

Nuclear and Radiation Physics (PHY2005)

Lecture 4

D. Margarone

2021-2022



QUEEN'S
UNIVERSITY
BELFAST

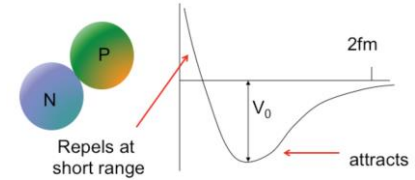
Recap & Learning Goals

Summary of Lecture 3 (Chap.2)

- Properties of the nuclear force
 - ✓ nucleon-nucleon interaction
 - ✓ spin-orbit interaction
- Exchange force model
 - ✓ virtual meson exchange

$$\begin{aligned} V(r) &= +\infty & r < R_{core} & \quad (R_{core} \sim 0.5 \text{ fm}) \\ V(r) &= -V_0 & R_{core} \leq r \leq R \\ V(r) &= 0 & r > R \end{aligned}$$

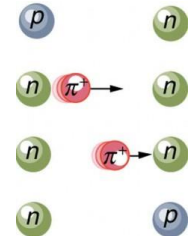
$$P = \frac{N(\uparrow) - N(\downarrow)}{N(\uparrow) + N(\downarrow)}$$



$$R = \frac{200 \text{ MeV} \cdot \text{fm}}{m_\pi c^2}$$

n-p scattering experiments
($T \sim 100 \text{ MeV}$)

$$\theta = \frac{V_0}{2T} \quad (\theta \leq 10^\circ)$$



Learning goals of of Lecture 4 (Chap.3)

- Knowing the terminology and notation of the nuclear *Shell model*
- Understanding physical reasoning behind the *Shell model*

3. Nuclear Models

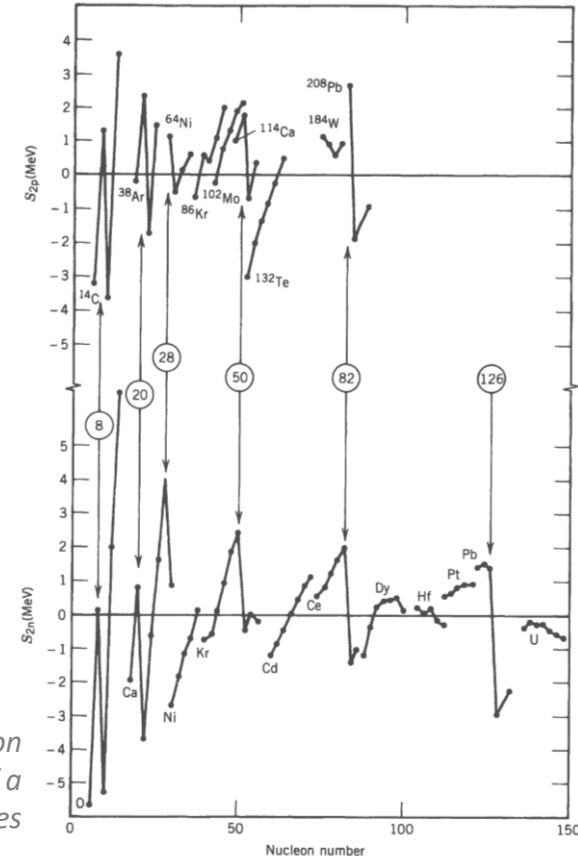
3.1. The Shell model I

Nuclear Shell model

- heavy nuclei (many-body problem) 😞 → simplified nuclear models 😊
- nuclear Shell model → valence nucleons experience a potential created by the nucleons themselves
- nucleons are large (size of nucleus) → collisions (?)
- experimental evidence → nuclear shells
- magic numbers → filled major shells
(Z or $N = 2, 8, 20, 28, 50, 82, \text{ and } 126$)
- nucleons orbit as if they were transparent to one another



*two-proton/neutron
separation energies of a
sequence of isotones/isotopes*



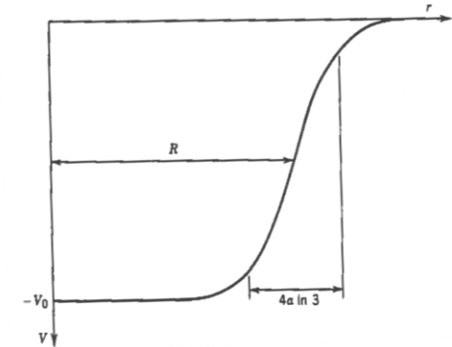
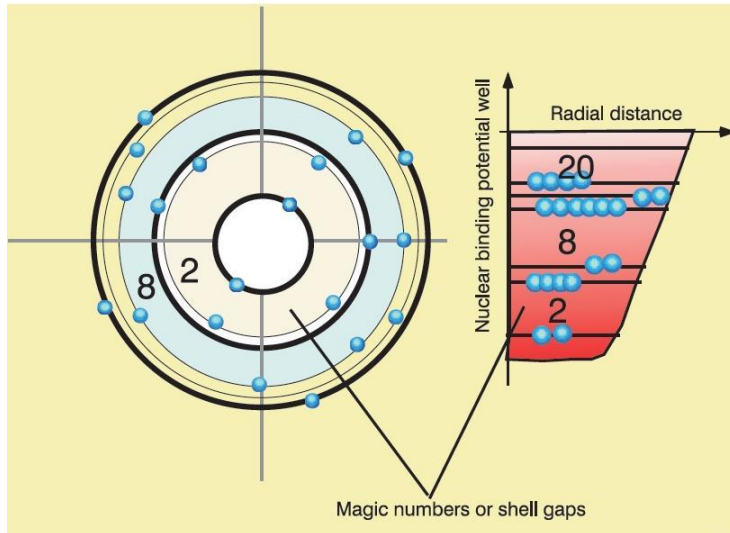
3. Nuclear Models

3.1. The Shell model II

Nuclear Shell model

- *Shell model potential* → Woods-Saxon potential

$$V(r) = \frac{-V_0}{1 + \exp\left(\frac{r-R}{a}\right)}$$

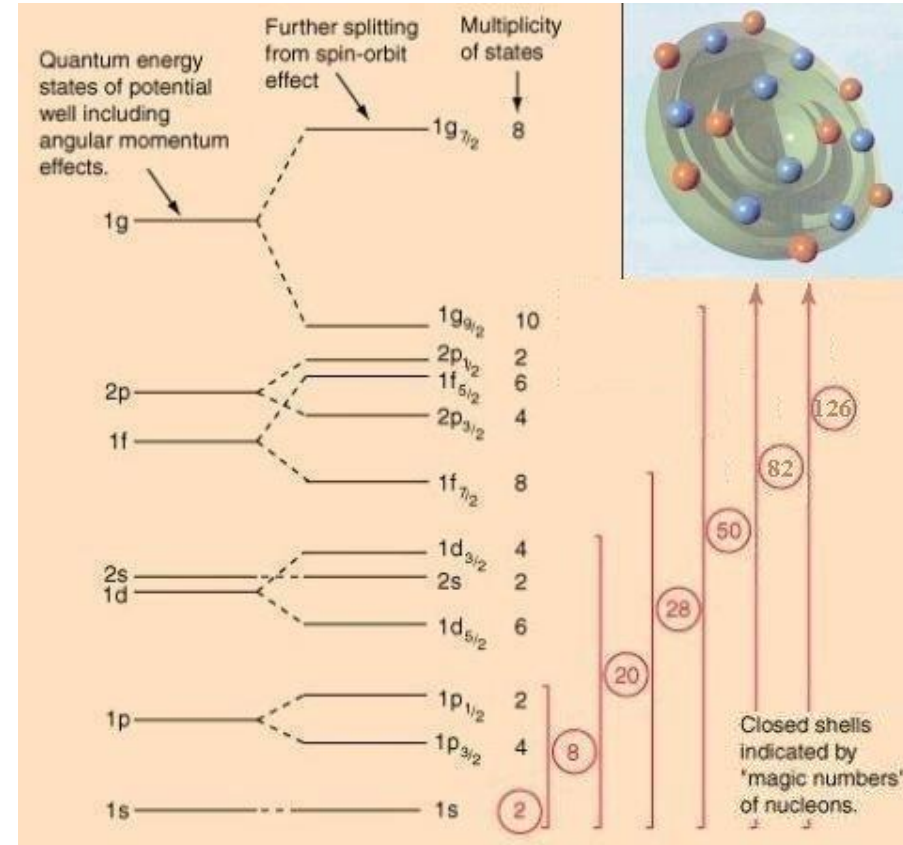


3. Nuclear Models

3.1. The Shell model III

Quantum states and magic numbers

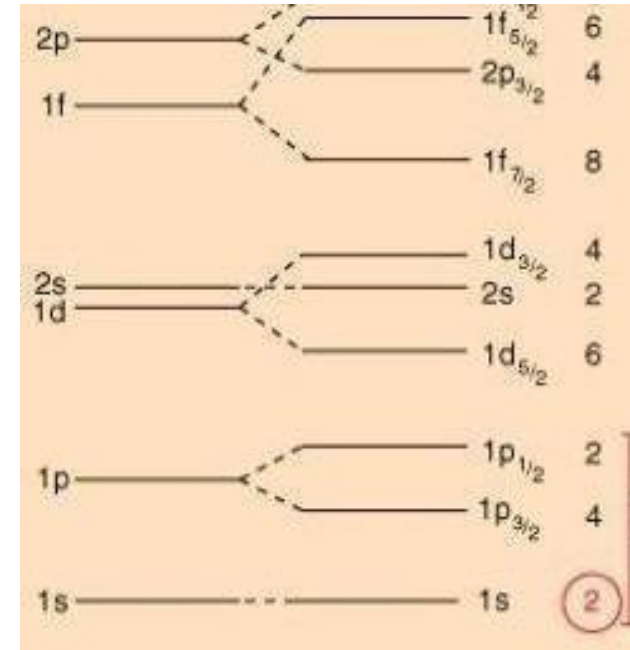
- level degeneracy (Pauli principle) $\rightarrow 2(2L + 1)$
 $m_L \rightarrow (2L + 1)$; $m_S \rightarrow 2$
- Nuclear spectroscopic notation $\rightarrow n$ is not the principal quantum number (!)
- OK only for magic numbers 2, 8, and 20 ☹
- spin-orbit potential \rightarrow OK for all magic numbers! 😊
- $\mathbf{J} = \mathbf{L} + \mathbf{S} \rightarrow \mathbf{J} = \mathbf{L} \pm \frac{1}{2}$ (for a single nucleon $s = \frac{1}{2}$)
- ... except for $L = 0 \rightarrow$ only $J = 1/2$ is allowed
- level degeneracy (spin-orbit) $\rightarrow (2J + 1)$



3. Nuclear Models

3.1. The Shell model IV

- ✓ **1s** ($L = 0$)
 - L degeneracy (no spin-orbit): $2(2L + 1) = 2$
 - possible J values ($L \pm \frac{1}{2}$): $1/2 \rightarrow 1s_{1/2}$
 - J degeneracy (spin-orbit): $2J + 1 = 2$
- ✓ **1p** ($L = 1$)
 - L degeneracy (no spin-orbit): $2(2L + 1) = 6$
 - possible J values ($L \pm \frac{1}{2}$): $3/2, 1/2 \rightarrow 1p_{3/2}; 1p_{1/2}$
 - J degeneracy (no spin-orbit): $2J + 1 = 4; 2$
- ✓ **1d** ($L = 2$)
 - L degeneracy (no spin-orbit): $2(2L + 1) = 10$
 - possible J values ($L \pm \frac{1}{2}$): $5/2, 3/2 \rightarrow 1d_{5/2}; 1d_{3/2}$
 - J degeneracy (spin-orbit): $2J + 1 = 6; 4$
- ✓ **1f** ($L = 3$)
 - L degeneracy (no spin-orbit): $2(2L + 1) = 14$
 - possible J values ($L \pm \frac{1}{2}$): $7/2, 5/2 \rightarrow 1f_{7/2}; 1f_{5/2}$
 - J degeneracy (spin-orbit): $2J + 1 = 8; 6$



3. Nuclear Models

3.1. The Shell model V

Magnetic Dipole Moments (Shell model)

- Shell model → not exact agreement with measured magnetic dipole moments ☹
- expectation value of magnetic moment operator of odd-A nucleon (state with max z-projection of L):

$$g_L = +1 \text{ (proton)}$$

$$g_L = 0 \text{ (neutron)}$$

$$g_S = +5.6 \text{ (proton)}$$

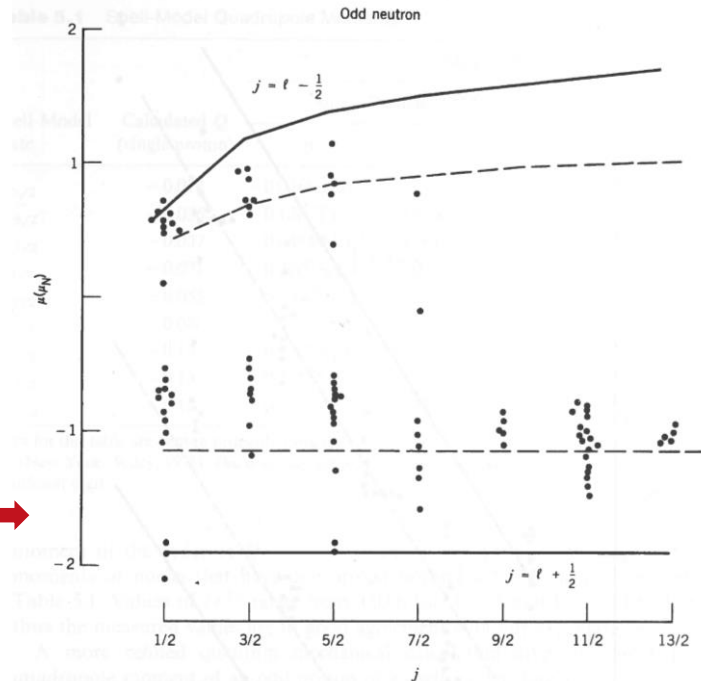
$$g_L = -3.8 \text{ (neutron)}$$

$$\mu = \frac{\mu_N(g_L L_z + g_S S_z)}{\hbar}$$

- even-A nuclei (paired) → $\langle \mu \rangle = 0$
- odd-A nuclei (unpaired) → $\langle \mu \rangle \neq 0$

- *theoretical vs. experimental data of odd-A nuclei*
- *exp. data are smaller in magnitude ☹*
- *wrong assumption on g_s (μ -nucleon = μ -nucleus)*
- *meson cloud of bounded nucleon \neq meson cloud of free nucleon*
- *the shell-model theory oversimplifies μ calculation!*

experimental vs. calculated values of the magnetic moments of odd-A nuclei



3. Nuclear Models

3.1. The Shell model VI

Electric Quadrupole Moments (*Shell model*)

- electric quadrupole moment $\rightarrow 3z^2 - r^2$ operator
- expectation value of electric quadrupole operator of odd-A nuclei (state with max z-projection of L, i.e. $m_J = J$)
- single-particle electric quadrupole moment of an odd proton (state J):

$$\langle r^2 \rangle = 3/5 R^2 = 3/5 R_0^2 A^{2/3}$$

$$\langle Q_{sp} \rangle = -\frac{2J-1}{2(J+1)} \langle r^2 \rangle$$

- electric quadrupole moment (subshell with more than a single-particle):

n: number of nucleons in the subshell
($1 \leq n \leq 2J$)

$$\langle Q \rangle = \langle Q_{sp} \rangle \left[1 - 2 \frac{n-1}{2J-1} \right]$$

- $n = 2J$ (subshell lacking only one nucleon) $\rightarrow Q = -Q_{sp}$ ("hole" states)
- Experiments \rightarrow electric quadrupole moment of an odd-n nucleus $\neq 0$!!! ☹
- Shell model fails to predict Q of heavy nuclei ☹
- *the shell-model theory oversimplifies μ calculation!*

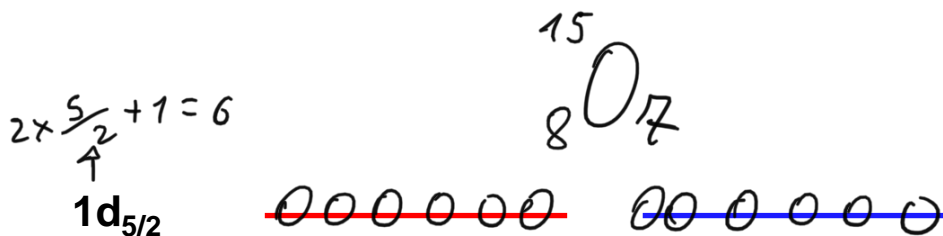
3. Nuclear Models

Example 3.1

Fill the given quantum states with protons and neutrons for the two isotopes ^{15}O and ^{17}O , according to the Nuclear Shell model rules for the level degeneracy in each state.

$$\text{degeneracy} \rightarrow 2J + 1$$

$$J = L \pm \frac{1}{2}$$



$$2$$

$$1p_{1/2}$$



$$1p_{3/2}$$



$$2 \times \frac{3}{2} + 1 = 4$$

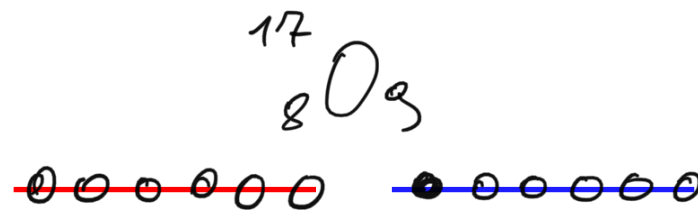
$$1s_{1/2}$$



$$2 \times \frac{1}{2} + 1 = 2$$

protons

neutrons



protons

neutrons

3. Nuclear Models

Example 3.2

Fill the given quantum states with protons and neutrons for the two nuclei ^{15}N and ^{17}F , and specify their nuclear spin and parity.

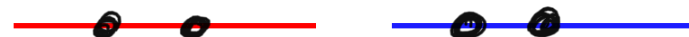
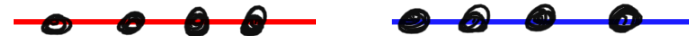
$$\pi \rightarrow (-1)^L$$

$$^{15}_7\text{N}_8 \rightarrow \frac{1}{2}^-$$

$$\text{degeneracy} \rightarrow 2J + 1$$

$$J = L \pm \frac{1}{2}$$

$$^{17}_9\text{F}_8 \rightarrow \frac{5}{2}^+$$



protons

neutrons

protons

neutrons

3. Nuclear Models

Example 3.3

Calculate the theoretical electric quadrupole moments of the nuclei ^{15}O , ^{17}O , ^{15}N , and ^{17}F , according to the Nuclear Shell model assumptions.

$$\langle Q_{sp} \rangle = -\frac{2J-1}{2(J+1)} \langle r^2 \rangle$$

$$\langle r^2 \rangle = \frac{3}{5} R_0^2 A^{2/3}$$

$$\langle Q \rangle = \langle Q_{sp} \rangle \left[1 - 2 \frac{n-1}{2J-1} \right]$$

$^{15}_8\text{O}_7$ protons in complete sub-shells $\rightarrow \langle Q \rangle = 0$

$^{17}_8\text{O}_9$ neutrons do not contribute $\rightarrow \langle Q \rangle = 0$

$^{15}_7\text{N}_8$ 1 proton in $1p_{1/2} \rightarrow \langle Q_{sp} \rangle = -\frac{2 \cdot \frac{1}{2} - 1}{2(\frac{1}{2} + 1)} \langle r^2 \rangle = 0$ ($L=0$!)

$^{17}_9\text{F}_8$ 1 proton in $1d_{5/2} \rightarrow \langle Q_{sp} \rangle = -\frac{2 \cdot \frac{5}{2} - 1}{2(\frac{5}{2} + 1)} \langle r^2 \rangle =$

$$= -\frac{4}{7} \times \frac{3}{5} \times (1.2 \text{ fm})^2 \times (10^{-15} \text{ s})^2 \times 17^{2/3} = 3.26 \times 10^{-30} \text{ C} \times \text{m}^2$$

