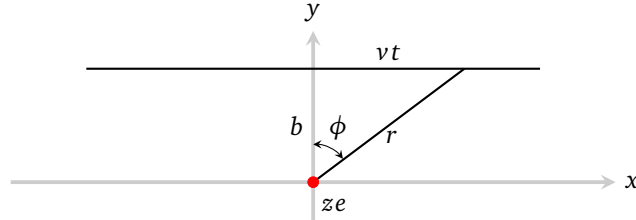


1. In the rest frame of the heavy particle, the electron is coming in with impact parameter  $b$  and velocity  $v$ . The change in momentum is given by the impulse

$$\Delta \mathbf{p} = \int_{-\infty}^{\infty} -\nabla V(r) dt = \int_{-\infty}^{\infty} -\nabla \left[ \frac{ze^2 e^{-k_D r}}{r} \right] dt = ze^2 \int_{-\infty}^{\infty} \frac{e^{-k_D r}}{r^2} (1 + k_D r) \hat{\mathbf{r}} dt \quad (1)$$

Transforming back to the lab frame, we can obtain the energy loss

$$\Delta E = \frac{|\Delta \mathbf{p}|^2}{2m} \quad (2)$$



As shown in the diagram above, the longitudinal component ( $x$ ) of the integral in (1) vanishes, giving

$$\Delta E = \frac{z^2 e^4}{2m} I^2 \quad \text{where} \quad I = \int_{-\infty}^{\infty} \frac{e^{-k_D r}}{r^2} (1 + k_D r) \cos \phi dt \quad (3)$$

For a fast velocity, the trajectory is approximated as a straight line (small deflection angle), so we have

$$r = \sqrt{b^2 + v^2 t^2} = b \sqrt{1 + u^2} \quad \text{where} \quad u \equiv \frac{vt}{b} \quad \cos \phi = \frac{b}{r} \quad (4)$$

The integral  $I$  becomes

$$\begin{aligned} I &= \int_{-\infty}^{\infty} \frac{e^{-k_D b \sqrt{1+u^2}}}{bv(1+u^2)^{3/2}} (1 + k_D b \sqrt{1+u^2}) du && \text{let } u = \sinh w \\ &= \frac{1}{bv} \int_{-\infty}^{\infty} \frac{e^{-k_D b \cosh w}}{\cosh^3 w} (1 + k_D b \cosh w) \cosh w dw \\ &= \frac{1}{bv} \int_{-\infty}^{\infty} \frac{e^{-k_D b \cosh w}}{\cosh^2 w} (1 + k_D b \cosh w) dw \end{aligned} \quad (5)$$

Let  $\kappa \equiv k_D b$ , then

$$\frac{d(bvI)}{d\kappa} = \int_{-\infty}^{\infty} e^{-\kappa \cosh w} \left[ \frac{-\cosh w (1 + \kappa \cosh w) + \cosh w}{\cosh^2 w} \right] dw = -\kappa \int_{-\infty}^{\infty} e^{-\kappa \cosh w} dw \quad (6)$$

Recall the integral representation of  $K_\nu(\kappa)$ , [DLMF 10.32.E9](#)

$$K_\nu(\kappa) = \int_0^\infty e^{-\kappa \cosh t} \cosh(\nu t) dt \quad (7)$$

and the recurrence relation [DLMF 10.29.E2](#)

$$(-1)^\nu K'_\nu(\kappa) = (-1)^{\nu-1} K_{\nu-1}(\kappa) - \frac{\nu}{\kappa} (-1)^\nu K_\nu(\kappa) \quad (8)$$

we see that

$$\frac{d(bvI)}{d\kappa} = -2\kappa K_0(\kappa) = 2[\kappa K'_1(\kappa) + K_1(\kappa)] = 2 \frac{d[\kappa K_1(\kappa)]}{d\kappa} \implies I = \frac{2}{bv} [\kappa K_1(\kappa)] + C \quad (9)$$

With [DLMF 10.25.E3](#)

$$K_\nu(\kappa) \rightarrow \sqrt{\frac{\pi}{2\kappa}} e^{-\kappa} \quad \text{as } \kappa \rightarrow \infty \quad (10)$$

matching the behavior of  $I$  as  $\kappa \rightarrow \infty$  establishes  $C = 0$ . Putting  $I$  back to  $\Delta E$  yields

$$\Delta E = \frac{2z^2 e^4}{mv^2} k_D^2 K_1^2(k_D b) \quad (11)$$

2. The energy loss per unit distance can be obtained by the weighted integral over all possible impact parameters (see Jackson (13.7))

$$\frac{dE}{dx} = 2\pi NZ \int_{b_{\min}}^{\infty} \Delta E(b) b db = \frac{4\pi NZ z^2 e^4}{mv^2} \int_{b_{\min}}^{\infty} k_D^2 K_1^2(k_D b) b db = \frac{4\pi NZ z^2 e^4}{mv^2} \int_{\kappa_{\min}}^{\infty} K_1^2(\kappa) \kappa d\kappa \quad (12)$$

The integral can be looked up from integration tables (e.g., see 5.54.2 from <sup>1</sup>)

$$\int \kappa K_1^2(\kappa) d\kappa = \underbrace{\frac{\kappa^2}{2} [K_1^2(\kappa) - K_0(\kappa) K_2(\kappa)]}_{F(\kappa)} + C \quad (13)$$

From the asymptotic form (10), we see that  $F(\infty) = 0$ . In small argument  $\kappa \rightarrow 0$ , we have (see Jackson (3.103))

$$K_1(\kappa) \rightarrow \frac{1}{\kappa} \quad K_0(\kappa) \rightarrow -\left[\ln\left(\frac{\kappa}{2}\right) + \gamma\right] \quad K_2(\kappa) \rightarrow \frac{2}{\kappa^2} \quad (14)$$

thus

$$\begin{aligned} F(\kappa) &\rightarrow \frac{\kappa^2}{2} \left[ \frac{1}{\kappa^2} + \left[ \ln\left(\frac{\kappa}{2}\right) + \gamma \right] \cdot \frac{2}{\kappa^2} \right] \\ &= \frac{1}{2} + \ln\left(\frac{\kappa}{2}\right) + \gamma \\ &= \ln\left(\frac{\kappa e^{\gamma+1/2}}{2}\right) \\ &\approx \ln(1.47\kappa) \end{aligned} \quad (15)$$

The energy loss per unit distance is then

$$\begin{aligned} \frac{dE}{dx} &\approx \frac{4\pi NZ z^2 e^4}{mv^2} \ln\left(\frac{1}{1.47 k_D b_{\min}}\right) && \text{by (7.60) in Gaussian units } \omega_p^2 \equiv \frac{4\pi NZ e^2}{m} \\ &= \frac{z^2 e^2}{v^2} \omega_p^2 \ln\left(\frac{1}{1.47 k_D b_{\min}}\right) \end{aligned} \quad (16)$$

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<sup>1</sup>I. S. Gradshteyn and I. M. Ryzhik, *Table of Integrals, Series and Products Eighth Edition*.