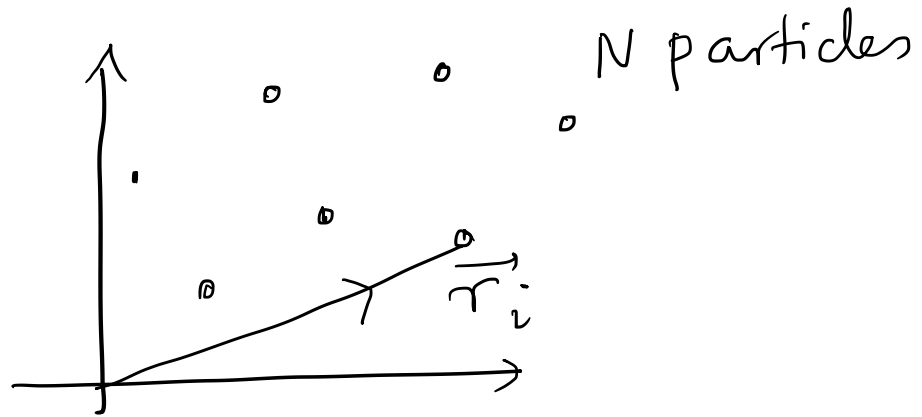


# Physics I

## Lecture 23

# Many particle system dynamics



Newton's Law  $i^{\text{th}}$  particle

$$\vec{p}_i = \sum_{\substack{j, \\ i \neq j}} \vec{F}_{ji} + \vec{F}_i^{(e)} \quad \text{--- (1)}$$

Assume Newton's 3rd. Law (weak form)  
not necessarily  
acting along line  
joining particles

$$\vec{F}_{ij} = -\vec{F}_{ji} \quad \text{--- (2)}$$

Sum (1) over all particles.

$$\sum_{i=1}^N \vec{p}_i = \sum_i \sum_{\substack{j, \\ i \neq j}} \underbrace{\vec{F}_{ji}}_{=0} + \underbrace{\sum_i \vec{F}_i^{(e)}}_{\vec{F}^{(e)}} \quad \text{--- (3)}$$

$\vec{F}^{(e)} = \text{total ext force}$

$$\vec{P} = \sum_i \vec{p}_i$$

$$\frac{d\vec{P}}{dt} = \vec{F}^{(e)} \quad \text{--- (3)}$$

$$\frac{d^2}{dt^2} \sum_i m_i \vec{r}_i = \vec{F}^{(e)} \quad \text{--- (4)}$$

→ centre of mass coordinate .

$$\text{Define } \vec{R} = \frac{\sum_i m_i \vec{r}_i}{\sum_i m_i} = \frac{\sum_i m_i \vec{r}_i}{M} \quad \text{--- (5)}$$

(4) reduces to

$$\boxed{M \frac{d^2 \vec{R}}{dt^2} = \vec{F}^{(e)}} \quad \text{--- (6)}$$

→ purely internal forces have no effect on motion of CM .

Total linear momentum

$$\vec{P} = \sum_i m_i \frac{d\vec{r}_i}{dt} = M \frac{d\vec{R}}{dt} \quad (7)$$

If  $\vec{F}^{(e)} = 0$ , total linear momentum is conserved

## Angular Momentum

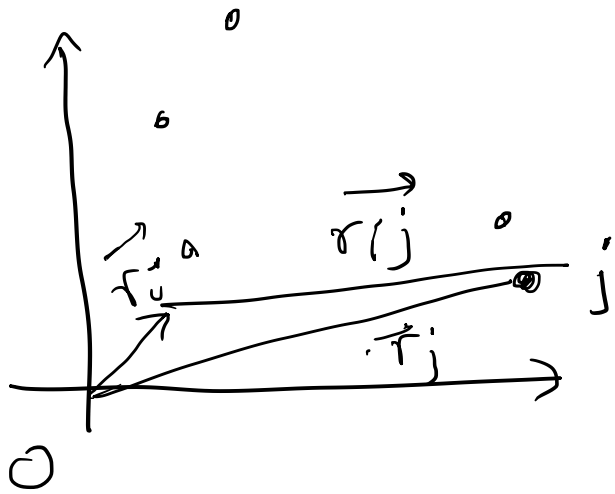
$$\vec{L}_{\text{tot}} = \sum_i (\vec{r}_i \times \vec{p}_i) = \vec{L}$$

$$\dot{\vec{L}} = \sum_i \dot{\vec{r}}_i \times \dot{\vec{p}}_i = \sum_i \dot{\vec{r}}_i \times F_i^{(e)} + \sum_i \sum_{\substack{j \\ i \neq j}} \dot{\vec{r}}_i \times \vec{F}_{ji} \quad \text{--- (8)}$$

Last term can be considered as sum of pairs of the following form

$$\dot{\vec{r}}_i \times \vec{F}_{ji} + \dot{\vec{r}}_j \times \vec{F}_{ij} = (\dot{\vec{r}}_i - \dot{\vec{r}}_j) \times \vec{F}_{ji} \quad \text{--- (9)}$$

vanishes  $\vec{r}_{ij} = \vec{r}_i - \vec{r}_j$   
in the direction  
of  $\vec{F}_{ji}$

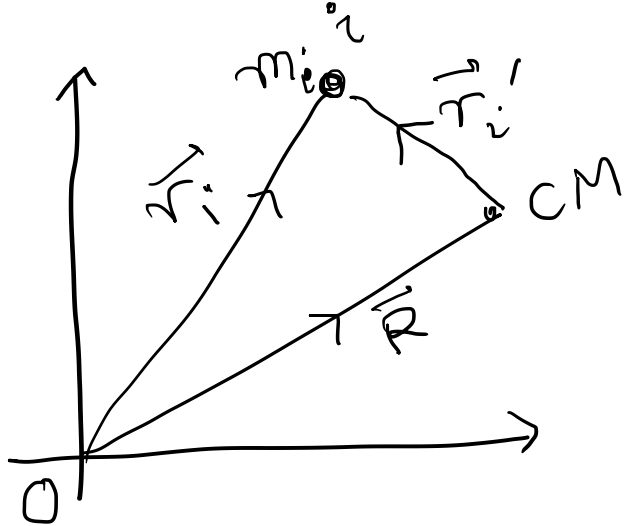


So if strong form  
of 3rd Law holds

$$(\vec{r}_i - \vec{r}_j) \times \vec{F}_{ji} = 0$$

$$\frac{d\vec{L}}{dt} = \sum_i \vec{r}_i \times \vec{F}^{(e)} = \vec{N}^{(e)} \quad \text{--- (10) external torque}$$

$$\hookrightarrow \vec{N}^{(e)} = 0 \implies \vec{L} \text{ is conserved}$$



Angular momentum about origin

$$\vec{L} = \sum_i \vec{r}_i \times \vec{p}_i$$

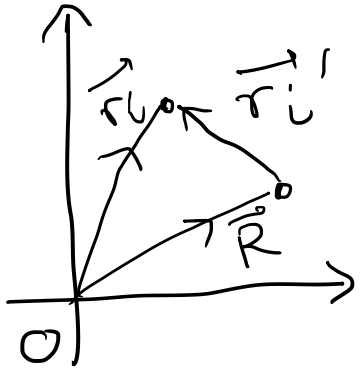
$$\left. \begin{aligned} \vec{r}_i &= \vec{r}_i' + \vec{R} \\ \vec{v}_i &= \vec{v}_i' + \vec{V} \end{aligned} \right\} \text{--- (11)}$$

$$\boxed{\begin{aligned} \vec{V} &= \dot{\vec{R}} \\ \vec{v}_i' &= \dot{\vec{r}}_i' \\ \vec{v}_i &= \dot{\vec{r}}_i \end{aligned}}$$

$$\begin{aligned} \vec{L} &= \sum_i \vec{R} \times m_i \vec{v}_i + \sum_i \vec{r}_i' \times m_i \vec{v}_i' \\ &+ \sum_i \vec{r}_i' \times m_i \vec{V} + \sum_i \vec{R} \times m_i \vec{v}_i' \end{aligned} \text{--- (12)}$$

Rewriting (12)

$$\vec{L} = \sum_i \vec{R} \times m_i \vec{V} + \sum_i \vec{r}_i' \times m_i \vec{v}_i' + \underbrace{\left( \sum_i m_i \vec{r}_i' \right)}_{=0} \times \vec{V} + \vec{R} \times \underbrace{\frac{d}{dt} \left( \sum_i m_i \vec{r}_i' \right)}_{=0}$$




$$\boxed{\vec{L} = \vec{R} \times M \vec{V} + \sum_i \vec{r}_i' \times \vec{p}_i'} \quad (13)$$

If the C.M. is at rest w.r.t  $O$ , angular momentum will be independent of point of ref.



# Energy

$$W_{12} = \sum_i \int_1^2 \vec{F}_i \cdot d\vec{s}_i$$


$$\sum_i \int_1^2 \vec{F}_i \cdot d\vec{s}_i = \sum_i \int_1^2 \vec{F}_i^{(e)} \cdot d\vec{s}_i + \sum_{\substack{j \\ i \neq j}} \int_1^2 \vec{F}_{ji} \cdot d\vec{s}_i$$

Using eqns of motion

$$\sum_i \int_1^2 \vec{F}_i \cdot d\vec{s}_i = \sum_i \int_1^2 m \dot{\vec{v}}_i \cdot \vec{v}_i dt = \sum_i \int_1^2 d\left(\frac{1}{2} m v_i^2\right)$$

$$W_{12} = T_2 - T_1, \text{ where } T = \frac{1}{2} \sum_i m_i v_i^2$$

Making use of transfr. to CM coordinates.

$$T = \frac{1}{2} \sum_i m_i (\vec{v}_i' + \vec{V}) \cdot (\vec{v}_i' + \vec{V})$$

$$= \frac{1}{2} \sum_i m_i V^2 + \frac{1}{2} \sum_i m_i v_i'^2 + \vec{V} \cdot \frac{d}{dt} \left( \underbrace{\sum m_i \vec{r}_i'}_{=0} \right)$$

$$T = \frac{1}{2} M V^2 + \frac{1}{2} \sum_i m_i v_i'^2 \quad (13)$$

↙ K.E of  
CM

↘ K.E of motion  
about the C.M.

RHS -

$$\sum_i \int_1^2 \vec{F}_i^{(e)} \cdot d\vec{s}_i$$

If ext force conservative

$$= \sum_i \int_1^2 -\vec{\nabla}_i U_i \cdot d\vec{s}_i = - \sum_i \int_1^2 dU_i = - \sum_i U_i \Big|_1^2$$

If internal forces also conservative

$\vec{F}_{ij}$  can be derived from a potential  $U_{ij}$

$U_{ij} = U_{ij}(|\vec{r}_i - \vec{r}_j|) \rightarrow$  to satisfy 3<sup>rd</sup>

Law

$$\vec{F}_{ij} = -\vec{\nabla}_i U_{ij} = +\vec{\nabla}_j U_{ij} = -\vec{F}_{ji} \quad (14)$$

$$-\vec{F}_{ij} = -\vec{\nabla}_i U_{ij}(\vec{r}_i - \vec{r}_j) = (\vec{r}_i - \vec{r}_j) f$$

Please fill in steps.

$\underbrace{\quad}_{\rightarrow}$  scalar fn.  
in the direction of  
line joining two  
particles.

↓ finally, we find that

Total potential energy

$$U = \sum_i U_i + \frac{1}{2} \sum_{\substack{i,j \\ i \neq j}} U_{ij}$$

Consequence  $\left[ T + U \Rightarrow \text{conserved} \right]$