



### 1. General setup

The relevant vectors are shown in the diagram above, where the initial polarization  $\epsilon_0$  is either  $\epsilon_0^{(E)} = \hat{x}$  for  $E$  plane or  $\epsilon_0^{(H)} = \hat{y}$  for  $H$  plane. The outgoing polarization  $\epsilon$  is spanned by

$$\epsilon_{\parallel} = \cos \theta \hat{x} - \sin \theta \hat{z} \quad \epsilon_{\perp} = \hat{y} \quad (1)$$

The total differential scattering cross section is obtained by

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma_{\parallel}}{d\Omega} + \frac{d\sigma_{\perp}}{d\Omega} = \frac{r^2 \frac{1}{2Z_0} \left( |\epsilon_{\parallel}^* \cdot \mathbf{E}_{sc}|^2 + |\epsilon_{\perp}^* \cdot \mathbf{E}_{sc}|^2 \right)}{\frac{1}{2Z_0} |\epsilon_0^* \cdot \mathbf{E}_{inc}|^2} = \frac{|\epsilon_{\parallel}^* \cdot \mathbf{F}_{sc}|^2 + |\epsilon_{\perp}^* \cdot \mathbf{F}_{sc}|^2}{|\epsilon_0^* \cdot \mathbf{E}_{inc}|^2} \quad (2)$$

For plane incident wave  $\mathbf{E}_{inc} = \epsilon_0 E_0 e^{ik_0 \cdot \mathbf{x}}$ , the above becomes

$$\frac{d\sigma}{d\Omega} = \frac{1}{E_0^2} \left( |\epsilon_{\parallel}^* \cdot \mathbf{F}_{sc}|^2 + |\epsilon_{\perp}^* \cdot \mathbf{F}_{sc}|^2 \right) \quad (3)$$

With  $ka \gg 1$ , we can use (10.127) and (10.132) to find  $\epsilon^* \cdot \mathbf{F}_{sc}$  (where  $\epsilon$  is either  $\epsilon_{\parallel}$  or  $\epsilon_{\perp}$ )

$$\begin{aligned} \epsilon^* \cdot \mathbf{F}_{sc} &= \epsilon^* \cdot \mathbf{F}_{sh} + \epsilon^* \cdot \mathbf{F}_{ill} \\ &= ika^2 E_0 (\epsilon^* \cdot \epsilon_0) \frac{J_1(ka \sin \theta)}{ka \sin \theta} + E_0 \frac{a}{2} e^{-2ika \sin(\theta/2)} (\epsilon^* \cdot \epsilon_r) \end{aligned} \quad (4)$$

where

$$\epsilon_r = -\epsilon_0 + 2(\mathbf{n}_r \cdot \epsilon_0) \mathbf{n}_r \quad (5)$$

$$\mathbf{n}_r = \frac{\mathbf{k} - \mathbf{k}_0}{|\mathbf{k} - \mathbf{k}_0|} = \frac{\sin \theta \hat{x} - (1 - \cos \theta) \hat{z}}{2 \sin(\theta/2)} = \cos \frac{\theta}{2} \hat{x} - \sin \frac{\theta}{2} \hat{z} \quad (6)$$

### 2. $E$ plane

For this part,  $\epsilon_0 = \epsilon_0^{(E)} = \hat{x}$ , so

$$\epsilon_r = -\hat{x} + 2 \cos \frac{\theta}{2} \left( \cos \frac{\theta}{2} \hat{x} - \sin \frac{\theta}{2} \hat{z} \right) = \cos \theta \hat{x} - \sin \theta \hat{z} \quad (7)$$

we see from (4) that  $\epsilon_{\perp}^* \cdot \mathbf{F}_{sc} = 0$ , so the total scattering cross section is

$$\begin{aligned} \frac{d\sigma}{d\Omega} &= \frac{d\sigma_{\parallel}}{d\Omega} = \left| ika^2 \cos \theta \frac{J_1(ka \sin \theta)}{ka \sin \theta} + \frac{a}{2} e^{-2ika \sin(\theta/2)} \right|^2 \\ &= \left| \frac{a}{2} \cos \left( 2ka \sin \frac{\theta}{2} \right) + i \left[ a \cot \theta J_1(ka \sin \theta) - \frac{a}{2} \sin \left( 2ka \sin \frac{\theta}{2} \right) \right] \right|^2 \\ &= \frac{a^2}{4} \left[ 4 \cot^2 \theta J_1^2(ka \sin \theta) + 1 - 4 \cot \theta J_1(ka \sin \theta) \sin \left( 2ka \sin \frac{\theta}{2} \right) \right] \end{aligned} \quad (8)$$

### 3. $H$ plane

In this case  $\epsilon_0 = \epsilon_0^{(H)} = \hat{\mathbf{y}}$ , giving  $\epsilon_r = -\hat{\mathbf{y}}$ . Thus  $\epsilon_{\parallel}$  has no contribution in the total differential cross section (4), and

$$\begin{aligned}
 \frac{d\sigma}{d\Omega} &= \frac{d\sigma_{\perp}}{d\Omega} = \left| ika^2 \frac{J_1(ka \sin \theta)}{ka \sin \theta} - \frac{a}{2} e^{-2ika \sin(\theta/2)} \right|^2 \\
 &= \left| -\frac{a}{2} \cos\left(2ka \sin \frac{\theta}{2}\right) + i \left[ a \csc \theta J_1(ka \sin \theta) + \frac{a}{2} \sin\left(2ka \sin \frac{\theta}{2}\right) \right] \right|^2 \\
 &= \frac{a^2}{4} \left[ 4 \csc^2 \theta J_1^2(ka \sin \theta) + 1 + 4 \csc \theta J_1(ka \sin \theta) \sin\left(2ka \sin \frac{\theta}{2}\right) \right] \quad (9)
 \end{aligned}$$