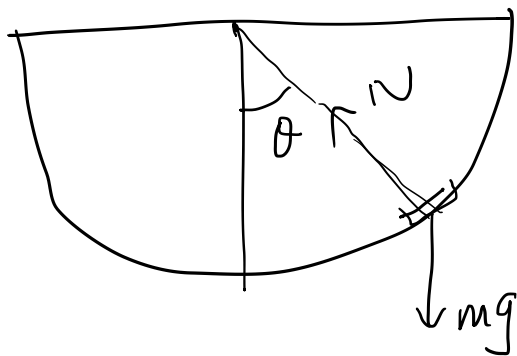


Physics I

Lecture 4

Recap

$$\vec{F} = m\vec{a}$$



$$\ddot{\theta} = -\frac{g}{R} \sin \theta$$

$$\ddot{\theta} \approx -\frac{g}{R} \theta$$

equilibrium position

$$\ddot{\theta} = 0, \theta = 0$$

$$\theta > 0, \ddot{\theta} < 0$$

$$\theta < 0, \ddot{\theta} > 0$$

$$\ddot{\theta} + \frac{g}{R} \theta = 0$$

$$\frac{g}{R} = \omega^2$$

$$\theta(t) = A \sin \omega t + B \cos \omega t$$

$$t=0, \theta = \theta_0, \dot{\theta} = 0$$

$$\boxed{\theta(t) = \theta_0 \cos \omega t}$$

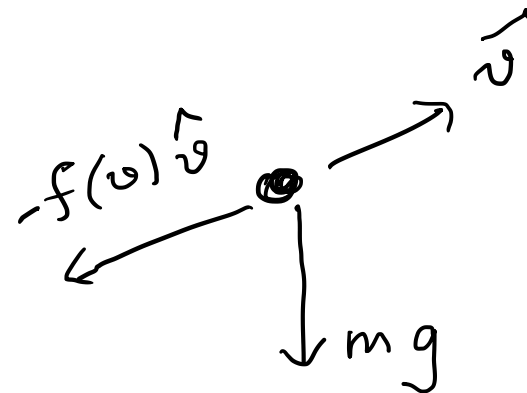
$$\vec{F} = m\vec{a}$$

$$\hookrightarrow \vec{F}(\vec{v}, \vec{r}, t)$$

Projectile motion with air resistance

Retarding forces

$$\vec{F} = \vec{F}(\vec{v}) = -f(v)\hat{v}$$



At low speeds

$$f(v) = bv + cv^2 = f_{lin} + f_{quad}$$

$f_{lin} \rightarrow$ viscous drag of medium \propto viscosity of medium
depends on the size of the particle.

Stokes Law $f = 6\pi r \eta v$

$f_{quad} \rightarrow$ projectiles need to accelerate mass of air which they are in contact with, continuously colliding
 \propto density of medium and cross sectional area.

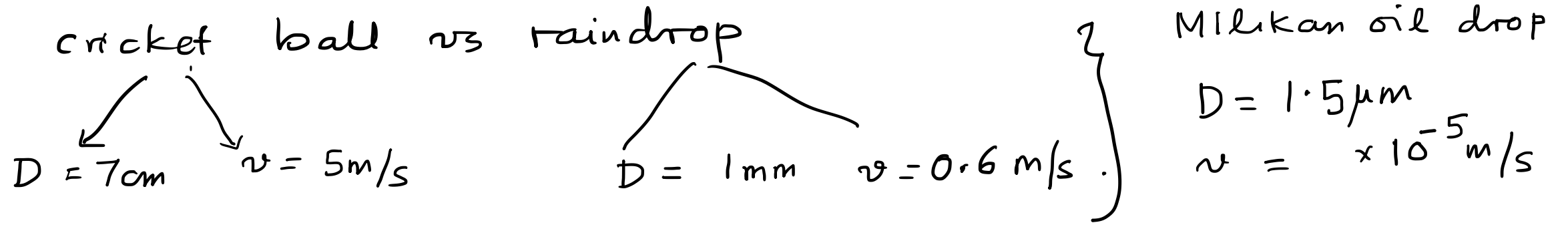
for a spherical projectile

$$b = \beta D \overset{\text{diameter}}{\quad}, \quad = \gamma D^2$$

for air $\beta = 1.6 \times 10^{-4} \frac{\text{N} \cdot \text{s}}{\text{m}^2}$

$$\gamma = 0.25 \frac{\text{N} \cdot \text{s}}{\text{m}^2}$$

$$\frac{f_{quad}}{f_{lin}} = \frac{c v^2}{b v} = \frac{\gamma D}{\beta} v = \left(1.6 \times 10^3 \frac{\text{s}}{\text{m}^2} \right) D v$$



$$\frac{f_{\text{quad}}}{f_{\text{lin}}} \approx 600 \text{ cricket ball}$$

$$\frac{f_q}{f_e} \sim 1 \text{ raindrop.}$$

$$\frac{f_q}{f_e} \approx 10^{-7} \quad \vec{f} = -b\vec{v}$$

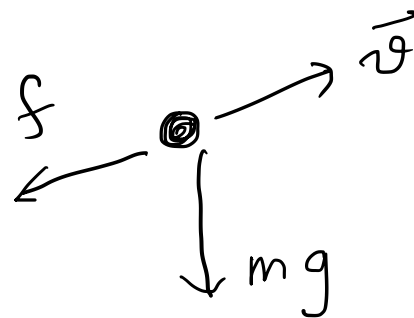
Linear air resistance

$$\vec{F} = m\vec{g} - b\vec{v} \quad (1)$$

$$m\ddot{\vec{r}} = m\vec{g} - b\vec{v} \quad (2)$$

$$m\dot{v}_x = -bv_x \quad (3)$$

$$m\dot{v}_y = mg - bv_y \quad (4)$$



for quadratic drag . $\vec{f} = -c v^2 \hat{v} = -c v \vec{v}$

$$m\dot{v}_x = -c \sqrt{v_x^2 + v_y^2} v_x$$

$$m\dot{v}_y = mg - c \sqrt{v_x^2 + v_y^2} v_y$$

Horizontal motion with linear drag

at $t=0$, $x=0$, $v_x = v_{x0}$

$$m \dot{v}_x = -b v_x$$

$$\dot{v}_x = -k v_x$$

$$\boxed{k = \frac{b}{m}}$$

$$\frac{\dot{v}_x}{v_x} = -k$$

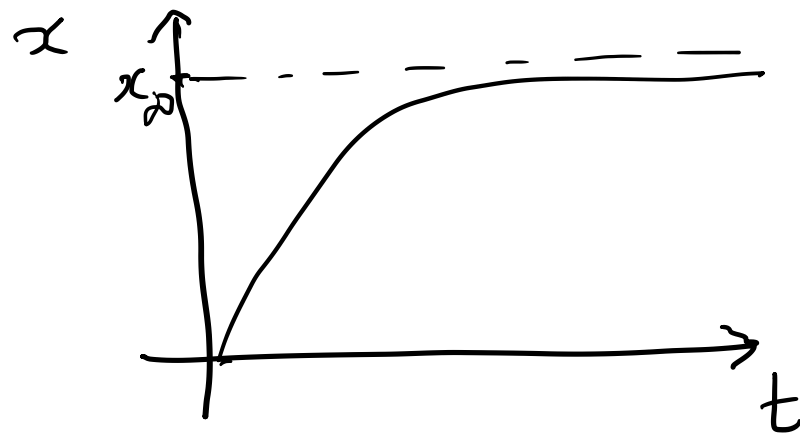
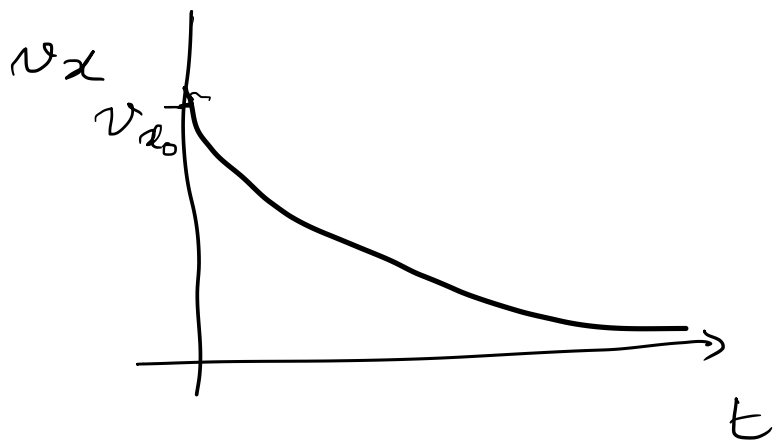
$$\frac{dv}{v} = -k dt$$

$$\rightarrow \boxed{v_x = A e^{-kt}} + \text{initial conditions}$$

$$\boxed{v_x = v_{x0} e^{-kt} = v_{x0} e^{-t/\tau}}$$

$$\begin{aligned}\tau &= \frac{1}{k} \\ &= \frac{m}{b}\end{aligned}$$

$$v_x(t \rightarrow \infty) = 0$$



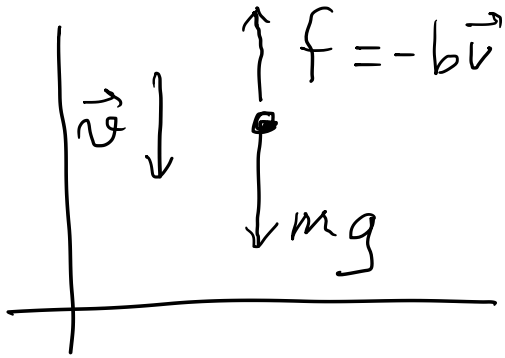
$$\frac{dx}{dt} = v_{x0} e^{-kt}$$

$$x = -\frac{v_{x0}}{k} e^{-kt} + C, \quad t=0, x=0 \Rightarrow C = \frac{v_{x0}}{k}$$

$$\boxed{x = \frac{v_{x0}}{k} (1 - e^{-t/\tau})}$$

$$\left. \begin{aligned} x_{\infty} &= \frac{v_{x0}}{k} \end{aligned} \right\} \begin{array}{l} \text{short time limit} \\ x \simeq \frac{v_{x0}}{k} (1 - 1 + \frac{t}{\tau}) \\ \simeq v_{x0} t \end{array}$$

Vertical motion with linear drag



$$m \dot{v}_y = mg - bv_y$$

$$(v_y > 0)$$

retarding force
upward.

$$\begin{aligned} \text{when } mg - bv_y &= 0 \\ v_y &= \frac{mg}{b} = v_{ter} \end{aligned}$$

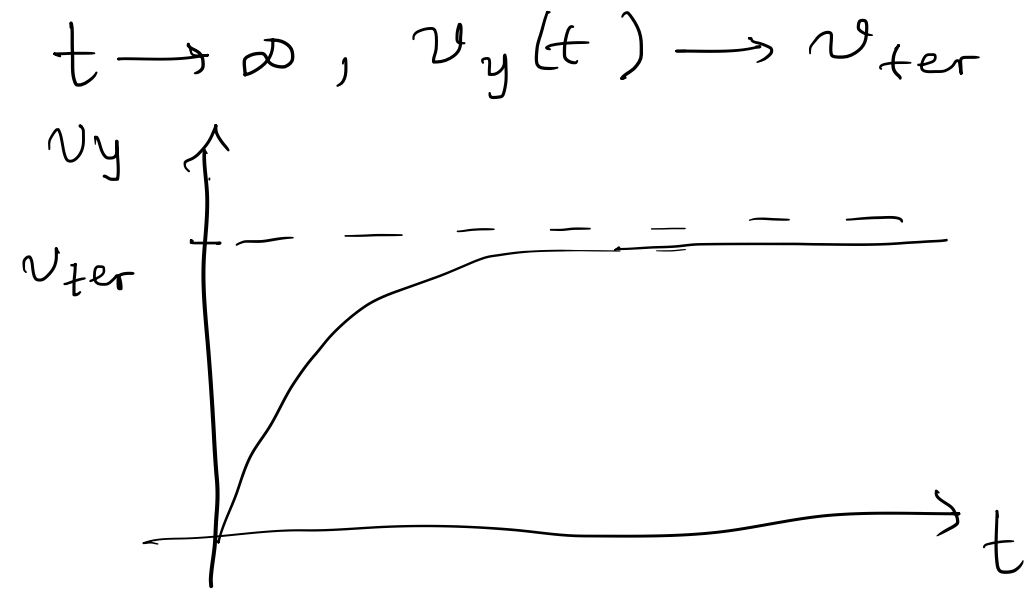
$$m \dot{v}_y = -b(v_y - v_{ter})$$

$$\frac{dv_y}{v_y - v_{ter}} = -\frac{b}{m}$$

$$\hookrightarrow v_y - v_{ter} = A e^{-b/mt} \quad ; \quad \begin{aligned} t=0, v_y &= v_{y0} \\ \boxed{A = v_{y0} - v_{ter}} \end{aligned}$$

$$v_y(t) = v_{ter} + (v_{y_0} - v_{ter})e^{-t/\tau}$$

$$= v_{y_0}e^{-t/\tau} + v_{ter}(1 - e^{-t/\tau})$$



Short time approx.

$$v_y(t) \simeq$$

$$v_{y_0}(1 - t/\tau) + v_{ter} t/\tau$$

$$\simeq v_{y_0} + (v_{ter} - v_{y_0})t/\tau$$

$$v_{y_0} = 0$$

$$v_y(t) \simeq v_{ter} t/\tau$$

Next integration

$$v_y(t) = v_{ter} (1 - e^{-t/\tau})$$

↓ $y(0) = 0$

$$y(t) = v_{ter} t + (v_{y0} - v_{ter}) \tau (1 - e^{-t/\tau})$$

$$x(t) = x_{\infty} (1 - e^{-t/\tau}) \quad x_{\infty} = v_{x0} \tau$$

→ orbit of a projectile subject to linear drag.

