

# Marlics: A finite difference liquid crystal simulation package

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## Abstract

In this paper we present Marlics to the world. Marlics is a software written in C++ to solve the Beris-Edwards equation of nematodynamics without flow for both nematic and cholesteric liquid crystals. The program takes as input an descriptive file giving the simulations parameters and initial conditions generating a series of different snapshots. The code is organized in class modules which can be modified by the user base to attend their further needs. **Review the abstract after the paper is finished**

**Keywords:** Liquid crystals, Landau-de Gennes, finite differences.

## PROGRAM SUMMARY

*Program Title: MarLicS*

*Licensing provisions: GNU General Public License v3.0 (GPL)*

*Programming language: C++*

*Supplementary material: Complete instructions about program usage can be found in the user-guide.*

*Nature of problem: Marlics was developed to simulate liquid crystal devices via solution of the Beris-Edwards system of differential equations without flow.*

*Solution method: The system of equations is solved using finite differences in both time and space. The time integration is performed using an explicit integrator with or without variable time-step.*

*External routines: The code needs the GSL (Gnu scientific library), an implementation of the CBLas library and an implementation of the OpenMp library(optional).*

*Running time: From minutes to hours depending on the problem size.*

*Computer: Single or multi-core processor with shared memory.*

*RAM: From hundreds of megabytes to gigabytes depending on the problem size.*

*Restrictions: The code is parallelized using OpenMp, consequently it can only be run in parallel with shared memory processors.*

*Additional comments: The source code comes with a Mathematica notebook with can be used to aid the user to implement additional interactions, boundary conditions or other situations situation not covered by the current software.*

## 1. Introduction

### Misses an introduction to optical devices or light modulating devices.

Liquid crystals are excellent materials to perform these functions, since it can selectively reflect or transmit incoming light depending on its state, which can be controlled by an external perturbation. In displays, one of the most common application, it is used electric fields to control the liquid crystals state [1]. There are many types of LCDs current available, being the major difference among them given by the distribution of electrodes in the device, and the inclusion of protrusions in the boundaries of the confining surface [2].

The design of new devices requires a great amount of experimentation and empirical knowledge. Performing every trial with a real apparatus would require a unpractical number of prototypes to construct and tested. As an alternative, the researcher can turn to modeling softwares to aid in the discovery process. Before prototyping a device it can be simulated in a package in many forms before being constructed saving time and resources.

There are some commercial available softwares for simulation of liquid crystals devices, for instance LC3D[3] and LCD master[4]. These softwares provides out of the box capabilities for simulation and visualization, however it lacks the possibility for extension by the user. Open sources alternatives have been appearing to fill this gap, Licra [5] is one of this attempt. The licra software is written in C an support a few modes of use, however, its implementation is very monolithic, and if the user wants to change a simulation parameter, or a mode of operation, he/she has to code it directly to the source code and recompile the executable. Recently, another open source liquid crystal simulation software has been released [6]. The program named openQmin, performs the minimization of the liquid crystal energy, searching for its minimum energy state. The program works in many situations and can be fine tuned by the user via a script or a graphical interface. It has many desirable capabilities, and performs very well, but the program is focused

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in static problems, and in some situations the dynamic operation of the device is a matter of great importance.

In this paper we present *marlics*, it is an open source code designed to simulate the dynamics of liquid crystal order parameters. The code is written in C++ using an independent system of classes, which provide the building blocks for the most common cases, and provide a framework where the user can develop their specific applications. The mode of operation is defined by a script which is passed to the software.

## 2. Theoretical background

Liquid crystal are anisotropic liquids which have a certain degree of long range ordering. Nematic LCs presents orientational order, but no positional one. Its state can be described by a combination of scalar and vectorial fields. The preferred direction can be represented by a vector  $\vec{n}$ , and due to the phase organization  $\vec{n}$  and  $-\vec{n}$  are equivalent. In some circumstances, it happens that the liquid crystal presents a second preferred direction, which it is also described by a vector  $\vec{l}$  called co-director. The degree which the system is ordered in the direction of the director is given by the scalar order parameter  $S$ , and by the scalar order parameter  $P$  in the co director direction. In this way, when  $S = 1$  the system is perfectly oriented along  $\vec{n}$ , while  $S = 0$  gives an isotropic liquid (no preferred orientation). It is possible to have  $S < 0$ , in this case the molecules of the liquid crystal are oriented in average perpendicular to the director  $\vec{n}$ .

The density of energy  $f_s(S)$  associated with the scalar order parameter  $S$  is given by:

$$f_L(S) = \frac{1}{2}a(T - T^*)S^2 - \frac{1}{3}S^3 + \frac{1}{4}cS^4 + \frac{1}{2}L(\nabla S)^2 \quad (1)$$

where  $\{a, b, c, L\}$  are thermodynamics constants. The energy associated with variation of  $\vec{n}$  in space is given by the Frank density of energy:

$$f_f(\vec{n}) = \frac{1}{2}K_{11}(\nabla \cdot \vec{n})^2 + \frac{1}{2}K_{22}(\vec{n} \cdot \nabla \times \vec{n} + q_0)^2 + \frac{1}{2}K_{33}|\vec{n} \times \nabla \times \vec{n}|^2 + \frac{1}{2}K_{24}\nabla \cdot (\vec{n} \times \nabla \times \vec{n} + \vec{n} \cdot \nabla \vec{n}) \quad (2)$$

being  $\{K_{11}, K_{22}, K_{33}, K_{24}\}$  elastic constants with units of force.

The scalar order parameters  $\{S, P\}$  and the director and co-director  $\{\vec{n}, \vec{l}\}$  can be combined in a unique order parameter  $\mathbf{Q}$ , which is a second rank tensor whose elements are given by:

$$Q_{ij} = \frac{1}{2}S(3n_i n_j - \delta_{ij}) + \frac{P}{2}(l_i l_j - m_i m_j) \quad (3)$$

where  $i = 1, 2, 3$  and  $j = 1, 2, 3$ . This order parameter is symmetric and traceless, therefore only 5 independent elements, for example  $\{Q_{11}, Q_{12}, Q_{13}, Q_{22}, Q_{23}\}$  are necessary to fully determine it.

Deviations from the equilibrium value of the scalar order parameter, or spatial variations of the director has associated energy density given by Landau- de Gennes energy density

$(f_{LDG}(\mathbf{Q}))$ :

$$f_{LDG}(\mathbf{Q}) = \frac{a}{2}(T - T^*)\text{Tr}(\mathbf{Q}^2) + \frac{B}{3}\text{Tr}(\mathbf{Q}^3) + \frac{C}{4}\text{Tr}(\mathbf{Q}^2)^2 + \frac{1}{2}L_1(\partial_i Q_{jk})(\partial_i Q_{jk}) + \frac{1}{2}L_2(\partial_i Q_{ji})(\partial_k Q_{jk}) + \frac{1}{2}L_3 Q_{ij}(\partial_i Q_{kl})(\partial_j Q_{kl}) + \frac{4\pi}{P_0}L_q \epsilon_{ijk} Q_{ij}(\partial_j Q_{ik}) \quad (4)$$

where  $Q_{ij,k} = \partial Q_{ij}/\partial x_k$ ,  $\epsilon_{ijk}$  is the Levi-Civita tensor,  $\{L_1, L_2, L_3, L_q, L_s\}$  are the elastic constants,  $\{a, B, C\}$  are thermodynamic constants related to the nematic isotropic transition and  $T$  and  $T^*$  are the system temperature and the virtual nematic-isotropic phase transition temperature, respectively. Here we used Einstein summation convention in repeated indexes.

The elastic constants of the Frank energy  $\{K_{11}, K_{22}, K_{33}, K_{24}\}$  and the Landau-De Gennes ones  $\{L_1, L_2, L_3, L_q, L_s\}$  are related by the expressions:

$$\begin{aligned} L_1 &= \frac{2.0(k_{33} - k_{11} + 3.0k_{22})}{(27.0S_{eq}^2)} \\ L_2 &= \frac{4.0(k_{11} - k_{22} - k_{24})}{(9.0S_{eq}^2)} \\ L_3 &= \frac{4.0(k_{33} - k_{11})}{(27.0S_{eq}^3)} \\ L_q &= \frac{2.0(k_{22})}{(9.0S_{eq}^2)} \\ L_s &= \frac{4(k_{24})}{(9.0S_{eq}^2)} \end{aligned} \quad (5)$$

The dielectric energy density is given by:

$$f_e(\mathbf{Q}) = -\frac{1}{3}\epsilon_0 \Delta \epsilon^m E_i E_j Q_{ij} + \frac{\epsilon_0}{2} \mathbf{E} \cdot \mathbf{E}, \quad (6)$$

where  $\epsilon_0$  is the vacuum dielectric constant,  $\Delta \epsilon = \epsilon_{||} - \epsilon_{\perp}$  is the dielectric anisotropy which measures the difference between the dielectric constant parallel ( $\epsilon_{||}$ ) and perpendicular ( $\epsilon_{\perp}$ ) to the liquid crystal director. The volume energy density ( $f_v$ ) will be given by the sum of all energy terms being considered:

$$f_v(\mathbf{Q}) = f_{ldg}(\mathbf{Q}) + f_e(\mathbf{Q}) \quad (7)$$

The liquid crystal also interacts with the confining surfaces, which can induce an order parameter at the surface and a preferred direction for the direction, which is called surface easy axis  $n_0$ . One of the simplest form of the surface energy density between a liquid crystal and a surface is given by the Rapini-Papoular (also called Nobili-Durant) which is given by:

$$f_{rp}(\mathbf{Q}) = \frac{1}{2}W_1(Q_{ij} - Q_{ij}^0)(Q_{ij} - Q_{ij}^0) \quad (8)$$

where  $Q_{ij}^0$  is the surface induced order parameter, which can be given in its tensorial form, or constructed using the induced scalar order parameters  $\{S^0, P^0\}$  and the induced easy axis  $\{\vec{n}, \vec{l}\}$  using expression 3.

When the liquid crystals boundary is a liquid or gas, instead of inducing a preferred direction the surface may induces a preferred plane of orientation perpendicular to the surface. Any variation of the director inside this plane gives the same energy. This type of surface energy is described by the Fournier-Galatola anchoring energy given by [7]:

$$f_{FG}(\mathbf{Q}) = W \left( \tilde{Q}_{ij}(\mathbf{Q}) - \tilde{Q}_{ij}^\perp(\mathbf{Q}) \right) \left( \tilde{Q}_{ij}(\mathbf{Q}) - \tilde{Q}_{ij}^\perp(\mathbf{Q}) \right) \quad (9)$$

where  $W$  is the anchoring strength constant,  $\tilde{Q}_{ij}(\mathbf{Q}) = Q_{ij} + \frac{S_0}{3} \delta_{ij}$  and  $\tilde{Q}_{ij}^\perp(\mathbf{Q}) = (\delta_{ik} - v_i v_k) Q_{kl} (\delta_{lj} - v_l v_j)$ . Here  $\vec{v} = \{v_1, v_2, v_3\}$  is the normal surface vector.

The time evolution of the order parameter is given by the Beris-Edwards set of equations. If we neglect the liquid crystal flow, the time evolution of  $Q_{ij}$  in the bulk will be given by:

$$\frac{\partial Q_{ij}}{\partial t} = -\frac{1}{\mu} \left( \frac{\partial f_v(\mathbf{Q})}{\partial Q_{ij}} - \frac{\partial}{\partial x_k} \frac{\partial f_v(\mathbf{Q})}{\partial Q_{ijk}} \right) = F_{ij}(\mathbf{Q}), \quad (10)$$

where  $\mu$  is the bulk viscosity and  $\partial Q_{ij} \partial Q_{kl} = (\delta_{ik} \delta_{jl} + \delta_{jk} \delta_{il} - 2\delta_{ij} \delta_{kl}/3)$ . Meanwhile the dynamics in the bulk will be given by

$$\frac{\partial Q_{ij}}{\partial t} = -\frac{1}{\mu_s} \left( v_k \frac{\partial f_{LDG}(\mathbf{Q})}{\partial Q_{ijk}} - \frac{\partial f_{pen}(\mathbf{Q})}{\partial Q_{ij}} \right) = F_{ij}^s(\mathbf{Q}), \quad (11)$$

here  $\mu_s$  is the surface viscosity.

We solve the system of equations using the method of lines [8]. In this method the spatial and time discretization of the governing equations are performed separated and independently, being the spatial dimensions of the equations discretized first. We discretize the R.H.S of (10) and (11) by finite differences. In the bulk we have taken centered differences for both first order and second order derivatives. In the surface we have taken centered differences for derivatives perpendicular to the normal and first order to derivatives parallel to the surface normal.

In the methods of lines the temporal discretization depends on the kind of integrator intended to propagate the solution. In the current version of marlics, we have implemented only explicit integrators, thus we have approximated the time derivative by forward differences. As an example we will take the Euler method; assuming  $Q_{ij}^t$  is the value  $Q_{ij}$  at time  $t$ , the value of  $Q_{ij}^{t+\Delta t}$  can be calculated by:

$$Q_{ij}^{t+\Delta t} = Q_{ij}^t + \Delta t F_{ij}(\mathbf{Q}) \quad (12)$$

Although the Euler method is convergent and can be used in some cases, it has its drawbacks. The value of  $\Delta t$  is fixed during the simulation, and the allowed timestep size is too small for some applications. As an alternative we have provided another 2 explicit integrator with the program: explicit second order Runge-Kutta, which also has a fixed timestep but is more stable and the Dormand-Prince 5(4), which has a self adaptive timestep. We implemented the time adaption as proposed in reference [9].

### 3. Software Usage:

Here we present the basic information necessary to install and use the software. The complete information about software usage can be found in the supplementary material "Userguide".

#### 3.1. Installation:

The installation of Marlics in Unix systems is very straightforward. The source code comes with a *makefile* to help user compile it in its computer. Actually the makefile provides an automatic installation for two compilers (being one of them free). The user will need just to assure he/she has the *make* software and the necessary external libraries and their respective developer files installed. These libraries are: GSL (Gnu Scientific library), OpenMp(optional, but highly recommended) and a CBLAS implementation (you can use the GSL implementation for example). If everything is present, the user just need to open a terminal in the program folder and type:

```
make
```

to compile the program. Once the compilation is done, the user will find a executable named *marlics* in the installation folder. For ease of use, this executable can be added to one of the system search paths for binaries files, or the user can add the installation folder to the list of search-able paths. Its is also possible to run the simulations in the same path that the software is installed, but we strong recommend against, since it can be become very confuse.

#### 3.2. Simulation set up and Execution:

To execute marlics you must call the program passing an input file which sets the simulation parameter as follows:

```
marlics input_file
```

being *input\_file* the file containing the simulation parameters.

We already provide input files for some situations that the user can use to test the program, or as a base to their own simulations. All the parameters necessary to set up the simulation must be passed to the program via the input file. An entry in the input file is set by passing the parameter name followed by the parameter required values:

```
parameter value
```

Comments can be placed in the input file starting a line with "#", everything in this line will be ignored by marlics. Also, everything following the required parameters will also be ignored by marlics. We found it quite useful to let the parameters units after its values. Some parameters must be set, while others can have standard values associated with them. Whenever marlics use an standard value, it inform the user in the standard output which value was used. The parameters standard values and its units can be found in the table Appendix B.

The simulation constants must be filled with a real number. The user has two options to pass the elastic constants: the user can pass the  $L_i$  values, which are the constants actually used in the calculations, or pass the  $k_{ii}$  values and let marlics calculate

188  $L_i$ . Also the chirality power can be passed in two different ways,240  
189 as  $p_0$ , i.e the helix passe, or as  $q_0$ , i.e the helix vector.241

190 We provide the most common initial conditions: random, ho-242  
191 mogeneous oriented along a direction  $\vec{v}$ . We have also provide243  
192 an initial condition set as `read_from_file` and pass an file con-244  
193 taining the initial condition. The refereed file must be formatted  
194 as the standard output file presented in section 3.3. The remain-245  
195 ing initial co nditions can be checked in the user manual.

196 We provided 3 differnt geometries in marlics: bulk, sphere246  
197 and slab. Each geometry has its number of boundary con-247  
198 ditions, in the case, 0, 1 and 2 respectively. To define a  
199 boundary the user must start a line with the *boundary* key-248  
200 word followed by the boundary name and its number. There249  
201 is 4 forms of boundary conditions implemented in marlics:250  
202 Rapinni-Papoular (in Nobili-Durant) form, Fournier-Galatola,251  
203 strong boundaries and homeotropic.

204 For more details see the supplementary file “Userguide”. For252  
205 reference, we also include an example of input file in the ap-  
206 pendix Appendix C.253

### 207 3.3. Output files:

208 The software produces two outputs: an log of the program257  
209 execution, and a series of files containing the spatial distribu-258  
210 tion of the order parameters. The log has the function of in-259  
211 forming the user about the parameters read by the program, so260  
212 it can be used as reference in the future, and informing the cur-261  
213 rent state of the simulation. The log is printed in the standard262  
214 output, which can aid the preparation of the input file, but we263  
215 strong recommend redirecting it to a separate file for reference264  
216 in future.265

217 The main output of the software are the files containing the266  
218 spatial distribution of the order parameter: the main LC direc-267  
219 tor director  $\vec{n}$ , the co-director  $\vec{l}$ , the uniaxial order parameter  $S$ 268  
220 and biaxial order parameter  $P$ . We decided to output the order269  
221 parameters in this form instead of the elements of  $Q_{ij}$ , since the270  
222 former are more ready to use and interpret than the later. We271  
223 also preferred to refer to the position in space using the lattice272  
224 numbers instead of the space position in the Cartesian frame.273  
225 The actual position can be easily repried multiplying the col-274  
226 umn by its referred grid spacing ( $dx$ ,  $dy$  or  $dz$ ). Although every275  
227 output file has associate with it an time  $t$ , this number is not276  
228 output in the file. Instead the output file number and is referred  
229 output time is informed in the log output.

230 An example of a truncated output file can be viewed in ap-277  
231 pendix Appendix D. More information can be found the in the  
232 supplementary material “Userguide”.

## 233 4. Test problems:

234 To validate our code we performed a few standard simula-  
235 tions. Even though we are presenting a source code which can  
236 be used as a framework for user implemented situations, here  
237 we wanted to show out of the box capabilities. In this way  
238 we chose some of well documented scenarios performed in the  
239 marlics framework.

### 4.1. Bulk Nematic:

### 4.2. Cholesteric Slab:

It is know that a cholesteric liquid crystal with pitch  $p_0$   
placed inside slab with planar anchoring in both substrates will  
organize itself with the profile

### 4.3. Nematic sphere with strong anchoring

### 4.4. Cholesteric sphere with weak anchoring

## 5. Conclusions

In conclusion, we have presented marlics. Marlics is open  
source and offers some of the most common cases for device  
simulation. Moreover it provides a class organized framework  
where the user can program its on cases if necessary.

## Appendix A. Software implementation and extension:

Marlics is subdivided in a series of C++ classes. There are  
4 main super classes which confers most of marlics usability :  
energy, boundary, geometry and integrator. Each specific func-  
tionality must be implemented deriving one of these classes, **for  
instance, the Landau-de Gennes energy is implemented through  
a class derived from the energy super class.**

Boundary is a derivation of energy and contains the same  
functions plus some members containing information about  
which boundary it is associated.

Integrator contains the routine to evolve the system of equa-  
tions.

Geometry contain the information about the geometry of the  
system and their boundary conditions. Here the user must pro-  
vide a way to calculate the boundaries normals, and the pa-  
rameter field derivatives. With all inofrmation present in the  
geometry class, we put the calculation of RHS in it, therefore,  
there is a routine to calculate the evolution RHS. To do this the  
geoemtry has as member an energy instance and a linked list  
containing the boundaries pointers.

There is an final class called driver which parses the input file  
and assembly the simulation pieces, except for the boundaries,  
it setups all the other classes. In the appendix ?? we presents  
a short description showing how the user can implement their  
new situations (boundaries, integrators and geometries).

## Appendix B. List of available parameters and its units:

Parameter name	variable type	units	mandatory/standard value
Geometry	string		Yes
Nx	integer		Yes
Ny	integer		Yes
Nz	integer		Yes
dx	real	nm	Yes
dy	real	nm	Yes
dz	real	nm	Yes
integrator	string		Yes
facmin	real		0.4
facmax	real		3
prefac	real		0.8
Atol	real		0.001
Rtol	real		0.001
a	real	$\text{MJ}/(m^2 K)$	Yes
b	real	$\text{MJ}/m^2$	Yes
c	real	$\text{MJ}/m^2$	Yes
K1	real	pN	See 3
K2	real	pN	See 3
K3	real	pN	See 3
L1	real	pN	See 3
L2	real	pN	No
L3	real	pN	No
Ls	real	pN	No
Lq	real	pN/m	No
p0 or q0	real	nm or 1/nm	No
T	real	K	No
Mu or gamma	real	Pa/s	Yes
Mu_s or gamma_s	real	nm Pa m/s	See 3
ti	real	$\mu\text{s}$	0.0
tf	real	$\mu\text{s}$	Yes
dt	real	$\mu\text{s}$	Tf/1e6
timeprint	real	$\mu\text{s}$	Tf/20
timeprint_type	string		Linear
timeprint_increase_factor	real		Tf/20
output_folder	string		.
output_fname	string		director_field_\$.csv
initial_output_file_number	int		0
initial_conditions	string		yes
initial_file_name	string		See 3
theta_i	real	degrees	See 3
phi_i	real	degrees	See 3
anchoring_type	int + string		See 3
Wo1	int + real		See 3
theta_0	int + real		See 3
phi_0	int + real		See 3

279 **Appendix C. Input File example:**

280     Here you can find the complete input file:

```

281
282 #Geometry Parameters:
283 geometry slab
284 Nx 200 /* grid size */
285 Ny 200 /* grid size */
286 Nz 100 /* grid size */
287 dx 10.0 /* 10^-9 m */
288 dy 10.0 /* 10^-9 m */
289 dz 10.0 /* 10^-9 m */
290
291
292 #Integrator parameters:
293 integrator DP5
294 atol 0.005
295 rtol 0.005
296 facmax 3.0
297 facmin 0.4
298 prefac 0.8
299
300
301 #Liquid crystal parameters:
302 a 0.182
303 b -2.12
304 c 1.73
305 T -1 Kelvin
306 k11 16.7 /* pN */
307 k22 7.8 /* pN */
308 k33 18.1 /* pN */
309 k24 0 /* pN */
310 p0 500
311 mu_1 0.3 /* Pa s */
312 mu_1_s 30.0 /* Pa nm s */
313
314
315 #Time parameters:
316 dt 0.001 /* 10^-6 s */
317 ti 0.0 /* 10^-6 s */
318 tf 5000.0 /* 10^-6 s */
319
320 #Output Parameters:
321 time_print_type logarithmic
322 timeprint 50. /* 10^-6 s */
323 timeprint_increase_factor 1.16
324 output_folder .
325 output_fname output_$.csv
326 initial_output_file_number 0
327
328 #Initial conditions:
329 initial_conditions random
330
331
332 #Boundaries conditions
333
334 #Bottom boundaries:
335 anchoring_type 0 Rapini-Papoular
336 Wol 0 1000.0
337 theta_0 0 45.0

```

```
338     phi_0          0  45.0
339
340 #Top boundaries:
341 anchoring_type 1  Fournier-Galatola
342 Wol            1  1000.0
```



## Appendix D. Output file:

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