

1. The canonical momentum of the particle is

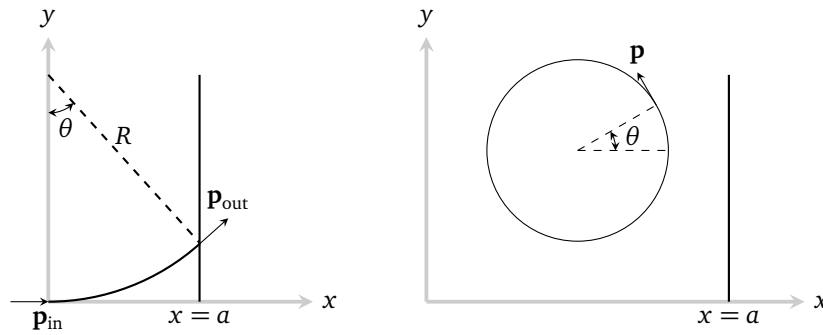
$$\mathbf{G} = \gamma m \mathbf{v} + \frac{q}{c} \mathbf{A} \quad (1)$$

Only knowing  $\mathbf{B} = -B\hat{\mathbf{z}}$  is not sufficient to determine  $\mathbf{A}$ . For example, up to a gauge transformation, we can write

$$\mathbf{A} = -xB\hat{\mathbf{y}} + \nabla\chi \quad \Rightarrow \quad \mathbf{G} = \gamma m \mathbf{v} - \frac{qxB}{c}\hat{\mathbf{y}} + \frac{q}{c}\nabla\chi \quad (2)$$

2. In this part, the radius of the particle's movement inside the magnetic field is

$$R = \frac{cP}{qB} \quad (3)$$



With the help of the diagram on the left, we see that when the particle emerges out of the magnetic field, the components of its mechanical momentum are

$$p_{\text{out},y} = p \sin \theta = \frac{pa}{R} = \frac{qBa}{c} \quad (4)$$

$$p_{\text{out},x} = p \cos \theta \quad (5)$$

With the choice of gauge  $\chi = 0$ , we see that

$$G_{\text{in},x} = p \quad G_{\text{in},y} = 0 \quad G_{\text{in},z} = 0 \quad (6)$$

$$G_{\text{out},x} = p \cos \theta \quad G_{\text{out},y} = \frac{qBa}{c} - \frac{qBa}{c} = 0 \quad G_{\text{out},z} = 0 \quad (7)$$

Clearly the  $y, z$  component of the canonical momentum are conserved, but not its  $x$  component.

This because the Euler-Lagrange equation of motion

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} = \frac{\partial L}{\partial q_i} \quad (8)$$

indicates that for the canonical momentum  $G_i = \partial L / \partial \dot{q}_i$  to be conserved, the Lagrangian must not have explicit dependence on the coordinate  $q_i$ . Our particular choice of gauge  $\chi = 0$  made it so that  $L$  has dependence on  $x$ , but not  $y, z$ , therefore  $G_y, G_z$  are conserved, but not  $G_x$ .

3. In this part, first let's keep using the zero gauge function  $\chi = 0$ , which gives

$$\begin{aligned} \mathbf{G}(\theta) &= \mathbf{p}(\theta) - \frac{qxB}{c}\hat{\mathbf{y}} \\ &= \mathbf{p}(\theta) - \frac{qRB}{c} \cos \theta \hat{\mathbf{y}} \\ &= p(-\sin \theta \hat{\mathbf{x}} + \cos \theta \hat{\mathbf{y}}) - p \cos \theta \hat{\mathbf{y}} \\ &= -p \sin \theta \hat{\mathbf{x}} \end{aligned} \quad (9)$$

of which the  $y$  component is conserved, but not  $x$ .

But if we choose gauge function to be

$$\chi = \frac{xyB}{2} \quad (10)$$

then we have

$$\mathbf{A} = -xB\hat{\mathbf{y}} + \nabla\chi = -xB\hat{\mathbf{y}} + \frac{yB}{2}\hat{\mathbf{x}} + \frac{xB}{2}\hat{\mathbf{y}} = -\frac{B}{2}(x\hat{\mathbf{y}} - y\hat{\mathbf{x}}) \quad (11)$$

Then the canonical momentum becomes

$$\begin{aligned} \mathbf{G}(\theta) &= \mathbf{p}(\theta) - \frac{qB}{2c}(x\hat{\mathbf{y}} - y\hat{\mathbf{x}}) \\ &= \mathbf{p}(\theta) - \frac{qBR}{2c}(\cos\theta\hat{\mathbf{y}} - \sin\theta\hat{\mathbf{x}}) \\ &= p(-\sin\theta\hat{\mathbf{x}} + \cos\theta\hat{\mathbf{y}}) - \frac{p}{2}(\cos\theta\hat{\mathbf{y}} - \sin\theta\hat{\mathbf{x}}) \\ &= \frac{p}{2}(-\sin\theta\hat{\mathbf{x}} + \cos\theta\hat{\mathbf{y}}) \end{aligned} \quad (12)$$

which is not a conserved quantity for either  $x$  or  $y$ . This is indicated from our choice of  $\chi$  which makes  $\mathbf{A}$ , hence  $L$ , depend on both  $x$  and  $y$ .