from} MD simulation ~~trajectories from~~ ^{of pure} the explicitly disordered lattice supercells, the ~~mode to total and kinetic energy autocorrelation functions~~ do not always follow the simple functional forms, ^{as shown in} see Fig 10 (a) and (b). By calculating the ~~phonon mode spectral energy~~ ¹⁰ Φ , which is a ~~phonon mode~~ energy method in the frequency-domain,²⁴ artifacts such as multiple peaks

? (i.e, a gamma-mode)

be specific about 31

what is plotted, there is no Fig 10(b)!

are

~~in a given mode's energy spectrum Φ can be observed for high concentrations, $c = 0.5$, Fig. 10 (a).~~

These artifacts ~~in the mode energy spectrum and autocorrelations~~ are not surprising given two considerations: (i) the MD simulations contain explicit disorder which influences the atomic trajectories, (ii) the VC normal modes are not the exact normal modes ~~and~~ of the explicitly disordered lattice supercells. ~~It is important to remember that the VC normal modes are exact in the limit $c \rightarrow 0$.~~

~~In the case of multiple peaks in the VC phonon mode spectrum Φ , the choice of which peak to fit to predict the phonon properties can be ambiguous.~~ An effective lifetime can be predicted ~~unambiguously~~ using Eq. (11) because the VC ^{total} phonon mode ^{energy} autocorrelations still decay to zero in a finite time, ~~Fig. 10 (b).~~ This results is to be expected given that the atomic trajectories contain information about the lattice energy, which from general statistical physics principles will have exponential relaxation behavior in an equilibrium ensemble.⁷²⁻⁷⁴

↓ not totally satisfying

not relevant to your analysis