

1. The force can be written as

$$\mathbf{F} = \int_V (\nabla \times \mathbf{M}) \times \mathbf{B}_e d^3x + \oint_S (\mathbf{M} \times \mathbf{n}) \times \mathbf{B}_e da \quad (1)$$

With vector identity (see Jackson inner cover)

$$\nabla (\mathbf{a} \cdot \mathbf{b}) = (\mathbf{a} \cdot \nabla) \mathbf{b} + (\mathbf{b} \cdot \nabla) \mathbf{a} + \mathbf{a} \times (\nabla \times \mathbf{b}) + \mathbf{b} \times (\nabla \times \mathbf{a}) \quad (2)$$

the first integral in (1) can be written as

$$\begin{aligned} \int_V (\nabla \times \mathbf{M}) \times \mathbf{B}_e d^3x &= - \int_V \mathbf{B}_e \times (\nabla \times \mathbf{M}) d^3x \\ &= \int_V (\mathbf{M} \cdot \nabla) \mathbf{B}_e d^3x + \int_V (\mathbf{B}_e \cdot \nabla) \mathbf{M} d^3x + \int_V \mathbf{M} \times \overbrace{(\nabla \times \mathbf{B}_e)}^0 d^3x - \int_V \nabla (\mathbf{M} \cdot \mathbf{B}_e) d^3x \\ &= \int_V (\mathbf{M} \cdot \nabla) \mathbf{B}_e d^3x + \int_V (\mathbf{B}_e \cdot \nabla) \mathbf{M} d^3x - \int_V \nabla (\mathbf{M} \cdot \mathbf{B}_e) d^3x \end{aligned} \quad (3)$$

Again with vector identities and vector calculus

$$\mathbf{a} \times (\mathbf{b} \times \mathbf{c}) = (\mathbf{a} \cdot \mathbf{c}) \mathbf{b} - (\mathbf{a} \cdot \mathbf{b}) \mathbf{c} \quad (4)$$

$$\int_V \nabla \psi d^3x = \oint_S \psi \mathbf{n} da \quad (5)$$

the second integral in (1) can be written as

$$\begin{aligned} \oint_S (\mathbf{M} \times \mathbf{n}) \times \mathbf{B}_e da &= \oint_S \mathbf{B}_e \times (\mathbf{n} \times \mathbf{M}) da \\ &= \oint_S (\mathbf{B}_e \cdot \mathbf{M}) \mathbf{n} da - \oint_S (\mathbf{B}_e \cdot \mathbf{n}) \mathbf{M} da \\ &= \int_V \nabla (\mathbf{B}_e \cdot \mathbf{M}) d^3x - \oint_S (\mathbf{B}_e \cdot \mathbf{n}) \mathbf{M} da \end{aligned} \quad (6)$$

With (3) and (6) inserted into (1), we have

$$\mathbf{F} = \int_V (\mathbf{M} \cdot \nabla) \mathbf{B}_e d^3x + \int_V (\mathbf{B}_e \cdot \nabla) \mathbf{M} d^3x - \oint_S (\mathbf{B}_e \cdot \mathbf{n}) \mathbf{M} da \quad (7)$$

Now for arbitrary vector field \mathbf{P}, \mathbf{Q} , let

$$\mathbf{R} = \int_V (\mathbf{P} \cdot \nabla) \mathbf{Q} d^3x \quad (8)$$

we see (with Einstein summation convention)

$$\begin{aligned} R_j &= \int_V \left(P_i \frac{\partial}{\partial x_i} \right) Q_j d^3x = \int_V P_i \frac{\partial Q_j}{\partial x_i} d^3x \\ &= \int_V \left[\frac{\partial (P_i Q_j)}{\partial x_i} - Q_j \frac{\partial P_i}{\partial x_i} \right] d^3x \\ &= \int_V [\nabla \cdot (Q_j \mathbf{P}) - Q_j (\nabla \cdot \mathbf{P})] d^3x \\ &= \oint_S (Q_j \mathbf{P}) \cdot \mathbf{n} da - \int_V Q_j (\nabla \cdot \mathbf{P}) d^3x \\ &= \oint_S Q_j (\mathbf{P} \cdot \mathbf{n}) da - \int_V Q_j (\nabla \cdot \mathbf{P}) d^3x \end{aligned} \quad (9)$$

Assembling R_j 's back into \mathbf{R} ,

$$\int_V (\mathbf{P} \cdot \nabla) \mathbf{Q} d^3x = \mathbf{R} = \oint_S \mathbf{Q} (\mathbf{P} \cdot \mathbf{n}) da - \int_V \mathbf{Q} (\nabla \cdot \mathbf{P}) d^3x \quad (10)$$

Applying (10) to (7) gives

$$\begin{aligned} \mathbf{F} &= \left[\oint_S \mathbf{B}_e (\mathbf{M} \cdot \mathbf{n}) da - \int_V \mathbf{B}_e (\nabla \cdot \mathbf{M}) d^3x \right] + \left[\oint_S \mathbf{M} (\mathbf{B}_e \cdot \mathbf{n}) da - \int_V \mathbf{M} (\overbrace{\nabla \cdot \mathbf{B}_e}^0) d^3x \right] - \oint_S (\mathbf{B}_e \cdot \mathbf{n}) \mathbf{M} da \\ &= \oint_S \mathbf{B}_e (\mathbf{M} \cdot \mathbf{n}) da - \int_V \mathbf{B}_e (\nabla \cdot \mathbf{M}) d^3x \end{aligned} \quad (11)$$

as desired.

2. This is a brute force calculation. The magnetization is uniform

$$\mathbf{M}(r, \theta, \phi) = M \begin{bmatrix} \sin \theta_0 \cos \phi_0 \\ \sin \theta_0 \sin \phi_0 \\ \cos \theta_0 \end{bmatrix} \quad (12)$$

so the second term in (11) vanishes. The force only has contribution from the surface integral, which is

$$\mathbf{F} = \oint_S \mathbf{B}_e (\mathbf{M} \cdot \mathbf{n}) da \quad (13)$$

where

$$\mathbf{B}_e(\theta, \phi) = B_0 \begin{bmatrix} 1 + R\beta \sin \theta \sin \phi \\ 1 + R\beta \sin \theta \cos \phi \\ 0 \end{bmatrix} \quad (14)$$

$$\begin{aligned} \mathbf{M} \cdot \mathbf{n} &= M \begin{bmatrix} \sin \theta_0 \cos \phi_0 \\ \sin \theta_0 \sin \phi_0 \\ \cos \theta_0 \end{bmatrix} \cdot \begin{bmatrix} \sin \theta \cos \phi \\ \sin \theta \sin \phi \\ \cos \theta \end{bmatrix} \\ &= M (\sin \theta_0 \cos \phi_0 \sin \theta \cos \phi + \sin \theta_0 \sin \phi_0 \sin \theta \sin \phi + \cos \theta_0 \cos \theta) \end{aligned} \quad (15)$$

Thus

$$\mathbf{F} = B_0 M R^3 \int_0^{2\pi} d\phi \int_0^\pi \sin \theta d\theta \mathbf{f}(\theta, \phi) \quad (16)$$

where

$$f_1(\theta, \phi) = \beta \sin \theta \sin \phi (\sin \theta_0 \cos \phi_0 \sin \theta \cos \phi + \sin \theta_0 \sin \phi_0 \sin \theta \sin \phi + \cos \theta_0 \cos \theta) \quad (17)$$

$$f_2(\theta, \phi) = \beta \sin \theta \cos \phi (\sin \theta_0 \cos \phi_0 \sin \theta \cos \phi + \sin \theta_0 \sin \phi_0 \sin \theta \sin \phi + \cos \theta_0 \cos \theta) \quad (18)$$

$$f_3(\theta, \phi) = 0 \quad (19)$$

where we have ignored the "1" in (14) since it has no net contribution to the force.

While performing the integral, notice the terms with $\sin \phi$, $\cos \phi$, $\sin \phi \cos \phi$ will all vanish after the $d\phi$ integral, so the result is actually easy to obtain, which is

$$\mathbf{F} = \beta B_0 M \left(\frac{4\pi R^3}{3} \right) (\sin \theta_0 \sin \phi_0 \hat{\mathbf{x}} + \sin \theta_0 \cos \phi_0 \hat{\mathbf{y}}) \quad (20)$$