

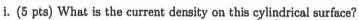
1. (25 pts) According to Ampere's Law, the magnetic H field generated by a cylindrically symmetric current distribution along the z-axis points in the azimuthal $(\tilde{\phi})$ direction, and its magnitude varies with distance r from the z-axis in terms of the $+\hat{z}$ -directed current I_{encl} passing through a circular contour centered on the z-axis:

$$H_{\phi}(r) = \frac{I_{encl}}{2\pi r} = \frac{B_d}{M}$$

- a) PART ONE: Consider a volumetric current density $J = 3\hat{z} \text{ A/m}^2$ embedded in free space which is symmetric around the z-axis and uniformly distributed across a cross sectional area of radius
 - i. (3 pts) In terms of cartesian coordinate unit directions, what is B at (x, y, z) = (0, 2, 2)? r = 2 (distance along $\frac{2}{3}$ is invelovant in cylindrical coords) $I_{end} = \int_{0}^{2} \vec{J} \cdot d\vec{s} = \int_{0}^{2} 3 ds = 3 \left(\pi a^{2} \right) = I_{total}$ $\vec{B} = -\hat{x} \mu_0 H_0 = -\hat{x} \mu_0 \frac{3\pi a^2}{2\pi (2)} = -\hat{x} \frac{3}{4} \mu_0 a^2 \left[\frac{Wb}{m^2} \right]$
 - ii. (3 pts) In terms of cartesian coordinate unit directions, what is **B** at $(x, y, z) = (-\frac{1}{2}, 0, 0)$?

$$r = \frac{1}{2}$$
 $\frac{1}{2}$
 $r = \frac{1}{2}$ $\frac{1}{2}$
 $r = \frac{1}{2}$
 $r = \frac{1}{$

b) PART TWO: Now consider that the same total current I (in A) is instead carried as a surface current on a hollow cylindrical shell having the same radius a = 1 m as the thick wire described in PART ONE above.



i. (5 pts) What is the current density on this cylindrical surface? Fortal current
$$I = 3(\pi a^2)[A]$$
 is spread around the circumference of the Shell = $2\pi a$.

So $J_s = \frac{2}{3} \frac{3\pi a^2}{2\pi a} = \frac{2}{3} \frac{3}{2} a \left[\frac{A}{m}\right]$

ii. (5 pts) How would your answers in PART ONE above change for this new current configuration?

Outside the shell, I end is the same. so No CHANGE in B(0,2,2)

Inside the shell,
$$Jend = 0$$
.
So $\vec{B}(\frac{1}{2},0,0) = 0$ Field components

cancel by symmetry

- 2. (30 pts) Short answer questions, each is independent of the others.
 - a) (5 pts) A surface current density $J = f(t)\hat{z}$ on the y = 0 plane generates a TEM wave that propagates into free space on both sides. The following expression is NOT the wave equation for the plane wave just described:

$$\nabla^2 \mathbf{E} = \frac{\partial^2 E_x}{\partial y^2} + \frac{\partial^2 E_x}{\partial z^2} = \frac{\partial^2 E_x}{\partial t^2}$$

Identify and explain the errors in this wave equation and write the correct expression below.

Source @ $y=0 \Rightarrow \hat{v}_p = \pm \hat{y}$ so \vec{E} only depends on $y \cdot \frac{\partial}{\partial z} \rightarrow \vec{p}$ \vec{J} along $\hat{z} \Rightarrow \vec{E}$ along \hat{z} so $\vec{E} = E_z(y, \pm)$. Not E_x .

$$\nabla^2 \vec{E} = \frac{\partial^2 E_7}{\partial y^2} = \frac{\partial^2 E_7}{\partial t^2}$$

- b) (5 pts) Which statement best explains the boundary condition between two perfect dielectric media having different magnetic permeabilities μ_1 and μ_2 ?
 - X. All components of the B and H fields are continuous across the boundary.
 - ii. The normal and the tangential components of the **B** field are continuous across the boundary.
 - The normal and the tangential components of the H field are continuous across the boundary.

no free

charges to carry current at the

boundary: J = 0

- The normal component of B and the tangential components of H are continuous across the boundary.
 - ✓. Only the normal component of B is continuous across the boundary.
 - yf. None of the above.

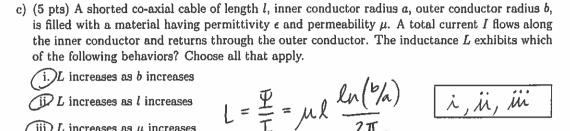
boundary condition equations:

$$\hat{N} \cdot (\hat{B}^{\dagger} - \hat{B}^{-}) = 0$$
 so $\hat{B}_{n}^{\dagger} = \hat{B}_{n}^{-}$

normal B continuous

tangential H continuous.

Since
$$\overline{H} = \overline{B}/\mu$$
, normal \overline{H} and tangential \overline{B} will be discontinuous since $\mu + \mu$



iii)
$$L$$
 increases as μ increases \mathcal{K} . L increases as I increases

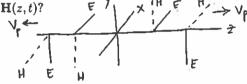
d) (5 pts) If $J = 2\sin(z)\hat{z}$ A/m², then the charge density at the origin is

$$-\frac{\partial f}{\partial t} = \nabla \cdot \vec{J} = \frac{\partial J_z}{\partial \vec{z}} = 2\cos(z)$$

at
$$z=0$$
 $\frac{\partial f}{\partial t}=-2$

e) (5 pts) A TEM wave propagating in free space is known to have the following electric field waveform: $\mathbf{E}(z,t) = u(t\pm\frac{z}{c})\hat{x} - u(t\pm\frac{z}{c})\hat{y}$, where u(t) is the unit step function. What is the waveform of the associated magnetic field $\mathbf{H}(z,t)$? $|E| = \eta_o \qquad \hat{e} \times \hat{h} = \hat{V}_p$

$$\frac{|E|}{|H|} = \eta_0 \quad \hat{e} \times \hat{h} =$$



$$\vec{H}(z,t) = \frac{1}{\eta_0} u(t + \vec{z})(\pm \hat{y}) + \frac{1}{\eta_0} u(t + \vec{z})(\pm \hat{x})$$

f) (5 pts) A pair of infinite current sheets are positioned on the x = -2 m plane and x = +2 m plane and carry equal and opposite current density of constant magnitude J_0 . The region between the sheets is filled with a perfect dielectric material having permittivity $9\epsilon_0$ and permeability $4\mu_0$. The magnetic field in the region between the current sheets is known to be $\mathbf{B} = 8\mu_0 \hat{z} \text{ Wb/m}^2$. What is the magnitude and direction of the surface current on the x = -2 m plane? You are not required to derive the field strength of a single current sheet using Ampere's Law to answer direction = -y

$$\vec{B} = 0$$

Nere

 $\langle s \leftarrow | 0^B \rangle$

$$|B| = M|J_s| = 4 M_0 J_0 = 8 M_0$$

 $S_0 J_0 = 2 A/m$

$$|\vec{H}| = \frac{|J_s|}{2} \text{ due to one sheet.}$$

$$-or - \hat{n} \times (\vec{H}^+ - \vec{H}^-) = \vec{J}_s = -\hat{\chi} \times (0 - |B|/\hat{z}) = -\hat{y}J_o$$

$$-\hat{y}|B|_{L} = -\hat{y}J_o \Rightarrow J_o = 2^{A/m}$$

Fall 2017 329 EXAM 2 3. Conducting wive shown as: Embedded in $\overline{B} = 5y\cos(\omega t)\hat{x} + 5x\sin(\omega t)\hat{y} + 10\hat{e}\hat{z}$ Wb/m2 The position of the loop and orientation of ds (RH rule of C) is critical to answers for parts (a)-(c). Full credit was given to the following answers: · CASE 1: here, ds = + ŷdxdz only By component contributes to magnetic flux through the loop a) At t=0, By = $5 \times \sin(wt)$ is negative $(-1 < \times < 0)$ and increasing in magnitude d/sin(wt) >0 Induced current will counteract this change in flux by generating B' with a positive By component. This current must flow along & => TRUE. (Therefore, I expect derived & > 0 below) b) $\underline{\mathbf{F}}(t) = \int \vec{\mathbf{B}} \cdot d\hat{\mathbf{s}} = \iint \mathbf{B}_{\mathbf{y}} \hat{\mathbf{y}} \cdot \hat{\mathbf{y}} d\mathbf{x} d\mathbf{z}$ $\int \hat{\mathbf{y}} \cdot \hat{\mathbf{y}} = +1$ = $\int dz \int 5 \times \sin(\omega t) dx$

= (1)(5 sin(wt)) $\left[\frac{1}{2}x^2\right]_{-1}$

 $=-\frac{5}{2}\sin(wt)$ [wb]

c)
$$\varepsilon = -\frac{d\Psi}{dt} = -\frac{d}{dt}(-\frac{5}{2}\sin(wt)) = \frac{5}{2}w\cos(wt)$$
 [v] (indeed $\varepsilon > 0$ at $t = 0$)

· CASE 2: NOTE: by moving
$$\hat{x}$$
, coordinate

System is now left-handed, so

this solution is physically wrong.

a) loop $ds = +\hat{y}dxdz$ as before. but now By is positive and increasing at t=0.

=> current will how opposite to C to viduce B, < 0 (lenz law)

b)
$$\Psi(t) = \int_{0}^{t} \left(B_{y} \hat{y} \cdot \hat{y} dx dz \right) = \int_{0}^{t} dz \int_{0}^{t} 5 \times \sin(\omega t) dx$$

$$=+\frac{5}{2}\sin(\omega t)$$
 [Wb]

c)
$$\varepsilon = \frac{-dI}{dt} = \frac{-\frac{5}{2}\omega\cos(\omega t)}{\left(v\right)} \left(v\right) \left(\varepsilon\cos(\omega t) + v\right)$$

Consistent $\psi(z)$

For RH coordinate system, reverse & and ŷ.

a) now ds = -ijdxdq. By is shill positive f increasing at t=0, so current must generate regardle By. \Rightarrow along $C \Rightarrow TRUE$

b)
$$\Psi(t) = \iint By\hat{y} \cdot (-\hat{y}) dxdz$$

$$= -\frac{5}{2} \sin(\omega t) [\omega b].$$

d) Regardless of coordinate system, consider changing material around loop to $\mu = 1.5 \, \mu_0$. Assuming \vec{B} is the same as above, $\Xi(t)$ stays the same, and $-d\vec{\Psi} = \varepsilon$ also stays the same.

So current direction stays the same. Induced current magnitude I = E/R stays the same. Note! I is driven by an electric field in the loop frame: $E = \oint \vec{E} \cdot d\vec{l}$, generates \vec{B} induced that mil change magnitude in the magnetic material.

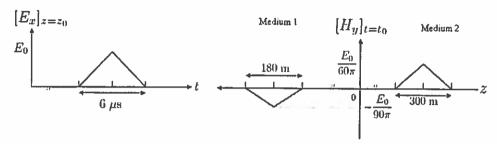
e) One test for \vec{B} natisfy is whether it satisfies maxwellegn $\nabla \cdot \vec{B} = 2Bx + 2By + 2Bz = 0$. This field does.

It is also realistic to have temporal variation that is co-sinusoidal in one direction and decaying in another - fields from different sources superpose.

However, this field strength $\rightarrow \infty$ for increasing distances from origin. This is not physically realistic unless the expression for B is only valid locally.

4. (25 pts) An infinite plane current sheet of uniform density $J_s = -J_s(t)\hat{x}$ A/m is sandwiched on the z=0 plane between two perfect dielectric media, having material parameters ϵ_1 and μ_1 in Region 1 (on the $-\hat{z}$ side) and ϵ_2 and μ_2 in Region 2 (on the $+\hat{z}$ side). $J_s(t)$ is a triangular pulse which begins at t=0 and lasts a duration of 6 μ s.

The figure below depicts the associated wavefield $E_x(z_0, t)$ at some location $z = z_0 > 0$ and the associated wavefield $H_y(z, t_0)$ at some time $t = t_0 > 0$.



Based on the information in the figures, answer the following questions:

a) (4 pts) What are the intrinsic impedances η_1 and η_2 in medium 1 and medium 2, respectively?

$$\frac{|E|}{|H|} = \eta \quad \text{by vispection:} \quad |H_1| = \frac{|E_0|}{90\pi} \quad \text{so} \quad \eta_1 = 90\pi \quad \Omega$$

$$|H_2| = \frac{|E_0|}{60\pi} \quad \text{so} \quad \eta_2 = 60\pi \quad \Omega$$

b) (6 pts) What are the propagation velocities v_{p1} and v_{p2} in medium 1 and medium 2, respectively? Pulse dwarion of 6 MS Spans $\Delta \tau$ where $v_p = \frac{\Delta \tau}{6} \mu s$.

$$V_{P_1} = \frac{180 \, \text{m}}{6 \, \mu \text{s}} = 30 \, \text{m/\mu s}$$

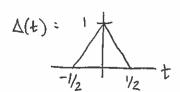
$$V_{P_2} = \frac{300 \, \text{m}}{6 \, \mu \text{s}} = 50 \, \text{m/\mu s}$$

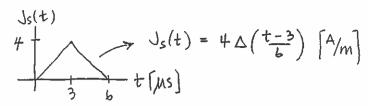
c) (3 pts) If the maximum magnitude of $J_s(t)$ is 4 A/m, what is E_0 ?

$$|E_0| = \frac{\gamma}{2} |J_0| \left(\frac{\forall}{m}\right)$$

$$= \begin{cases} \frac{\gamma_1(4)}{2} = 180\pi \left[\frac{\forall}{m}\right] = 2<0\\ \frac{\gamma_2(4)}{2} = 120\pi \left[\frac{\forall}{m}\right] = 2>0 \end{cases}$$

Can also relate $|J_{S}|$ to |H| via boundary condition eqn: $(\hat{J}_{S} = \hat{n} \times (H^{+} - H^{-})^{6})_{max} = 4 = \frac{E_{o}}{b0\pi} + \frac{E_{o}}{90\pi} \Rightarrow E_{o} = 144\pi$ but this approach assumes $E_{o}(2>0) = E_{o}(2<0)$, which is not true (Full credit given regardless)





d) (7 pts) Write the general expression for the propagating electric field E(z,t) in terms of an appropriately scaled and shifted triangular pulse function $\Delta(t)$. You should express your answer in terms of E_0 , v_{p1} , etc., rather than using the numerical values you found above.

$$\overrightarrow{E}(z,t) = \begin{cases} \int_{0}^{\infty} \frac{\eta_{1}}{2} \Delta\left(\frac{t+\frac{2}{2}/v_{p1}-3}{6}\right) (+\hat{x}) & [v/m] & z < 0 \\ \frac{J_{0} \eta_{2}}{2} \Delta\left(\frac{t-\frac{2}{2}/v_{p2}-3}{6}\right) (+\hat{x}) & [v/m] & z > 0 \end{cases}$$

e) (5 pts) Is it possible to determine the material parameters μ_1 and ϵ_1 from the information given? If so, how? If not, what other information is needed?

YES.
$$\eta = \sqrt{\frac{M_I}{E_I}} \quad V_{P_I} = \frac{1}{\sqrt{M_I E_I}}$$
So
$$\xi = \frac{1}{\gamma_I V_{P_I}}$$

$$M_I = \frac{\gamma_I}{V_{P_I}}$$