

1. From Prob 3.23, we have the following alternative forms of the Green function

$$G(\mathbf{x}, \mathbf{x}') = \frac{4}{a} \sum_{m=-\infty}^{\infty} e^{im(\phi-\phi')} \sum_{n=1}^{\infty} \frac{J_m\left(\frac{x_{mn}\rho}{a}\right) J_m\left(\frac{x_{mn}\rho'}{a}\right)}{J_{m+1}^2(x_{mn}) x_{mn} \sinh\left(\frac{x_{mn}L}{a}\right)} \sinh\left(\frac{x_{mn}z_{<}}{a}\right) \sinh\left[\frac{x_{mn}(L-z_{>})}{a}\right] \quad (1)$$

$$G(\mathbf{x}, \mathbf{x}') = \frac{4}{L} \sum_{m=-\infty}^{\infty} e^{im(\phi-\phi')} \sum_{n=1}^{\infty} \sin\left(\frac{n\pi z}{L}\right) \sin\left(\frac{n\pi z'}{L}\right) \frac{I_m\left(\frac{n\pi\rho_{<}}{L}\right)}{I_m\left(\frac{n\pi a}{L}\right)} \\ \times \left[I_m\left(\frac{n\pi a}{L}\right) K_m\left(\frac{n\pi\rho_{>}}{L}\right) - I_m\left(\frac{n\pi\rho_{>}}{L}\right) K_m\left(\frac{n\pi a}{L}\right) \right] \quad (2)$$

$$G(\mathbf{x}, \mathbf{x}') = \frac{8}{La^2} \sum_{m=-\infty}^{\infty} e^{im(\phi-\phi')} \sum_{k=1}^{\infty} \sin\left(\frac{k\pi z}{L}\right) \sin\left(\frac{k\pi z'}{L}\right) \sum_{n=1}^{\infty} \frac{J_m\left(\frac{x_{mn}\rho}{a}\right) J_m\left(\frac{x_{mn}\rho'}{a}\right)}{\left[\left(\frac{x_{mn}}{a}\right)^2 + \left(\frac{k\pi}{L}\right)^2\right] J_{m+1}^2(x_{mn})} \quad (3)$$

Given the setup of this problem, any interior point's potential is given by

$$\Phi(\rho, \phi, z) = -\frac{1}{4\pi} \oint_S \Phi(\mathbf{x}') \frac{\partial G}{\partial n'} da' \\ = -\frac{1}{4\pi} \int_0^{2\pi} d\phi' \int_0^b V \rho' d\rho' \left(\frac{\partial G}{\partial z'} \right) \Big|_{z'=L} \quad (4)$$

For all 3 alternative forms, the integration in $d\phi'$ will eliminate all m s but $m = 0$. Carrying out the remaining steps gives us the following

(a) For form (1):

$$\Phi(\rho, \phi, z) = -\frac{1}{4\pi} \cdot \frac{4}{a} \cdot 2\pi \int_0^b V \rho' d\rho' \sum_{n=1}^{\infty} \frac{J_0\left(\frac{x_{0n}\rho}{a}\right) J_0\left(\frac{x_{0n}\rho'}{a}\right)}{J_1^2(x_{0n}) x_{0n} \sinh\left(\frac{x_{0n}L}{a}\right)} \sinh\left(\frac{x_{0n}z}{a}\right) \left(-\frac{x_{0n}}{a}\right) \cosh 0 \\ = \frac{2V}{a^2} \sum_{n=1}^{\infty} \frac{J_0\left(\frac{x_{0n}\rho}{a}\right)}{J_1^2(x_{0n})} \frac{\sinh\left(\frac{x_{0n}z}{a}\right)}{\sinh\left(\frac{x_{0n}L}{a}\right)} \int_0^b J_0\left(\frac{x_{0n}\rho'}{a}\right) \rho' d\rho' \quad (\text{use } \int J_0(x) x dx = x J_1(x)) \\ = \frac{2bV}{a} \sum_{n=1}^{\infty} \frac{J_0\left(\frac{x_{0n}\rho}{a}\right) J_1\left(\frac{x_{0n}b}{a}\right)}{J_1^2(x_{0n}) x_{0n}} \frac{\sinh\left(\frac{x_{0n}z}{a}\right)}{\sinh\left(\frac{x_{0n}L}{a}\right)} \quad (5)$$

(b) For form (2), denote

$$g(\rho, \rho') \equiv \frac{I_0\left(\frac{n\pi\rho_{<}}{L}\right)}{I_0\left(\frac{n\pi a}{L}\right)} \left[I_0\left(\frac{n\pi a}{L}\right) K_0\left(\frac{n\pi\rho_{>}}{L}\right) - I_0\left(\frac{n\pi\rho_{>}}{L}\right) K_0\left(\frac{n\pi a}{L}\right) \right] \quad (6)$$

then,

$$\Phi(\rho, \phi, z) = -\frac{1}{4\pi} \cdot \frac{4}{L} \cdot 2\pi \int_0^b V \rho' d\rho' \sum_{n=1}^{\infty} \sin\left(\frac{n\pi z}{L}\right) \left(\frac{n\pi}{L}\right) \cos(n\pi) g(\rho, \rho') \\ = -\frac{2\pi V}{L^2} \sum_{n=1}^{\infty} (-1)^n n \sin\left(\frac{n\pi z}{L}\right) \underbrace{\int_0^b g(\rho, \rho') \rho' d\rho'}_I \quad (7)$$

The evaluation of integral I depends on whether $\rho \leq b$ or $\rho > b$. If $\rho \leq b$, the integral must be done in two segments $\int_0^\rho + \int_\rho^b$. We will not derive the general form here but will calculate in the next part with specific values of ρ, z .

(c) For form (3):

$$\begin{aligned}
\Phi(\rho, \phi, z) &= -\frac{1}{4\pi} \cdot \frac{8}{La^2} \cdot 2\pi \int_0^b V \rho' d\rho' \sum_{k=1}^{\infty} \sin\left(\frac{k\pi z}{L}\right) \left(\frac{k\pi}{L}\right) \cos(k\pi) \sum_{n=1}^{\infty} \frac{J_0\left(\frac{x_{0n}\rho}{a}\right) J_0\left(\frac{x_{0n}\rho'}{a}\right)}{\left[\left(\frac{x_{0n}}{a}\right)^2 + \left(\frac{k\pi}{L}\right)^2\right] J_1^2(x_{0n})} \\
&= -\frac{4\pi V}{L^2 a^2} \sum_{k=1}^{\infty} (-1)^k k \sin\left(\frac{k\pi z}{L}\right) \sum_{n=1}^{\infty} \frac{J_0\left(\frac{x_{0n}\rho}{a}\right)}{\left[\left(\frac{x_{0n}}{a}\right)^2 + \left(\frac{k\pi}{L}\right)^2\right] J_1^2(x_{0n})} \int_0^b J_0\left(\frac{x_{0n}\rho'}{a}\right) \rho' d\rho' \\
&= -\frac{4b\pi V}{L^2 a} \sum_{n=1}^{\infty} \frac{J_0\left(\frac{x_{0n}\rho}{a}\right) J_1\left(\frac{x_{0n}b}{a}\right)}{J_1^2(x_{0n}) x_{0n}} \sum_{k=1}^{\infty} \frac{(-1)^k k \sin\left(\frac{k\pi z}{L}\right)}{\left(\frac{x_{0n}}{a}\right)^2 + \left(\frac{k\pi}{L}\right)^2}
\end{aligned} \tag{8}$$

2. For the given parameters $\rho = 0, z = L/2, b = L/4 = a/2$, we can simplify the potential expression as the following

(a) For form (5):

$$\Phi(\rho = 0, \phi, z = a) = V \sum_{n=1}^{\infty} \frac{J_1\left(\frac{x_{0n}}{2}\right)}{J_1^2(x_{0n}) x_{0n}} \frac{\sinh(x_{0n})}{\sinh(2x_{0n})} \tag{9}$$

(b) For form (7), first note

$$g(\rho = 0, \rho') = K_0\left(\frac{n\pi\rho'}{L}\right) - \frac{K_0\left(\frac{n\pi}{2}\right)}{I_0\left(\frac{n\pi}{2}\right)} I_0\left(\frac{n\pi\rho'}{L}\right) \tag{10}$$

which gives the integral

$$I = \int_0^b K_0\left(\frac{n\pi\rho'}{L}\right) \rho' d\rho' - \frac{K_0\left(\frac{n\pi}{2}\right)}{I_0\left(\frac{n\pi}{2}\right)} \int_0^b I_0\left(\frac{n\pi\rho'}{L}\right) \rho' d\rho' \tag{11}$$

We will make use of equation 10.43.1 on nist.gov

$$\int z^{v+1} Z_v(z) dz = z^{v+1} Z_{v+1}(z) \quad \text{for } Z_v(z) = I_v(z) \text{ or } e^{i\nu\pi} K_\nu(z) \tag{12}$$

which gives

$$I = T_1(\rho) \Big|_{\rho=0}^{\rho=b} - T_2(\rho) \Big|_{\rho=0}^{\rho=b} \quad \text{where}$$

$$T_1(\rho) = \left(\frac{L}{n\pi}\right)^2 \left(\frac{n\pi\rho}{L}\right) \left[-K_1\left(\frac{n\pi\rho}{L}\right)\right] \tag{13}$$

$$T_2(\rho) = \frac{K_0\left(\frac{n\pi}{2}\right)}{I_0\left(\frac{n\pi}{2}\right)} \left(\frac{L}{n\pi}\right)^2 \left(\frac{n\pi\rho}{L}\right) I_1\left(\frac{n\pi\rho}{L}\right) \tag{14}$$

While we can easily see $T_2(0) = 0$, an easy *mistake* is to declare $T_1(0) = 0$ due to its ρ factor. Because $K_1(x)$ diverges at 0, we must evaluate $T_1(0)$ with a limiting procedure. Denote $x = n\pi\rho/L$, by the asymptotic form (3.103), we have

$$T_1(0) = \left(\frac{L}{n\pi}\right)^2 [-xK_1(x)] \longrightarrow -\left(\frac{L}{n\pi}\right)^2 \quad \text{as } x \rightarrow 0 \tag{15}$$

Now we can evaluate I correctly by

$$\begin{aligned}
I &= T_1(b) - T_1(0) - T_2(b) \\
&= \left(\frac{L}{n\pi}\right)^2 - \frac{L^2}{4n\pi} \left[K_1\left(\frac{n\pi}{4}\right) + \frac{K_0\left(\frac{n\pi}{2}\right)}{I_0\left(\frac{n\pi}{2}\right)} I_1\left(\frac{n\pi}{4}\right) \right]
\end{aligned} \tag{16}$$

Plugging this into (7), we end up with

$$\begin{aligned}
\Phi(\rho = 0, \phi, z = L/2) &= -\frac{2\pi V}{L^2} \sum_{n=1}^{\infty} (-1)^n n \sin\left(\frac{n\pi}{2}\right) \cdot I \quad (\text{only odd } ns \text{ remain}) \\
&= \frac{2\pi V}{L^2} \sum_{n \text{ odd}} n \sin\left(\frac{n\pi}{2}\right) \left\{ \left(\frac{L}{n\pi}\right)^2 - \frac{L^2}{4n\pi} \left[K_1\left(\frac{n\pi}{4}\right) + \frac{K_0\left(\frac{n\pi}{2}\right)}{I_0\left(\frac{n\pi}{2}\right)} I_1\left(\frac{n\pi}{4}\right) \right] \right\} \\
&= \frac{2V}{\pi} \sum_{n \text{ odd}} \frac{1}{n} \sin\left(\frac{n\pi}{2}\right) + \frac{V}{2} \sum_{n \text{ odd}} \sin\left(\frac{n\pi}{2}\right) \left[K_1\left(\frac{n\pi}{4}\right) + \frac{K_0\left(\frac{n\pi}{2}\right)}{I_0\left(\frac{n\pi}{2}\right)} I_1\left(\frac{n\pi}{4}\right) \right] \\
&= \frac{V}{2} + \frac{V}{2} \sum_{n=0}^{\infty} (-1)^n \left\{ K_1\left[\frac{(2n+1)\pi}{4}\right] + \frac{K_0\left[\frac{(2n+1)\pi}{2}\right]}{I_0\left[\frac{(2n+1)\pi}{2}\right]} I_1\left[\frac{(2n+1)\pi}{4}\right] \right\} \quad (17)
\end{aligned}$$

(c) For form (8):

$$\begin{aligned}
\Phi(\rho = 0, \phi, z = a) &= -\frac{2\pi V}{L^2} \sum_{n=1}^{\infty} \frac{J_1\left(\frac{x_{0n}}{2}\right)}{J_1(x_{0n}) x_{0n}} \sum_{k=1}^{\infty} \frac{(-1)^k k \sin\left(\frac{k\pi}{2}\right)}{\left(\frac{2x_{0n}}{L}\right)^2 + \left(\frac{k\pi}{L}\right)^2} \\
&= V \sum_{n=1}^{\infty} \frac{J_1\left(\frac{x_{0n}}{2}\right)}{J_1^2(x_{0n}) x_{0n}} \sum_{k=1}^{\infty} \frac{(-2\pi)(-1)^k k \sin\left(\frac{k\pi}{2}\right)}{(2x_{0n})^2 + (k\pi)^2} \quad (18)
\end{aligned}$$

Again, the relation between (18) and (9) is given by the Fourier transform

$$\frac{\sinh(x_{0n})}{\sinh(2x_{0n})} = \sum_{k=1}^{\infty} \frac{(-2\pi)(-1)^k k \sin\left(\frac{k\pi}{2}\right)}{(2x_{0n})^2 + (k\pi)^2} \quad (19)$$

which is provable from equation (33) of [my notes for Prob 3.23](#), by taking the derivative with respect to z' at $z' = L$.

Apparently form (18) will converge much slower than (9) or (17) since it involves the extra infinite sum in k .

A simple [python script](#) is used to calculate (9) and (17) for 10 iterations. The result is shown below.

Iteration	$\Phi(\rho, \phi, z)$ by (9)	$\Phi(\rho, \phi, z)$ by (17)
0	0.0689326815	0.0328306826
1	0.0715821867	0.0773298612
2	0.0715328168	0.0705415934
3	0.0715293025	0.0717072322
4	0.0715293676	0.0714963211
5	0.0715293730	0.0715356408
6	0.0715293729	0.0715281674
7	0.0715293729	0.0715296072
8	0.0715293729	0.0715293270
9	0.0715293729	0.0715293819