

Incoherent Thomson Scattering (ITS) applied to low temperature plasma sources



electrons

laser beam

ITS on electrons

probed wavevector

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• Introduction •

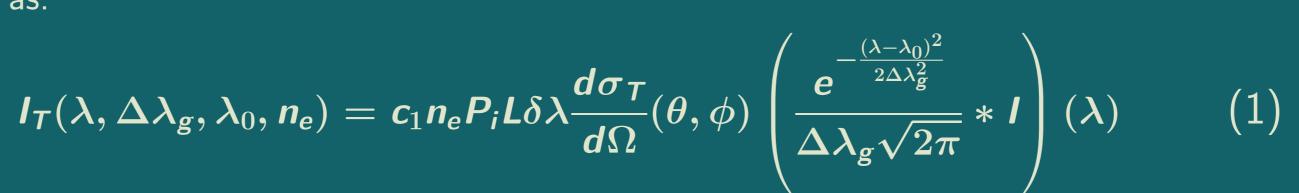
Low-temperature plasma sources have wide ranges of applications such as in electric propulsion, thin film deposition and particle sources for accelerators. As the behavior of such plasmas is strongly correlated to electron properties, reliable diagnostics able to probe electron properties are needed. Access to these information would help to increase our understanding of the physics of such complex plasma sources and validate predictive simulations under development $^{[1]}$.

O Aim O

In this work, our objective is to develop a new highly-sensitive and compact Incoherent Thomson Scattering (ITS) diagnostic for the measurement of electron properties in different low temperature plasma sources with the spatial and temporal resolution required.

Method

ITS methods involves the analysis of scattered photons by free charged particles (electrons in our study) at length scales smaller than the Debye length. Averaged over the scattering volume and the solid angle of observation, the scattered photon spectral intensity measured in the case of Maxwell-Boltzmann Electron Velocities Distribution Function (EVDF) can be expressed as:



With c_1 the total transmission factor obtained after a Raman calibration, $P_i L \delta \lambda$ some known experimental parameters, $\frac{d\sigma_T}{d\Omega}(\theta,\phi)=0$ $r_e^2(1-sin^2 heta\cos^2\phi)$ the differential Thomson scattering cross section along the observation direction and $I(\lambda)$ the normed instrument function. From a fitting of the scattered photons spectral intensity the electron density (n_e) is directly obtained, $\Delta \lambda_g$ gives the electron temperature (T_e) and $\lambda_0 - \lambda_i$ the electron drift velocity $(v_{e,drift})$.

In case of Thomson signals with high signal to noise ratio, the Electron Energy Distribution Function (EEDF) can be extracted from the normed derivative of the spectral intensity^[2]: $f_{E}(E) \propto \frac{dl}{d\lambda}$ with $E = \frac{m_e}{2} \cdot \frac{c(\lambda - \lambda_0)}{2\lambda_i sin(\theta/2)}$

• Experimental setup (THETIS^[3]) •

Transmission branch:

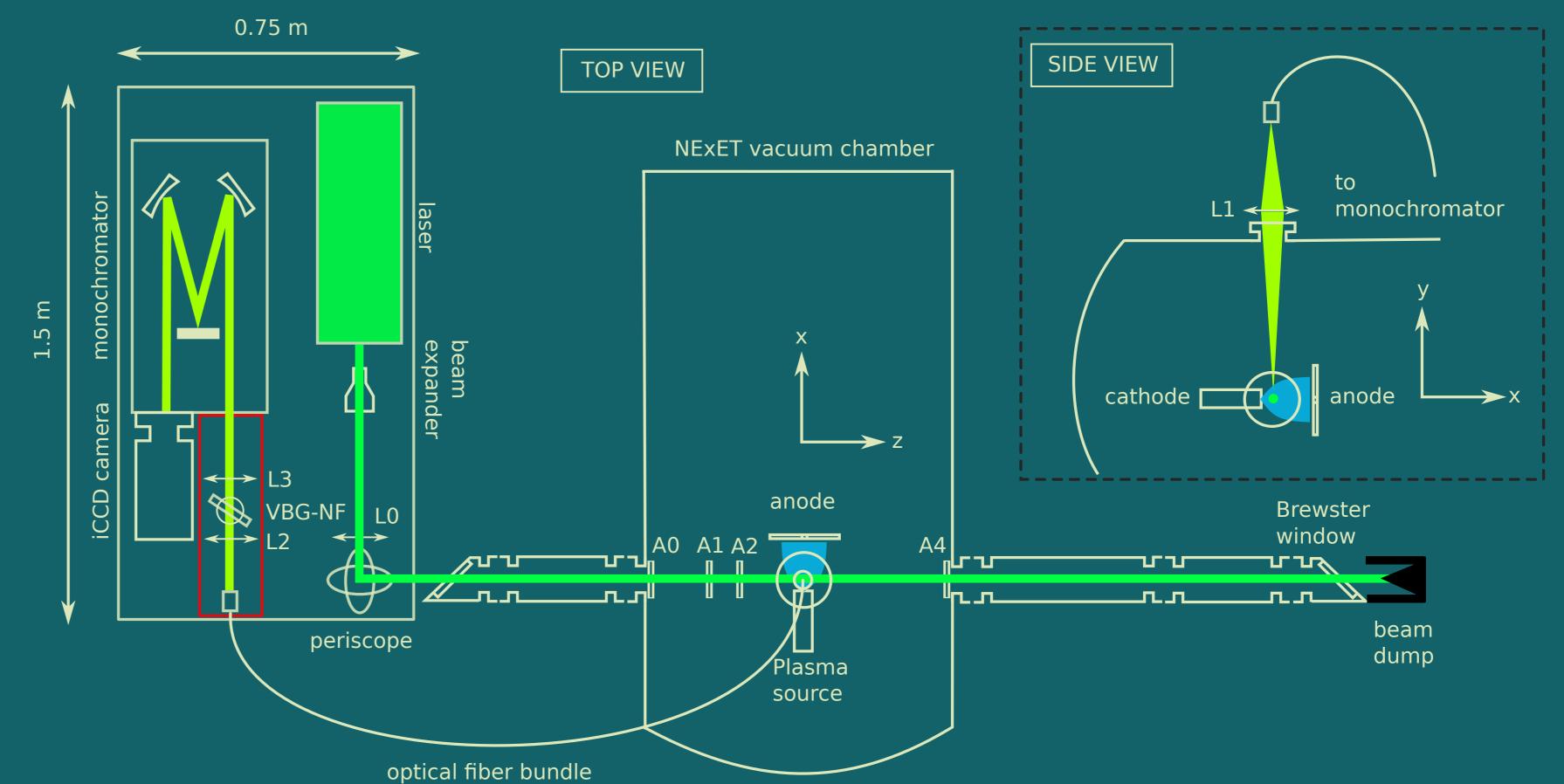
- ightharpoonup Q-switch Nd:YAG laser ($\lambda_i = 532 \; nm$; $au = 5 \; ns$; $f = 10 \; Hz$; E = 0.43 J
- ightharpoonup 2 m focal lens \Rightarrow laser beam waist $w_0 \approx 0.3 mm$
- \triangleright Brewster windows, apertures and large aperture beam dump \Rightarrow reduce stray light propagation

Detection branch:

- \blacktriangleright Fiber bundle $(5 \times 3, 0.3 \mu m) \Rightarrow$ increase the etendue collected
- ightharpoonup Volume Bragg Grating based Notch Filter (VBG-NF) \Rightarrow filter laser stray light contribution [4] with fewer losses than Triple Gratings Spectrometer (TGS)
- ightharpoonup Acton SP-2750 spectrometer \Rightarrow disperse the collected light
- \blacktriangleright ICCD PI-MAX 5 camera (Gen II intensifier) \Rightarrow gated detection to reduce plasma stray light contribution and obtain temporal resolution

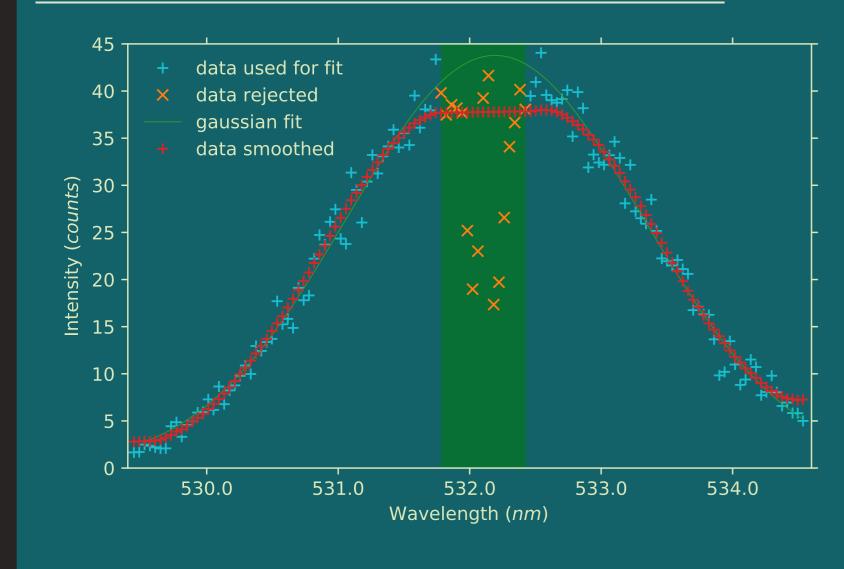
Plasma sources under study:

- ► Hall thruster and its electron source (thermo-emissive cathode)
- ► Planar magnetron (plasma assisted thin film deposition)
- ECR source (ion source for particle accelerators)

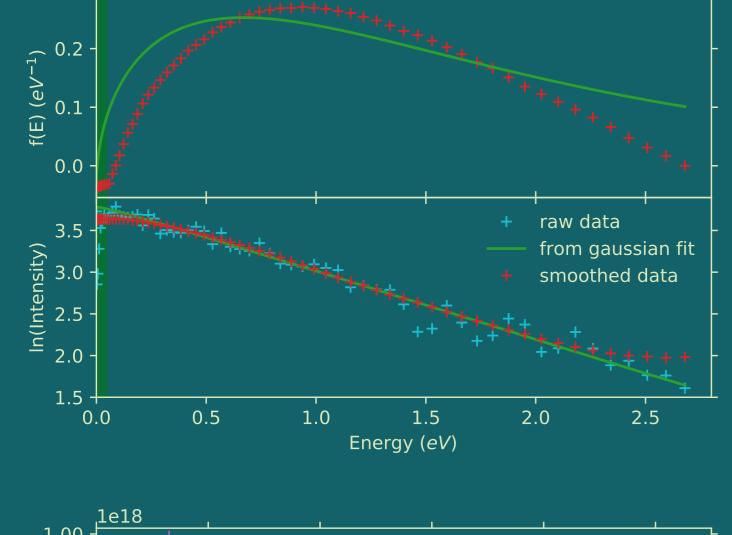


• Results •

Cathode sources (spatial probing and EEDF):

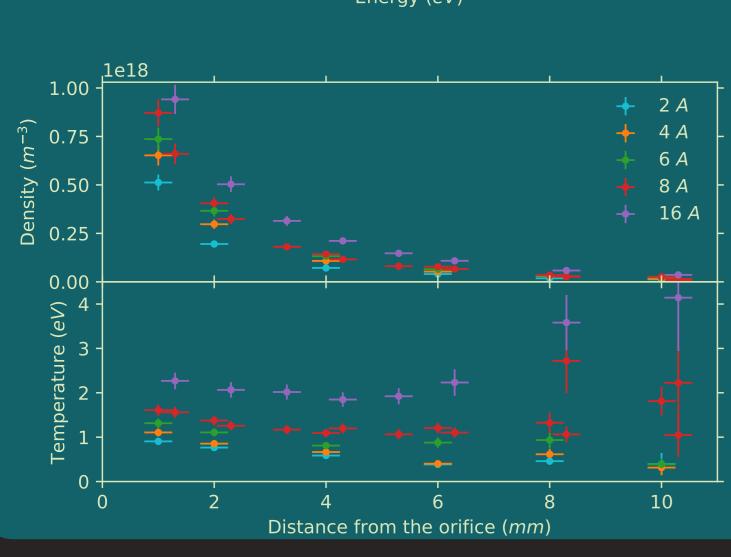


Thomson spectrum obtained $1.3\,$ mm from the cathode orifice ($I_{discharge} = 16 \ A$; $D_{Xe} = 0.8 \text{ mg.s}^{-1}$; $P_{chamber} = 10^{-3} Pa$; $N_{averaged} = 6000 \text{ pulses}$).



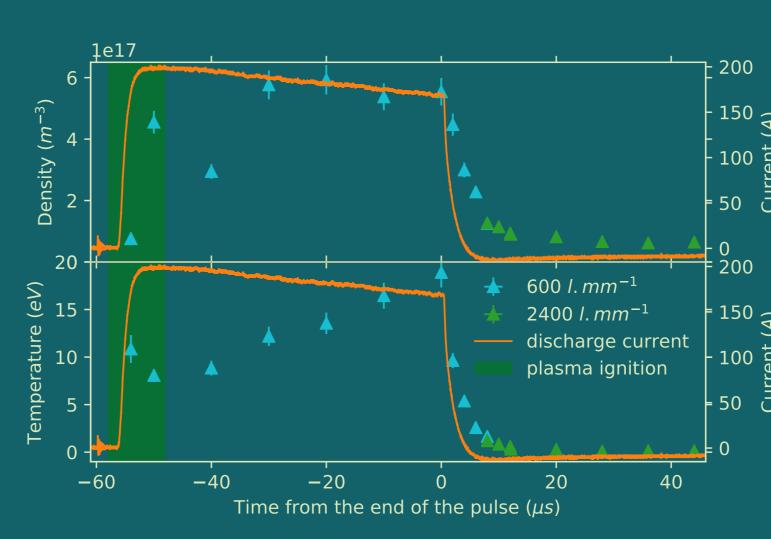
EEDF extracted from the Thomson spectrum derivative.

⇒ Weak deviations from ideal distribution coming from either the VBG-NF spectral distortion or Maxwell-Boltzmann distribution deviations (NB: the commonly used $ln(I_T)vs\Delta\lambda^2$ representation is not sensitive to such deviations).



 T_e and n_e profiles obtained under the hypothesis of Maxwell-Boltzmann EVDF (i.e. Gaussian Thomson spectrum) for various discharge current and $D_{Xe} = 0.8 \ mg.s^{-1}$. \Rightarrow Expected increase of T_e and n_e with discharge current.

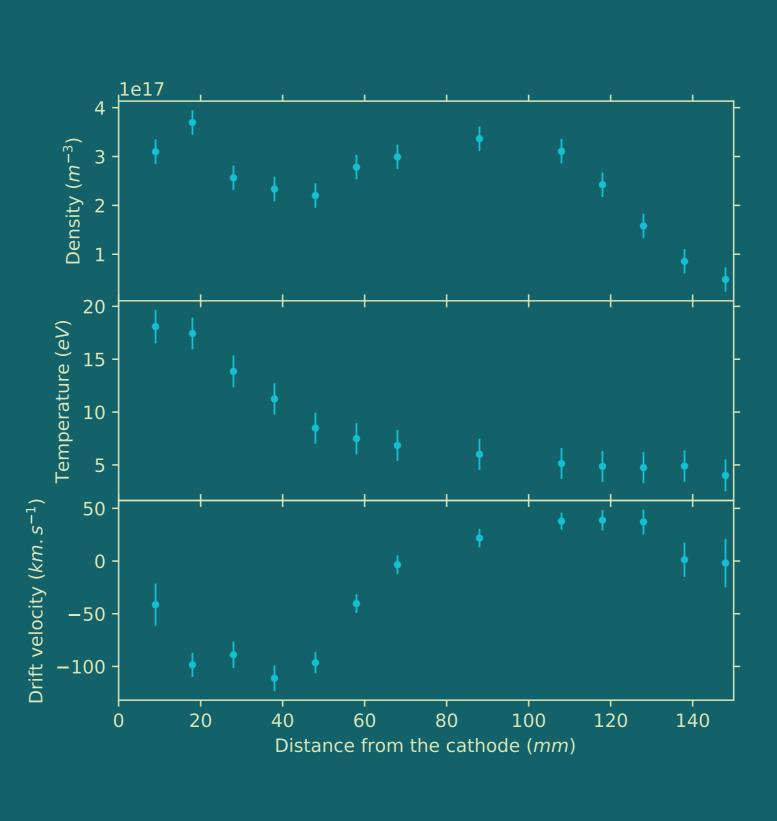
 \Rightarrow Unexpected change of trend of T_e at $\approx 5 \ mm$ (visible for high discharge current). Magnetron source (spatio-temporal probing):



Temporal profile of T_e and n_e at 9 mm from 150 $_{\widehat{\$}}$ the cathode of the magnetron, with a $10~\emph{Pa}$ 100 background pressure of Helium. \rightarrow Except during the unstable plasma

ignition, precise temporal measurement of electron properties is achieved through synchronization of the discharge with the - 50 5 laser and $extbf{ extit{N}}_{avaeraged}=6000$ $extbf{ extit{pulses}}$.

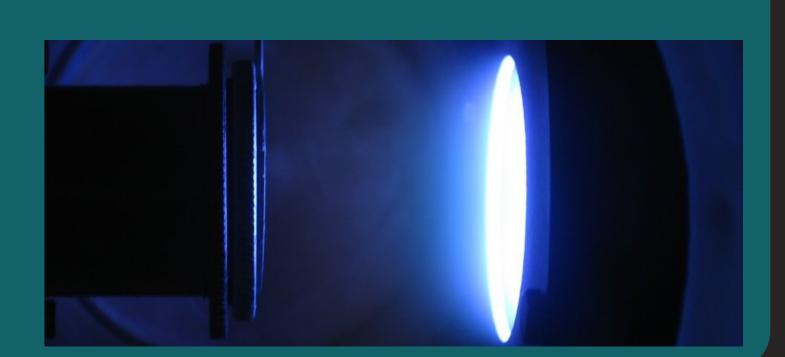
 \Rightarrow Steady $n_e~(pprox 5 imes 10^{17}~m^{-3})$ during the pulse and a linear increase of T_e (up to $\approx 20eV$).



Axial profiles of T_e , n_e and $v_{e,drift}$ obtained at $0 \ \mu s$ from the end of the pulse.

 \Rightarrow Monotonous decrease of T_e (possibly due to energy diffusion).

 \Rightarrow Non-monotonous behavior of n_e and $v_{e,drift} \Rightarrow$ Possibly due to electric field inversion^[5] and/or wave energization^[6].



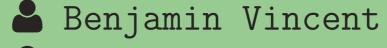
Conclusion

A new highly-sensitive ITS diagnostic has been successfully developed and applied for electron property measurement in low-temperature plasmas. Electron properties measurements by ITS have been obtained for the first time in a Hall thruster cathode and a planar magnetron with spatial resolution of $0.3 \mu m$ and temporal resolution of 15~ns.

• References •

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