#### Example 6-1:

How to use (6.2) to find  $\mathcal{L}\{\sin(\omega_o t)u(t)\}$ .

$$\begin{split} \mathcal{L} \left\{ \sin\left(\omega_{o}t\right) u(t) \right\} &= \int_{0^{-}}^{\infty} \sin\left(\omega_{o}t\right) \exp\left(-st\right) dt \\ &= \frac{1}{2j} \int_{0^{-}}^{\infty} \left[ \exp\left(j\omega_{o}t\right) - \exp\left(-j\omega_{o}t\right) \right] \exp\left(-st\right) dt \\ &= \frac{1}{2j} \int_{0^{-}}^{\infty} \left[ \exp\left((j\omega_{o}-s)t\right) - \exp\left(-(j\omega_{o}+s)t\right) \right] dt \\ &= \frac{1}{2j} \left[ \frac{\exp\left((j\omega_{o}-s)t\right)}{(j\omega_{o}-s)} - \frac{\exp\left(-(j\omega_{o}+s)t\right)}{-(j\omega_{o}+s)} \right]_{0^{-}}^{\infty} \\ &= \frac{1}{2j} \left[ -\frac{1}{(\mathbf{j}\boldsymbol{\omega}_{o}-s)} - \frac{1}{(\mathbf{j}\boldsymbol{\omega}_{o}+s)} \right] \cdots \underbrace{\frac{\operatorname{Re}[s] > 0}{\operatorname{Region of convergence}}}_{\operatorname{Region of convergence}} \\ &= \frac{\omega_{o}}{s^{2} + \omega_{o}^{2}} \end{split}$$

The range of values of s for which the Laplace transform converges is called the region of convergence (ROC).

Can you find  $\mathcal{L}\{\cos(\omega_o t)u(t)\}$ ?

#### Multiple Pole Case

If D(s) has multiple roots, i.e., it contains factors of the form  $(s+p_n)^r$ , we say that  $-p_n$  is a multiple pole of F(s) with multiplicity r. The expansion of F(s) will consist of terms of the form

$$\frac{\gamma_{1}}{s+p_{n}} + \frac{\gamma_{2}}{(s+p_{n})^{2}} + \dots + \frac{\gamma_{r}}{(s+p_{n})^{r}}$$
 (6.19)

where

$$\gamma_{r-k} = \frac{1}{k!} \frac{d^k}{ds^k} \left[ \left( s + p_n \right)^r F(s) \right]_{s=-p_n}; \quad k = 0, 1, \dots, r-1.$$
(6.20)

### Example 6.4:

$$F(s) = \frac{s^2 + 2s + 5}{(s+3)(s+5)^2} = \frac{\alpha_1}{(s+3)} + \frac{\gamma_1}{(s+5)} + \frac{\gamma_2}{(s+5)^2}$$

$$\text{Using } (\mathbf{6.18}) : \alpha_1 = (s+3)F(s)\big|_{s=-3} = 2$$

$$\text{Using } (\mathbf{6.20}) : \gamma_2 = (s+5)^2F(s)\big|_{s=-5} = -10$$

$$\text{Using } (\mathbf{6.20}) : \gamma_1 = \frac{d}{ds}\Big[(s+5)^2F(s)\Big]\big|_{s=-5} = \frac{d}{ds}\Big[\frac{s^2 + 2s + 5}{s+3}\Big]\bigg|_{s=-5} = \frac{s^2 + 6s + 1}{(s+3)^2}\bigg|_{s=-5} = -1$$

$$\therefore F(s) = \frac{2}{(s+3)} - \frac{1}{(s+5)} - \frac{10}{(s+5)^2} \implies \left(f(t) = \mathcal{L}^{-1}\{F(s)\} = \left(2e^{-3} - e^{-5} - 10te^{-5t}\right)u(t)\right)$$

# **B.** F(s) is an Improper Rational Function $(M \ge N)$

If  $M \geq N$  , we can apply long division to express F(s) in the form

$$F(s) = \frac{N(s)}{D(s)} = Q(s) + \frac{R(s)}{D(s)}$$

$$(6.21)$$

Errata

 $\text{such that the } \begin{cases} \text{Quotient} &: Q(s) \text{ is a polynomial in } s \text{ with degree } (M \text{-} N), \\ \text{Remainder} &: R(s) \text{ is a polynomial in } s \text{ with degree strictly less than } N. \end{cases}$ 

The inverse Laplace transform of R(s)/D(s), which is now a **proper rational function**, can be computed by first expanding into partial fractions.

The inverse Laplace transform of  $oldsymbol{Q}(s)$  can be computed by using

$$\mathcal{L}^{-1}\left\{s^{k}\right\} = \frac{d^{k}}{dt^{k}}\delta(t); \quad k = 0, 1, 2, \dots$$
(6.22)

Example 6.5:

$$F(s) = \underbrace{\frac{2s^2 + 10s + 10}{(s+1)(s+3)}}_{\text{by long division}} = 2 + \underbrace{\frac{2s+4}{(s+1)(s+3)}}_{\text{by long division}}$$
$$\therefore F(s) = 2 + \underbrace{\frac{1}{s+1}}_{+1} + \underbrace{\frac{1}{s+3}}_{+3} \implies \left( f(t) = \mathcal{L}^{-1} \left\{ F(s) \right\} = 2\delta(t) + e^{-t}u(t) + e^{-3t}u(t) \right)$$

#### 6.4 Relationship between the Fourier Transform and the Laplace Transform

Fourier Transform 
$$\left. : \Im \left\{ f(t) \right\} = \int_{-\infty}^{\infty} f(t) \exp(-j\omega t) dt \right.$$
 (6.23)

Bilateral Laplace Transform 
$$\left. \right\}$$
:  $\tilde{F}(s) = \int_{-\infty}^{\infty} f(t) \exp(-st) dt$  (6.24)

Unilateral Laplace Transform 
$$\left. \right\}$$
:  $F(s) = \int_{0^{-}}^{\infty} f(t) \exp(-st) dt$ 



• From (6.23) and (6.24), we see that the Fourier transform is a special case of the **bilateral** Laplace transform in which  $\mathbf{s} = \mathbf{j}\boldsymbol{\omega}$ , that is

$$\Im\{f(t)\} = \tilde{F}(s)\Big|_{s=j\omega} \tag{6.26}$$

• Setting  $\mathbf{s} = \boldsymbol{\sigma} + \boldsymbol{j}\boldsymbol{\omega}$  in (6.24), we have

$$\tilde{F}(\sigma + j\omega) = \int_{-\infty}^{\infty} f(t) \exp(-(\sigma + j\omega)t) dt = \int_{-\infty}^{\infty} [f(t) \exp(-\sigma t)] \exp(-j\omega t) dt$$

$$= \Im\{f(t) \exp(-\sigma t)\} \tag{6.27}$$

which shows that the **bilateral** Laplace transform of f(t) can be viewed as the Fourier transform of  $f(t)e^{-\sigma t}$ .

• Considering the Laplace transform as a generalization of the Fourier transform where the frequency variable is generalized from  $j\omega$  to  $s = \sigma + j\omega$ , the complex variable s is often referred to as the complex frequency.

• It should not be automatically assumed that the Fourier transform of f(t) is the Laplace transform with s replaced by  $j\omega$ . If f(t) is absolutely integrable, i.e.  $\int_{-\infty}^{\infty} |f(t)| \, dt < \infty \quad \text{, then the}$ 

Fourier transform of f(t) can be obtained from its bilateral Laplace transform with s replaced by  $j\omega$ . This is not generally true if f(t) is not absolutely integrable.

The above relationship between the Fourier Transform and the *Bilateral* Laplace Transform extends fully to *Unilateral* Laplace transform if f(t) is, in addition, a right-sided function.

Example 6.6:

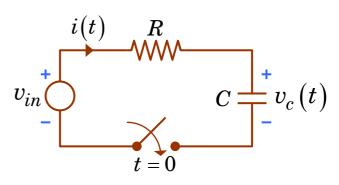
	f(t)	F(s)	$\Im\{f(t)\}$	Right-sided?	Absolutely Integrable?	$\Im\{f(t)\} = F(j\omega)$
Unit Impulse	$\delta(t)$	1	1	Yes	Yes	Yes
Unit Step	u(t)	$\frac{1}{s}$	$\pi\delta(\omega) + \frac{1}{j\omega}$	Yes	No	No
Exponential	$\exp(- \alpha t)u(t)$	$\frac{1}{s+ \alpha }$	$rac{1}{j\omega+ lpha }$	Yes	Yes	Yes

#### Example 6.7: Series RC Circuit

For the RC circuit shown, find the voltage  $v_c\left(t
ight)$  across the capacitor C .

$$RC\frac{dv_{c}\left(t\right)}{dt} + v_{c}\left(t\right) = v_{in}$$

$$sRCV_{c}\left(s\right) - RCv_{c}\left(0^{-}\right) + V_{c}\left(s\right) = \frac{v_{in}}{s}$$



$$V_{c}\left(s\right) = \frac{RCv_{c}\left(0^{-}\right)}{sRC+1} + \frac{v_{in}}{s\left(sRC+1\right)} = \frac{RCv_{c}\left(0^{-}\right)}{sRC+1} + \frac{v_{in}}{s} - \frac{v_{in}RC}{sRC+1} \quad (\clubsuit)$$

$$\begin{aligned} v_c\left(t\right) &= \mathcal{L}^{-1}\left\{V_c\left(s\right)\right\} = \mathcal{L}^{-1}\left\{\frac{v_c\left(0^-\right) - v_{in}}{s + \frac{1}{RC}}\right\} + \mathcal{L}^{-1}\left\{\frac{v_{in}}{s}\right\} \end{aligned}$$

$$= \left[v_c\left(0^-\right) - v_{in}\right] \exp\left(-\frac{t}{RC}\right) + v_{in}$$

From  $(\clubsuit)$  and  $(\blacktriangledown)$ , we observe that:  $\lim_{t\to\infty} v_c(t) = \lim_{s\to 0} sV_c(s) = v_{in}$ 

## RETURN 1

# Unilateral Laplace Transform : $X(s) = \int_{0^{-}}^{\infty} x(t) \exp(-st) dt$

ROC added

	x(t)	X(s)	ROC
Unit Impulse	$\delta(t)$	1	All $s$
Unit Step	u(t)	1/s	$\operatorname{Re}[s] > 0$
Ramp	tu(t)	$1/s^{2}$	$\operatorname{Re}[s] > 0$
n <sup>th</sup> order Ramp	$t^n u(t)$	$\frac{n!}{s^{n+1}}$	$\operatorname{Re}[s] > 0$

	x(t)	X(s)	ROC
Exponential	$\exp(-\alpha t)u(t)$	$1/(s+\alpha)$	$\operatorname{Re}[s] > -\operatorname{Re}[\alpha]$
Damped Ramp	$t\exp(-lpha t)u(t)$	$1/(s+\alpha)^2$	$\operatorname{Re}[s] > -\operatorname{Re}[\alpha]$
Cosine	$\cos(\omega_0 t)u(t)$	$s/(s^2+\omega_o^2)$	$\operatorname{Re}[s] > 0$
Sine	$\sin(\omega_0 t)u(t)$	$\omega_o/(s^2+\omega_o^2)$	$\operatorname{Re}[s] > 0$

	x(t)	X(s)	ROC
Damped Cosine	$\exp(-lpha t)\cos(\omega_{o}t)u(t)$	$(s+\alpha)/[(s+\alpha)^2+\omega_o^2]$	$\operatorname{Re}[s] > -\operatorname{Re}[\alpha]$
Damped Sine	$\exp(-\alpha t)\sin(\omega_o t)u(t)$	$\omega_o/[(s+\alpha)^2+\omega_o^2]$	$\operatorname{Re}[s] > -\operatorname{Re}[\alpha]$

	x(t)	X(s)	
Step response of 1 <sup>st</sup> order system	$\Big[1-\exp\!\left(-at\right)\Big]u(t)$	$rac{a}{sig(s+aig)}$	$\omega_d$
Step response of 2 <sup>nd</sup> order underdamped system	$Kiggl\{1-rac{\exp\left(-\omega_{n}\zeta t ight)}{\sqrt{1-\zeta^{2}}}\siniggl[\omega_{n}\sqrt{1-\zeta^{2}}t+\phiiggr]iggr\}u(t)$	$\frac{K\omega_n^2}{s\left(s^2 + 2\zeta\omega_n s + \omega_n^2\right)}$	$\phi \wedge \phi \wedge$
	$K \left\{ 1 - \frac{\sqrt{\sigma^2 + \omega_d^2} \exp(-\sigma t)}{\omega_d} \sin[\omega_d t + \phi] \right\} u(t)$	$\frac{K\omega_d^2}{s\Big[\big(s+\sigma\big)^2+\omega_d^2\Big]}$	$-\omega_n$