#### PROJECT

Development of low-order shock-capturing scheme for discontinuous Galerkin method

### Sana Amri Research Master Student

#### MIDTERM EVALUATION TEAM MEETING

18<sup>th</sup> March 2019



## von Karman Institute for Fluid Dynamics Aeronautics & Aerospace department

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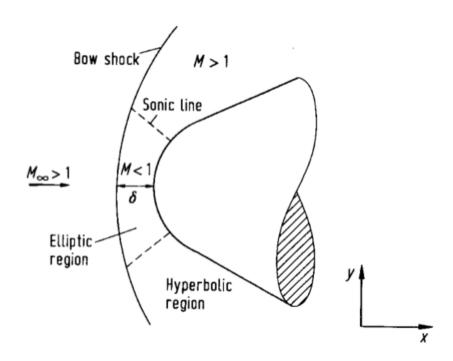
## OUTLINE

- 1 Motivation
- Bibliographic research approach
- Aims of the project
- Progress up to today
- 5 Conclusion & future work

#### **CHALLENGES**

- Measure the aerodynamic forces and heat flux on a body:
  - → Experiments in wind tunnels
  - → Simulation by means of CFD
- Hypersonic flow: → Multi-physics phenomena
  - → High temperature and compressible effects

"The role of CFD in engineering predictions has become so strong that today it can be viewed as a new 'third dimension' in fluid dynamics, the other two dimensions being the classical cases of pure experiment and pure theory."



## OUTLINE

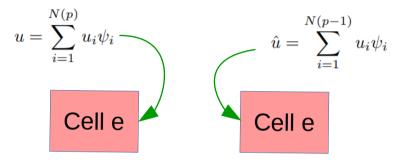
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#### Shock detector and artificial viscosity

**Shock detectors** are part of most shock capturing techniques and have for purpose to identify the location of discontinuities in the computational domain

Artificial viscosity makes shocks broader, so that they can be resolved over several grid cells.

#### Method:



**IF**(the polynomials shapes are similar) **Then** the solution on the cell e is smooth **Else** the solution contains discontinuities

Sensor of the discontinuity on the cell e:

$$S_e = \frac{(u - \hat{u}, u - \hat{u})_e}{(u, u)_e}$$

 $(\cdot,\cdot)_e$  inner product in  $L_2(\Omega_e)$ 

#### Inputs defined by the user:

- $\rightarrow$  The threshold  $s_0 \sim 1/p^4$
- $\rightarrow$  The interval  $\kappa$
- ightharpoonup The amount of artificial viscosity  $arepsilon_0 \sim h/p_0$

#### Artificial viscosity for each cell e:

$$\varepsilon_e = \begin{cases} 0 & \text{if } s_e < s_0 - \kappa \\ \frac{\varepsilon_0}{2} \left( 1 + \sin \frac{\pi(s_e - s_0)}{2\kappa} \right) & \text{if } s_0 - \kappa \le s_e \le s_0 + \kappa \\ \varepsilon_0 & \text{if } s_e > s_0 + \kappa \end{cases}$$

with  $s_e = \log_{10} S_{e}$ .

#### Main weakness of sensors

- The quantity needs to be compared to a threshold defined arbitrarily
- The method can not be used for p=0

#### The Navier-Stokes equations

$$oldsymbol{w} = egin{pmatrix} 
ho oldsymbol{u} \\ 
ho E \end{pmatrix} \in \Omega_{adm}$$

Let 
$$\Omega_{adm} = \left\{ oldsymbol{w} \in \mathbb{R}^4 : 
ho > 0, oldsymbol{u} \in \mathbb{R}^2, e > 0 
ight\}$$

The conservative form of the NS equations in 2D for a perfect calorically and thermally perfect gas are defined as

$$\partial_t \boldsymbol{w} + \nabla . \mathbb{F}_c(\boldsymbol{w}) - \nabla . \mathbb{F}_v(\boldsymbol{w}, \nabla \boldsymbol{w}) = 0, \ \forall (\boldsymbol{x}, t) \in \Omega \times ]0; +\infty[$$
 
$$\boldsymbol{w}(\boldsymbol{x}, 0) = \boldsymbol{w}_{t=0}(\boldsymbol{x}), \ \forall \boldsymbol{x} \in \Omega$$
 
$$p = \rho RT$$

Physical boundary conditions (defined on the following slides): 
$$\Gamma = \Gamma_{in} \cup \Gamma_{out} \cup \Gamma_{wall}$$

$$\mathbb{F}_c(\boldsymbol{w}) = \begin{pmatrix} \rho \boldsymbol{u}^T \\ \rho \boldsymbol{u} \otimes \boldsymbol{u} + p \mathbb{I} \\ (\rho E + p) \boldsymbol{u}^T \end{pmatrix}_{4 \times 2} \qquad \mathbb{F}_v(\boldsymbol{w}) = \begin{pmatrix} \boldsymbol{0}_{\mathbb{R}^2}^T \\ \boldsymbol{\tau} \\ \boldsymbol{u}^T . \boldsymbol{\tau} - \boldsymbol{q}^T \end{pmatrix}_{4 \times 2}$$

- Two important vector spaces:  $H^1(\Omega_h)=\left\{f\in L^2(\Omega_h):\phi_{|K}\in H^1(K),\; \forall K\in\Omega_h
  ight\}$   $v_h^p=\left\{f\in L^2(\Omega_h):\phi_{|K}\in\mathcal{P}^p(K),\; \forall K\in\Omega_h
  ight\}$
- $\mathcal{P}^p(K)$ : Polynomial of degree inferior or equal to p on K
- Each cell K is approximate by a polynomial that results from a linear combination of  $N_p$  basis functions  $\phi_i^K$  that are polynomial in the cell K and equal to zero in  $\Omega_h\backslash \mathrm{K}$

Variational Formulation: Find  $\boldsymbol{w}_h \in (v_h^p)^4$  such that  $\forall \phi_i^K, \ i \in \llbracket 1, N_p \rrbracket, \boldsymbol{w}_h$  is solution of

$$\sum_{K \in \Omega_h} \int_K \phi_i^K \partial_t \boldsymbol{w}_h dV - \sum_{K \in \Omega_h} \int_K \left( \mathbb{F}_c(\boldsymbol{w}_h) - \mathbb{F}_v(\boldsymbol{w}_h, \Upsilon_h) \right) \nabla \phi_i^K dv + \sum_{K \in \Omega_h} \oint_{\partial K} \phi_i^K \left( \hat{\mathbb{F}}_c - \hat{\mathbb{F}}_v \right) dS = 0$$

•  $\Upsilon_h$  is an auxiliary variable that represents the gradient  $abla w_h$  in the space of discretisation

$$dim(\mathcal{P}^p) := \prod_{i=1}^d \frac{(p+i)}{i} = N_p \quad dim(v_h^p) := N \times N_p$$

## Generic Form of the Variational Formulation

$$\begin{aligned} & \text{Generic Form of the Variational} \\ & \text{Formulation} \\ & \sum_{K \in \Omega_h} \int_K \phi_i^K \partial_t \boldsymbol{w}_h dV + \mathcal{L}_c(\boldsymbol{w}_h, \phi^K) + \mathcal{L}_v(\boldsymbol{w}_h, \nabla \boldsymbol{w}_h, \phi^K) = 0 \end{aligned} \qquad \begin{aligned} & z^{\pm}(x_e) = \lim_{y \to x_e, \ y \in K^{\pm}} z(y) \\ & \mathbb{I}[z] = \begin{cases} z^+ - z^- & \text{if } e \in \epsilon_i \\ z^+ - z_{wall} & \text{if } e \in \epsilon_{wall} \end{cases} \\ & \{z\} = \begin{cases} \frac{z_+ + z_-}{2} & \text{if } e \in \epsilon_i \\ z_{wall} & \text{if } e \in \epsilon_{wall} \end{cases} \end{aligned}$$

$$z^{\pm}(x_e) = \lim_{y \to x_e, \ y \in K^{\pm}} z(y)$$

$$[z] = \begin{cases} z^+ - z^- & \text{if } e \in \epsilon_i \\ z^+ - z_{wall} & \text{if } e \in \epsilon_{wall} \end{cases}$$

$$\{z\} = \begin{cases} \frac{z_{+} + z_{-}}{2} & if \ e \in \epsilon_{i} \\ z_{wall} & if \ e \in \epsilon_{wall} \end{cases}$$

$$\sum_{K \in \Omega_h} \int_K \phi_i^K \partial_t \boldsymbol{w}_h dV + \mathcal{L}_c(\boldsymbol{w}_h, \phi^K) + \mathcal{L}_v(\boldsymbol{w}_h, \nabla \boldsymbol{w}_h, \phi^K) = 0$$

## Lax Friedrich Riemann Solver

$$\hat{\mathbb{F}}_c(\boldsymbol{w}_h^+, \boldsymbol{w}_h^-, \boldsymbol{n}) = \{\mathbb{F}_c(\boldsymbol{w}_h)\} \, \boldsymbol{n} + \frac{a}{2} \, [\![\boldsymbol{w}_h]\!]$$

$$a = max \left\{ \left| \frac{\partial \mathbb{F}_c \boldsymbol{n}}{\partial \boldsymbol{w}} \right| : \boldsymbol{w} = \boldsymbol{w}_h^{\pm} \right\}$$

$$\mathcal{L}_{c}(\boldsymbol{w}_{h}, \phi^{K}) = -\sum_{K \in \Omega_{h}} \int_{K} \nabla(\phi^{K}) \mathbb{F}_{c}(\boldsymbol{w}_{h})$$

$$+ \sum_{e \in \epsilon_{i}} \int_{e} \llbracket \phi^{K} \rrbracket \, \hat{\mathbb{F}}_{c}(\boldsymbol{w}_{h}^{+}, \boldsymbol{w}_{h}^{-}, \boldsymbol{n})$$

$$+ \sum_{e \in \epsilon_{wall}} \int_{e} \phi^{+} \mathbb{F}_{c}(\boldsymbol{w}_{wall})$$

$$z^{\pm}(x_e) = \lim_{y \to x_e, \ y \in K^{\pm}} z(y)$$

$$[z] = \begin{cases} z^+ - z^- & \text{if } e \in \epsilon_i \\ z^+ - z_{wall} & \text{if } e \in \epsilon_{wall} \end{cases}$$

$$z^{\pm}(x_e) = \lim_{y \to x_e, y \in K^{\pm}} z(y)$$

$$\begin{bmatrix} z \end{bmatrix} = \begin{cases} z^+ - z^- & \text{if } e \in \epsilon_i \\ z^+ - z_{wall} & \text{if } e \in \epsilon_{wall} \end{cases}$$

$$\{z\} = \begin{cases} \frac{z_+ + z_-}{2} & \text{if } e \in \epsilon_i \\ z_{wall} & \text{if } e \in \epsilon_{wall} \end{cases}$$

$$\sum_{K \in \Omega_h} \int_K \phi_i^K \partial_t \boldsymbol{w}_h dV + \mathcal{L}_c(\boldsymbol{w}_h, \phi^K) + \mathcal{L}_v(\boldsymbol{w}_h, \nabla \boldsymbol{w}_h, \phi^K) = 0$$

Internal penalty inconsistent at p=0

$$z^{\pm}(x_e) = \lim_{y \to x_e, \ y \in K^{\pm}} z(y)$$

$$[z] = \begin{cases} z^+ - z^- & \text{if } e \in \epsilon_i \\ z^+ - z_{wall} & \text{if } e \in \epsilon_{wall} \end{cases}$$

$$z^{\pm}(x_e) = \lim_{y \to x_e, y \in K^{\pm}} z(y)$$

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$$\sum_{K \in \Omega_h} \int_K \phi_i^K \partial_t \boldsymbol{w}_h dV + \mathcal{L}_c(\boldsymbol{w}_h, \phi^K) + \mathcal{L}_v(\boldsymbol{w}_h, \nabla \boldsymbol{w}_h, \phi^K) = 0$$

The second scheme of Bassi Rebay

$$z^{\pm}(x_e) = \lim_{y \to x_e, \ y \in K^{\pm}} z(y)$$

$$[z] = \begin{cases} z^+ - z^- & \text{if } e \in \epsilon_i \\ z^+ - z_{wall} & \text{if } e \in \epsilon_{wall} \end{cases}$$

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$$\begin{cases} \Upsilon_h = \nabla \boldsymbol{w}_h \\ \sum_{K \in \Omega_h} \int_K \phi_i^K \partial_t \boldsymbol{w}_h dV + \mathcal{L}_c(\boldsymbol{w}_h, \phi^K) + \mathcal{L}_v(\boldsymbol{w}_h, \Upsilon_h, \phi^K) = 0 \end{cases}$$

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$$\mathbb{F}_v(\boldsymbol{w}_h, \nabla \boldsymbol{w}_h) := (G(\boldsymbol{w}_h) \Upsilon_h)$$

Find  $(\Upsilon_h, \boldsymbol{w}_h)$  in  $(\upsilon_h^p)^{4\times 2} \times (\upsilon_h^p)^4$  such that for all  $(\alpha^K, \phi^K)$  in  $(\upsilon_h^p)^2 \times \upsilon_h^p$ ,

$$\begin{cases} \sum_{K \in \Omega_{h}} \int_{K} \Upsilon_{h} \alpha^{K} dV + \sum_{K \in \Omega_{h}} \int_{K} (\nabla . \alpha^{K}) \boldsymbol{w}_{h} dV \\ - \sum_{e \in \epsilon_{i}} \int_{e} (\{\boldsymbol{w}_{h}\} \otimes \boldsymbol{n}) \left[\!\left[\alpha^{K}\right]\!\right] dS - \sum_{e \in \epsilon_{wall}} \int_{e} (\boldsymbol{w}_{wall} \otimes \boldsymbol{n}) \alpha^{K,+} dS = 0 \\ \sum_{K \in \Omega_{h}} \int_{K} \phi_{i}^{K} \partial_{t} \boldsymbol{w}_{h} dV + \mathcal{L}_{c}(\boldsymbol{w}_{h}, \phi^{K}) + \sum_{K \in \Omega_{h}} \int_{K} \nabla \phi^{K} (G\Upsilon_{h}) dV \\ - \sum_{e \in \epsilon_{i}} \int_{e} \left[\!\left[\phi^{K}\right]\!\right] \{G\Upsilon_{h}\} \boldsymbol{n} dS - \sum_{e \in \epsilon_{wall}} \int_{e} \phi^{K,+} \mathbb{F}_{v}(\boldsymbol{w}_{wall}, \nabla \boldsymbol{w}_{wall}) \boldsymbol{n} dS = 0 \end{cases}$$

$$\begin{cases} \Upsilon_h = \nabla \boldsymbol{w}_h \\ \sum_{K \in \Omega_h} \int_K \phi_i^K \partial_t \boldsymbol{w}_h dV + \mathcal{L}_c(\boldsymbol{w}_h, \phi^K) + \mathcal{L}_v(\boldsymbol{w}_h, \Upsilon_h, \phi^K) = 0 \end{cases}$$

Find 
$$(\Upsilon_h, \boldsymbol{w}_h)$$
 in  $(v_h^p)^{4\times 2} \times (v_h^p)^4$  such that for all  $(\alpha^K, \phi^K)$  in  $(v_h^p)^2 \times v_h^p$ ,

$$z^{\pm}(x_e) = \lim_{y o x_e, \ y \in K^{\pm}} z(y)$$
 $\llbracket z 
Vert = \begin{cases} z^+ - z^- & if \ e \in \epsilon_i \ z^+ - z_{wall} & if \ e \in \epsilon_{wall} \end{cases}$ 
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abla oldsymbol{1}_h \ Global \ lifting \ operator: \ \Upsilon_h \coloneqq 
abla oldsymbol{w}_h + R_h (\llbracket oldsymbol{w}_h 
Vert)$ 

Local lifting operator:

 $\{G\Upsilon_h\} = \{G\nabla \boldsymbol{w}_h\} + \{Gr_h^e\}$ 

$$\begin{cases} \sum_{K \in \Omega_{h}} \int_{K} \Upsilon_{h} \alpha^{K} dV + \sum_{K \in \Omega_{h}} \int_{K} (\nabla . \alpha^{K}) \boldsymbol{w}_{h} dV \\ - \sum_{e \in \epsilon_{i}} \int_{e} (\{\boldsymbol{w}_{h}\} \otimes \boldsymbol{n}) \left[\!\left[\alpha^{K}\right]\!\right] dS - \sum_{e \in \epsilon_{wall}} \int_{e} (\boldsymbol{w}_{wall} \otimes \boldsymbol{n}) \alpha^{K,+} dS = 0 \\ \sum_{K \in \Omega_{h}} \int_{K} \phi_{i}^{K} \partial_{t} \boldsymbol{w}_{h} dV + \mathcal{L}_{c}(\boldsymbol{w}_{h}, \phi^{K}) + \sum_{K \in \Omega_{h}} \int_{K} \nabla \phi^{K} (\boldsymbol{G} \Upsilon_{h}) dV \\ - \sum_{e \in \epsilon_{i}} \int_{e} \left[\!\left[\phi^{K}\right]\!\right] \left(\!\left[G \Upsilon_{h}\right]\!\right) \boldsymbol{n} dS - \sum_{e \in \epsilon_{wall}} \int_{e} \phi^{K,+} \mathbb{F}_{v}(\boldsymbol{w}_{wall}, \nabla \boldsymbol{w}_{wall}) \boldsymbol{n} dS = 0 \end{cases}$$

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scheme of

Bassi Rebay

$$\mathcal{L}_{v}(\boldsymbol{w}_{h}, \Upsilon_{h}, \phi^{K}) = -\sum_{K \in \Omega_{h}} \int_{K} \nabla \phi^{K} G(\boldsymbol{w}_{h}) \left( \nabla \boldsymbol{w}_{h} + R_{h} \left( \llbracket \boldsymbol{w}_{h} \rrbracket \right) \right) dV$$

$$+ \sum_{e \in \epsilon_{i}} \int_{r} \int_{r} \mathcal{L}^{K} \left[ \left\{ G(\boldsymbol{w}_{h}) \left( \nabla \boldsymbol{w}_{h} + \eta_{r} r_{h}^{e} \left( \llbracket \boldsymbol{w}_{h} \rrbracket \right) \right) \right\} \boldsymbol{n} dS$$

$$+ \sum_{e \in \epsilon_{wall}} \int_{e} \phi^{K,+} \mathbb{F}_{v}(\boldsymbol{w}_{wall}, \nabla \boldsymbol{w}_{wall} + \eta_{r} r_{h}^{e} \left( \llbracket \boldsymbol{w}_{h} \rrbracket \right) \right) \boldsymbol{n} dS$$

$$z^{\pm}(x_e) = \lim_{y o x_e, \ y \in K^{\pm}} z(y)$$
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Global lifting operator:
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Local lifting operator:

 $\{G\Upsilon_h\} = \{G\nabla \boldsymbol{w}_h\} + \{Gr_h^e\}$ 

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#### Aim of the project

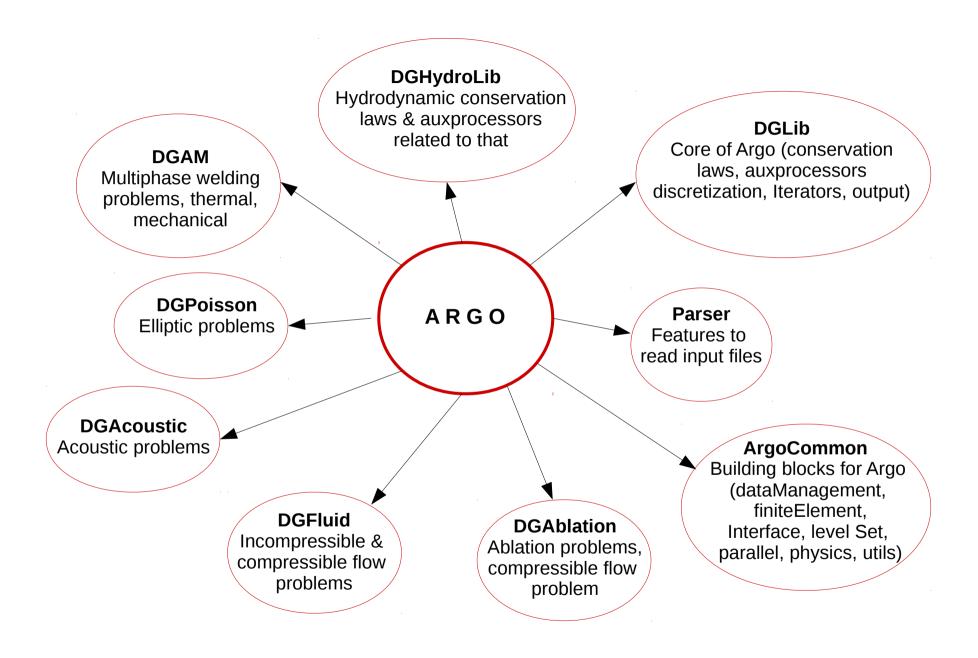
Develop a robust shock capturing scheme for the DG method

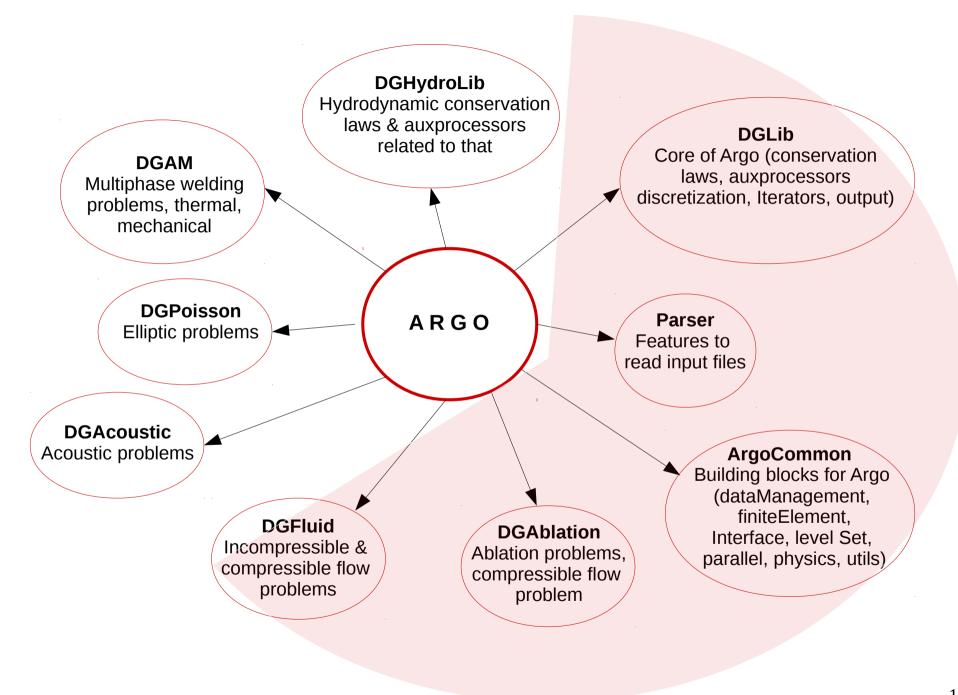
#### Main tasks

- 1. Review the hypersonic flow challenges for high order numerical simulation
- 2. Identify the limit of the artificial viscosity into the Argo platform
- **3.** Optimize and adapt a hybrid solver for hypersonic applications which uses a **degraded order scheme** (BR2) in the shock region and a high order method (DG) everywhere else
- 4. Validation on a specific test case

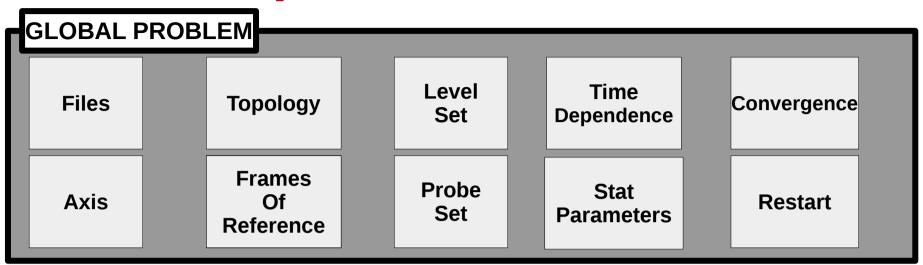
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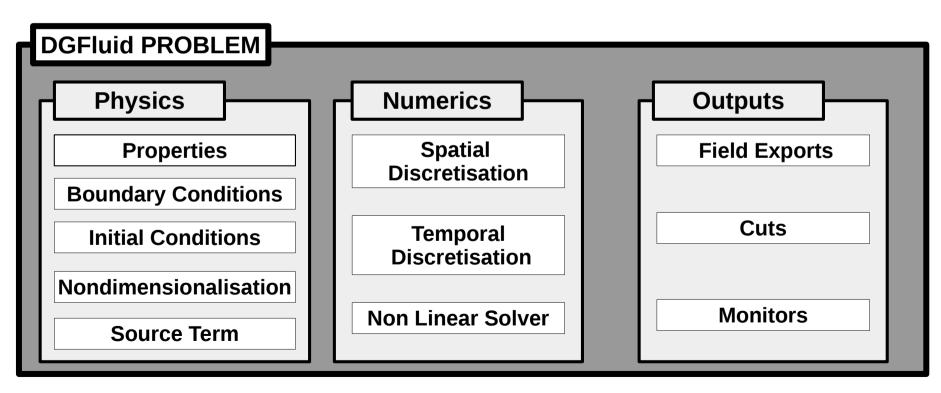
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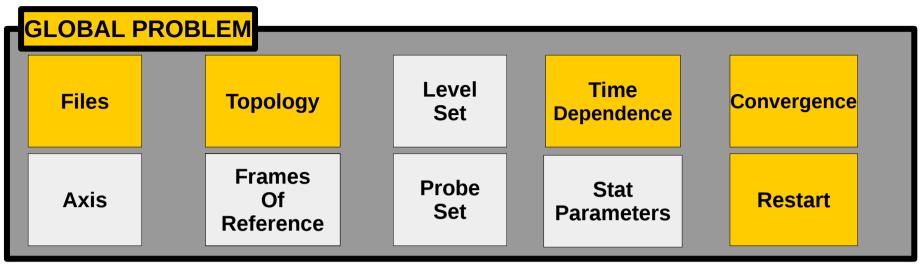


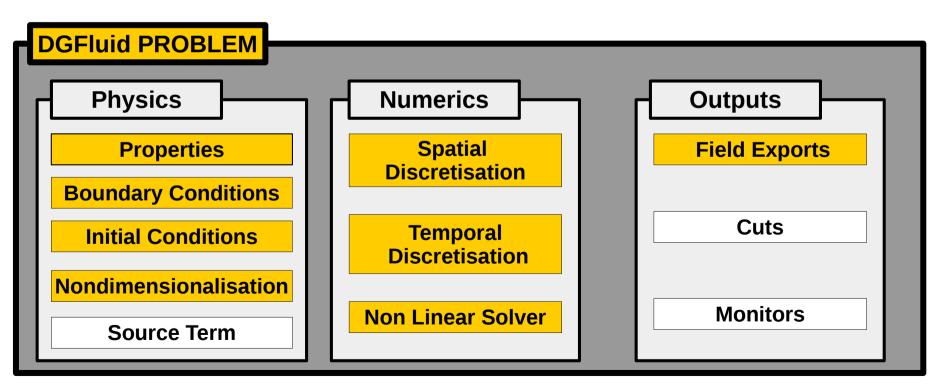
#### Structure of the input file



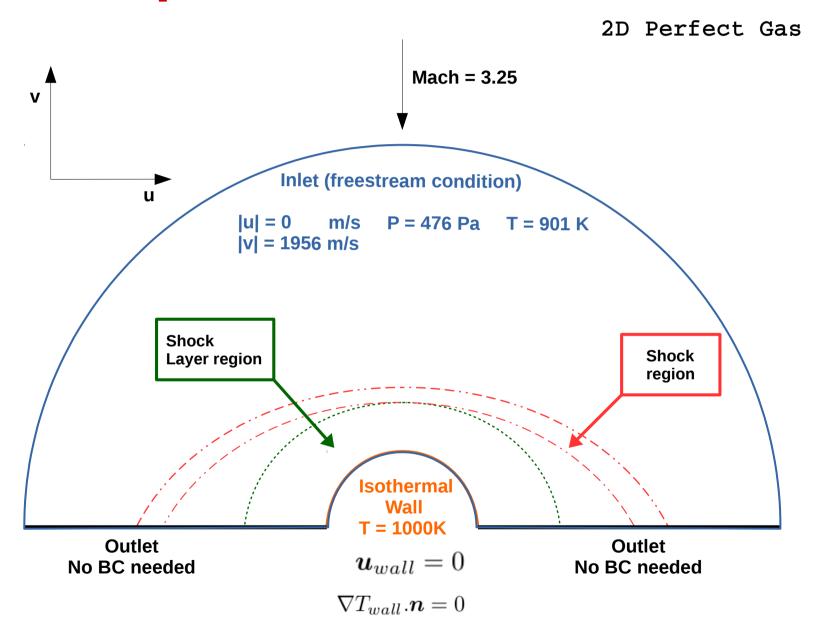


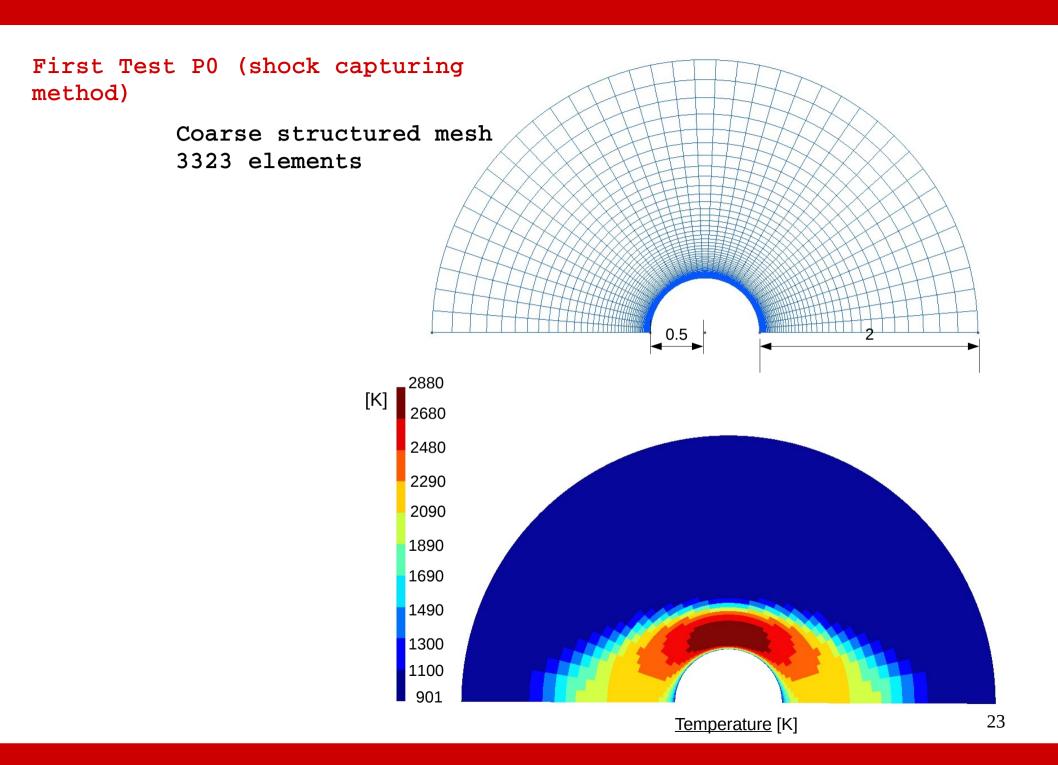
### Structure of the input file

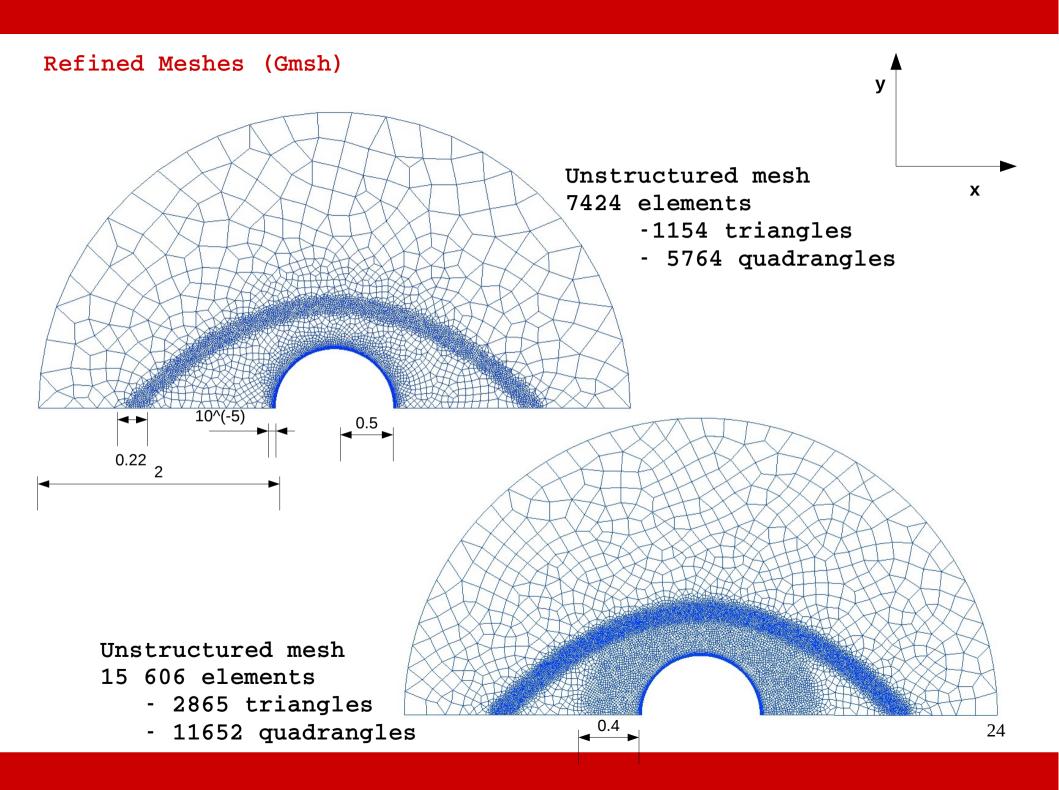


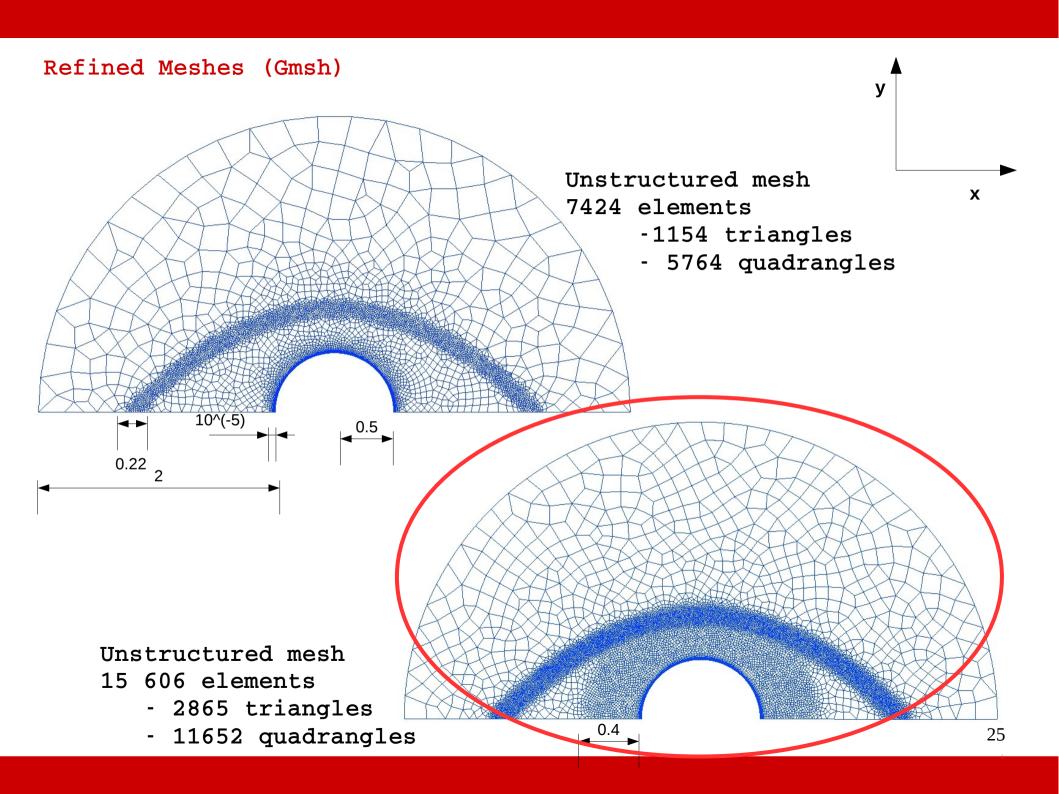


## Scheme of the Physical model

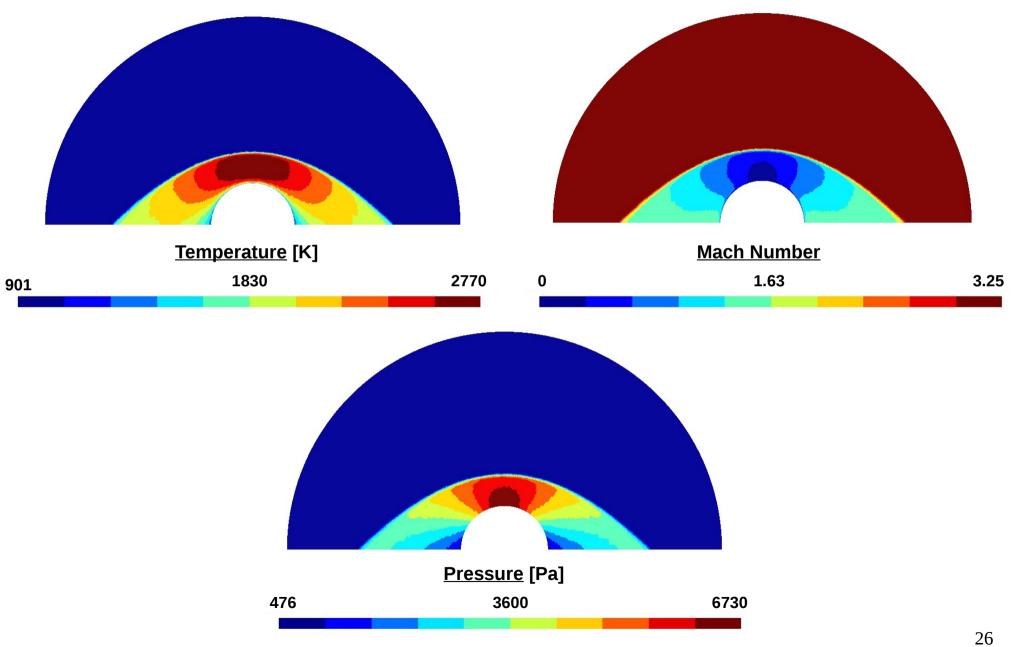




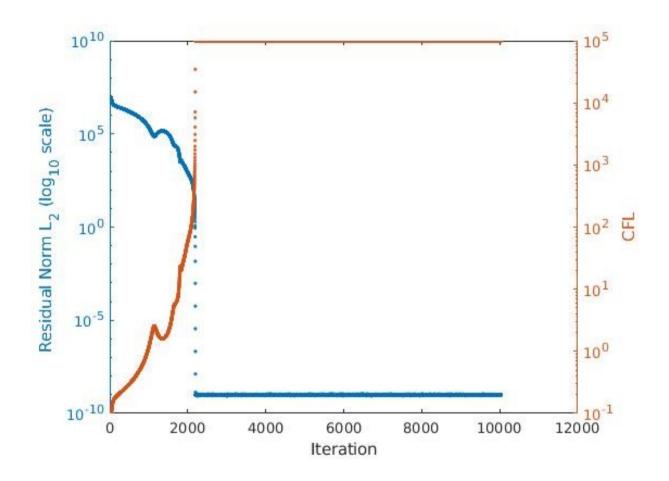


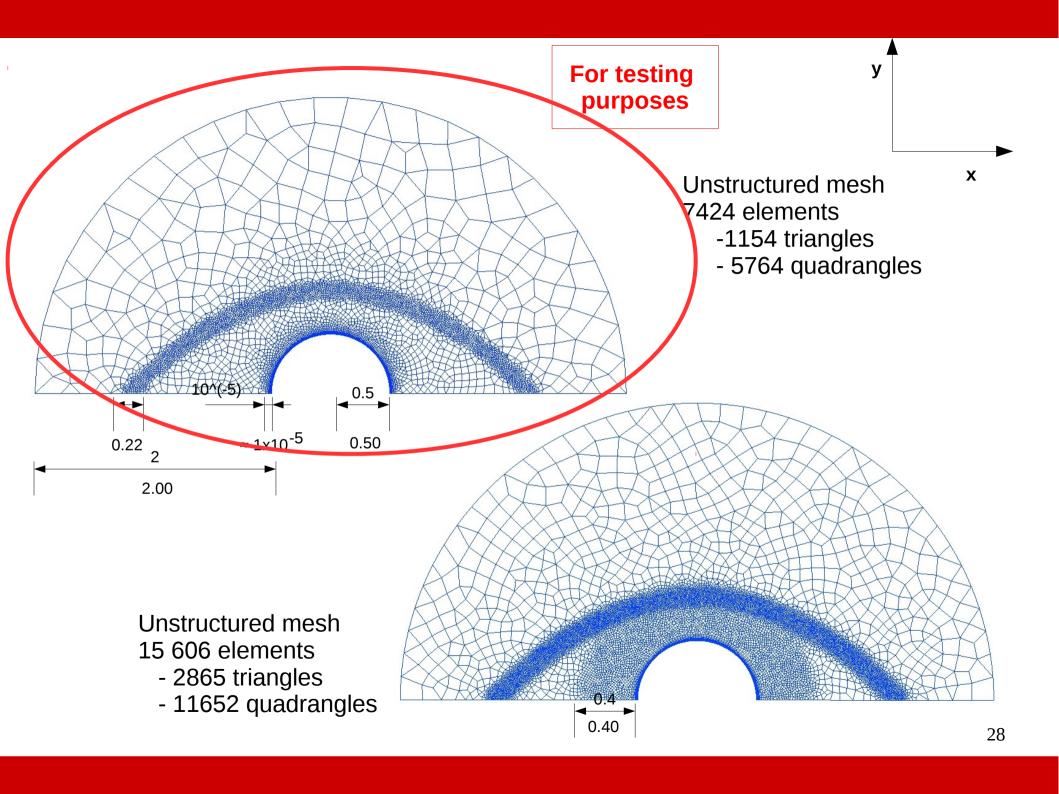


## Solution on Refined 15606 elements mesh PO

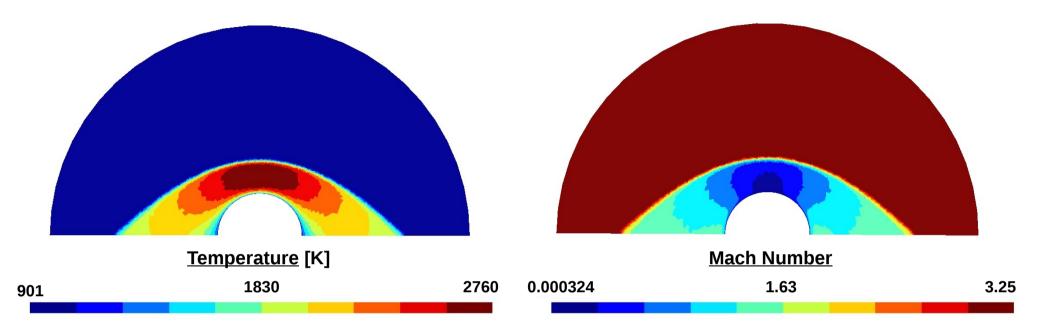


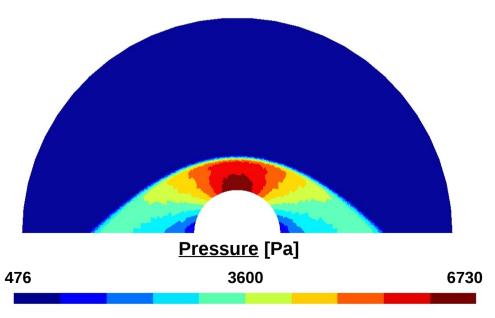
## Solution on Refined 15606 elements mesh PO





## Solution on Refined 7424 element mesh PO





Solution along the wall on Refined 7424 element mesh PO

0.5

1.5

1

7000

6000

5000

4000

3000

2000

1000

0.025

0.02

0.015

0.01

0.005

-1.5

Density [kg·m<sup>-2</sup>]

-1.5

-1

-0.5

-0.5

-1

0

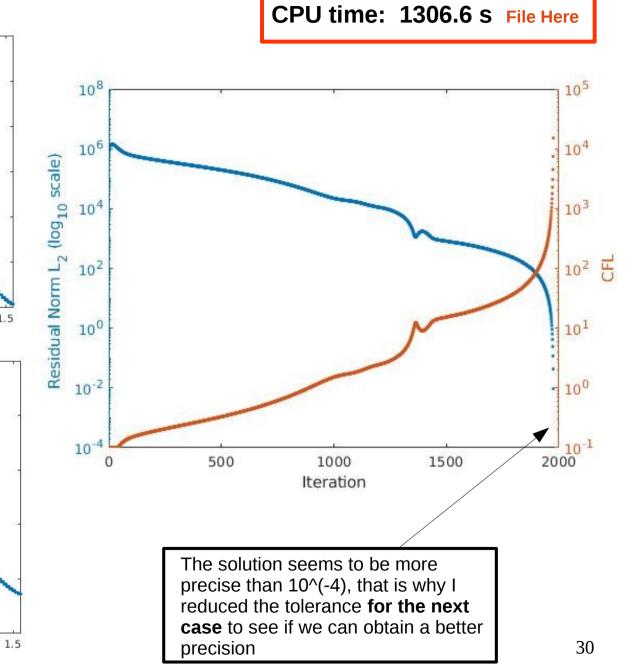
Angle [radian]

0.5

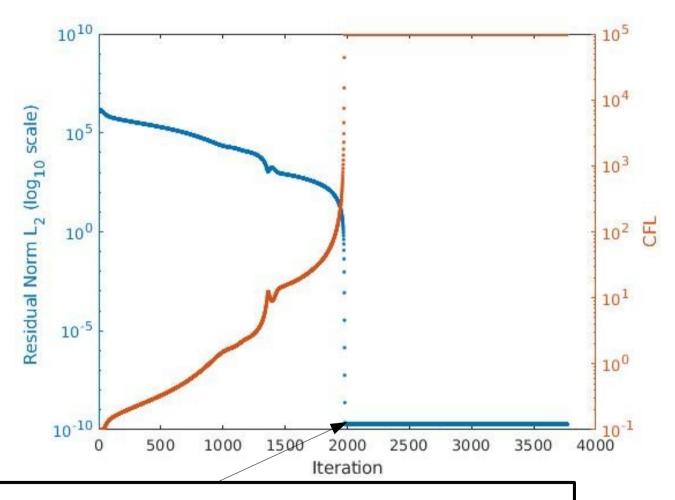
1

Angle [radian]

Total Pressure [Pa]

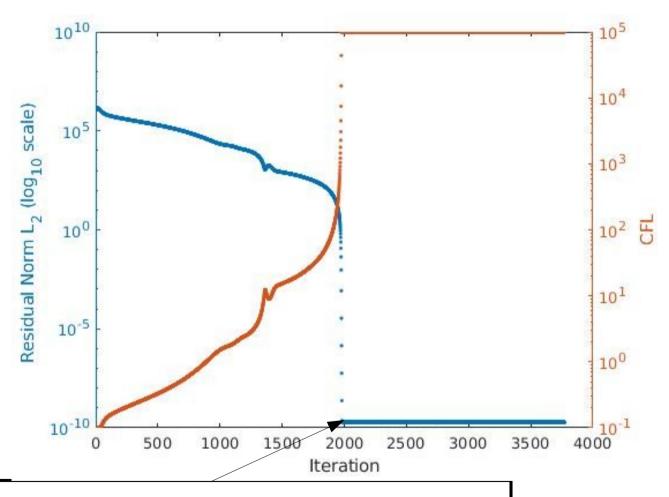


## Residual on refined 7424 elements mesh P0 (tol=10^(-16))



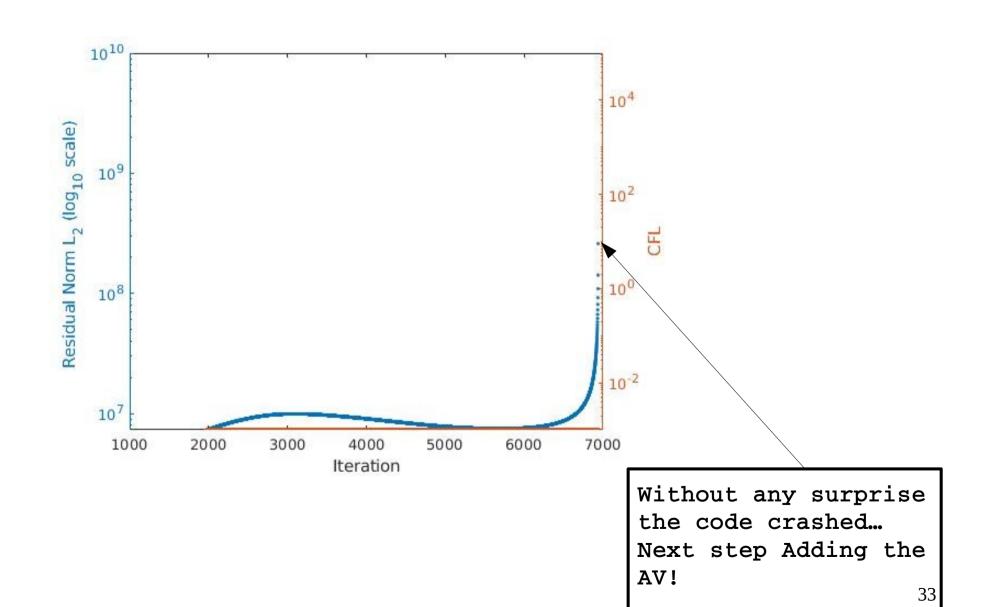
We reached a precision of 10^(-10) that seems to be the best result (I had to stop the code manually at one point)

## Residual on refined 7424 elements mesh P0 (tol=10^(-16))

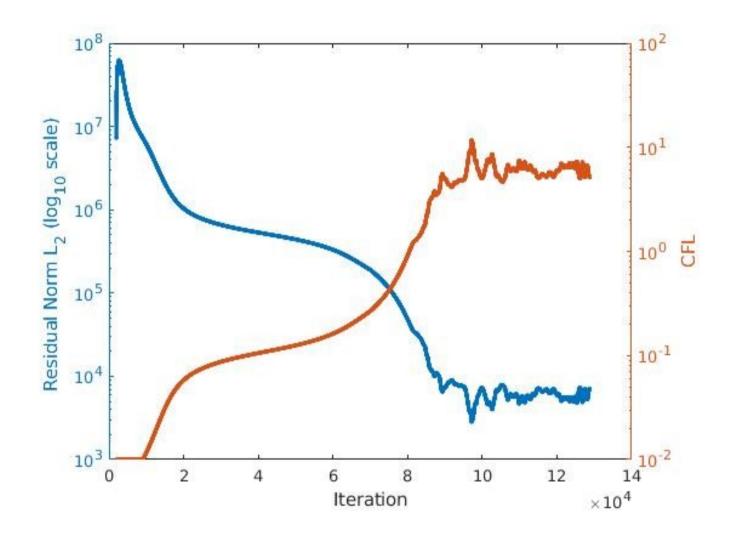


How can we improve the convergence ? By increasing the order of the DGM !

# Solution along the wall on Refined 7424 element mesh P1 without Artificial Viscosity



# Solution along the wall on Refined 7424 element mesh P1 WITH Artificial Viscosity



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#### Conclusion & Futur Work

• The artificial viscosity and the CFL are the two main components that allow the convergence for the Galerkin Method in high order in the Argo platform.

#### • Suggestion of future work:

• Adapt the the modified Localized Laplacian Artificial Viscosity (LLAV) applied in the flux reconstruction method in COOLFluiD from the following article.

Ray Vandenhoeck, "Massively Parallel and Robust High-Order Methods for Transitional Hypersonic Flow Modelling on Unstructured Grids"

## BACKUP SLIDES

## **Local lifting**

$$\int_{K^L \cup K^R} \phi^K r_h^e \left( \llbracket \boldsymbol{w}_h \rrbracket \right) dV = \begin{cases} -\int_e \left\{ \phi^K \right\} \llbracket \boldsymbol{w}_h \rrbracket \otimes \boldsymbol{n} dS & if \ e \in \epsilon_i \\ -\int_e \phi^{K,+} \llbracket \boldsymbol{w}_h \rrbracket \otimes \boldsymbol{n} dS & if \ e \in \epsilon_{wall} \end{cases}$$

## **Global Lifting**

$$\sum_{K \in \Omega_h} \int_K \phi^K R_h \left( \llbracket \boldsymbol{w}_h \rrbracket \right) dV = -\sum_{e \in \epsilon_i} \int_e \left\{ \phi^K \right\} \llbracket \boldsymbol{w}_K \rrbracket \otimes \boldsymbol{n} ds + \sum_{e \in \epsilon_{wall}} \int_e \phi^{K,+} \llbracket \boldsymbol{w}_K \rrbracket \otimes \boldsymbol{n} ds$$

## <u>Discretisation of the diffusive term</u>: The Interior penalty method limitation

Gouverning PDE's expressed in conservative and hypervectorial form:

$$\frac{\partial \mathbf{U}}{\partial t} + \nabla \cdot \mathbf{F}(\mathbf{U}) = \nabla \cdot \mathbf{F}_{v}(\mathbf{U}, \nabla \mathbf{U}) + \mathbf{S}(\mathbf{U}, \nabla \mathbf{U})$$

Weak form of the equation (DG method):

$$\sum_{\Omega_{e}} \int_{\Omega_{e}} v \frac{\partial u}{\partial t} d\Omega_{e} + \sum_{\Omega_{e}} \int_{\Omega_{e}} v \frac{\partial F(u)}{\partial x} d\Omega_{e} = \sum_{\Omega_{e}} \int_{\Omega_{e}} v \frac{\partial}{\partial x} \left( D \frac{\partial u}{\partial x} \right) d\Omega_{e}$$

$$[v] = v^{-} n^{-} + v^{+} n^{+} n^{+} n^{-} + v^{-} n^{-} +$$

Rewritting the diffusive term as:

$$\sum_{\Omega_{e}} \int_{\Omega_{e}} v \frac{\partial}{\partial x} \left( D \frac{\partial u}{\partial x} \right) d\Omega_{e} = \sum_{\Omega_{e}} \int_{\Omega_{e}} \frac{\partial v}{\partial x} \left( D \frac{\partial u}{\partial x} \right) d\Omega_{e}$$

$$- \sum_{\partial \Omega_{e}} \int_{\partial \Omega_{e}} \left\langle D(u) \frac{\partial u}{\partial x} \right\rangle [v] dS - \theta \sum_{\partial \Omega_{e}} \int_{\partial \Omega_{e}} \left\langle D(u) \frac{\partial v}{\partial x} \right\rangle [u] dS + \alpha \sum_{\partial \Omega_{e}} \int_{\partial \Omega_{e}} [v] [u] dS$$

The interior penalty coefficient  $\theta$  can take the value -1, 0 or 1.