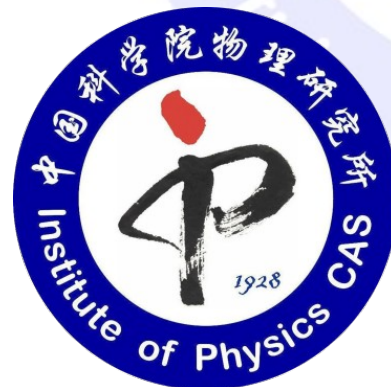


Determinantal Quantum Monte Carlo: Algorithm and Measurements

Zi Yang Meng

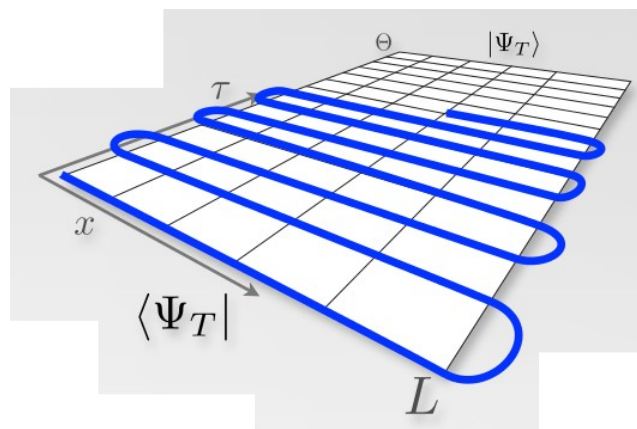
(孟子杨)

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Quantum Monte Carlo

■ Determinantal QMC for fermions

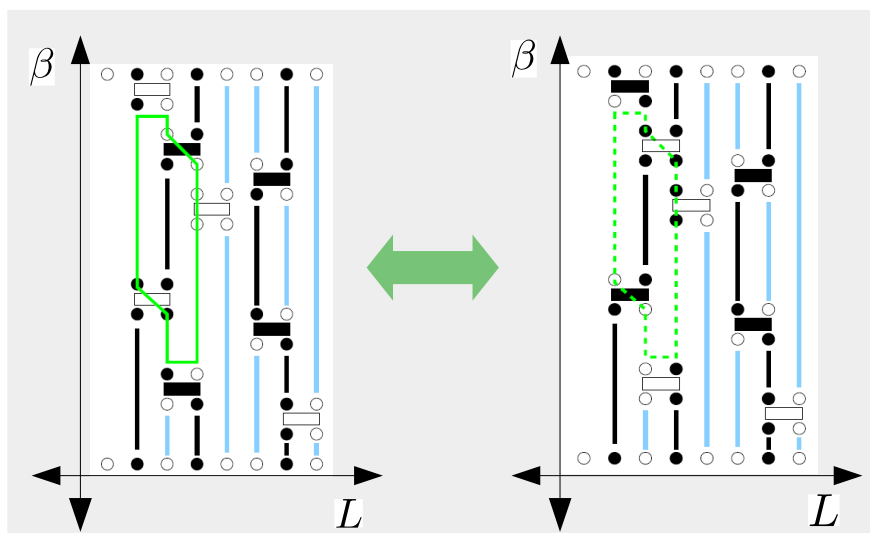


Hubbard model:

- Metal-Insulator transition
- Unconventional superconductivity
- Non-Fermi-liquid
- Interaction effects on topological state of matter

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■ World-line QMC for bosons

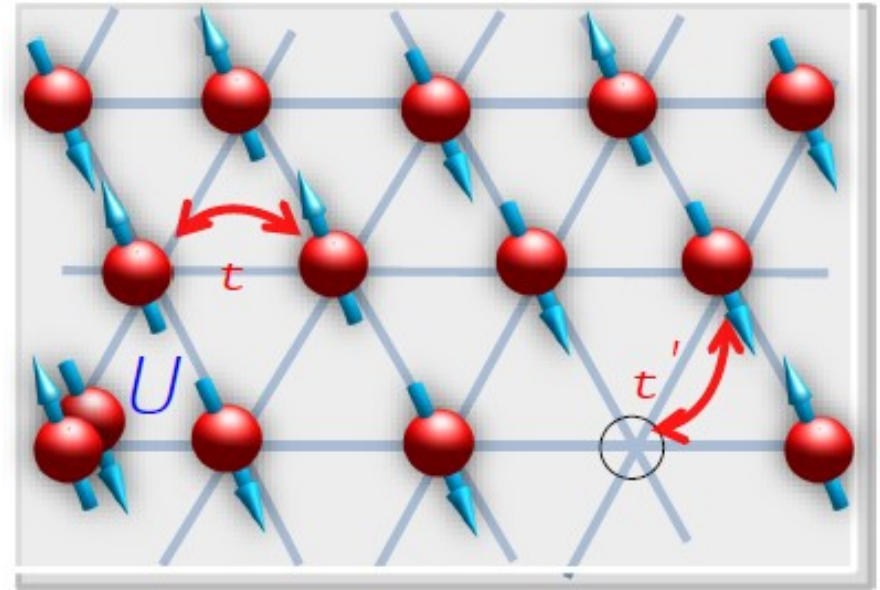
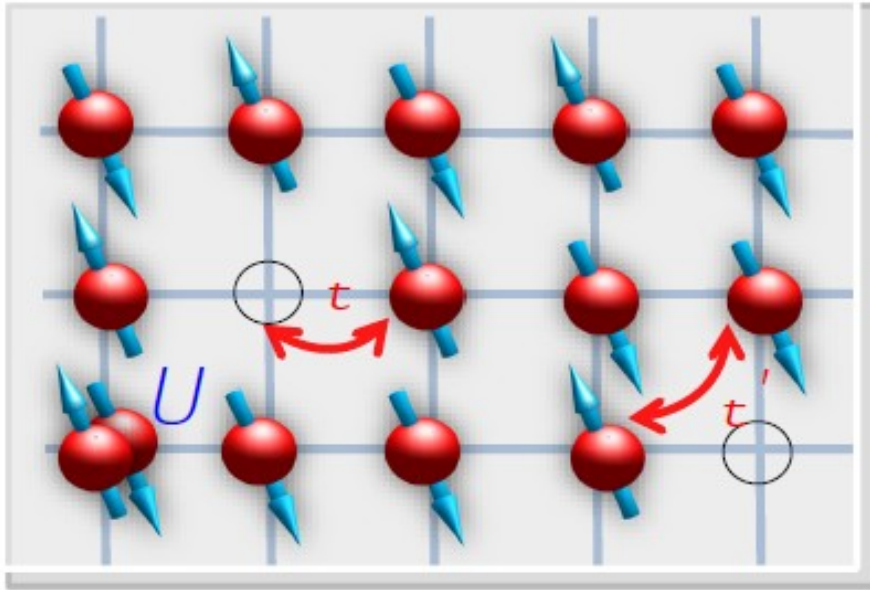


Heisenberg model:

- Quantum magnetism
- Phase transition and critical phenomena
- Quantum spin liquids
- Quantum spin ice

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Basic problem



Partition function:
$$Z = \text{Tr} [e^{-\beta(\hat{H} - \mu\hat{N})}] = \sum_n \langle n | e^{-\beta(\hat{H} - \mu\hat{N})} | n \rangle$$

Observables :
$$\langle \hat{A} \rangle = \frac{\text{Tr} [\hat{A} e^{-\beta(\hat{H} - \mu\hat{N})}]}{\text{Tr} [e^{-\beta(\hat{H} - \mu\hat{N})}]} = \frac{\sum_n \langle n | \hat{A} e^{-\beta(\hat{H} - \mu\hat{N})} | n \rangle}{\sum_n \langle n | e^{-\beta(\hat{H} - \mu\hat{N})} | n \rangle}$$

Fock space :
$$\{|n\rangle\} \sim e^{N_e \ln(2)} (2^{N_e}) \quad \text{or} \quad e^{N_e \ln(4)} (4^{N_e})$$

Determinantal quantum Monte Carlo

$$H = -t \sum_{\langle i,j \rangle, \sigma}^N (c_{i,\sigma}^\dagger c_{j,\sigma} + \text{h.c.}) + U \sum_{i=1}^N (n_{i,\uparrow} - \frac{1}{2})(n_{i,\downarrow} - \frac{1}{2})$$

■ Path-integral & Trotter-Suzuki decomposition

$$Z = \text{Tr}[e^{-\beta \hat{H}}] \approx \text{Tr}\left[\prod_{l=1}^m e^{-\Delta\tau \hat{H}_t} e^{-\Delta\tau \hat{H}_U}\right] \quad \Delta\tau = \frac{\beta}{m}, m \rightarrow \infty$$

■ Free fermion (Slater) determinant

$$\text{Tr}\left[e^{-\sum_{i,j} c_i^\dagger A_{i,j} c_j}\right] = \text{Det}[\mathbf{1} + e^{-\mathbf{A}}]$$

Determinantal quantum Monte Carlo

$$H = -t \sum_{\langle i,j \rangle, \sigma}^N (c_{i,\sigma}^\dagger c_{j,\sigma} + \text{h.c.}) + U \sum_{i=1}^N (n_{i,\uparrow} - \frac{1}{2})(n_{i,\downarrow} - \frac{1}{2})$$

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■ Discrete Hubbard-Stratonovich transformation

$$e^{-\Delta\tau U \sum_{i=1}^N (n_{i,\uparrow} - \frac{1}{2})(n_{i,\downarrow} - \frac{1}{2})} = C \sum_{s_1, \dots, s_N = \pm 1} e^{\alpha \sum_{i=1}^N s_i (n_{i,\uparrow} - n_{i,\downarrow})} \quad (C, \alpha)(U, N, \Delta\tau)$$

$$\int_{-\infty}^{\infty} dx e^{-\frac{a}{2}x^2} = \sqrt{\frac{2\pi}{a}}$$

- Blankenbecler et. al., Phys. Rev. D 24, 2278 (1981)
- Hirsch, Phys. Rev. B 28, 4059(R) (1983)
- Hirsch, Phys. Rev. B 31, 4403 (1985)
- Assaad and Evertz, Lec. Notes. In Phys. 739 (2008)

$$\int_{-\infty}^{\infty} dx e^{-\frac{a}{2}x^2 + bx} = \int_{-\infty}^{\infty} dx e^{-\frac{a}{2}(x - \frac{b}{a})^2 + \frac{b^2}{2a}} = \sqrt{\frac{2\pi}{a}} e^{\frac{b^2}{2a}}$$

Determinantal quantum Monte Carlo

$$H = -t \sum_{\langle i,j \rangle, \sigma}^N (c_{i,\sigma}^\dagger c_{j,\sigma} + \text{h.c.}) + \frac{U}{2} \sum_{i=1}^N (n_{i,\uparrow} + n_{i,\downarrow} - 1)^2$$

■ Path-integral & Trotter-Suzuki decomposition

$$Z = \text{Tr}[e^{-\beta \hat{H}}] \approx \text{Tr}\left[\prod_{l=1}^m e^{-\Delta\tau \hat{H}_t} e^{-\Delta\tau \hat{H}_U}\right] \quad \Delta\tau = \frac{\beta}{m}, m \rightarrow \infty$$

■ Discrete Hubbard-Stratonovich transformation

$$e^{-\Delta\tau \frac{U}{2} (n_{i,\uparrow} + n_{i,\downarrow} - 1)^2} = \frac{1}{4} \sum_{s_1, s_2, \dots, s_N = \pm 1, \pm 2} \gamma(s_i) e^{i\sqrt{\Delta\tau \frac{U}{2}} \eta(s_i) (n_{i,\uparrow} + n_{i,\downarrow} - 1)} + O[(\Delta\tau)^4]$$

$$\int_{-\infty}^{\infty} dx e^{-\frac{a}{2} x^2} = \sqrt{\frac{2\pi}{a}}$$

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Determinantal quantum Monte Carlo

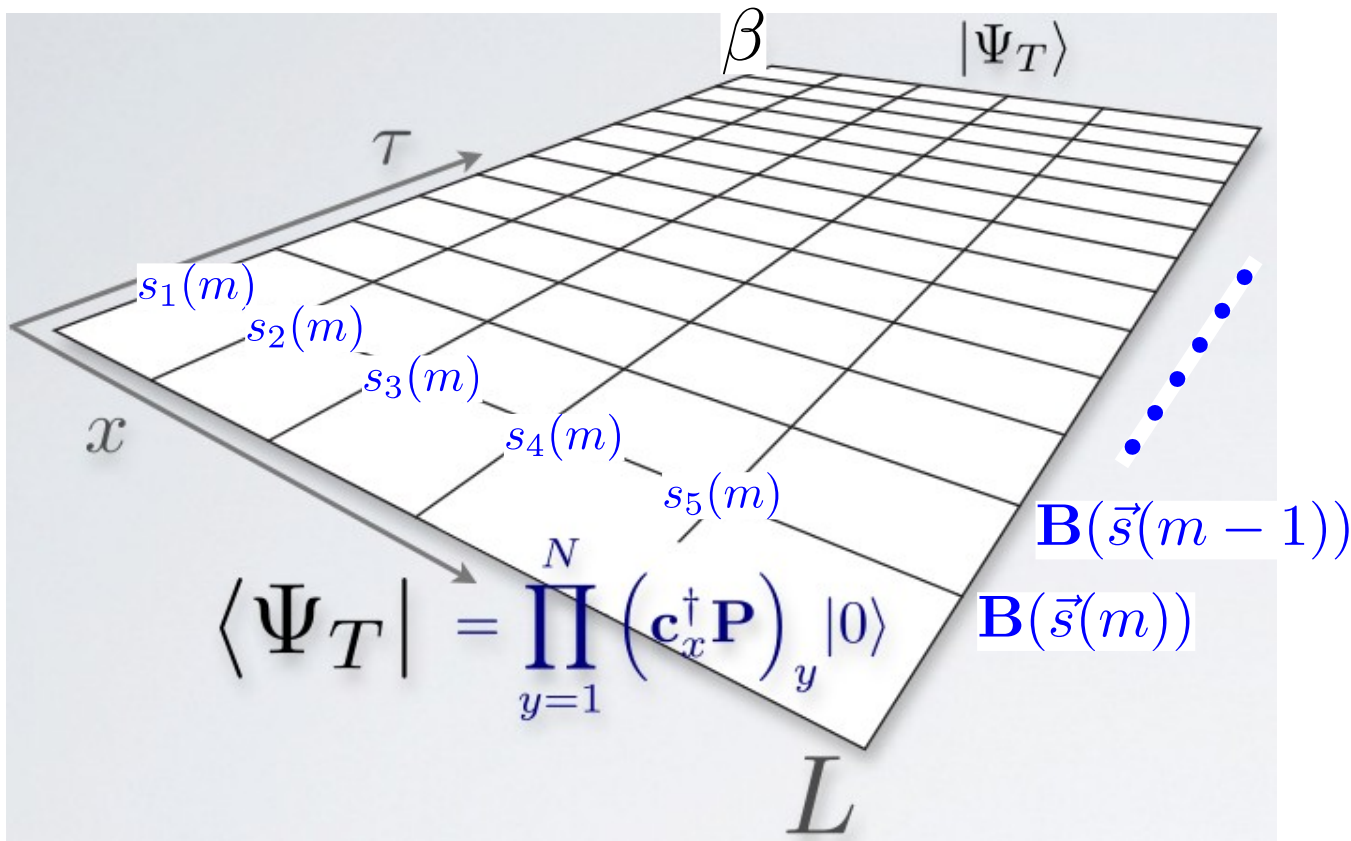
- Write Path-integral into determinant

$$Z = \text{Tr} \left[\prod_{l=1}^m e^{-\Delta\tau \hat{H}_t} e^{-\Delta\tau \hat{H}_U} \right] = C^m \sum_{\vec{s}_1, \vec{s}_2, \dots, \vec{s}_m} \det[1 + \mathbf{P}^\dagger \mathbf{B}_m \mathbf{B}_{m-1} \dots \mathbf{B}_1 \mathbf{P}]$$

$$\mathbf{B}_l = e^{-\Delta\tau \mathbf{H}_t} e^{-\Delta\tau \mathbf{H}_U(\vec{s}(l))}$$

- Monte Carlo sampling in configuration space

$$\mathbf{H}_U(\vec{s}(l)) \propto \alpha \vec{s}(l)$$



Quantum Monte Carlo

■ Hubbard-Stratonovich Transformation

$$\exp\left\{-\Delta\tau\frac{U}{2}(n_{i,\uparrow} + n_{i,\downarrow} - 1)^2\right\} = \frac{1}{4} \sum_{s_i=\pm 1, \pm 2} \gamma(s_i) e^{i\sqrt{\Delta\tau\frac{U}{2}}\eta(s_i)(n_{i,\uparrow} + n_{i,\downarrow} - 1)} + O[(\Delta\tau)^4]$$

$$\exp\left\{-\Delta\tau\frac{J}{8}(D_{i,j} - D_{i,j}^\dagger)^2\right\} = \frac{1}{4} \sum_{t_{i,j}=\pm 1, \pm 2} \gamma(t_{i,j}) e^{i\sqrt{\Delta\tau\frac{J}{8}}\eta(t_{i,j})(D_{i,j} - D_{i,j}^\dagger)} + O[(\Delta\tau)^4]$$

■ QMC measurements

$$\langle H \rangle = \langle H_{KM} \rangle + \langle H_U \rangle + \langle H_J \rangle \qquad G(k, \tau) \propto Z_{sp}(k) e^{-\Delta_{sp}(k)\tau} \quad k = K, M$$

$$\frac{S^{xy}(Q_{AF})}{L^2} = m_S^2 + \frac{a_1}{L} + \frac{a_2}{L^2} + \dots \qquad S^{xy}(k, \tau) \propto Z_S(k) e^{-\Delta_S(k)\tau} \quad k = Q_{AF} = \Gamma$$

Quantum Monte Carlo

- Computation effort scales linearly with βN^3

System sizes: $N_s = 2 \times L^2$ $L = 3, 6, 9, 12, 15, 18$

Time discretization: $\beta t \propto L$, $\Delta\tau t = 0.05$

Parallelization: $\sim 10^3$ CPUs, $\sim 10^6$ CPU hours



国家超级计算天津中心
National Supercomputer Center in Tianjin

Tianhe-1



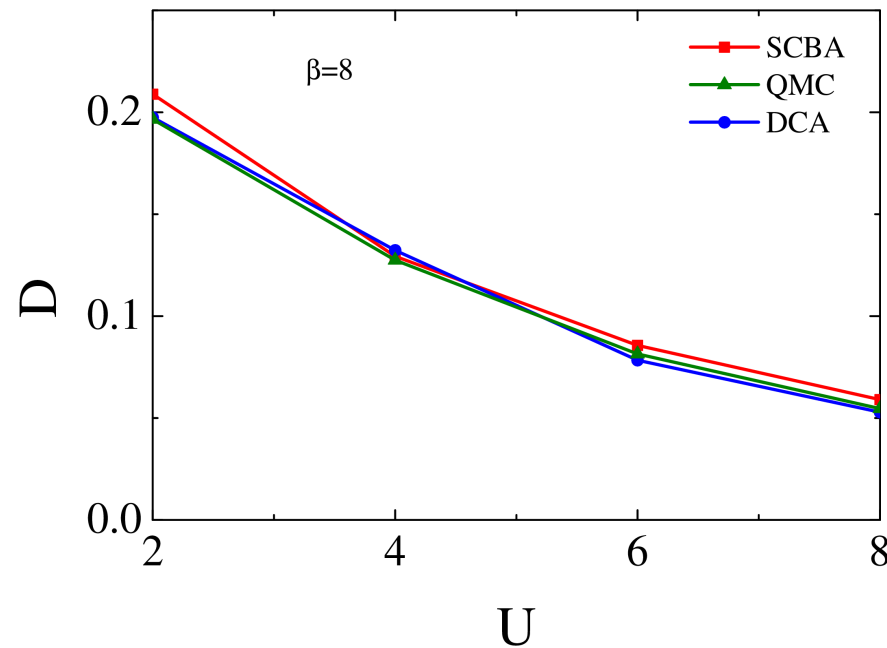
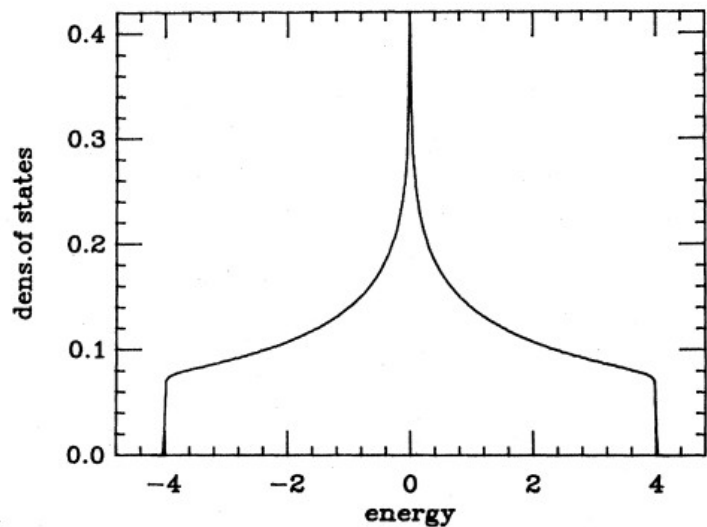
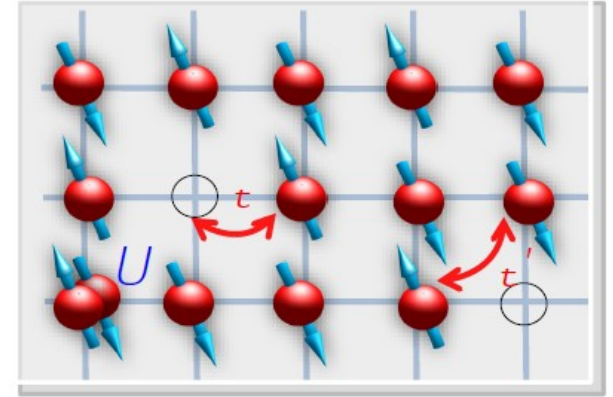
国家超级计算广州中心
NATIONAL SUPERCOMPUTER CENTER IN GUANGZHOU

Tianhe-2



Square lattice Hubbard model

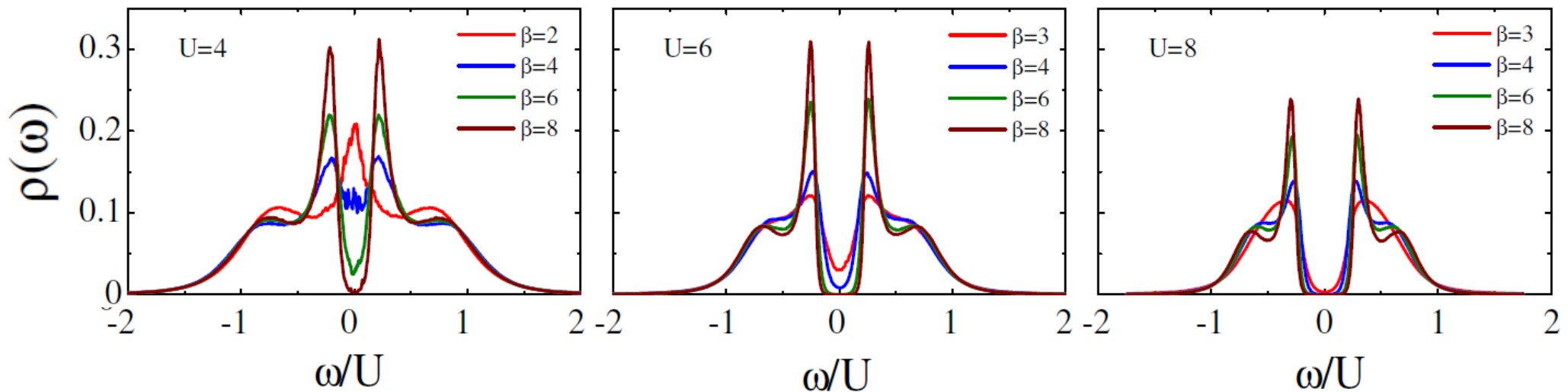
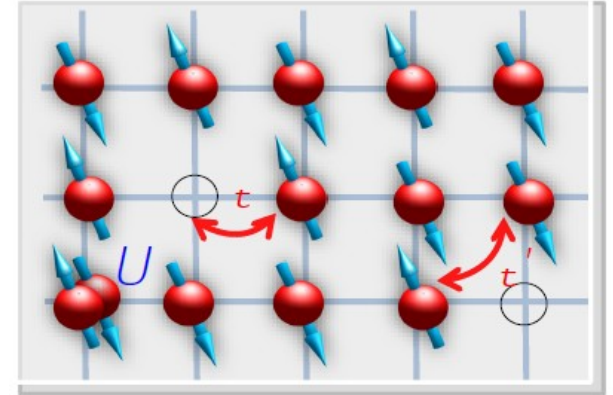
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➤ C. Chen, Bachelor Thesis (2016)

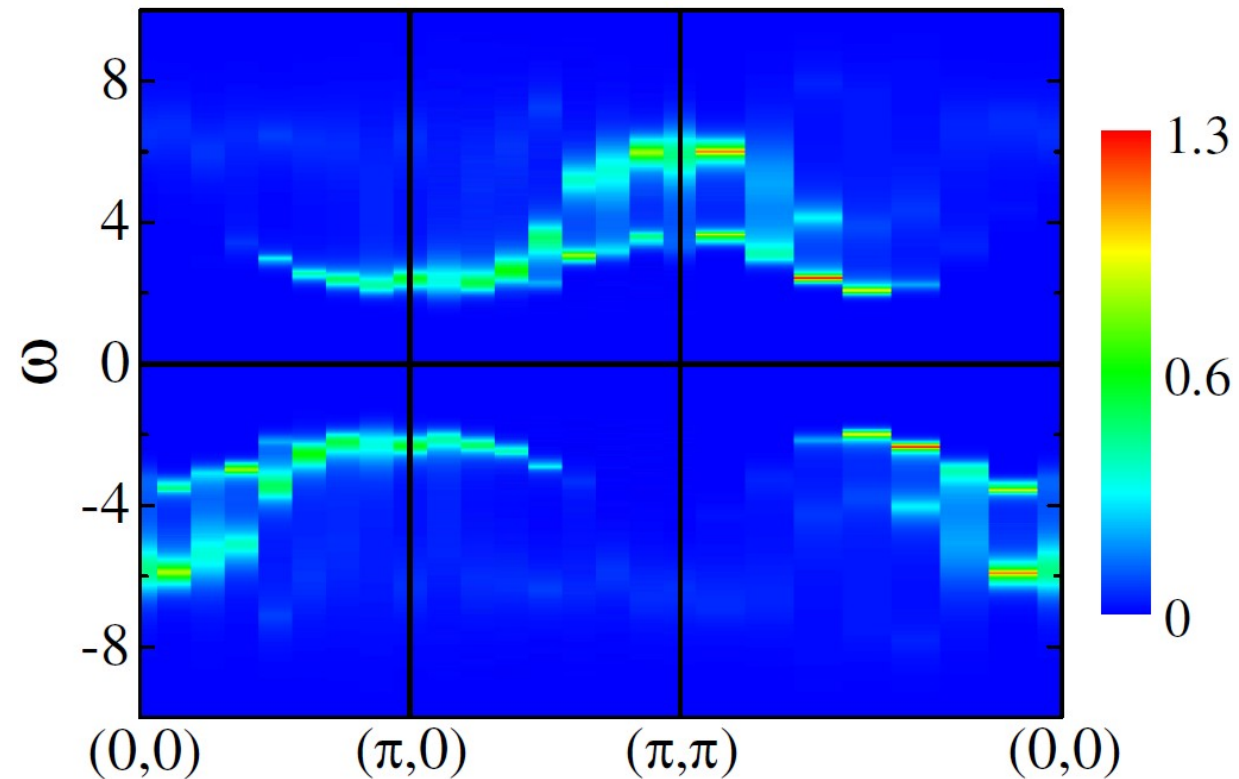
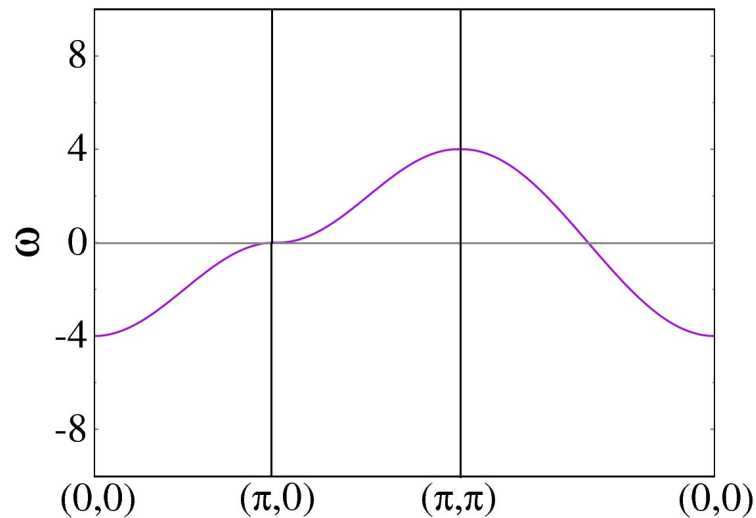
Square lattice Hubbard model

$$H = -t \sum_{\langle i,j \rangle, \sigma} (c_{i,\sigma}^\dagger c_{j,\sigma} + \text{h.c.}) + U \sum_{i=1}^N (n_{i,\uparrow} - \frac{1}{2})(n_{i,\downarrow} - \frac{1}{2})$$



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Square lattice Hubbard model



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