

## Lecture 14:

### Solved problems on angular momentum

**Exercise 1:** Demonstrate that  $L^2 = L_-L_+ + L_3^2 + L_3$

**Solution**

We know that:

a)  $L^2 = L_1^2 + L_2^2 + L_3^2$

b)  $L_- = L_1 - iL_2$

c)  $L_+ = L_1 + iL_2$

d)  $[L_+, L_-] = 2L_3$

Let us calculate  $L_-L_+ = (L_1 - iL_2)(L_1 + iL_2) = L_1^2 + L_2^2 + i(L_1L_2 - L_2L_1)$ . The term in brackets is the commutator  $[L_1, L_2]$  which, neglecting  $\hbar$  (we consider here to have  $\hbar = 1$ ), is equal to  $[L_1, L_2] = iL_3$ .

We thus have  $L_-L_+ = L_1^2 + L_2^2 - L_3$ . If we add  $L_3 + L_3^2$ , we obtain:  $L_-L_+ + L_3 + L_3^2 = L_1^2 + L_2^2 + L_3^2 = L^2$ , which concludes our demonstration. ■

**Exercise 2:** Demonstrate that, if  $J^2$  has an eigenvalue  $j(j+1)$ , there exists  $2j+1$  states that are eigenfunctions of  $J_3$  with eigenvalues  $j, j-1, j-2, \dots, -j$ .

**Solution**

First of all, we demonstrate that there exists an eigenfunction of  $J^2$  with eigenvalue  $j(j+1)$  which has the maximum eigenvalue for  $J_3$ , and that this is equal to  $j$ . In other words, there exists a state  $\psi$  such as  $J^2\psi = j(j+1)\psi$  and, simultaneously,  $J_3\psi = j\psi$ . If the state  $\psi$  has the maximum eigenvalue of  $J_3$ , it immediately follows that  $J_+\psi = 0$ . So, remembering that  $J^2 = J_-J_+ + J_3^2 + J_3$ , we have:

$$J^2\psi = J_-J_+\psi + J_3^2\psi + J_3\psi$$

The first term is equal to 0 ( $\psi$  is the state with maximum eigenvalue for  $J_3$  implying that  $J_+\psi = 0$ ), the second is equal to  $j^2\psi$  and the third is equal to  $j\psi$  giving:  $J^2\psi = j(j+1)\psi$ .

On the other hand, there must be also a state with the *minimum* eigenvalue for  $J_3$  meaning that I can only apply the operator  $J_-$  a certain amount of times (which I will denote as  $n$ ).

We will have:

$$J_3J_-\psi = (j-1)J_-\psi, \quad J_3J_-^2\psi = (j-2)J_-^2\psi, \quad J_3J_-^3\psi = (j-3)J_-^3\psi, \quad \dots, \quad J_3J_-^n\psi = 0$$

For this state of minimum eigenvalue of  $J_3$  we have:

$$J^2 J_-^n \psi = (J_+ J_- + J_3^2 - J_3) J_-^n \psi = ((j-n)^2 - (j-n)) J_-^n \psi$$

However, also the state  $J_-^n \psi$  is eigenfunction of  $J^2$  with eigenvalue  $j(j+1)$ . We thus have the equality:

$$j(j+1) = (j-n)^2 - (j-n)$$

which implies  $n = 2j$ . However,  $n$  is a positive integer number (it represents the times I can apply  $J_-$  to the state with maximum eigenvalue of  $J_3$ ) implying that there are  $2j+1$  states all with the same eigenvalue for  $J^2$ . ■

**Exercise 3:** Calculate the spherical harmonics  $Y_{\ell,m}$  associated with: a)  $\ell=0, m=0$ , a)  $\ell=1, m=0$ , a)  $\ell=1, m=1$ .

**Solution**

We know that a general spherical harmonic can be written as:

$$Y_{\ell}^m(\theta, \phi) = (-)^{(m+|m|)/2} i^{\ell} \sqrt{\frac{(2\ell+1)(\ell-|m|)!}{4\pi(\ell+|m|)!}} P_{\ell}^m(\cos \theta) e^{im\phi} \quad (1)$$

where the associated *Legendre* polynomials  $P_{\ell}^m$  are defined as:

$$P_{\ell}^m(x) = (x^2 - 1)^{m/2} \frac{d^m}{dx^m} P_{\ell}(x) \quad (2)$$

and the *Legendre* polynomials  $P_{\ell}$  are defined as:

$$P_{\ell}(x) = \frac{1}{2^{\ell} \ell!} \frac{d^{\ell}}{dx^{\ell}} (x^2 - 1)^{\ell} \quad (3)$$

We then start by calculating  $P_0(x)$ :

$$P_0(x) = 1 \quad (4)$$

and  $P_1(x)$ :

$$P_1(x) = \frac{1}{2} \frac{d}{dx} (x^2 - 1) = x \quad (5)$$

Then,  $P_0^0(x)$ :

$$P_0^0(x) = 1 \quad (6)$$

and  $P_1^0(x)$ :

$$P_1^0(x) = x \quad (7)$$

and  $P_1^1(x)$ :

$$P_1^1(x) = (x^2 - 1)^{1/2} \frac{d}{dx} x = \sqrt{x^2 - 1} \quad (8)$$

We now have all the ingredients to calculate the various spherical harmonics (bear in mind that  $x = \cos \theta$ ):

$$a) \quad Y_{0,0} = \frac{1}{\sqrt{4\pi}} \quad (9)$$

$$b) \quad Y_{1,0} = i \frac{3}{\sqrt{4\pi}} \cos \theta \quad (10)$$

$$c) \quad Y_{1,1} = -i \frac{3}{\sqrt{8\pi}} \sin \theta e^{i\phi} \quad (11)$$

**Exercise 4:** Consider the Hamiltonian:

$$\hat{H} = g \hat{L} \cdot \hat{s}$$

Here  $\hat{L}$  is the angular momentum operator,  $\hat{s}$  is the spin of the electron and  $g$  is a coupling constant.

**a** Making use of the commutation rule between  $\hat{L}$  and  $\hat{s}$ , demonstrate that, the square of total angular momentum is given by:

$$\hat{J}^2 = \hat{L}^2 + \hat{s}^2 + 2\hat{L} \cdot \hat{s}$$

**b** If we have an electron with  $\ell = 1$  and  $s = 1/2$ , what are the possible values of the eigenvalues of  $\hat{J}_z$  and  $\hat{J}^2$ ?

**c** Still considering  $\ell = 1$  and  $s = 1/2$ , find the eigenvalues of the Hamiltonian (energy) for all the possible configurations of the total angular momentum.

**Solution**

$J^2 = (L + s)^2 = L^2 + s^2 + Ls + sL$ . However  $[L, s] = 0$  implying that  $J^2 = L^2 + s^2 + 2Ls$ . The Hamiltonian can thus be written as:

$$H = g(J^2 - L^2 - s^2)/2 \quad (12)$$

The possible values for the total angular momentum  $j$  are  $j = \ell \pm s = 3/2$  or  $1/2$ . We thus have two possible energy eigenvalues:

$$a) \quad j = 3/2 \quad \ell = 1 \quad s = 1/2 \quad \rightarrow \quad E = g/2 \quad (13)$$

or

$$b) \quad j = 1/2 \quad \ell = 1 \quad s = 1/2 \quad \rightarrow \quad E = -g/2 \quad (14)$$