

Existence and Stability of Travelling Waves

Björn Sandstede
(March 2007)

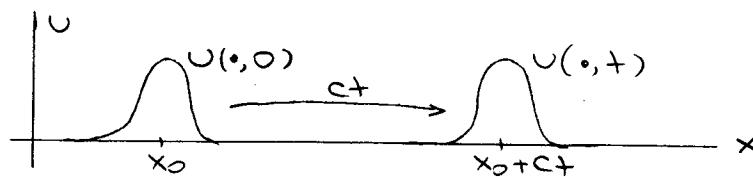
§1 Introduction

Travelling waves :

$$U(x,t) = \underbrace{U_*(x-ct)}_{\text{constant velocity } c \text{ / same shape \& profile}}$$

constant velocity c / same shape & profile

x position
 t time
 c speed



Applications :

- nonlinear optics (electric field)
- nerve impulses (voltage & ion concentrations)
- chemical waves (concentrations)
- flame fronts (temperature & concentrations)
- buckled structures (displacement)
- fluids (surface elevation)

Often modelled by partial differential equations (PDEs) :

- Reaction-diffusion eqns. $u_t = D u_{xx} + f(u)$
- Hamiltonian PDEs :
 - Korteweg-de Vries $u_t + u_{xxx} + (u^2)_x = 0$
 - nonlinear Schrödinger $i u_t + u_{xx} + |u|^2 u = 0 \quad (u \in \mathbb{C})$
- Conservation Laws $u_t = f(u)_x$ or $u_t = \underbrace{u_{xx}}_{\text{viscous}} + f(u)_x$
- Fourth-order PDEs $u_t + u_{xxxx} = f(u, u_x, u_{xx})$
- Lattice equations $\partial_t u_n = \alpha(u_{n+1} - 2u_n + u_{n-1}) + f(u_n) \quad (n \in \mathbb{Z})$
- Multidimensional PDEs $u_t = D \Delta u + f(u)$
with $x \in \mathbb{R} \times \Omega$ and $\Omega \subset \mathbb{R}^d$

Concentrate on reaction-diffusion equations

$$(1) \quad u_t = D u_{xx} + f(u) \quad x \in \mathbb{R}, t > 0, u \in \mathbb{R}^n$$

$$D = \text{diag}(d_j) \text{ with } d_j > 0$$

Travelling wave: $u(x, t) = u_* (\underbrace{x - ct}_{\xi = x - ct} \text{ comoving frame})$

→ substitute into (1) to get

$$(2) \quad -c u_\xi = D u_{\xi\xi} + f(u) \quad \xi \in \mathbb{R}, c \in \mathbb{R}$$

2nd-order ordinary differential equation (ODE)
for profile u_* travelling with speed c

Alternatively, use $(x, t) \rightarrow (\xi, t) = (x - ct, t)$ which transforms (1)

via $u(x, t) = \tilde{u}(x - ct, t) = \tilde{u}(\xi, t)$ into

$$(3) \quad u_t = D u_{\xi\xi} + c u_\xi + f(u) \quad \xi \in \mathbb{R}, t > 0$$

Travelling wave
of (1) with speed c_*



ODE solution
of (2) for $c = c_*$



stationary solution
of (3) for $c = c_*$

↑
offer two complementary
viewpoints that can be exploited

• PDEs on \mathbb{R} : travelling waves correspond to ODE solutions

• PDEs on $\mathbb{R} \times \Omega$: $u_t = D(u_{xx} + \Delta_y u) + f(u) \quad y \in \Omega \quad (\text{PDE})$

$D u_{xx} + c u_x + D \Delta_y u + f(u) = 0 \quad (\text{TW})$

travelling-wave equation is still a PDE

• Lattice equations: $\dot{u}_n = \alpha(u_{n+1} - 2u_n + u_{n-1}) + f(u_n) \quad (n \in \mathbb{Z})$

travelling wave: $u_n(t) = \varphi(n - ct)$

$-c \varphi_\xi(\xi) = \alpha(\varphi(\xi+1) - 2\varphi(\xi) + \varphi(\xi-1)) + f(\varphi(\xi))$

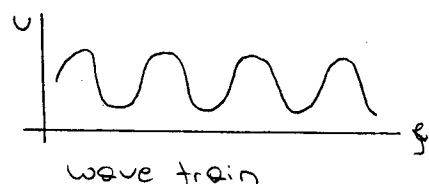
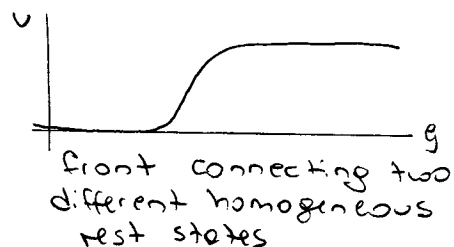
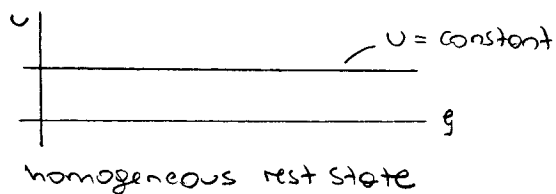
travelling-wave equation is an
advanced-retarded functional differential eqn.

"Interesting" travelling waves:

PDE

equilibria of

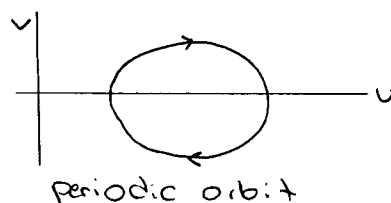
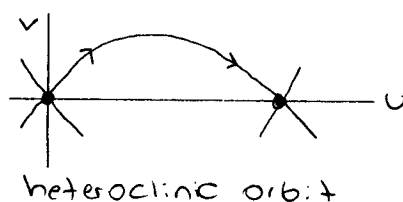
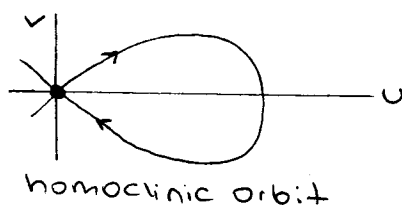
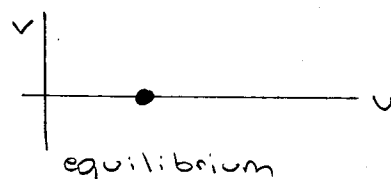
$$v_t = D u_{\xi\xi} + c u_{\xi} + f(u)$$



ODE

solution of

$$\begin{pmatrix} u \\ v \end{pmatrix}_{\xi} = \begin{pmatrix} v \\ -D^{-1}[cv + f(u)] \end{pmatrix}$$



Of interest to us : Connection between PDE and ODE

-
- 1) What does PDE stability of a wave tell us about its properties as a solution to the ODE?
 - 2) What can the ODE tell us about PDE stability?
 - 3) Bifurcations of waves :
Should we analyse them using (PDE) or (ODE)?
 - 4) Can these ideas be applied to multidimensional PDEs and to lattice equations?

Stability of travelling waves u_* with speed c_* :

→ equilibria of PDE (3) for $c=c_*$

→ Linearise right-hand side of (3) to get linear operator

$$\mathcal{L}_* = D\partial_{\xi\xi} + c_* \partial_{\xi} + f_u(u_*(\xi))$$

Function spaces for admissible perturbations :

(i) $X = L^2(\mathbb{R}, \mathbb{R}^n)$ square-integrable functions

$$\|u\|_{L^2} = \left(\int_{\mathbb{R}} |u(x)|^2 dx \right)^{1/2} \quad (\text{Hilbert space})$$

(ii) $X = C_{\text{unif}}^0(\mathbb{R}, \mathbb{R}^n)$ bounded, uniformly continuous functions

$$\|u\|_{C^0} = \sup_{x \in \mathbb{R}} |u(x)| \quad (\text{Banach space: no scalar product!})$$

Spectral stability : calculate spectrum of $\mathcal{L}_* : X \rightarrow X$
defined on domain $\mathcal{D}(\mathcal{L}_*) \subset X$.

• $\lambda \in \rho(\mathcal{L}_*)$ resolvent set iff $\exists K : \forall h \in X \exists ! u \in X : (\mathcal{L}_* - \lambda)u = h$
and $\|u\|_X \leq K \|h\|_X$

• Spectrum of \mathcal{L}_* : $\text{spec}(\mathcal{L}_*) := \mathbb{C} \setminus \rho(\mathcal{L}_*)$

• λ is an eigenvalue of \mathcal{L}_* iff $\exists u \in X \setminus \{0\} : \mathcal{L}_* u = \lambda u$

• We write

$$\text{spec}(\mathcal{L}_*) = \Sigma_{\text{pt}} \cup \Sigma_{\text{ess}} \quad \begin{array}{l} \text{essential spectrum:} \\ \text{defined as } \text{spec}(\mathcal{L}_*) \setminus \Sigma_{\text{pt}} \end{array}$$

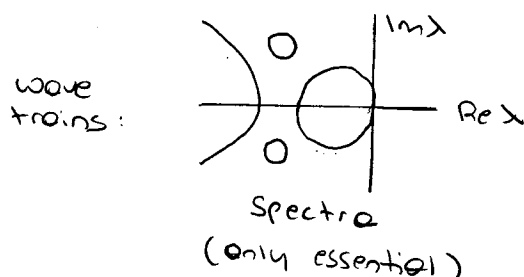
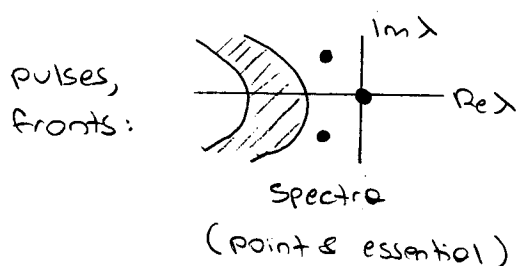
point spectrum :
defined as the union of all
isolated eigenvalues with
finite multiplicity

Example $\mathcal{L} : \ell^\infty \rightarrow \ell^\infty, (a_0, a_1, \dots) \mapsto (0, a_0, a_1, \dots)$

→ $\lambda=0$ is not an eigenvalue but $0 \in \text{spec}(\mathcal{L})$

since $\mathcal{L}u = (1, 0, 0, \dots)$ does not have a solution in ℓ^∞ .

Expectation



Example If $u_*' \neq 0$ and $u_*' \in X$, then $0 \in \text{spec}(\mathcal{L}_*)$

Proof

The travelling wave $u_*(\xi)$ satisfies

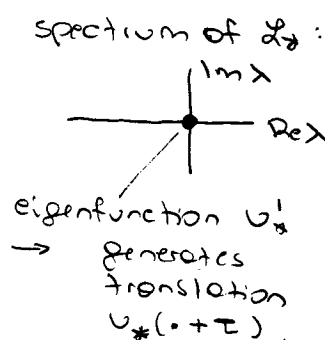
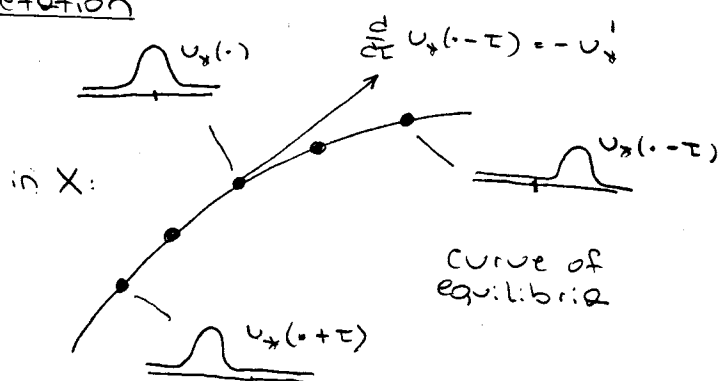
$$D u_*''(\xi) + c_* u_*'(\xi) + f(u_*(\xi)) = 0 \quad \forall \xi \in \mathbb{R}$$

Taking a further derivative $\frac{d}{d\xi}$, we find

$$D u_*'''(\xi) + c_* u_*''(\xi) + f_u(u_*(\xi)) u_*'(\xi) = 0 \quad \forall \xi \in \mathbb{R}$$

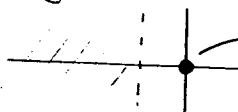
and therefore $\mathcal{L}_* u_*' = 0$ □

Interpretation

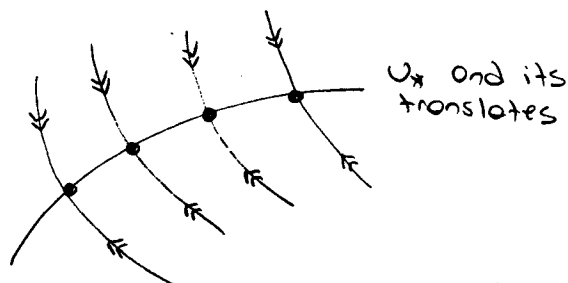


Nonlinear stability

Theorem: If u_* is a travelling wave of $u_t = D u_{xx} + f(u)$ on X with spectrum



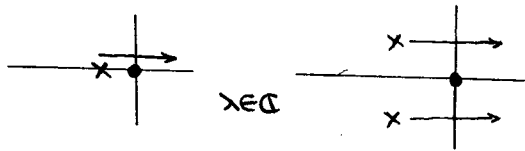
is nonlinearly asymptotically stable: $\forall \varepsilon_* > 0 \exists \delta_* > 0$:
 $\forall u_0 \in X$ with $\|u_0 - u_*\|_X < \delta_*$, we have $\|u(\cdot, t) - u_*\|_X < \varepsilon_*$
 $\forall t \geq 0$ and $\exists \xi_* \in \mathbb{R}$ with $\|u(\cdot, t) - u_*(\cdot - \xi_*)\|_X \rightarrow 0$
 exponentially as $t \rightarrow \infty$.



Remark: Theorem is also true if some of the diffusion coefficients vanish or for higher-order PDEs.

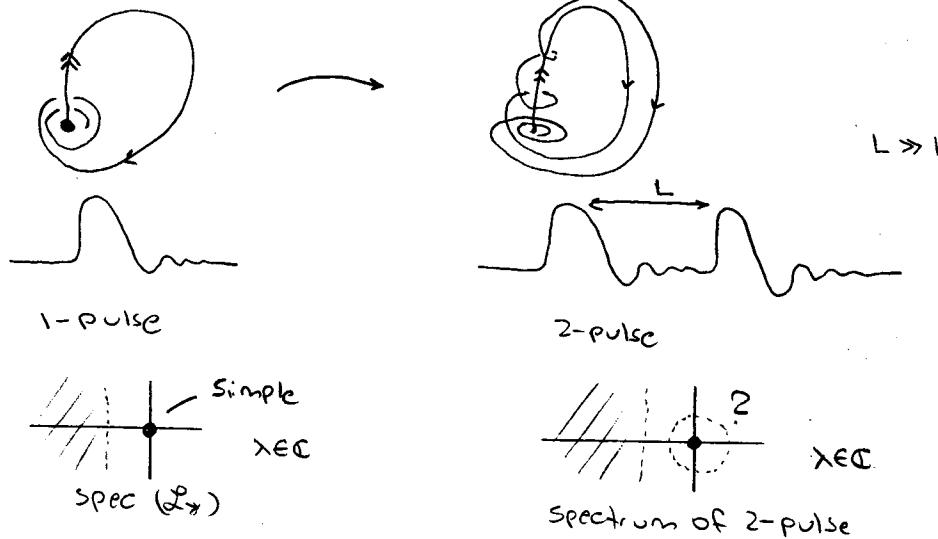
Bifurcations of travelling waves:

(i) Saddle-node and Hopf:

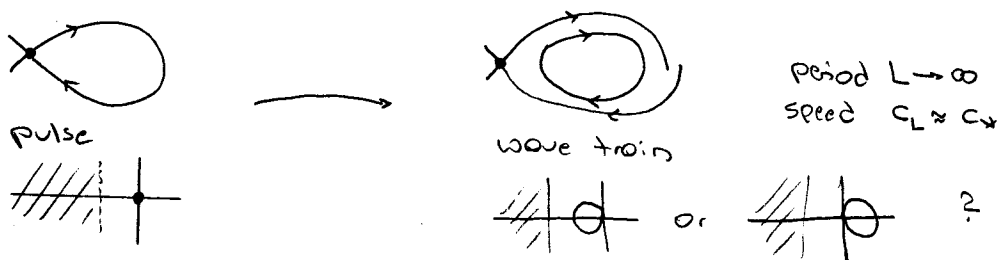


as a parameter μ is varied

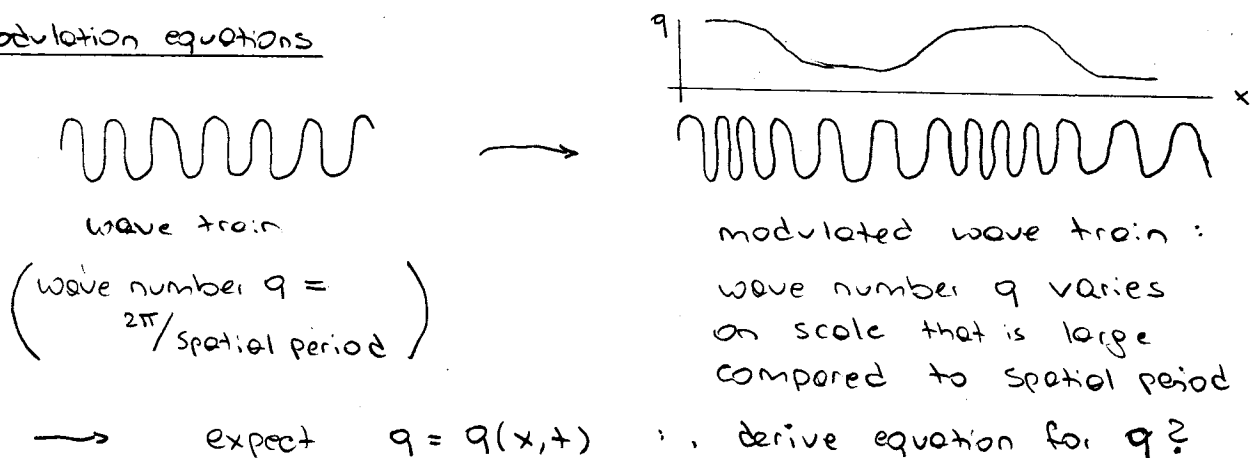
(ii) multi-pulses:



(iii) wave trains



Modulation equations



§2 Spectra

Recall $\mathcal{L}_* = D\partial_{xx} + c_* \partial_x + f_u(u_*(x))$ on $X = L^2$ or $X = C^0_{\text{unit}}$

$$\lambda \in \mathbb{C} \setminus \text{spec } \mathcal{L}_* \text{ iff } \exists k > 0 : \forall h \in X \exists ! u \in X : (\mathcal{L}_* - \lambda)u = h, \|u\|_X \leq k \|h\|_X$$

§2.1 Homogeneous rest states $u_*(x) = u_0$ constant

$$(4) \quad u_t = D u_{xx} + c_* u_x + f_u(u_0) u \quad \text{constant coefficients}$$

$$u(x, t) = e^{\lambda t + \nu x} v_0 \quad \text{for some } \nu \in \mathbb{C}, v_0 \in \mathbb{C}^n \setminus \{0\}$$

$$\rightarrow \lambda v_0 = [D\nu^2 + c_* \nu + f_u(u_0)] v_0$$

$$\rightarrow d(\lambda, \nu) := \det [D\nu^2 + c_* \nu + f_u(u_0) - \lambda] = 0 \quad \begin{array}{l} \text{necessary \& sufficient} \\ \text{for having an eigenmode} \\ e^{\lambda t + \nu x} v_0 \end{array}$$

↳ linear dispersion relation

intuition: want $e^{\lambda t + ikx} v_0$ for $k \in \mathbb{R}$ to satisfy (4)

$$\rightarrow \text{spec } \mathcal{L}_* = \left\{ \lambda \in \mathbb{C} \mid d(\lambda, ik) = \det (-k^2 D + ikc_* + f_u(u_0) - \lambda) = 0 \text{ for some } k \in \mathbb{R} \right\}$$

Theorem Assume that $u_*(x) = u_0$, then

$$\text{spec } \mathcal{L}_* = \left\{ \lambda \in \mathbb{C} \mid d(\lambda, ik) = 0 \text{ for some } k \in \mathbb{R} \right\} \text{ on } L^2 \text{ and } C^0_{\text{unit}}$$

Proof Fix $\lambda \in \mathbb{C}$.

(i) Assume $d(\lambda, ik) \neq 0 \forall k \in \mathbb{R}$. Pick $h \in X$ and consider

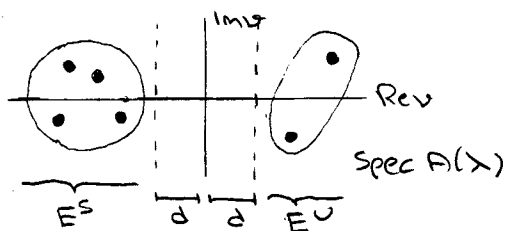
$$D u_{xx} + c_* u_x + f_u(u_0) u = h(x) \quad \text{constant coefficients}$$

which we write as

$$(5) \quad \begin{pmatrix} u \\ v \end{pmatrix}_x = \underbrace{\begin{pmatrix} 0 & 1 \\ D^{-1}[\lambda - f_u(u_0)] & -c_* D^{-1} \end{pmatrix}}_{=: A(\lambda)} \begin{pmatrix} u \\ v \end{pmatrix} + \begin{pmatrix} 0 \\ h(x) \end{pmatrix}$$

$$\text{key} \quad \det [A(\lambda) - \nu] = d(\lambda, \nu) \frac{1}{\det D}$$

$\rightarrow A(\lambda)$ is hyperbolic since $d(\lambda, ik) \neq 0 \forall k \in \mathbb{R}$



2n spatial eigenvalues
in $\text{spec}(A(\lambda))$
for each $\lambda \in \mathbb{C}$

\rightarrow generalized stable and unstable eigenspaces with
spectral projections $P^s(\lambda)$ and $P^u(\lambda)$

$$E^S(\lambda) \oplus E^U(\lambda) = \mathbb{C}^{2n}, \quad P^U(\lambda) + P^S(\lambda) = 1_{\mathbb{C}^{2n}}$$

$$e^{A(\lambda)x} = e^{A(\lambda)x} P^S(\lambda) + e^{A(\lambda)x} P^U(\lambda) \quad \text{with}$$

$$\begin{cases} \|e^{A(\lambda)x} P^S(\lambda)\| \leq \tilde{K} e^{-d|x|} & x \geq 0 \\ \|e^{A(\lambda)x} P^U(\lambda)\| \leq \tilde{K} e^{-d|x|} & x \leq 0 \end{cases}$$

→ Unique bounded solution of (5) is given by

$$\begin{aligned} \begin{pmatrix} u \\ v \end{pmatrix}(x) &= \int_{-\infty}^x e^{A(\lambda)(x-y)} P^S(\lambda) \begin{pmatrix} 0 \\ h(y) \end{pmatrix} dy + \int_{\infty}^x e^{A(\lambda)(x-y)} P^U(\lambda) \begin{pmatrix} 0 \\ h(y) \end{pmatrix} dy \\ &=: \int_{\mathbb{R}} G(x-y) \begin{pmatrix} 0 \\ h(y) \end{pmatrix} dy \quad \text{with } \|G(x)\| \leq \tilde{K} e^{-d|x|} \end{aligned}$$

Using Young's inequality

$$\|G * H\|_{L^p} \leq \|G\|_{L^1} \|H\|_{L^p} \quad \text{for } 1 \leq p \leq \infty$$

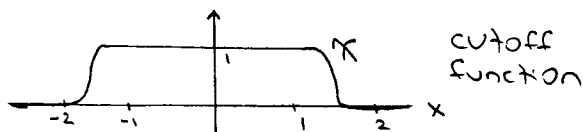
we find

$$\|u\|_{C^0} \leq \frac{2\tilde{K}}{d} \|h\|_{C^0}, \quad \|u\|_{L^2} \leq \frac{2\tilde{K}}{d} \|h\|_{L^2}$$

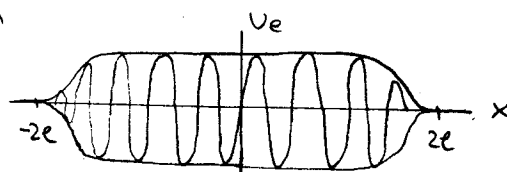
(ii) Assume $d(\lambda, i k_0) = 0$ for some $k_0 \in \mathbb{R}$.

We shall prove that $\exists u_e \in X : \|(\mathcal{L}_\lambda - \lambda) u_e\|_X \leq \frac{1}{e} \|u_e\|_X$ as $e \rightarrow \infty$

which shows that $\lambda \in \text{spec } \mathcal{L}_*$ (otherwise, $\exists K : \|u_e\|_X \leq K \|(\mathcal{L}_\lambda - \lambda) u_e\|_X \leq \frac{K}{e} \|u_e\|_X \rightarrow 0$)



$$\text{Set } u_e(x) = \chi\left(\frac{x}{e}\right) e^{i k_0 x} v_0$$



$$\text{where } v_0 \in \mathbb{C}^n \setminus \{0\} : (-k_0^2 D + i k_0 c_* + f_v(v_0) - \lambda) v_0 = 0$$

$$\rightarrow \|u_e\|_{L^2}^2 \geq 2e \|v_0\|_\infty^2 \quad (\|v_0\|_\infty = \max_j |v_0^{(j)}|)$$

$$\|(\mathcal{L}_\lambda - \lambda) u_e\|_{L^2}^2 \leq \| [D(\frac{2i k_0}{e} \chi' + \frac{1}{e^2} \chi'') + \frac{c_*}{e} \chi'] e^{i k_0 x} v_0 \|_{L^2}^2 \leq \frac{\text{const.}}{e}$$

□ 2

Conclusion v_0 homogeneous rest state : $A(\lambda) = \begin{pmatrix} 0 & 1 \\ D^{-1}[\lambda - f_v(v_0)] & -c_* D^{-1} \end{pmatrix}$

$$\lambda \in \text{Spec } \mathcal{L}_* \iff \text{spec } A(\lambda) \cap i\mathbb{R} \neq \emptyset$$

temporal PDE
Spectrum

spatial ODE
Spectrum

Note

travelling-wave ODE

$$D u_{xx} + c_* u_x + f(u) = 0 \quad \text{or}$$

$$\begin{pmatrix} u \\ v \end{pmatrix}_x = \begin{pmatrix} v \\ -D^{-1}[c_* v + f(u)] \end{pmatrix}$$

$$\text{Linearize about equilibrium } (u, v) = (u_0, 0) \rightarrow A(0) = \begin{pmatrix} 0 & 1 \\ -D^{-1} f_u(u_0) & -c_* D^{-1} \end{pmatrix}$$

Properties of $\text{Spec } \mathcal{L}_*$

$$u_t = \underbrace{Du_{xx} + c_* u_x + f_u(u_*) u}_{= \mathcal{L}_* u}$$

$$u(x,t) = e^{\lambda t + v x} v_0 \quad \text{Soln for some } v_0 \neq 0 \iff$$

$$d(\lambda, v) = \det [Dv^2 + c_* v + f_u(u_*) - \lambda] = 0$$

$$\begin{aligned} \text{Spec}(\mathcal{L}_*) &= \{ \lambda \mid d(\lambda, ik) = 0 \text{ for some } k \in \mathbb{R} \} \\ &= \{ \lambda \mid \text{Spec } A(\lambda) \cap i\mathbb{R} \neq \emptyset \} \end{aligned}$$

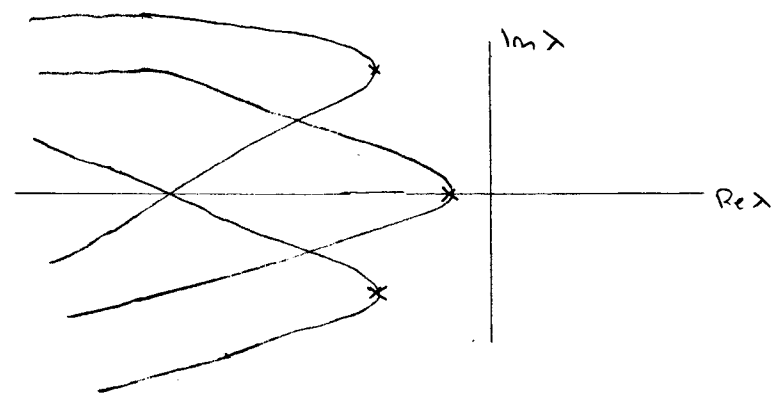
$$A(\lambda) = \begin{pmatrix} 0 & 1 \\ D^{-1}(\lambda - f_u(u_*)) & -c_* D^{-1} \end{pmatrix}$$

$A(0)$ ODE Linearization about $(u_*, 0)$

Structure of $\text{Spec } \mathcal{L}_*$

- all eigenvalues of $f_u(u_*)$ lie in $\text{Spec } \mathcal{L}_*$: set $k=0$
- if $d(\lambda_0, ik_0) = 0$, $d_\lambda(\lambda_0, ik_0) \neq 0$, then \exists curve $\lambda_*(ik)$ defined for $k \neq k_0$ with $\lambda_*(ik_0) = \lambda_0$ so that $\lambda_*(ik) \in \text{Spec}(\mathcal{L}_*)$ for all k
- As $|k| \rightarrow \infty$, we have $\text{Re } \lambda \rightarrow -\infty$.

Assume that $d(\lambda, ik) = 0$ implies $d_\lambda(\lambda, ik) \neq 0$, then



$n=3$

\times eigenvalues of $f_u(u_*)$

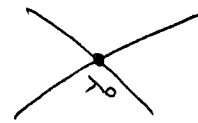
Genericity

- Do we expect to encounter points (λ_0, ik_0) with $d(\lambda_0, ik_0) = d_\lambda(\lambda_0, ik_0) = 0$?
 $\begin{pmatrix} d(\lambda, v) \\ d_\lambda(\lambda, v) \end{pmatrix} = 0$: 2 eqns for 2 unknowns \rightarrow expect finite set of solns (λ_j, v_j) (true when $d_i \neq d_j \forall i \neq j$)

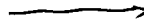
However, "typically", we have $\text{Re } v_j \neq 0 \forall j$

\Rightarrow We do not expect to encounter points where the implicit function theorem fails when continuing λ in \mathbb{C} .

- If $d(\lambda_0, i\omega_0) = d_\lambda(\lambda_0, i\omega_0) = 0$ and $d_{\lambda\lambda}(\lambda_0, i\omega_0) \neq 0$, then



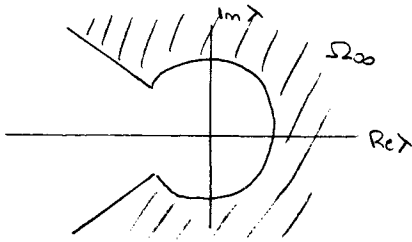
Structurally
unstable



as a
system
parameter
is varied

Implications

(i) Spectrum lies in sector:



$$\lambda = \frac{e^{i\varphi}}{\varepsilon^2} \quad \text{where } 0 < \varepsilon \ll 1, \quad |\varphi| \leq \pi - \delta \quad (\delta > 0 \text{ fixed})$$

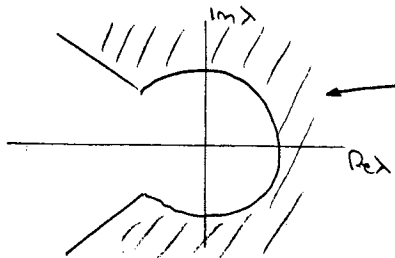
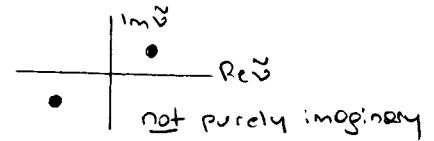
$v \in \text{spec } A(\lambda)$ iff

$$\det [Dv^2 + C_\star v + f_v(u_0) - \lambda] = 0$$

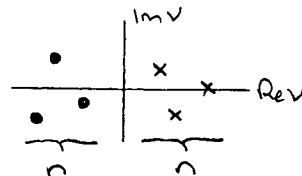
$$\rightarrow \det (v^2 D + v C_\star + f_v(u_0) - \frac{e^{i\varphi}}{\varepsilon^2}) \stackrel{v = \frac{\tilde{v}}{\varepsilon}}{=} \frac{1}{\varepsilon^{2n}} \det (\tilde{v}^2 D + \tilde{v} C_\star \varepsilon + \varepsilon^2 f_v(u_0) - e^{i\varphi}) = 0$$

$$\varepsilon = 0: \det (\tilde{v} D - e^{i\varphi}) = 0 \rightarrow \tilde{v}_j^0 = \pm \sqrt{\frac{e^{i\varphi}}{d_j}}$$

$\varepsilon > 0$: eigenvalues $\tilde{v}_j = \tilde{v}_j^0 + O(\varepsilon)$: hyperbolic



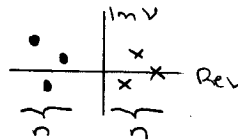
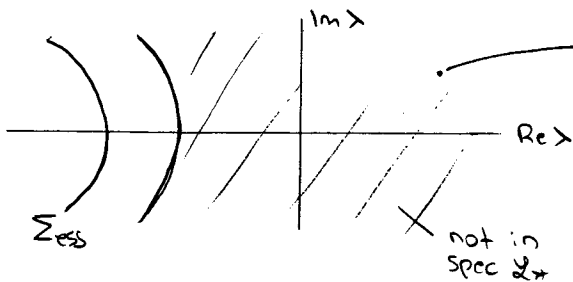
$\text{Spec } A(\lambda)$



$\begin{cases} n \text{ stable} \\ n \text{ unstable} \end{cases}$ Spatial eigenvalues

\rightarrow numbers can change only if $\text{spec } A(\lambda) \cap i\mathbb{R} \neq \emptyset$
ie when $\lambda \in \text{spec } \mathcal{L}_\star$

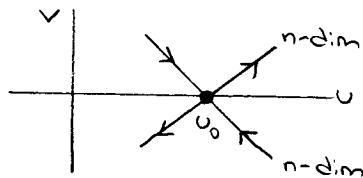
\rightarrow if u_0 is PDE stable, then



$\text{Spec } A(\lambda)$

applied to $A(0) = \text{ODE Linearization about } \begin{pmatrix} u_0 \\ 0 \end{pmatrix}$:

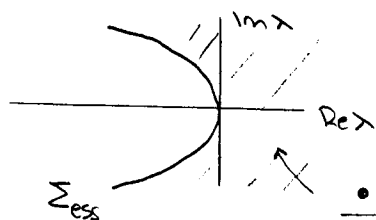
$\begin{pmatrix} u_0 \\ 0 \end{pmatrix}$ saddle with n -dimensional stable and unstable manifolds



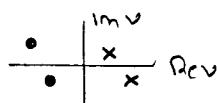
Provided u_0 is PDE stable

"PDE stability \Rightarrow ODE hyperbolicity with $n|n$ split of spatial eigenvalues"

(ii) Suppose

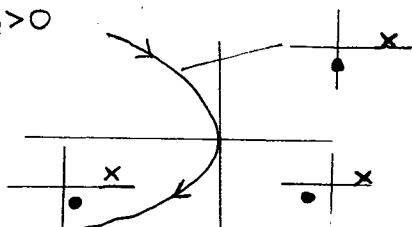


Curve $\lambda = \lambda(i\kappa)$ for $\kappa \in \mathbb{R}$
with $\lambda(0) = 0$, $\lambda'(0) =: -c_g \neq 0$.

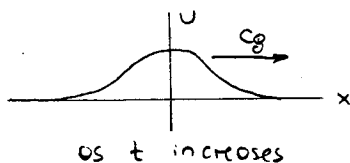


$$\lambda(v) = \lambda(0) + \lambda'(0)v + O(|v|^2) = -c_g v + O(|v|^2)$$

① $c_g > 0$

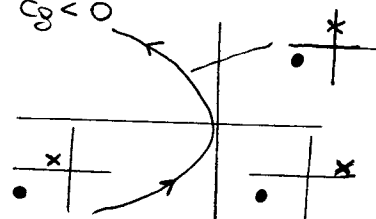


$$\lambda(i\kappa) = -i\kappa c_g$$

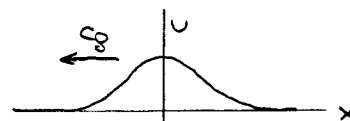


as t increases

② $c_g < 0$



$$\lambda(i\kappa) = i\kappa |c_g|$$

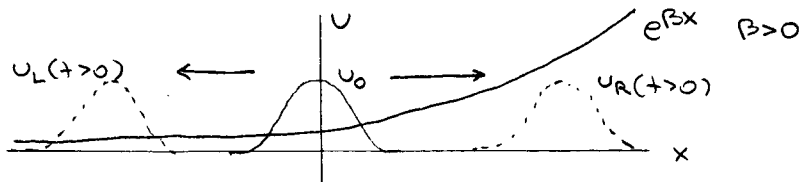


as t increases

$$u_t = \mathcal{L}_0 u$$

to validate this claim, introduce weighted norm

$$\|u\|_\beta := \sup_{x \in \mathbb{R}} e^{\beta x} |u(x)| \quad \text{for } \beta \in \mathbb{R} \text{ fixed}$$



$$u_t = \mathcal{L}_0 u$$

$$\beta > 0: \begin{cases} \|u_L(\cdot, t)\|_\beta & \text{decreases} \\ \|u_R(\cdot, t)\|_\beta & \text{increases} \end{cases}$$

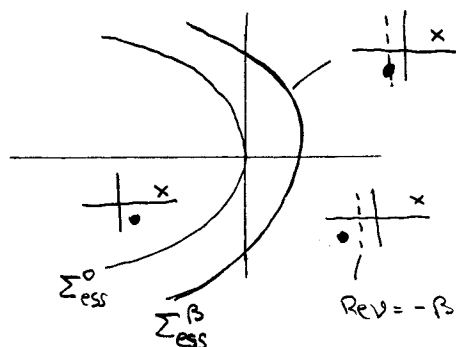
→ compute spectra in

$$\begin{aligned} X_\beta &= \{u \mid \|u\|_\beta < \infty\} \\ &= \{u = e^{-\beta x} v \mid \|v\|_{\ell^\infty} < \infty\} \end{aligned}$$

$$\begin{aligned} \Sigma_{\text{ess}}^\beta &= \{\lambda \in \mathbb{C} \mid d(\lambda, -\beta + i\kappa) = 0 \text{ for some } \kappa \in \mathbb{R}\} \\ &= \{\lambda \in \mathbb{C} \mid \text{spec } A(\lambda) \cap \{\text{Re } v = -\beta\} \neq \emptyset\}. \end{aligned}$$

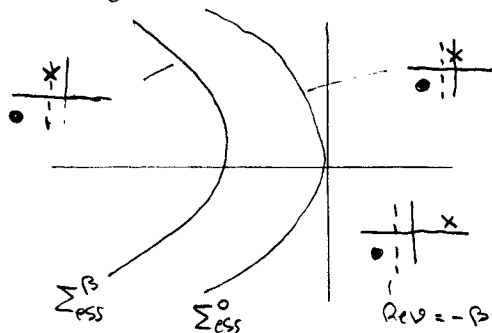
Continue with $0 < \beta \ll 1$:

① $c_0 > 0$



unstable in $\|\cdot\|_\beta$

② $c_0 < 0$



stable in $\|\cdot\|_\beta$

and analogously with $\beta < 0$.

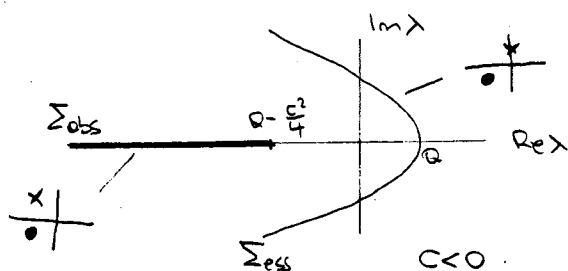
(iii) Absolute spectrum

Consider $\mathcal{L}_\lambda = \partial_{xx} + c\partial_x + a$ for $u \in \mathbb{R}$
on $(-L, L)$ with $u(\pm L) = 0$

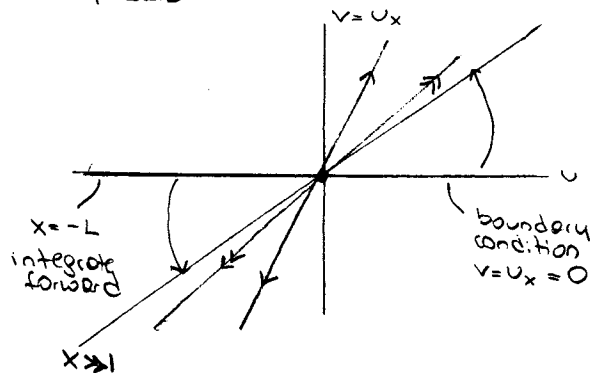
Need to solve $\begin{pmatrix} u \\ v \end{pmatrix}_x = \begin{pmatrix} 0 & 1 \\ \lambda - a & -c \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix} = A(x) \begin{pmatrix} u \\ v \end{pmatrix}$

$$\Sigma_{ESS} = \{ \lambda = -k^2 + ikc + a; k \in \mathbb{R} \}$$

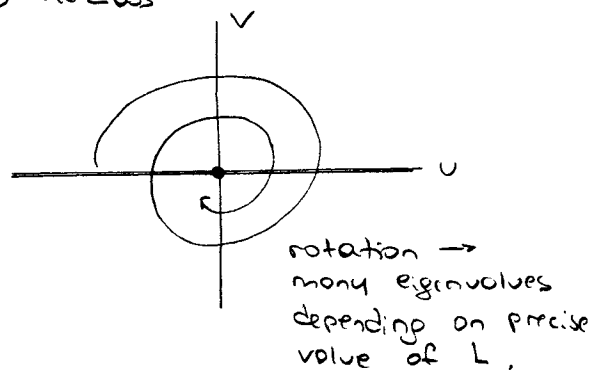
$$\Sigma_{OBS} = \{ \lambda = -k^2 + a - \frac{c^2}{4}; k \in \mathbb{R} \} = \{ \lambda \mid \text{eigenvalues of } A(x) \text{ have the same real part} \}$$



① $\lambda \notin \Sigma_{OBS}$



② $\lambda \in \Sigma_{OBS}$



rotation \rightarrow
many eigenvalues
depending on precise
value of L .

§2.2 Wave trains

$$U_*(x) = U_*(x + \frac{2\pi}{q}) \quad \forall x \in \mathbb{R}$$

Spectra of wave trains:

$$(\mathcal{L}_* - \lambda) u = D u_{xx} + c_* u_x + \underbrace{f_u(U_*(x))}_{\frac{2\pi}{q}\text{-periodic coefficients}} u - \lambda u = 0$$

becomes the ODE

$$(6) \quad \begin{pmatrix} u \\ v \end{pmatrix}_x = A(x; \lambda) \begin{pmatrix} u \\ v \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ D^{-1}[\lambda - f_u(U_*(x))] & -c_* D^{-1} \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix}$$

$$\text{with } A(x + \frac{2\pi}{q}; \lambda) = A(x; \lambda) \quad \forall x.$$

General solution of (6) \rightarrow Floquet theory:

$\exists R(x, \lambda)$: $\frac{2\pi}{q}$ -periodic in x with $R(0, \lambda) = 1$, and $B(\lambda)$:

$$\begin{pmatrix} u \\ v \end{pmatrix}(x) = R(x, \lambda) e^{B(\lambda)x} \begin{pmatrix} u_0 \\ v_0 \end{pmatrix} : \text{general solution of (6)}$$

Theorem Assume that $U_*(x + \frac{2\pi}{q}) = U_*(x) \quad \forall x$, then

$$\text{Spec } \mathcal{L}_* = \{ \lambda \in \mathbb{C} \mid \det(B(\lambda) - ik) = 0 \text{ for some } k \in \mathbb{R} \}$$

Proof Same proof as for homogeneous rest states:

Replace $e^{A(\lambda)x}$ by $R(x, \lambda) e^{B(\lambda)x}$ and use eigenspaces of $B(\lambda)$ instead of $A(\lambda)$

□
6)

Implications

(i) Eigenmodes are of the form $u(x) = u_{\text{per}}(x) e^{ikx}$ with $k \in \mathbb{R}$

where $u_{\text{per}}(x + \frac{2\pi}{q}) = u_{\text{per}}(x) \quad \forall x$ depends on $k \rightarrow$

$$\text{Spec } \mathcal{L}_* = \bigcup_{k \in [0, q)} \underbrace{\text{Spec } \mathcal{L}_*(ik)}_{\text{computed with periodic boundary conditions on } (0, \frac{2\pi}{q})}$$

$$\text{where } \mathcal{L}_*(v) = D(\partial_x + v)^2 + c_*(\partial_x + v) + f_u(U_*(x))$$

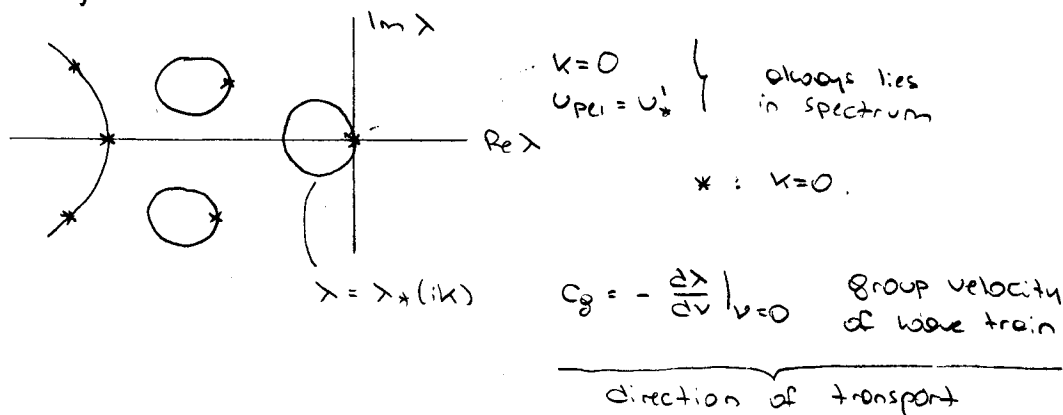
Proof Use $u(x) = v(x) e^{ikx}$ with $v(x + \frac{2\pi}{q}) = v(x) \quad \forall x$

\rightarrow Gives eigenvalue problem $\mathcal{L}_*(v) v = \lambda v$ for v □

(ii) Statements for } sectoriality of spectrum
weighted norms

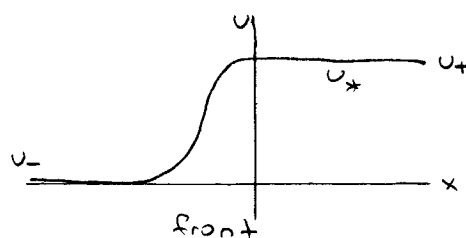
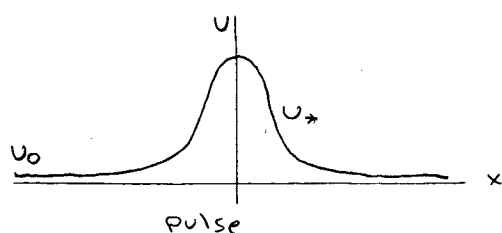
are also true for wave trains. Upon using $B(\lambda)$ instead of $A(\lambda)$

In particular,



§2.3 Pulses and fronts

$$u_*(x) \rightarrow u_{\pm} \text{ as } x \rightarrow \pm \infty$$



Notation

$$\mathcal{L}[u] = D \partial_{xx} + c \partial_x + f(u(x))$$

operator associated with wave u

$$\mathcal{L}_* = \mathcal{L}[u_*]$$

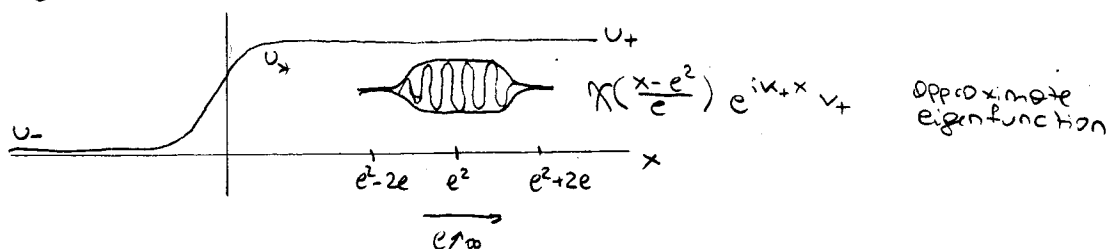
pulse or front

Theorem

If $\lambda \in \text{spec } \mathcal{L}[u_+] \cup \text{spec } \mathcal{L}[u_-]$, then $\lambda \in \text{spec } \mathcal{L}_*$

Proof

E.g. $\lambda \in \text{spec } \mathcal{L}[u_+]$, then



Pulses

Seek eigenvalues : $(\mathcal{L}_* - \lambda) u = 0$ ie bounded solution of

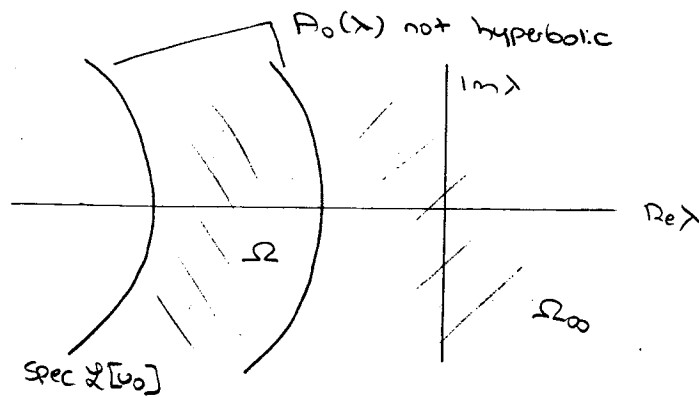
$$(7) \quad \begin{pmatrix} u \\ v \end{pmatrix}_x = \begin{pmatrix} 0 & 1 \\ D^{-1}[\lambda - f(u_*(x))] & -c_* D^{-1} \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix}$$

$$=: A(x, \lambda)$$

and $A(x, \lambda) \rightarrow A_0(\lambda)$ as $|x| \rightarrow \infty$.

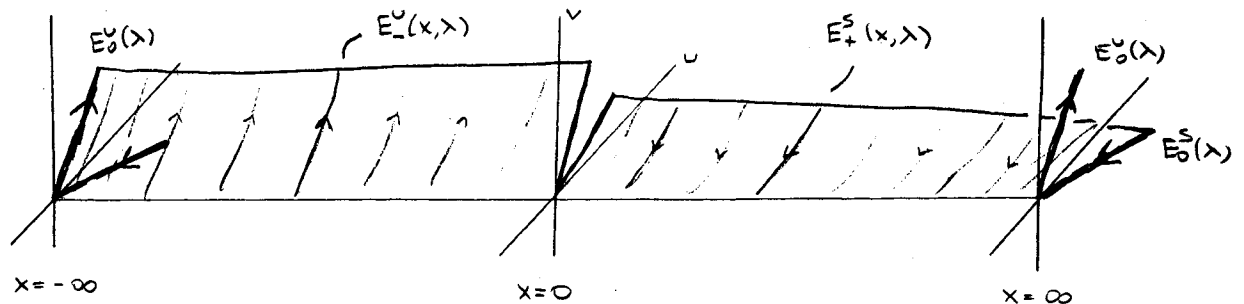
$$\text{where } A_0(\lambda) = \begin{pmatrix} 0 & 1 \\ D^{-1}[\lambda - f(u_0)] & -c_* D^{-1} \end{pmatrix}$$

corresponds to asymptotic rest state u_0



Fix connected component Ω of $\mathbb{C} \setminus \text{spec } \mathcal{L}[u_0]$

→ $A_0(\lambda)$ is hyperbolic $\forall \lambda \in \Omega$:



$$\begin{pmatrix} u \\ v \end{pmatrix}_x = A_0(\lambda) \begin{pmatrix} u \\ v \end{pmatrix}$$

$$(7) \quad \begin{pmatrix} u \\ v \end{pmatrix}_x = A(x, \lambda) \begin{pmatrix} u \\ v \end{pmatrix}$$

$$\begin{pmatrix} u \\ v \end{pmatrix}_x = A_0(\lambda) \begin{pmatrix} u \\ v \end{pmatrix}$$

$$E_+^s(x_0, \lambda) = \left\{ \begin{pmatrix} u_0 \\ v_0 \end{pmatrix} \in \mathbb{C}^{2n} \mid \begin{pmatrix} u \\ v \end{pmatrix}(x) \rightarrow 0 \text{ as } x \rightarrow \infty \text{ where } \begin{pmatrix} u \\ v \end{pmatrix}(x_0) = \begin{pmatrix} u_0 \\ v_0 \end{pmatrix} \right\}$$

$$E_-^u(x_0, \lambda) = \left\{ \begin{pmatrix} u_0 \\ v_0 \end{pmatrix} \in \mathbb{C}^{2n} \mid \begin{pmatrix} u \\ v \end{pmatrix}(x) \rightarrow 0 \text{ as } x \rightarrow -\infty \text{ where } \begin{pmatrix} u \\ v \end{pmatrix}(x_0) = \begin{pmatrix} u_0 \\ v_0 \end{pmatrix} \right\}$$

Note Due to hyperbolicity of the asymptotic equation, the only bounded solutions of (7) are, in fact, exponentially decaying as $|x| \rightarrow \infty$ and can therefore be found by seeking nontrivial intersections

$$(8) \quad E_-^u(0, \lambda) \cap E_+^s(0, \lambda) \neq \{0\}$$

Evans function Consider $\lambda \in \Omega$, then $E_-^u(0, \lambda)$ and $E_+^s(0, \lambda)$ depend analytically on λ .

→ Choose analytic bases $\{v_j^u(\lambda)\}_{j=1, \dots, k}$ and $\{v_j^s(\lambda)\}_{j=1, \dots, n-k}$ of $E_-^u(0, \lambda)$ and $E_+^s(0, \lambda)$, respectively.

→ Evans function :

$$D(\lambda) = \det [v_1^u(\lambda), \dots, v_k^u(\lambda), v_1^s(\lambda), \dots, v_{n-k}^s(\lambda)]$$

analytic in $\lambda \in \Omega$.

Theorem Let Ω be a connected component of $\mathbb{C} \setminus \text{spec } \mathcal{L}[U_0]$, then

$$(9) \quad \mathcal{D}(\lambda) = 0 \iff E_-^u(0, \lambda) \cap E_+^s(0, \lambda) \neq \{0\} \iff \lambda \text{ is an eigenvalue}$$

Furthermore, either

(i) $\Omega \cap \text{spec } \mathcal{L}^*$ is a discrete set of isolated eigenvalues with finite multiplicity: these eigenvalues correspond to roots of $\mathcal{D}(\lambda)$ and

$$(10) \quad \text{Order of roots of } \mathcal{D} = \text{PDE multiplicity of eigenvalues}$$

or else

(ii) $\Omega \subset \text{spec } \mathcal{L}^*$ and each $\lambda \in \Omega$ is an eigenvalue: $\mathcal{D}(\lambda) \equiv 0$ in Ω

For Ω_∞ , option (i) is the only possibility. □

Remark Case (ii) is ungeneric, and I do not know of any PDE example where it occurs.

Idea of Proof

We already proved (9). If $\mathcal{D}(\lambda) \equiv 0$ in Ω , then (ii) occurs. Otherwise, $\mathcal{D}(\lambda)$ has only a discrete set of roots, with finite order. I will not prove (10) \rightarrow see (9) for a sketch of proof.

It remains to show that $\lambda \notin \text{spec } \mathcal{L}^*$ when $\mathcal{D}(\lambda) \neq 0$. In this case, $E_-^u(x, \lambda) \oplus E_+^s(x, \lambda) = \mathbb{C}^{2n} \quad \forall x \in \mathbb{R}$. I claim that

$$(\mathcal{L}^* - \lambda)u = h$$

has a unique solution u with $\|u\|_{\mathcal{L}^*} \leq K \|h\|_X$.

Indeed, we have an x -dependent splitting into stable and unstable direction, given by $E_+^s(x, \lambda)$ and $E_-^u(x, \lambda)$, for $x \in \mathbb{R} \rightarrow$

$$\begin{pmatrix} u \\ v \end{pmatrix}(x) = \underbrace{\int_{-\infty}^x \Phi(x, y) P^s(y, \lambda) \begin{pmatrix} 0 \\ h(y) \end{pmatrix} dy}_{\text{evolution of } \begin{pmatrix} u \\ v \end{pmatrix}_x = A(x, \lambda) \begin{pmatrix} u \\ v \end{pmatrix}} + \underbrace{\int_{\infty}^x \Phi(x, y) P^u(y, \lambda) \begin{pmatrix} 0 \\ h(y) \end{pmatrix} dy}_{\text{Projections associated with } E_-^u(y, \lambda) \oplus E_+^s(y, \lambda) = \mathbb{C}^{2n}}$$

□

Fronts

$$\begin{pmatrix} u \\ v \end{pmatrix}_x = A(x, \lambda) \begin{pmatrix} u \\ v \end{pmatrix} \quad \text{and}$$

$$A(x, \lambda) \rightarrow A_{\pm}(\lambda) \quad \text{as } x \rightarrow \pm\infty$$

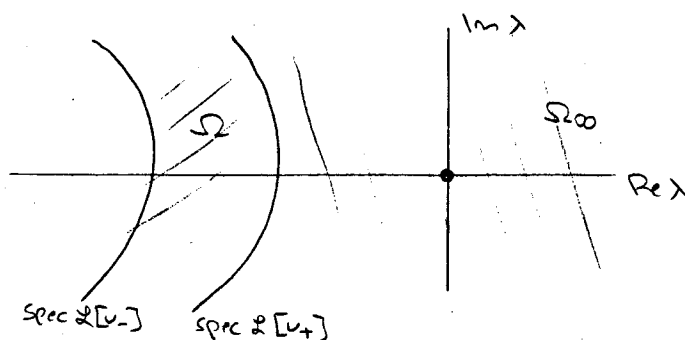
↑ associated with rest states u_{\pm} at $x = \pm\infty$.

Fix connected component Ω of $\mathbb{C} \setminus \text{spec } \mathcal{L}[u_-] \cup \text{spec } \mathcal{L}[u_+]$,

for $\lambda \in \Omega$, $A_+(\lambda)$ and $A_-(\lambda)$ are hyperbolic \rightarrow

Morse index: $i_{\pm}(\lambda) = \dim E_{\pm}^u(\lambda)$: dimension of unstable eigenspace of $A_{\pm}(\lambda)$.

- \rightarrow
- $i_{\pm}(\lambda)$ is constant for $\lambda \in \Omega$
 - $\lambda \in \Omega_{\infty}$ implies $i_+(\lambda) = i_-(\lambda) = n$.



(i) $i_+(\lambda) = i_-(\lambda) \quad \forall \lambda \in \Omega \rightarrow$ same situation as for pulses

Evans function well defined and analytic in $\lambda \rightarrow$

- either discrete set of isolated eigenvalues
- or else $\Omega \subset \text{spec } \mathcal{L}_*$ because $\mathcal{D}(\lambda) \equiv 0$.

(ii) $i_-(\lambda) > i_+(\lambda) \quad \forall \lambda \in \Omega \Rightarrow \Omega \subset \text{spec } \mathcal{L}_*$ consists entirely of eigenvalues:

$$\begin{cases} \dim E_-^u(0, \lambda) = \dim E_-^u(\lambda) = i_-(\lambda) \\ \dim E_+^s(0, \lambda) = \dim E_+^s(\lambda) = 2n - i_+(\lambda) \end{cases}$$

$$\begin{aligned} \text{sum of dimensions } -2n &= i_-(\lambda) + (2n - i_+(\lambda)) - 2n \\ &= i_-(\lambda) - i_+(\lambda) > 0 \end{aligned}$$

\rightarrow subspaces have intersection of dimension at least $i_-(\lambda) - i_+(\lambda) > 0$

$$\rightarrow E_-^u(0, \lambda) \cap E_+^s(0, \lambda) \neq \{0\}$$

$$\rightarrow \dim N(\mathcal{L}_* - \lambda) = \dim \{u \in X / (\mathcal{L}_* - \lambda)u = 0\} \geq i_-(\lambda) - i_+(\lambda)$$

$$(iii) \quad i_-(\lambda) < i_+(\lambda) \quad \forall \lambda \in \Omega \quad \Rightarrow \quad \Omega \subset \text{spec } \mathcal{L}_*$$

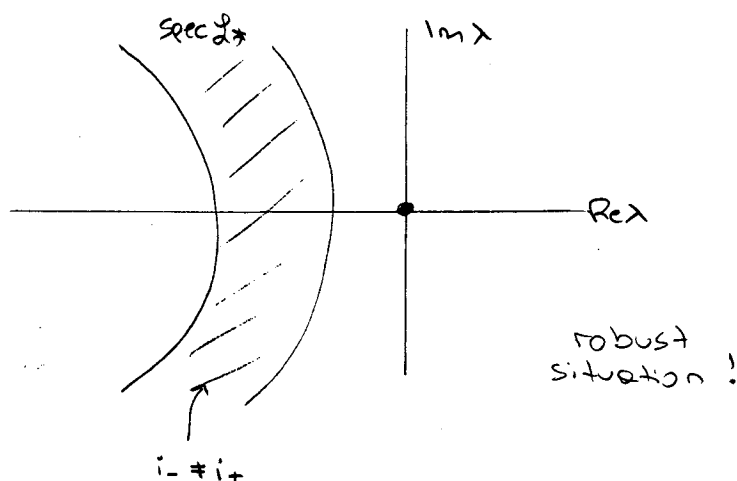
In fact, $(\mathcal{L}_* - \lambda)$ is not onto, and we have

$$\begin{aligned} \text{codim } \mathcal{R}(\mathcal{L}_* - \lambda) &= \text{codim } \{h \in X \mid \exists u \in X: (\mathcal{L}_* - \lambda)u = h\} \\ &\geq i_+(\lambda) - i_-(\lambda) \end{aligned}$$

The reason is that

$$\dim E_-^u(0, \lambda) + \dim E_+^s(0, \lambda) = 2n + i_-(\lambda) - i_+(\lambda) < 2n$$

→ do not have splitting into stable and unstable subspaces for $\lambda \in \mathbb{R}$



§2.4 Comments

• Spectrum in sector :

$$Du_{xx} + c_* u_x + (f_v(u_*(x)) - \lambda) u = 0$$

$$\lambda = \frac{e^{i\varphi}}{\varepsilon^2} \quad \rightarrow$$

$$Du_{xx} + c_* u_x + \left(f_v(u_*(x)) - \frac{e^{i\varphi}}{\varepsilon^2}\right) u = 0$$

$$x = \varepsilon y \quad \text{so that} \quad \frac{d}{dx} \rightarrow \frac{1}{\varepsilon} \frac{d}{dy} :$$

$$\frac{1}{\varepsilon^2} Du_{yy} + \frac{c_*}{\varepsilon} u_y + \left(f_v(u_*(\varepsilon y)) - \frac{e^{i\varphi}}{\varepsilon^2}\right) u = 0$$

$$\underbrace{Du_{yy} - e^{i\varphi} u}_{\text{constant-coefficients}} + \underbrace{\varepsilon c_* u_y + \varepsilon^2 f_v(u_*(\varepsilon y)) u}_{\text{small as } \varepsilon \rightarrow 0} = 0$$

$$\text{hyperbolic} \rightarrow \begin{pmatrix} u \\ v \end{pmatrix}_y = \underbrace{\begin{pmatrix} 0 & 1 \\ D^{-1}e^{i\varphi} & 0 \end{pmatrix}}_{\text{hyperbolic}} \begin{pmatrix} u \\ v \end{pmatrix}$$

→ no bounded nontrivial solutions for $0 < \varepsilon \ll 1$

- „Order of roots of $\mathcal{D}(\lambda)$ = algebraic PDE multiplicity“:

idea

$$(\mathcal{L}_* - \lambda) u_0 = 0$$

$$(\mathcal{L}_* - \lambda) u_1 = u_0$$

...

$$(\mathcal{L}_* - \lambda) u_j = u_{j-1}$$

Jordan chain
at $\lambda = \lambda_*$

Suppose $(\mathcal{L}_* - \lambda_*) u_0 = 0$ has nontrivial soln. u_0

We need to see whether $(\mathcal{L}_* - \lambda_*) u_1 = u_0$ has a soln. u_1 .

→ Solve $(\mathcal{L}_* - \lambda) u_0 = 0$ $\left\{ \begin{array}{l} \text{for } u_0(\lambda), u_1(\lambda), \text{ for } \lambda \approx \lambda_* \\ \text{separately on } \mathbb{R}^+ \text{ and } \mathbb{R}^- \\ \text{and look at jumps at } x=0 \end{array} \right.$

key $\frac{d}{d\lambda} [(\mathcal{L}_* - \lambda) u_0(\lambda) = 0]$ gives

$$(\mathcal{L}_* - \lambda) \underbrace{\partial_\lambda u_0(\lambda)}_{\text{should be } u_1} = u_0(\lambda)$$

→ derivatives of ODE solutions with respect to λ
are candidates for Jordan-chain solutions

→ this allows us (after much more algebra) to
relate derivatives of $\mathcal{D}(\lambda)$ to the length of
Jordan chains.

§3 Outlook and Guide to Literature

Nonlinear stability

(i) "Spectral stability \Rightarrow Linear stability"

We wish to see, whether $\text{spec } \mathcal{L}^* \subset \text{Re } \lambda < 0$ implies that solutions to

$$u_t = \mathcal{L}^* u, \quad u(t) = e^{\mathcal{L}^* t} u_0$$

decay to zero.

Spectral Mapping Theorem Under appropriate assumptions on \mathcal{L}^* ,

$$e^{\text{spec}(\mathcal{L}^*)t} = \text{spec}(e^{\mathcal{L}^* t})$$

[Pozny: §2.2], [Lunardi: Cor. 2.3.7].

(ii) "spectral (in)stability \Rightarrow nonlinear (in)stability"

dissipative systems (reaction-diffusion, even-order PDEs):

nonlinear stability: [Henry: Thm 6.2.1], [Henry: §5.1 exc.6]

nonlinear instability: [Henry: Thm. 5.1.5]

\mathcal{E}^0 -PDEs (some diffusion coefficients vanish)

[Bates & Jones: Dynamics Reported, 1989]

Hamiltonian PDEs:

[Grillakis, Shatah, Strauss: J. Funct. Anal. 1987+1990]

Merle, Weinstein

Conservation Laws:

[Zumbrun: In "Handbook of mathematical fluid dynamics III"
Elsevier, 2004]

Spectra

[Sondstedt]

[Henry: §5.4 + appendix to §5.4]

Multidimensional PDEs and Lattice equations

[Volpert³]

Evans function: [Deng & Nii: JDE 225 (2006)]

[Sondstedt & Scheel: Math. Nachr. 232 (2001)]

Exponential dichotomies:

[Mallet-Paret & Verduyn Lunel: JDE to appear]

[Härterich, S., Scheel: Indiana Univ. Math. J. 51 (2002)]

Bifurcations

n-pulses [Sondstedt: §5.2]

Center manifolds [Henry: Thm. 6.2.1]

References: Stability and bifurcations of travelling waves

- D. Henry. *Geometric theory of semilinear parabolic equations*. Lecture Notes **84**. Springer, 1981, 1989.
- A. Lunardi. *Analytic semigroups and optimal regularity in parabolic systems*. Birkhäuser 1995.
- A. Pazy. *Semigroups of linear operators and applications to partial differential operators*. Springer, 1983.
- A.I. Volpert, V.I. Volpert and V.I. Volpert. *Traveling wave solutions of parabolic systems*. AMS Monographs **140**, 1994. [Free download from AMS].
- B. Sandstede. "Stability of travelling waves". In: *Handbook of Dynamical Systems II* (B Fiedler, ed.). North-Holland (2002) 983-1055.

