

1. From Prob 3.26, we have the Green function

$$G(\mathbf{x}, \mathbf{x}') = \sum_{l=0}^{\infty} g_l(r, r') P_l(\cos \gamma) \quad \text{where}$$

$$g_l(r, r') = \begin{cases} \frac{r_{<}^l}{r_{>}^{l+1}} + \frac{1}{b^{2l+1} - a^{2l+1}} \left[\frac{l+1}{l} (rr')^l + \frac{l}{l+1} \frac{(ab)^{2l+1}}{(rr')^{l+1}} + a^{2l+1} \left(\frac{r^l}{r'^{l+1}} + \frac{r'^l}{r^{l+1}} \right) \right] & \text{for } l > 0 \\ \frac{1}{r_{>}} - \left(\frac{a^2}{a^2 + b^2} \right) \frac{1}{r'} + f(r) & \text{for } l = 0 \end{cases} \quad (1)$$

The potential with Neumann boundary condition is given by equation (1.46)

$$\Phi(\mathbf{x}) = \langle \Phi \rangle_S + \frac{1}{4\pi\epsilon_0} \int_V \rho(\mathbf{x}') G(\mathbf{x}, \mathbf{x}') d^3x' + \frac{1}{4\pi} \oint_S \frac{\partial \Phi}{\partial n'} G(\mathbf{x}, \mathbf{x}') da' \quad (2)$$

We can ignore the global constant $\langle \Phi \rangle_S$. Furthermore, for this problem, the volume integral is zero and the surface integral only has contribution from the outer sphere, where $\partial \Phi / \partial n' = -E_r = E_0 \cos \theta'$. This gives

$$\begin{aligned} \Phi(\mathbf{x}) &= \frac{1}{4\pi} \cdot b^2 \int_0^{2\pi} d\phi' \int_0^\pi \sin \theta' d\theta' E_0 \cos \theta' \sum_{l=0}^{\infty} g_l(r, b) P_l(\cos \gamma) \\ &= \frac{E_0 b^2}{4\pi} \sum_{l=0}^{\infty} g_l(r, b) \underbrace{\int_0^{2\pi} d\phi' \int_0^\pi \sin \theta' \cos \theta' d\theta' P_l(\cos \gamma)}_I \end{aligned} \quad (3)$$

By addition theorem

$$\begin{aligned} I &= \int_0^{2\pi} d\phi' \int_0^\pi \sin \theta' \cos \theta' d\theta' \left[\frac{4\pi}{2l+1} \sum_{m=-l}^l Y_{lm}^*(\theta, \phi) Y_{lm}(\theta', \phi') \right] \\ &= \frac{4\pi}{2l+1} \sum_{m=-l}^l Y_{lm}^*(\theta, \phi) \overbrace{\int_0^{2\pi} d\phi' e^{im\phi'}}^{2\pi\delta_{m0}} \int_0^\pi \sin \theta' \cos \theta' d\theta' \sqrt{\frac{2l+1}{4\pi}} P_l^m(\cos \theta') \\ &= \sqrt{\frac{4\pi}{2l+1}} Y_{l0}^*(\theta, \phi) 2\pi \underbrace{\int_0^\pi \sin \theta' \cos \theta' d\theta' P_l(\cos \theta')}_{\int_{-1}^1 x P_l(x) dx = \delta_{l1} \cdot 2/3} \\ &= \frac{4\pi}{3} \cos \theta \cdot \delta_{l1} \end{aligned} \quad (4)$$

Plugging (4) back into (3) yields

$$\Phi(\mathbf{x}) = \frac{E_0 b^2}{4\pi} \cdot g_1(r, b) \cdot \frac{4\pi}{3} \cos \theta = \frac{E_0 r \cos \theta}{1 - p^3} \left(1 + \frac{a^3}{2r^3} \right) \quad (5)$$

Simple calculation of spherical coordinate derivatives yields

$$E_r(r, \theta) = -\frac{\partial \Phi}{\partial r} = -\frac{E_0 \cos \theta}{1 - p^3} \left(1 - \frac{a^3}{r^3} \right) \quad E_\theta(r, \theta) = -\frac{1}{r} \frac{\partial \Phi}{\partial \theta} = \frac{E_0 \sin \theta}{1 - p^3} \left(1 + \frac{a^3}{2r^3} \right) \quad (6)$$

2. Omitted.