

1. For each of the following plane TEM waves in free space:

a) For plane waves propagating in free space, E and H are related by

$$H = \frac{1}{\eta_0} \hat{v} \times E,$$

where  $\hat{v}$  is the unit vector in the wave propagation direction and it is in the same direction as the Poynting vector. Therefore,

$$H_1 = \frac{1}{\eta_0} \hat{x} \times E_1 = \frac{1}{\eta_0} [\hat{z} 10 \cos(\omega t - \beta x)] \text{ (A/m)}.$$

The phasors are

$$\tilde{\mathbf{E}}_1 = \hat{y} 10 e^{-j\beta x} \text{ (V/m)},$$

$$\tilde{\mathbf{H}}_1 = \hat{z} \frac{10}{\eta_0} e^{-j\beta x} \text{ (A/m)}.$$

The instantaneous power across a  $1 \text{ m}^2$  area is simply the Poynting vector multiplied by the unit area

$$|S| \cdot 1 = |E \times H| = \left| E \times \left( \frac{1}{\eta_0} \hat{v} \times E \right) \right| = \frac{1}{\eta_0} |E|^2 \text{ (W)},$$

and we have

$$|S| \cdot 1 = \frac{1}{\eta_0} |E|^2 = \frac{1}{\eta_0} [100 \cos^2(\omega t - \beta x)] = \frac{100}{\eta_0} \cos^2(\omega t - \beta x) \text{ (W)}.$$

The time-average power across a  $1 \text{ m}^2$  area can be calculated as

$$\langle P_1 \rangle = \frac{1}{T} \int_0^T |S| \cdot 1 dt = \frac{50}{\eta_0} \text{ (W)}.$$

b) For plane waves propagating in free space, E and H are also related by

$$E = \eta_0 H \times \hat{v},$$

where  $\hat{v}$  is the unit vector in the wave propagation direction. Therefore,

$$E_2 = \eta_0 H \times (-\hat{z}) = \eta_0 [\hat{y} 5 \cos(\omega t + \beta x + \frac{\pi}{3}) - \hat{x} \sin(\omega t + \beta x - \frac{\pi}{6})] \text{ (V/m)}.$$

The phasors are

$$\tilde{\mathbf{E}}_2 = \eta_0 (\hat{y} 5 e^{j\beta z + j\frac{\pi}{3}} - \hat{x} 5 e^{j\beta z - j\frac{\pi}{2} - j\frac{\pi}{6}}) = \eta_0 (\hat{y} 5 e^{j\beta z + j\frac{\pi}{3}} - \hat{x} 5 e^{j\beta z - j\frac{2\pi}{3}}) \text{ (V/m)},$$

$$\tilde{\mathbf{H}}_2 = \hat{x} 5 e^{j\beta z + j\frac{\pi}{3}} + \hat{y} 5 e^{j\beta z - j\frac{\pi}{2} - j\frac{\pi}{6}} = \hat{x} 5 e^{j\beta z + j\frac{\pi}{3}} + \hat{y} 5 e^{j\beta z - j\frac{2\pi}{3}} \text{ (A/m)}.$$

Again, the instantaneous power across a  $1 \text{ m}^2$  area is the Poynting vector multiplied by the unit area

$$|S| \cdot 1 = |E \times H| = |(\eta_0 H \times \hat{v}) \times H| = \eta_0 |H|^2 \text{ (W)}.$$

Therefore,

$$\begin{aligned} |S_2| \cdot 1 &= \eta_0 |H_2|^2 = \eta_0 [25\cos^2(\omega t + \beta x + \frac{\pi}{3}) + 25\sin^2(\omega t + \beta x - \frac{\pi}{6})] \\ &= 50\eta_0 \cos^2(\omega t + \beta x + \frac{\pi}{3}) (W). \end{aligned}$$

Then, we obtain the time-average power across a  $1m^2$  area

$$\langle P_2 \rangle = \frac{1}{T} \int_0^T |S_2| \cdot 1 dt = 25\eta_0 (W).$$

c)

$$\begin{aligned} E_3 &= \eta_0 H_3 \times \hat{z} = \eta_0 [-\hat{y} \sin(\omega t - \beta z + \frac{\pi}{4}) - \hat{x} 2 \sin(\omega t - \beta z - \frac{3\pi}{4})] (V/m). \\ \tilde{\mathbf{E}}_3 &= \eta_0 (-\hat{y} e^{-j\beta z - j\frac{\pi}{2} + j\frac{\pi}{4}} - \hat{x} 2 e^{-j\beta z - j\frac{\pi}{2} - j\frac{3\pi}{4}}) = \eta_0 (-\hat{y} e^{-j\beta z - j\frac{\pi}{4}} - \hat{x} 2 e^{-j\beta z - j\frac{5\pi}{4}}) (V/m), \\ \tilde{\mathbf{H}}_3 &= \hat{x} e^{-j\beta z - j\frac{\pi}{2} + j\frac{\pi}{4}} - \hat{y} 2 e^{-j\beta z - j\frac{\pi}{2} - j\frac{3\pi}{4}} = \hat{x} e^{-j\beta z - j\frac{\pi}{4}} - \hat{y} 2 e^{-j\beta z - j\frac{5\pi}{4}} (A/m). \\ |S_3| \cdot 1 &= \eta_0 |H_3|^2 = \eta_0 [\sin^2(\omega t - \beta z + \frac{\pi}{4}) + 4\sin^2(\omega t - \beta z - \frac{3\pi}{4})] \\ &= 5\eta_0 \sin^2(\omega t - \beta z + \frac{\pi}{4}) (W). \\ \langle P_3 \rangle &= \frac{5}{2} \eta_0 (W). \end{aligned}$$

d)

$$\begin{aligned} H_4 &= \frac{1}{\eta_0} (-\hat{y}) \times E_4 = \frac{1}{\eta_0} [\hat{z} 2 \cos(\omega t + \beta y - \frac{\pi}{2}) + \hat{x} 2 \sin(\omega t + \beta y)] (A/m). \\ \tilde{\mathbf{E}}_4 &= \hat{x} 2 e^{+j\beta y - j\frac{\pi}{2}} - \hat{z} 2 e^{+j\beta y - j\frac{\pi}{2}} (V/m), \\ \tilde{\mathbf{H}}_4 &= \frac{1}{\eta_0} (\hat{z} 2 e^{+j\beta y - j\frac{\pi}{2}} + \hat{x} 2 e^{+j\beta y - j\frac{\pi}{2}}) (A/m). \\ |S_4| \cdot 1 &= \frac{1}{\eta_0} |E_4|^2 = \frac{1}{\eta_0} [4\cos^2(\omega t + \beta y - \frac{\pi}{2}) + 4\sin^2(\omega t + \beta y)] \\ &= \frac{8}{\eta_0} \sin^2(\omega t + \beta y) (W). \\ \langle P_4 \rangle &= \frac{4}{\eta_0} (W). \end{aligned}$$

e)

$$\begin{aligned} E_5 &= \eta_0 H_5 \times \hat{y} = \eta_0 [\hat{z} \cos(\omega t - \beta y) - \hat{x} \sin(\omega t - \beta y - \frac{\pi}{4})] (V/m). \\ \tilde{\mathbf{E}}_5 &= \eta_0 (\hat{z} e^{-j\beta y} - \hat{x} e^{-j\beta y - j\frac{\pi}{2} - j\frac{\pi}{4}}) = \eta_0 (\hat{z} e^{-j\beta y} - \hat{x} e^{-j\beta y - j\frac{3\pi}{4}}) (V/m), \\ \tilde{\mathbf{H}}_5 &= \hat{x} e^{-j\beta y} + \hat{z} e^{-j\beta y - j\frac{\pi}{2} - j\frac{\pi}{4}} = \hat{x} e^{-j\beta y} + \hat{z} e^{-j\beta y - j\frac{3\pi}{4}} (A/m). \\ |S_5| \cdot 1 &= \eta_0 |H_5|^2 = \eta_0 [\cos^2(\omega t - \beta y) + \sin^2(\omega t - \beta y - \frac{\pi}{4})] \\ &= \eta_0 \cos^2(\omega t - \beta y) + \eta_0 \cos^2(\omega t - \beta y - \frac{\pi}{4}) (W). \\ \langle P_5 \rangle &= \frac{\eta_0}{2} + \frac{\eta_0}{2} = \eta_0 (W). \end{aligned}$$

2. By comparing with the given expression and the general expression

$$\gamma\eta = j\omega\mu \text{ and } \frac{\gamma}{\eta} = \sigma + j\omega\epsilon$$

as well as

$$\mu = \frac{\gamma\eta}{j\omega}, \sigma = \operatorname{Re}\left\{\frac{\gamma}{\eta}\right\}, \epsilon = \frac{1}{\omega}\operatorname{Im}\left\{\frac{\gamma}{\eta}\right\}$$

Using these relations, for a plane wave propagating in a non-magnetic material ( $\mu = \mu_0$ ) with

$$\mathbf{H} = \hat{x}3e^y \cos(8\pi 10^6 t + \sqrt{3}y - \frac{\pi}{3}) \text{ A/m}$$

determine:

a) By comparing with the given expression and the general expression

$$\mathbf{H} = \hat{x}H_0 e^{\alpha y} \cos(\omega t + \beta y - \phi) \text{ (A/m)},$$

we find that  $\alpha = 1$  (Np/m),  $\beta = \sqrt{3}$  (rad/m),  $\gamma = \alpha + j\beta = 1 + j\sqrt{3}$ . As for the unit of  $\gamma$ , it should be a linear combination of (Np/m) and (rad/m).

b) The angular frequency  $\omega = 8\pi \times 10^6$  (rad/s), the frequency  $f = 4 \times 10^6$  (Hz), the wavelength  $\lambda = \frac{2\pi}{\beta} = \frac{2\pi}{\sqrt{3}}$  (m), and the phase velocity  $v_p = \frac{\omega}{\beta} = \frac{8\pi}{\sqrt{3}} \times 10^6$  (m/s).

c) From  $\gamma\eta = j\omega\mu$ , we have

$$\eta = \frac{j\omega\mu}{\gamma} = \frac{j \times 8\pi \times 10^6 \times 4\pi \times 10^{-7}}{1 + j\sqrt{3}} = \frac{4\pi^2}{5}(\sqrt{3} + j) \text{ (}\Omega\text{)}.$$

$$\therefore \frac{\gamma}{\eta} = \frac{1 + j\sqrt{3}}{\frac{4\pi^2}{5}(\sqrt{3} + j)} = 0.110 + j0.0633 \text{ (S/m)}.$$

$$\therefore \begin{cases} \epsilon &= \frac{1}{\omega} \operatorname{Im}\left(\frac{\gamma}{\eta}\right) = \frac{1}{8\pi \times 10^6} \times 0.0633 = 2.52 \times 10^{-9} \text{ (F/m)} \\ \sigma &= \operatorname{Re}\left(\frac{\gamma}{\eta}\right) = 0.110 \text{ (S/m)} \end{cases}$$

d)

$$\tilde{\mathbf{H}} = \hat{x}3e^{(1+j\sqrt{3})y-j\frac{\pi}{3}} \text{ (A/m)}.$$

$$\mathbf{E} = \eta \mathbf{H} \times \hat{v},$$

$$\tilde{\mathbf{E}} = \eta \tilde{\mathbf{H}} \times (-\hat{y}) = \left(\frac{8\pi^2}{5} \cdot e^{j\frac{\pi}{6}}\right) \cdot (-\hat{z}) \cdot 3e^{(1+j\sqrt{3})y-j\frac{\pi}{3}} = -\hat{z} \frac{24\pi^2}{5} e^{(1+j\sqrt{3})y-j\frac{\pi}{6}} \text{ (V/m)}.$$

e)

$$\langle \mathbf{E} \times \mathbf{H} \rangle = \frac{1}{2} \operatorname{Re}(\tilde{\mathbf{E}} \times \tilde{\mathbf{H}}^*)$$

$$= \frac{1}{2} \operatorname{Re}\left(-\hat{y} \frac{72\pi^2}{5} e^{2y+j\frac{\pi}{6}}\right)$$

$$= -\hat{y} \frac{36\pi^2}{5} e^{2y} \left(\frac{\sqrt{3}}{2}\right)$$

$$= -\hat{y} \frac{18\sqrt{3}\pi^2}{5} e^{2y} \text{ (W/m}^2\text{)}.$$

- f) For the amplitude of the electric field at a given distance  $y$  to reduce by a factor of  $1/e$  relative to its value at  $y = 0$ , we must solve the following expression for  $y$ :

$$\frac{|E(x)|}{|E(x=0)|} = \frac{\frac{24\pi^2}{5}e^y}{\frac{24\pi^2}{5}} = e^{-1}$$

so that

$$y = -1 \text{ (m)}.$$

Note that the absolute magnitude of this distance is always equal to  $1/\alpha$  and is known as the penetration depth.

- g) The total time-averaged power dissipated within a specified volume  $V$  (in W) is the volume integral of the time-averaged dissipated power density (in W/m<sup>3</sup>):

$$P_{dis} = \int_V \langle J \cdot E \rangle dV$$

Since  $J = \sigma E$ , the dissipated power density  $J \cdot E$  is easily calculated, but the volume integral of its time-average is not.

According to the integral form of the Poynting theorem (the law of electromagnetic energy conservation), the sum of the time-averaged dissipated power  $P_{dis}$  and the flux of time-averaged propagated power,  $\langle S \rangle = \langle E \times H \rangle$ , through the surface  $S$  bounding the volume  $V$  must equal zero (note that the divergence theorem is used to derive the integral form from the differential form):

$$\oint_s \langle E \times H \rangle \cdot dS + \int_V \langle J \cdot E \rangle dV = 0$$

so that

$$P_{dis} = - \oint_s \langle E \times H \rangle \cdot dS$$

Since  $\langle E \times H \rangle$  is along the  $-y$  direction, only the surfaces  $S_1 : y = 0$  and  $S_2 : y = 1$  have non-zero contributions to the total surface flux integral. Thus,

$$\begin{aligned} P_{dis} &= - \left[ \int_{s_1} \langle E \times H \rangle \cdot (-\hat{y}) dS + \int_{s_2} \langle E \times H \rangle \cdot (-\hat{y}) dS \right] \\ &= - \left[ \int_{s_1} \frac{18\sqrt{3}\pi^2}{5} dS - \int_{s_2} \frac{18\sqrt{3}\pi^2}{5} e^2 dS \right] \\ &= \frac{18\sqrt{3}\pi^2}{5} (e^2 - 1) \text{ (W)} \end{aligned}$$

Note that the value for  $P_{dis}$  is positive, indicating that, indeed, electromagnetic power is dissipated (via its conversion to kinetic energy of the free charges in the conducting material), rather than injected.

- h) Stored energy per unit volume is the stored energy density (in  $\frac{J}{m^3}$ ), which is given by  $\frac{1}{2}\epsilon E \cdot E + \frac{1}{2}\mu H \cdot H$ , such that the time rate change of this quantity  $\frac{d}{dt}[\frac{1}{2}\epsilon E \cdot E + \frac{1}{2}\mu H \cdot H]$  has units of  $\frac{W}{m^3}$  and represents stored power density. For this cosinusoidal wave, instantaneous stored energy density is proportional to  $\cos^2(\omega t + \beta x + \phi)$ , such that the instantaneous stored

power density is proportional to  $\cos(\omega t + \beta x + \phi) \sin(\omega t + \beta t + \phi)$ . The time average of this function will be 0 at all positions in space, including at (1,1,1).

Alternatively, one can express the time average of the stored power density at a given point as the time rate change of the time-averaged stored energy density, i.e. take the time average of  $[\frac{1}{2}\epsilon E \cdot E + \frac{1}{2}\mu H \cdot H]$  first and then the time derivative, which is thus the time average of a time-independent function (the time average of  $\cos^2(\omega t + \beta x + \phi)$  is  $\frac{1}{2}$ ), which also evaluates to 0.

3. We observe that

$$\frac{\sigma}{\omega\epsilon} = \frac{4}{2\pi \times 3 \times 10^4 \times 81 \times 8.85 \times 10^{-12}} \gg 1,$$

which means ocean water can be treated as a good conductor at 30 kHz.

a) In a good conductor

$$\alpha \approx \sqrt{\pi f \mu \sigma} = \sqrt{\pi \times 3 \times 10^4 \times 4\pi \times 10^{-7} \times 4} = 0.688 \text{ (Np/m)}$$

$$\beta \approx \sqrt{\pi f \mu \sigma} = \sqrt{\pi \times 3 \times 10^4 \times 4\pi \times 10^{-7} \times 4} = 0.688 \text{ (rad/m)}$$

$$|\eta| = \sqrt{\frac{\omega\mu}{\sigma}} = 0.243$$

because it is a good conductor,  $\tau = 45^\circ$

$$\eta = 0.243 \angle 45^\circ \text{ (}\Omega\text{)}$$

b) Assume the distance between the submarine and the ship is  $d$ , we want

$$e^{-\alpha d} \geq 1\% = 0.01$$

$$\therefore -\alpha d \geq \ln(0.01),$$

$$\therefore d \leq \frac{-\ln(0.01)}{\alpha} \approx 6.69 \text{ (m)}$$

c) Based on  $\beta$  the in Part a), the wavelength is

$$\lambda = \frac{2\pi}{\beta} = 9.13 \text{ (m)}$$

d) At  $f = 300 \text{ Hz}$ , since

$$\frac{\sigma}{\omega\epsilon} = \frac{4}{2\pi \times 300 \times 81 \times 8.85 \times 10^{-12}} = 2.96 \times 10^6 \gg 1,$$

we can still treat ocean water as a good conductor at 300 Hz. By using the general formula:

$$\alpha \approx \sqrt{\pi f \mu \sigma} = \sqrt{\pi \times 3 \times 10^2 \times 4\pi \times 10^{-7} \times 4} = 0.0688 \text{ (Np/m)}$$

$$\beta \approx \sqrt{\pi f \mu \sigma} = \sqrt{\pi \times 3 \times 10^2 \times 4\pi \times 10^{-7} \times 4} = 0.0688 \text{ (rad/m)}$$

$$|\eta| = \sqrt{\frac{\omega\mu}{\sigma}} = 0.0243$$

$$\eta = 0.0243 \angle 45^\circ \text{ } (\Omega)$$

$$d \leq \frac{-\ln(0.01)}{\alpha} \approx 66.9 \text{ } (m)$$

$$\lambda = \frac{2\pi}{\beta} = 91.3 \text{ } (m).$$

4.

a)

$$f = \frac{V_p \beta}{2\pi} = 3.18 \times 10^6 \text{ } (Hz)$$

b)

$$\lambda = \frac{2\pi}{\beta} = 31.4 \text{ } (m)$$

c)

$$\epsilon_r = \frac{1}{V_p^2 \epsilon_0 \mu} = \frac{9}{4} = 2.25$$

d)

$$\sigma = \frac{2\alpha}{\eta_0} \sqrt{\frac{\epsilon_r}{\mu_r}} = 3.98 \times 10^{-6} \text{ } (S/m)$$

e)

$$y = \frac{1}{\alpha} = 1000 \text{ } (m)$$

f)

$$\tau = \frac{\sigma}{2\omega\epsilon} = \frac{3.98 \times 10^{-6}}{2 \times 2\pi \times 3.18 \times 10^6 \times 2.25 \times 8.85 \times 10^{-12}} = 0.005$$

$$\begin{aligned} \tilde{\mathbf{H}} &= \frac{20}{|\eta|} e^{-(0.001+j0.2)y} e^{-j\tau} \hat{z} \\ &= \frac{1}{8\pi} e^{-(0.001+j0.2)y-j0.005} \hat{z} \text{ } (A/m). \end{aligned}$$