

The text has given the general solution of two dimensional problem using separation of variables (equation (2.71)):

$$\Phi(\rho, \phi) = a_0 + b_0 \ln \rho + \sum_{n=1}^{\infty} a_n \rho^n \sin(n\phi + \alpha_n) + \sum_{n=1}^{\infty} b_n \rho^{-n} \sin(n\phi + \beta_n) \quad (1)$$

When we consider the interior volume $\rho < b$, it includes the origin where $\rho = 0$. This requires all the $b_i, i = 0, 1, \dots$ to vanish.

Instead of having a_n and α_n as parameters, we convert (1) into an equivalent form

$$\Phi(\rho, \phi) = a_0 + \sum_{n=1}^{\infty} a_n \rho^n \sin n\phi + \sum_{n=1}^{\infty} c_n \rho^n \cos n\phi \quad (2)$$

For the surface $\rho = b$, it's clear that a_0 is the average potential on the surface, i.e.,

$$a_0 = \langle \Phi(b, \phi) \rangle = \frac{1}{2\pi} \int_0^{2\pi} \Phi(b, \phi) d\phi \quad (3)$$

Now the Fourier coefficients can be obtained readily

$$a_n = \frac{1}{\pi} \int_0^{2\pi} \frac{\Phi(b, \phi')}{b^n} \sin n\phi' d\phi' \quad (4)$$

$$c_n = \frac{1}{\pi} \int_0^{2\pi} \frac{\Phi(b, \phi')}{b^n} \cos n\phi' d\phi' \quad (5)$$

Inserting these into (2), we get

$$\begin{aligned} \Phi(\rho, \phi) &= a_0 + \sum_{n=1}^{\infty} \left[\frac{1}{\pi} \int_0^{2\pi} d\phi' \Phi(b, \phi') \frac{\rho^n}{b^n} \overbrace{(\sin n\phi \sin n\phi' + \cos n\phi \cos n\phi')}^{\cos n(\phi - \phi')} \right] \\ &= \frac{1}{2\pi} \int_0^{2\pi} d\phi' \Phi(b, \phi') \left[1 + 2 \sum_{n=1}^{\infty} \frac{\rho^n}{b^n} \cos n(\phi - \phi') \right] \end{aligned} \quad (6)$$

Let $\gamma = \phi - \phi'$, we recognize that the sum in (6) is the real part of the sum of the geometric series., which is

$$\operatorname{Re} \sum_{n=1}^{\infty} \left(\frac{\rho e^{i\gamma}}{b} \right)^n \quad (7)$$

Recall for $|x| < 1$,

$$\sum_{n=0}^{\infty} x^n = 1 + \sum_{n=1}^{\infty} x^n = \frac{1}{1-x} \quad (8)$$

This turns (7) into

$$\begin{aligned} \operatorname{Re} \left(\frac{1}{1 - \frac{\rho e^{i\gamma}}{b}} \right) - 1 &= \operatorname{Re} \left(\frac{\rho e^{i\gamma}}{b - \rho e^{i\gamma}} \right) \\ &= \operatorname{Re} \left(\frac{\rho \cos \gamma + i \rho \sin \gamma}{b - \rho \cos \gamma - i \rho \sin \gamma} \right) \\ &= \frac{\operatorname{Re} [(\rho \cos \gamma + i \rho \sin \gamma)(b - \rho \cos \gamma + i \rho \sin \gamma)]}{(b - \rho \cos \gamma)^2 + \rho^2 \sin^2 \gamma} \\ &= \frac{b \rho \cos \gamma - \rho^2}{b^2 + \rho^2 - 2b\rho \cos \gamma} \end{aligned} \quad (9)$$

Thus the bracket of (6) becomes

$$1 + 2 \cdot \frac{b \rho \cos \gamma - \rho^2}{b^2 + \rho^2 - 2b\rho \cos \gamma} = \frac{b^2 - \rho^2}{b^2 + \rho^2 - 2b\rho \cos \gamma} \quad (10)$$

which gives

$$\Phi(\rho, \phi) = \frac{1}{2\pi} \int_0^{2\pi} \Phi(b, \phi') \frac{b^2 - \rho^2}{b^2 + \rho^2 - 2b\rho \cos(\phi - \phi')} d\phi' \quad (11)$$

For the exterior volume of the cylinder, all of b_0, a_1, a_2, \dots must vanish to make $\rho \rightarrow \infty$ not diverge, which means the admissible solution will have similar form to (2) except all ρ^n will be replaced by ρ^{-n} . Then the same argument will lead to (6) with the $\rho \leftrightarrow b$ exchange, which then leads to (11) with the same exchange $\rho \leftrightarrow b$.