

## Black-Scholes equation.

Our next goal is to find the price of a call/put option. We will consider the call option, the function for the put option can be obtained using the Put-call parity eq-n, once the price of the call option is known.

Recall: if  $B$  is the balance in a bank at a risk-free rate  $r$ , then  $\frac{dB}{B} = r dt$ .

The price  $S$  of a stock satisfies the diff. eq-n:  $\frac{dS}{S} = \mu dt + \sigma dB(t)$ , where  $B(t)$  is the standard Brownian motion.

Rmk / defn. The parameter  $\mu$  is called the mean growth rate and  $\sigma$  is called volatility.

Let  $C(S, t)$  be the price of a call option for stock  $S$  at time  $t$ . Using Itô's formula (see page ) and eq-n (b) above, we get

$$dC(S, t) = \frac{\partial C}{\partial t} dt + \frac{\partial C}{\partial S} \cdot dS + \frac{1}{2} \frac{\partial^2 C}{\partial S^2} dS^2 =$$

$$\begin{aligned}
 &= \frac{\partial C}{\partial t} dt + \frac{\partial C}{\partial S} (\mu S dt + b S dB(t)) + \frac{\partial^2 C}{\partial S^2} \left( \frac{1}{2} b^2 S^2 dB(t)^2 + \text{O}(dt) \right) \\
 &= \left( \frac{\partial C}{\partial t} + \mu S \frac{\partial C}{\partial S} + \frac{b^2 S^2}{2} \frac{\partial^2 C}{\partial S^2} \right) dt + b S \frac{\partial C}{\partial S} dB(t) \quad (\text{I})
 \end{aligned}$$

On the other hand, we can represent  $C(S, t)$  as a linear combination of  $S$  and  $B$ , where the coefficients are functions in  $S$  and  $t$ :

$$C(S, t) = u(S, t)S - v(S, t)B$$

This expression gives

$$dC(S, t) = u dS + S du - v dB - v dB.$$

Notice that you need to borrow more money to buy more stock, i.e.  $S du = B dv$ , hence

$$dC(S, t) = u dS - v dB = u(\mu S dt + b S dB(t)) - v dB dt \quad (\text{II}).$$

Equating I and II, we get:

coefficients of  $dB(t)$ :  $u b S = b S \frac{\partial C}{\partial S}$  or, as  $b \neq 0$  and

$$S \neq 0, u = \frac{\partial C}{\partial S}$$

coefficients of  $dt$ :

$$\frac{\partial C}{\partial t} = -\frac{b^2 S^2}{2} \cdot \frac{\partial^2 C}{\partial S^2} - rS \frac{\partial C}{\partial S} + rC.$$

This equation is called the Black-Scholes differential equation.

Example. Suppose  $dS = 5dt + 6S^2 dB(t)$ . Find the corresponding Black-Scholes eq-n.

Sol-n: using Itô's formula, we get

$$\begin{aligned} dC(S, t) &= \frac{\partial C}{\partial t} dt + \frac{\partial C}{\partial S} dS + \frac{1}{2} \cdot \frac{\partial^2 C}{\partial S^2} dS^2 = \\ &= \frac{\partial C}{\partial t} dt + \frac{\partial C}{\partial S} (5dt + 6S^2 dB(t)) + \frac{1}{2} \frac{\partial^2 C}{\partial S^2} \cdot 36S^4 dB(t)^2 \\ &= \left( \frac{\partial C}{\partial t} + 5 \cdot \frac{\partial C}{\partial S} + 18S^4 \frac{\partial^2 C}{\partial S^2} \right) dt + 6S^2 \frac{\partial C}{\partial t} dB(t) \quad (\text{A}) \end{aligned}$$

Using the expression  $C(S, t) = u(S - rB)$ , obtain

$$\begin{aligned} dC(S, t) &= u dS - r dB = u(5dt + 6S^2 dB(t)) - rB dB dt = \\ &\approx (5u - rB) dt + 6S^2 dB(t). \quad (\text{**}) \end{aligned}$$

As  $(\text{**}) \approx (\text{A})$ ;

coeff-s of  $dB(t)$ :  $u = \frac{\partial C}{\partial S}$  which, using that  $r = \frac{uS - C(S, t)}{B}$ , gives  $r = \frac{\partial C}{\partial S} S - C(S, t)$ .

coeff-s of  $dt$ :  $\frac{\partial C}{\partial t} + 5 \cdot \cancel{\frac{\partial C}{\partial S}} + 18 \cdot \frac{B}{S^2} \cancel{\frac{\partial^2 C}{\partial S^2}} = \frac{5 \cdot \partial C}{\partial S} - rS \frac{\partial C}{\partial S}$

+  $rC$  or

$$\boxed{\frac{\partial C}{\partial t} + 18S^4 \cdot \frac{\partial^2 C}{\partial S^2} + rS \frac{\partial C}{\partial S} + rC = 0}$$

Fact. There is a change of variables after which the Black-Scholes equation becomes the heat equation:

$$\frac{\partial u(t, x)}{\partial t} = \frac{\partial^2 u(t, x)}{\partial x^2}.$$

Heat Equation.

Consider a thin (one-dimensional) bar. The solution of the heat eq-n identifies the temperature of the bar at location  $x$  at time  $t$ . We also need to know the initial temperature distribution (initial condition), i.e. the f-n  $u(x, 0) := \Psi(x)$ .

We need two more conditions:

$$u(L, t) = X_L(t) \text{ and } u(0, t) = X_R(t).$$



What is the solution of the heat eq-n?

Answer:  $u(x,t) = \frac{1}{\sqrt{2\pi t}} \int_{-\infty}^{\infty} \psi(s) e^{-\frac{(s-x)^2}{2t}} ds.$

Exercise: check it for  $\psi(x) = \text{const.}$

Rmk: the function  $u(x,t)$  for  $\psi(x) \equiv 1$  is the cdf for the distr-n  $N(x, t)$ .

Sol-ns of the Black-Scholes Eq-n.

Recall that the initial condition (value at time  $t=T$ ) for the call option is

$$C_E(S,T) = \max(0, S - E).$$

The sol-n is given by

$$C_E(S,t) = S \cdot \mathbb{Q}(d_1) - E e^{-r(T-t)} \mathbb{Q}(d_2), \text{ where } \\ \mathbb{Q}(x) \text{ is the cdf of } N(0,1) (\mathbb{Q}(x) = \int_{-\infty}^{x/\sqrt{t}} e^{-\frac{x^2}{2}} dx); \\ d_1 = \frac{\ln(S/E) + (T-t)(r + \frac{\sigma^2}{2})}{\sigma \sqrt{T-t}},$$

$$d_2 = \frac{\ln(S/E) + (T-t)(r - \frac{\sigma^2}{2})}{\sigma \sqrt{T-t}} = d_1 - \sigma \sqrt{T-t}.$$

## Analysis / Interpretations

① Consider the limit  $t \rightarrow T^-$  (from the left):

(a)  $S > E$ :  $\ln(S/E) > 0$ , so  $\frac{\ln(S/E)}{b\sqrt{T-t}} \rightarrow +\infty$  and

$$\frac{(T-t)(r^6/b^2)}{b\sqrt{T-t}} \rightarrow 0 \text{ (independent of the comparison of } S \text{ to } E), \text{ hence, } d_1, d_2 \rightarrow +\infty \text{ with } U(d_1) = U(d_2) \approx 0,$$

so  $C_E(S, T) \approx S - E$ .

(b)  $S < E$ :  $\ln(S/E) < 0$  and  $\frac{\ln(S/E)}{b\sqrt{T-t}} \rightarrow -\infty$  with  $d_1, d_2 \rightarrow -\infty$  and  $U(d_1) = U(d_2) \approx 0$ , so  $C_E(S, T) \approx 0$ .

(c)  $S = E$ :  $\ln(S/E) = \ln(1) = 0$  and  $d_1 = d_2 = 0$  with  $U(d_1) = U(d_2) = 1/b$ , giving  $C_E(S, T) \approx \frac{1}{2}S - \frac{1}{2}E = 0$ .

②  $S \rightarrow \infty$ :  $d_1 = d_2 \rightarrow \infty$ ,  $U(d_1) = U(d_2) \approx 1$  with  $C_E(S, T) \approx S - E = S$  ( $E$  is negligible)

③  $S \rightarrow 0$ :  $d_1 = d_2 \rightarrow -\infty$  with  $U(d_1) = U(d_2) \approx 0$  and  $C_E(S, T) \approx 0$ .

Rmk. As we already mentioned, the boundary condition for Black-Scholes eq-n for call options corresponding to  $\Psi(x)$  (for the heat eq-n) is

$$C_E(S, \Gamma) = \max(S - E, 0). \text{ This can be easily seen to be equivalent to the limit in } ①.$$

Notice that  $\Delta S = \mu S \Delta t + \sigma S \Delta B(t)$ , when  $S \rightarrow 0$  gives  $\Delta S \sim 0$ , giving the boundary condition  $C_E(0, t) \approx 0$ , which agrees with ③

①, ② and ③ form a complete set of boundary conditions and correspond to  $\Psi(x)$ ,  $\chi_{\ell}(t)$  and  $\chi_{\Gamma}(t)$  (boundary cond-ns for the heat eq-n with bar of infinite length, i.e.  $\ell = -\infty, \Gamma = +\infty$ ).

# Prices of stocks of some companies over last 5 years

