

1. Since $a \gg b$, we approximate the current distribution as if it is concentrated at the center of the cross section. Then by (5.37)

$$A_z = \frac{\mu_0}{4\pi} \frac{4Ia}{\sqrt{a^2 + r^2 + 2ar \sin \theta}} \left[\frac{(2 - k^2)K(k) - 2E(k)}{k^2} \right] \quad (1)$$

where

$$r^2 = (a + \rho \cos \phi)^2 + \rho^2 \sin^2 \phi \quad r \sin \theta = a + \rho \cos \phi \quad (2)$$

$$k^2 = \frac{4ar \sin \theta}{a^2 + r^2 + 2ar \sin \theta} = \frac{4a(a + \rho \cos \phi)}{4a^2 + \rho^2 + 4a\rho \cos \phi} \quad (3)$$

Note we have changed A 's subscript from ϕ to z which represents the direction perpendicular to the paper (as opposed to the direction of increasing ϕ in the problem's figure).

Thus

$$1 - k^2 = \frac{\rho^2}{4a^2 + \rho^2 + 4a\rho \cos \phi} \rightarrow \frac{\rho^2}{4a^2} \quad \text{as } \frac{\rho}{a} \rightarrow 0 \quad (4)$$

When $k \rightarrow 1$, the elliptic integrals can be approximated by (see [wolfram](https://www.wolfram.com/mathworld/elliptic-integral-k/), note on that website, $z = k^2$, see their definition page)

$$K(k) \rightarrow -\frac{1}{2} \ln(1 - k^2) + \ln 4 = \ln\left(\frac{8a}{\rho}\right) \quad (5)$$

$$E(k) \rightarrow 1 \quad (6)$$

Now (1) can be approximated as

$$A_z \rightarrow \frac{\mu_0}{4\pi} \frac{4Ia}{2a} \left[\ln\left(\frac{8a}{\rho}\right) - 2 \right] = \frac{\mu_0 I}{2\pi} \left[\ln\left(\frac{8a}{\rho}\right) - 2 \right] \quad \text{as } \frac{\rho}{a} \rightarrow 0 \quad (7)$$

2. Using the "straight wire" approximation, for $\rho < b$, the magnetic induction is given by

$$-\frac{\partial A_z}{\partial \rho} = B_\phi(\rho) = \frac{\mu_0 I \rho^2}{2\pi b^2 \rho} = \frac{\mu_0 I \rho}{2\pi b^2} \quad (8)$$

thus

$$A_z = -\frac{\mu_0 I \rho^2}{4\pi b^2} + C \quad (9)$$

C is determined by matching (7) and (9) at $\rho = b$:

$$C = \frac{\mu_0 I}{2\pi} \left[\ln\left(\frac{8a}{b}\right) - 2 \right] + \frac{\mu_0 I}{4\pi} \quad \Rightarrow \quad A_z = \frac{\mu_0 I}{4\pi} \left(1 - \frac{\rho^2}{b^2} \right) + \frac{\mu_0 I}{2\pi} \left[\ln\left(\frac{8a}{b}\right) - 2 \right] \quad (10)$$

(It doesn't seem necessary to impose the continuity of radial derivative at $\rho = b$ per the hint.)

3. From equation (5.149)

$$\begin{aligned} W &= \frac{1}{2} \int \mathbf{A} \cdot \mathbf{J} d^3x \approx \frac{1}{2} \cdot (2\pi a) \int_0^b \rho d\rho \int_0^{2\pi} d\phi \left\{ \frac{\mu_0 I}{4\pi} \left(1 - \frac{\rho^2}{b^2} \right) + \frac{\mu_0 I}{2\pi} \left[\ln\left(\frac{8a}{b}\right) - 2 \right] \right\} \cdot \frac{I}{\pi b^2} \\ &= \frac{\mu_0 a I^2}{2b^2} \int_0^b \left(1 - \frac{\rho^2}{b^2} \right) \rho d\rho + \frac{\mu_0 a I^2}{2} \left[\ln\left(\frac{8a}{b}\right) - 2 \right] \\ &= \frac{\mu_0 a I^2}{2} \left[\ln\left(\frac{8a}{b}\right) - \frac{7}{4} \right] \end{aligned} \quad (11)$$

which gives

$$L = \frac{2W}{I^2} = \mu_0 a \left[\ln\left(\frac{8a}{b}\right) - \frac{7}{4} \right] \quad (12)$$

If the current only flows on the surface, we will use (7) for the energy integration only, and

$$W = \frac{\mu_0 a I^2}{2} \left[\ln\left(\frac{8a}{b}\right) - 2 \right] \quad \Rightarrow \quad L = \mu_0 a \left[\ln\left(\frac{8a}{b}\right) - 2 \right] \quad (13)$$