

1. Consider the usual separation of variables in cylindrical coordinates, for potential at $z \rightarrow \infty$ to vanish, the z -dependency is going to be proportional to $e^{-k|z|}$ for k positive real.

Given the symmetry in ϕ , the ϕ -dependency $A_m \cos m\phi + B_m \sin m\phi$ must be reduced to a constant, i.e., $m = 0$.

For radial part, since $\rho = 0$ is in the range, we must only have J_m in the radial solution. In summary, the potential for $z > 0$ will have to have the form

$$\Phi(\rho, \phi, z) = \int_0^\infty \tilde{A}(k) J_0(k\rho) e^{-kz} dk \quad \text{where by (3.110)} \quad (1)$$

$$\tilde{A}(k) = k \int_0^\infty \Phi(\rho, \phi, 0) J_0(k\rho') \rho' d\rho' = kV \int_0^a J_0(k\rho') \rho' d\rho' \quad (2)$$

Plugging (2) into (1), and calculating for the axis point $\rho = 0$:

$$\begin{aligned} \Phi(0, \phi, z) &= \int_0^\infty \left[kV \int_0^a J_0(k\rho') \rho' d\rho' \right] J_0(0) e^{-kz} dk \\ &= V \int_0^a \rho' d\rho' \underbrace{\int_0^\infty J_0(k\rho') k e^{-kz} dk}_I \end{aligned} \quad (3)$$

where

$$\begin{aligned} I &= -\frac{d}{dz} \int_0^\infty J_0(k\rho') e^{-kz} dk \\ &= -\frac{d}{dz} \mathcal{L}\{J_0(k\rho')\}(z) \\ &= -\frac{d}{dz} \frac{1}{\sqrt{\rho'^2 + z^2}} = \frac{z}{\sqrt{\rho'^2 + z^2}^3} \end{aligned} \quad (4)$$

Thus

$$\Phi(0, \phi, z) = V \int_0^a \frac{z\rho' d\rho'}{\sqrt{\rho'^2 + z^2}^3} = V \left(1 - \frac{z}{\sqrt{a^2 + z^2}} \right) \quad (5)$$

2. Now we are looking for $\Phi(a, \phi, z)$. From (2)

$$\begin{aligned} \tilde{A}(k) &= kV \int_0^a J_0(k\rho') \rho' d\rho' \\ &= kV \int_0^a \sum_{j=0}^\infty \frac{(-1)^j}{j!j!} \left(\frac{k\rho'}{a} \right)^{2j} \rho' d\rho' \\ &= kV \sum_{j=0}^\infty \frac{(-1)^j}{j!j!} \left(\frac{k}{a} \right)^{2j} \int_0^a \rho'^{2j+1} d\rho' \\ &= kV \sum_{j=0}^\infty \frac{(-1)^j}{j!(j+1)!} \frac{k^{2j} a^{2j+2}}{2^{2j+1}} \\ &= kV \cdot \frac{a}{k} J_1(ka) = Va J_1(ka) \end{aligned} \quad (6)$$

This gives

$$\Phi(a, \phi, z) = Va \int_0^\infty J_1(ka) J_0(ka) e^{-kz} dk \quad (7)$$

The integral is non-trivial and can be looked up from the integral table (reference [Eduardo Kausel and Mirza M. Irfan Baig](#), in particular ENS-4.7), the result is

$$\int_0^\infty J_1(ka) J_0(ka) e^{-kz} dk = I_{10}^0 = -\frac{1}{\pi a} \frac{kz}{2\sqrt{a \cdot a}} K(k) - \frac{\text{sgn}(a-a)}{\pi a} \Lambda + \frac{1}{a} H(a-a) \quad (8)$$

where

$$k = \frac{2\sqrt{a \cdot a}}{\sqrt{(a+a)^2 + z^2}} = \frac{2a}{\sqrt{z^2 + 4a^2}} \quad (9)$$

and H is Heaviside step function which takes value of 1/2 at 0, and K the complete elliptic integral of the first kind. Λ is irrelevant because of the zero factor in front of it.

Plugging these all into (7) gives

$$\Phi(a, \phi, z) = \frac{V}{2} \left[1 - \frac{kz}{\pi a} K(k) \right] \quad (10)$$