

The emission rate of ionizing photons into the IGM per baryon can be written as:

$$\dot{n}_{\text{ion/b}} = \frac{d}{dt} \left[\bar{\rho}_m^{-1} \int_{M_{\text{min}}}^{\infty} dM_h \frac{dn}{d \ln M_h} f_b f_* N_{\gamma/b} f_{\text{esc}} \right]. \quad (80)$$

Here $\bar{\rho}_m$ is the mean matter density, and the average source (galaxy) properties are expressed in terms of:

- f_b – the baryon fraction of a halo, in units of the cosmic mean value, Ω_b/Ω_m
- f_* – the fraction of halo baryons ending up in stars
- $N_{\gamma/b}$ – the number of ionizing photons produced per stellar baryon
- f_{esc} – the fraction of produced ionizing photons which escape the galaxy into the IGM

and the lower limit of integration, M_{min} , corresponds to the minimum halo mass capable of hosting a star forming galaxy (i.e. the product $f_b f_* N_{\gamma/b} f_{\text{esc}} = 0$ for $M_h < M_{\text{min}}$).

Typically, these astrophysical parameters are combined into a single ionizing efficiency:

$$\zeta = 20 \left(\frac{N_{\gamma}}{4000} \right) \left(\frac{f_{\text{esc}}}{0.1} \right) \left(\frac{f_*}{0.05} \right) \left(\frac{f_b}{1} \right). \quad (81)$$

If ζ is independent of halo mass and time, eq. (80) simplifies to include the time derivative of the collapse fraction:

$$\dot{n}_{\text{ion/b}} = \zeta \frac{df_{\text{coll}}(> M_{\text{min}}, z)}{dt}. \quad (82)$$

More generally, one could parametrize ζ with a power-law scaling with the halo mass, $\zeta = \zeta_0 \left(\frac{M_h}{M_{\text{min}}} \right)^{\alpha}$, in which case eq. (80) becomes:

$$\dot{n}_{\text{ion/b}} = \frac{d}{dt} \left[\frac{\zeta_0}{\bar{\rho}_m} \int_{M_{\text{min}}}^{\infty} dM_h \frac{dn}{d \ln M_h} \left(\frac{M_h}{M_{\text{min}}} \right)^{\alpha} \right], \quad (83)$$

with $\alpha = 0$ reducing to eq. (82).

In principle, direct observations of galaxies can be used to constrain some of the above parameters. The abundance matching technique applied to $z \sim 8$ LBGs, discussed in the previous chapter can be used to motivate the scaling $f_b \propto M_h^{0.2-0.4}$, assuming a mass-independent mapping from the observed 1500 Å luminosity to the ionizing luminosity (e.g. Trenti et al. 2010; Greig & Mesinger 2015; Atek et al. 2015; Sun & Furlanetto 2015). However, f_{esc} likely has an opposite scaling, increasing towards smaller halo masses (e.g. Wise & Cen 2009; Paardekooper et al. 2015), perhaps even showing a non-monotonic evolution (e.g. Xu et al. 2016). Moreover, the IMF, which sets $N_{\gamma/b}$, could be more top-heavy in low-mass, poorly-enriched galaxies. For the purposes of EoR modeling, we only care about the product, ζ .