

Dielectric and Shear Mechanical Relaxation in Viscous Liquids: Are they Connected?

Kristine Niss¹ and Bo Jakobsen²
master students at Roskilde University
supervised by Niels Boye Olsen

<http://dirac.ruc.dk/~kniss/masterthesis/>

¹kniss@dirac.ruc.dk

²boj@dirac.ruc.dk

Methods of measurement



Dielectric: 22-layer gold platen capacitor with empty capacitance of 68pF . $10^{-3} - 10^6\text{Hz}$

Shear modulus: Piezoelectric shear modulus gauge (PSG)
[Christensen & Olsen, 1995] $10^{-3} - 10^{4.5}\text{Hz}$

Measurement: Standard equipment.
 $10^{-3} - 10^2\text{Hz}$: HP3458A multimeter in conjunction with a Keithley AWFG.

$10^2 - 10^6\text{Hz}$: HP 4284A LCR meter

Temperature: Nitrogen cooled cryostat.
Absolute temperature: better than 0.2K
Temperature stability: better than 20mK

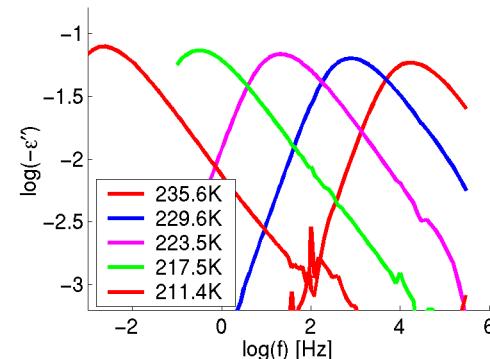
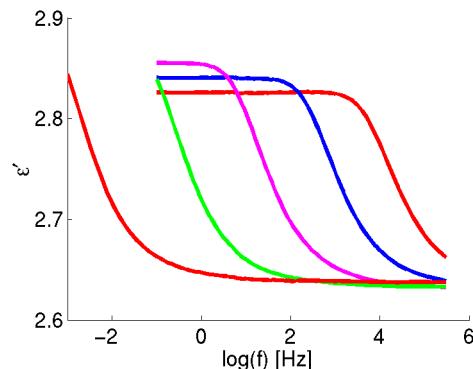
Dielectric relaxation

$$P_z e^{i(\omega t + \phi)} = \epsilon_0 \xi(\omega) E_0 e^{i\omega t} \quad \text{where} \quad \epsilon(\omega) = \xi(\omega) + 1$$

$$C(\omega) = \frac{A \epsilon_0 \epsilon(\omega)}{d} = C_0 \epsilon(\omega) \quad \text{where} \quad Q_0 e^{i(\omega t + \phi)} = C(\omega) U_0 e^{i\omega t}$$

$$P_z = (\epsilon - 1) \epsilon_0 E_m = N (E_i \alpha_i + \langle \mu_z \rangle) = N (\alpha_i E_i + \alpha_r E_d)$$

Typical spectrum (substance: DC704):



Debyes model

- Spherical non-interacting dipoles
- Surroundings can be described as a viscous continuum, with a frequency independent viscosity (η_0)
- No-slip boundary conditions
- Inertial effects are ignored

$$\frac{\partial f}{\partial t} + \nabla \cdot \mathbf{J} = 0 , \quad \mathbf{J} = -D_0 \nabla f + \mathbf{v} f$$

$$\zeta_0 = 8\pi r^3 \eta_0 , \quad \dot{\theta} = \frac{M}{\zeta_0}$$

$$\frac{\partial f}{\partial t} = \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left[\sin \theta \left(D_0 \frac{\partial f}{\partial \theta} - \frac{M}{\zeta_0} f \right) \right]$$

Microscopic DiMarzio-Bishop model

The Debye “rotational diffusion equation”:

$$\frac{\partial f}{\partial t} = \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left[\sin \theta \left(D_0 \frac{\partial f}{\partial \theta} - \frac{M}{\zeta_0} f \right) \right]$$

The generalized rotational diffusion equation by DiMarzio & Bishop [1974]:

$$\frac{\partial f}{\partial t} = \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left[\sin \theta \left(\frac{\partial}{\partial \theta} \int_{-\infty}^t D(t-\tau) f(\tau) d\tau - f \int_{-\infty}^t V(t-\tau) M(\tau) d\tau \right) \right]$$

The Stokes friction term is used

$$\zeta(\omega) = 8\pi r^3 \eta(\omega)$$

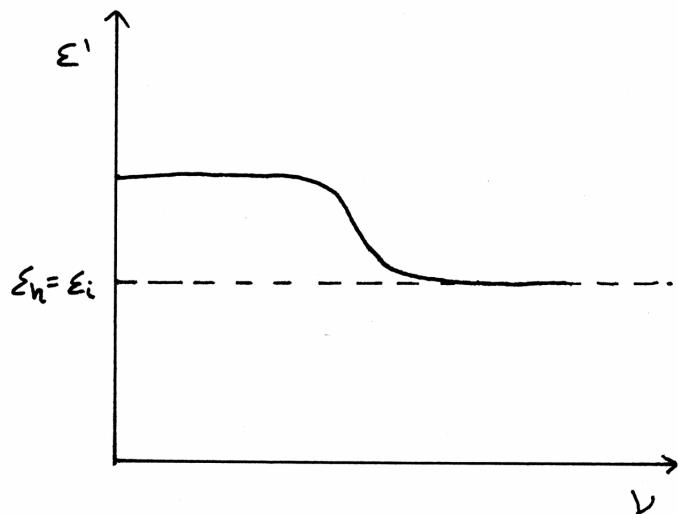
A first order solution is found

$$\alpha_r(\omega) = \frac{\mu^2}{3k_B T \left(1 + \left(\frac{4\pi r^3}{k_B T} \right) i\omega \eta(\omega) \right)} = \frac{\mu^2}{3k_B T \left(1 + \left(\frac{4\pi r^3}{k_B T} \right) G(\omega) \right)}$$

This microscopic polarizability has to be connected to macroscopic measurable quantities.

What controls the high frequency limit

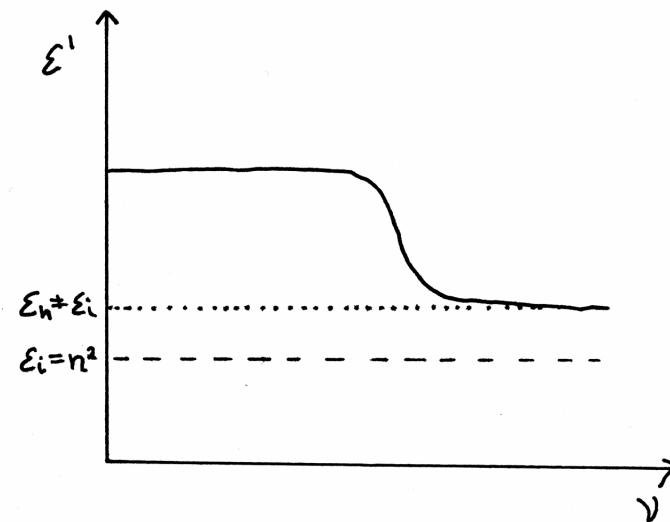
Original Debye model



$$\frac{\epsilon-1}{\epsilon+2} = \frac{N}{3\epsilon_0} \left[\alpha_i + \frac{\mu^2}{3k_B T \left(1 + \left(\frac{4\pi r^3}{k_B T} \right) i\omega \eta_0 \right)} \right]$$

$$\frac{\epsilon_h - 1}{\epsilon_h + 2} = \frac{N}{3\epsilon_0} \alpha_i$$

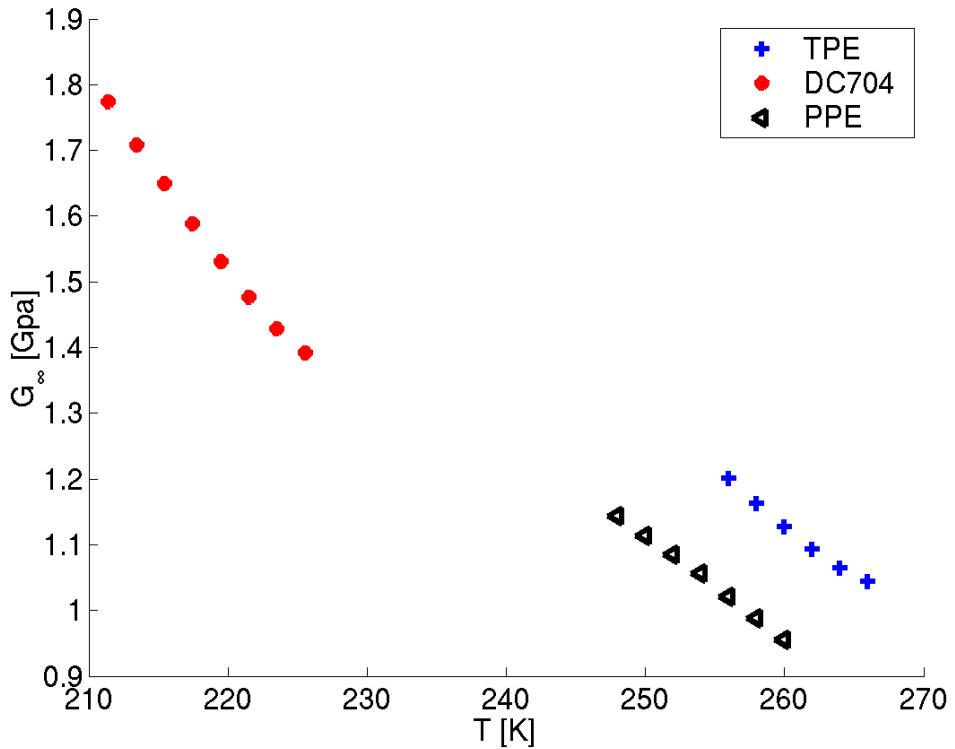
Generalized Debye model



$$\frac{\epsilon-1}{\epsilon+2} = \frac{N}{3\epsilon_0} \left[\alpha_i + \frac{\mu^2}{3k_B T \left(1 + \left(\frac{4\pi r^3}{k_B T} \right) G(\omega) \right)} \right]$$

$$\frac{\epsilon_h - 1}{\epsilon_h + 2} = \frac{N}{3\epsilon_0} \left[\alpha_i + \frac{\mu^2}{3k_B T \left(1 + \left(\frac{4\pi r^3}{k_B T} \right) G_\infty \right)} \right]$$

What controls the high frequency limit



Increasing T

Effects giving decreasing C_h

- $k_B T$ increases
- N decreases
- Spacing increases

Effect giving increasing C_h

- G_∞ decreases

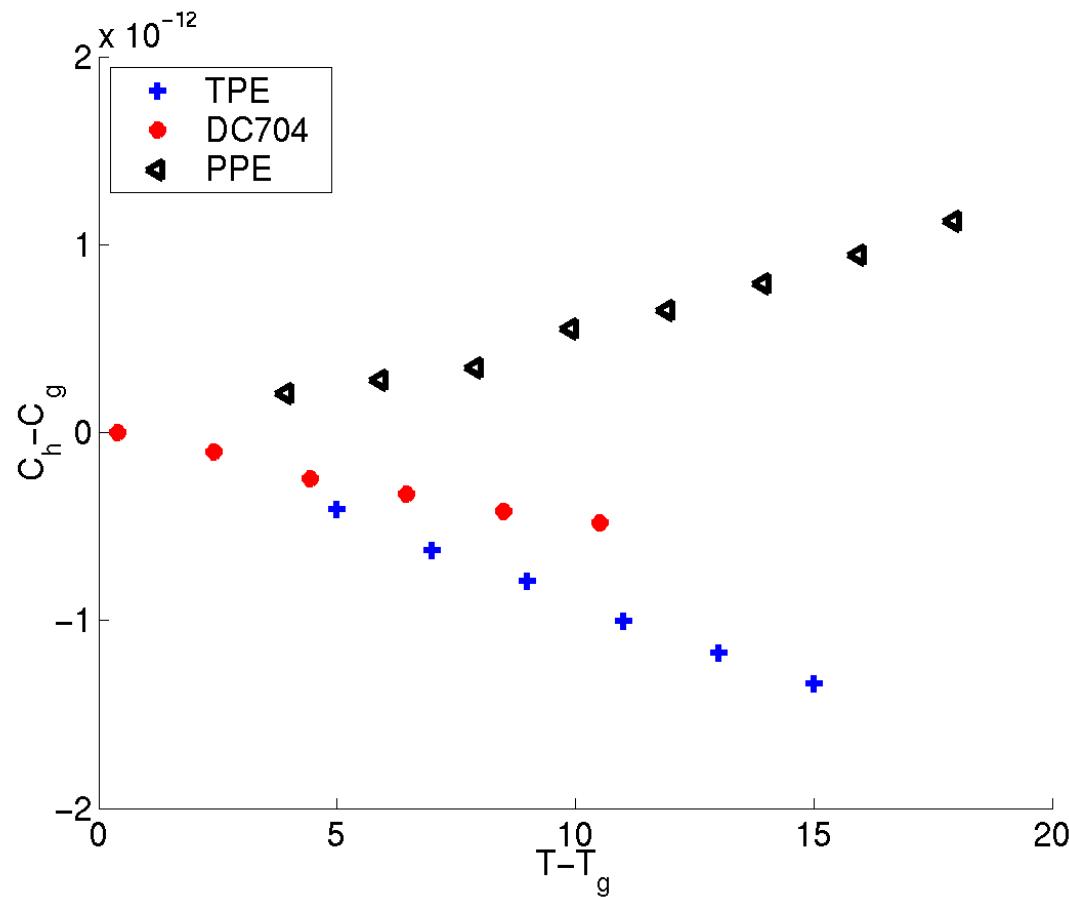
Original Debye model

$$\frac{\epsilon_h - 1}{\epsilon_h + 2} = \frac{N}{3\epsilon_0} \alpha_i$$

Generalized Debye model

$$\frac{\epsilon_h - 1}{\epsilon_h + 2} = \frac{N}{3\epsilon_0} \left[\alpha_i + \frac{\mu^2}{3k_B T \left(1 + \left(\frac{4\pi r^3}{k_B T} \right) G_\infty \right)} \right]$$

Temperature dependence C_h



Conclusions and questions

- There is an elastic contribution to ϵ_h .
- This is qualitatively in agreement with the generalized Debye model.
- Quantitative testing is difficult
- Other tests
 - The full spectrum, including the loss peak frequency.
 - The temperature dependence of fitted parameters.

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References

- Christensen, T. & Olsen, N. B. [1995]. A rheometer for the measurement of a high shear modulus covering more than seven decades of frequency below 50 kHz, *Review of scientific instruments* **66**(10): 5019.
- DiMarzio, E. A. & Bishop, M. [1974]. Connection between the macroscopic electrical and mechanical susceptibilities., *The Journal of Chemical Physics* **60**(10): 3802.