

# Chapter 1

## Soft Tissue Engineering

### 1.1 Stress & strain in large deformations

#### 1.1.1 Kinematics

- Mapping function

$$\vec{x} = \varphi(\vec{X}, t) \quad (1.1)$$

- Displacement function

$$\vec{x} = \vec{X} + \underline{U}(\vec{X}, t) \quad (1.2)$$

- The deformation gradient tensor

$$d\vec{x} = \mathbf{F}(X, t)d\vec{X} \Rightarrow \mathbf{F} = \frac{\partial \vec{x}}{\partial \vec{X}} = \frac{\partial \varphi}{\partial \vec{X}} \quad (1.3)$$

- Volume transformation

$$dv = \det(\mathbf{F})dV = JdV \quad (1.4)$$

- Surface transformation: Nanson's Formula

$$\vec{n}ds = J\mathbf{F}^{-\top}\vec{N}dS \quad (1.5)$$

- Polar decomposition

$$\mathbf{F} = \mathbf{R}\mathbf{U} = \mathbf{V}\mathbf{R} \quad (1.6)$$

$$\Rightarrow \mathbf{U}^2 = \mathbf{F}^\top \mathbf{F} = \mathbf{C} \quad (1.7)$$

$$\Rightarrow \mathbf{V}^2 = \mathbf{F}\mathbf{F}^\top = \mathbf{B} \quad (1.8)$$

#### 1.1.2 Strain

- Nominal strain, Biot strain or Engineering strain

$$\mathbf{e} = \mathbf{U} - \mathbf{I} \quad (1.9)$$

- Green-Lagrange strain

$$\mathbf{E} = \frac{1}{2}(\mathbf{F}^\top \mathbf{F} - \mathbf{I}) = \frac{1}{2}(\mathbf{U}^2 - \mathbf{I}) \quad (1.10)$$

- Euler-Almansi strain

$$\mathbf{A} = \frac{1}{2}(\mathbf{I} - \mathbf{V}^{-2}) \quad (1.11)$$

- True strain or logarithmic strain

$$\varepsilon = \ln(\mathbf{U}) \quad (1.12)$$

### 1.1.3 Stress

- True stress or Cauchy stress

$$\boldsymbol{\sigma}^\top = \frac{d\vec{f}}{\vec{n}dS} \quad (1.13)$$

- 1PK: first Piola Kirchoff stress (Nominal stress or Engineering stress)

$$\mathbf{P} = \frac{d\vec{f}}{\vec{N}dS_0} \quad (1.14)$$

- 2PK: second Piola Kirchoff stress

$$\mathbf{S}^\top = \frac{\mathbf{F}^{-1}d\vec{f}}{\vec{N}dS_0} \quad (1.15)$$

	$\boldsymbol{\sigma}$	$\mathbf{P}$	$\mathbf{S}$
• Switching between stresses:	$\boldsymbol{\sigma}$	$J^{-1}\mathbf{P}\mathbf{F}^\top$	$J^{-1}\mathbf{F}\mathbf{S}\mathbf{F}^\top$
	$\mathbf{P}$	$J\boldsymbol{\sigma}\mathbf{F}^\top$	$\mathbf{P}$
	$\mathbf{S}$	$J\mathbf{F}^{-1}\boldsymbol{\sigma}\mathbf{F}^{-\top}$	$\mathbf{F}^{-1}\mathbf{P}$

## 1.2 Constitutive modeling

The most general form:

$$\boldsymbol{\sigma}(t) = \mathcal{F}_{\tau \leq t}\{\mathbf{F}(\tau)\} \quad (1.16)$$

### 1.2.1 Hyperelasticity

Strain energy density function  $\Psi$  related to 2PK

$$\mathbf{S} = \frac{\partial \Psi}{\partial \mathbf{E}} = 2 \frac{\partial \Psi}{\partial \mathbf{C}} \quad (1.17)$$

Strain energy density function related to Cauchy stress

$$\boldsymbol{\sigma} = \frac{1}{J} \frac{\partial \Psi}{\partial \mathbf{F}} \mathbf{F}^\top \quad (1.18)$$

Small strain - linear elasticity (using Hooke's law)

$$\Psi_{\text{elastic}} = \frac{E\varepsilon^2}{2} \quad (1.19)$$

### 1.2.2 Isotropic hyperelastic materials

Right Cauchy-Green Tensor (with  $\lambda_i$  the PRICIPAL stretches)

$$\mathbf{C} = \begin{bmatrix} \lambda_1^2 & 0 & 0 \\ 0 & \lambda_2^2 & 0 \\ 0 & 0 & \lambda_3^2 \end{bmatrix} \quad (1.20)$$

with invariants:

$$I_1 = \text{tr}(\mathbf{C}) = \lambda_1^2 + \lambda_2^2 + \lambda_3^2 \quad (1.21)$$

$$I_2 = \frac{1}{2}[\text{tr}(\mathbf{C})^2 - \text{tr}(\mathbf{C}^2)] = \lambda_1^2\lambda_2^2 + \lambda_1^2\lambda_3^2 + \lambda_2^2\lambda_3^2 \quad (1.22)$$

$$I_3 = \det(\mathbf{C}) = \lambda_1^2\lambda_2^2\lambda_3^2 = J^2 \quad (1.23)$$

Some forms of isotropic SEDFs

- Generalized Ogden

$$\Psi = \sum_{p=1}^N \frac{\mu_p}{\alpha_p} (\lambda_1^{\alpha_p} + \lambda_2^{\alpha_p} + \lambda_3^{\alpha_p} - 3) \quad (1.24)$$

- Neo-Hookean ( $N = 1, \alpha_1 = 2$ )

$$\psi = c_1(I_1 - 3) \quad (1.25)$$

- Mooney-Rivlin ( $N = 2, \alpha_1 = 2, \alpha_2 = -2$ )

$$\psi = c_1(I_1 - 3) + c_2(\lambda_1^{-2} + \lambda_2^{-2} + \lambda_3^{-2} - 3) \quad (1.26)$$

### 1.2.3 The NeoHookean material formulation

$$\sigma_{\text{NH}} = \frac{2}{J} c_1 \mathbf{B} \quad (1.27)$$

### 1.2.4 Incompressibility

Consider as SEDF

$$\Psi_{\text{tot}} = \Psi(\mathbf{F}) - p(J - 1) \quad (1.28)$$

the 2PK stress becomes

$$\mathbf{S} = 2 \left( \frac{\partial \Psi}{\partial \mathbf{C}} - p \frac{J}{2} \mathbf{C}^{-\top} \right) \quad (1.29)$$

hence the Cauchy stress takes the following form

$$\sigma = \frac{2}{J} \frac{\partial \Psi}{\partial \mathbf{C}} \mathbf{F}^{\top} - p \mathbf{I} \quad (1.30)$$

Using the NeoHookean SEDF gives

$$\sigma = 2c_1 \mathbf{B} - p \mathbf{I} = -p_h \mathbf{I} \quad (1.31)$$

## 1.3 Mechanical testing & parameter fitting

### 1.3.1 Uniaxial tensile testing

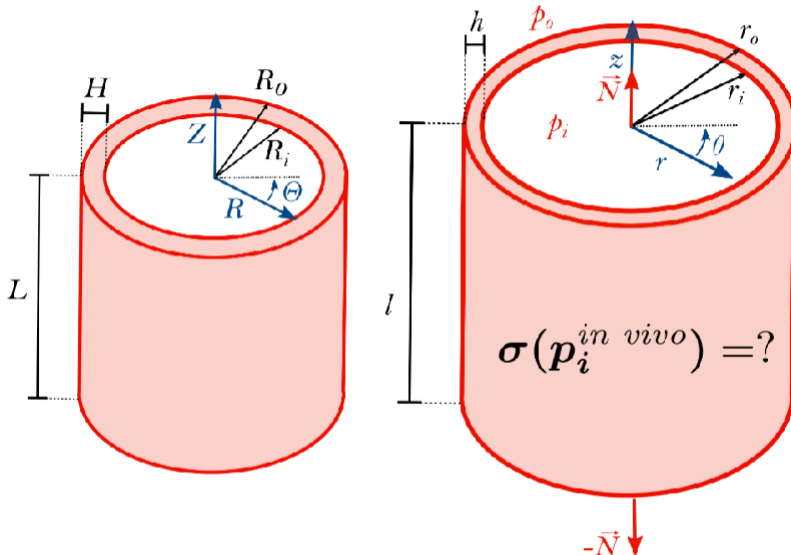
### 1.3.2 Biaxial tensile testing

### 1.3.3 Extension-inflation testing

### 1.3.4 Strain mapping

## 1.4 In vivo wall stress estimation

### 1.4.1 Laplace's law for thin-walled tubes



The deformation gradient tensor is given as

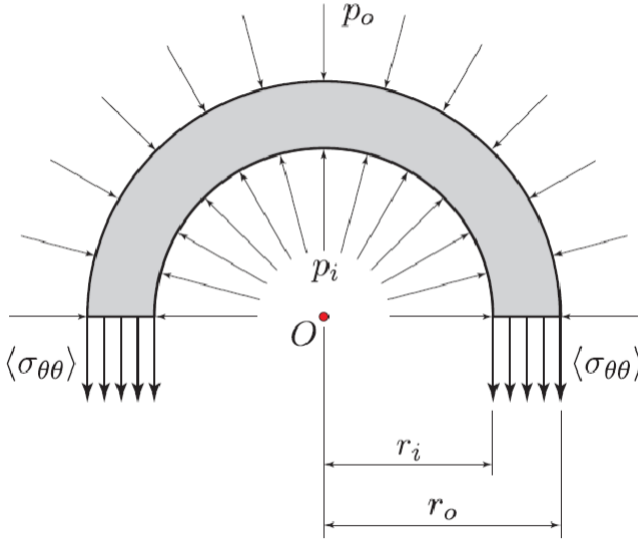
$$\mathbf{F} = \begin{bmatrix} 1/(\lambda_\theta \lambda_z) & 0 & 0 \\ 0 & \lambda_\theta & 0 \\ 0 & 0 & \lambda_z \end{bmatrix} \quad (1.32)$$

with corresponding stresses

$$\langle \sigma_{rr} \rangle = -\frac{P}{2} \quad (1.33)$$

$$\langle \sigma_{\theta\theta} \rangle = \frac{P}{h/r_i} \quad (1.34)$$

$$\langle \sigma_{zz} \rangle = \frac{\|\vec{N}\|}{\pi(r_o^2 - r_i^2)} \quad (1.35)$$



### 1.4.2 Analytical analysis for thick-walled cylinders

Stress equilibrium (Balance of Momentum)

$$\nabla \cdot \boldsymbol{\sigma} + \vec{f} = \rho \frac{D^2 \vec{u}}{Dt^2} \quad (1.36)$$

Stress equilibrium in a thick-walled cylinder

$$P = \int_{r_i}^{r_o} (\sigma_{\theta\theta} - \sigma_{rr}) \frac{dr}{r} \quad (1.37)$$

NeoHookean formulation for the Cauchy Stress in a thick-walled cylinder

$$\sigma_{rr} = 2c_1 \lambda_\theta^{-2} \lambda_z^{-2} - p \quad (1.38)$$

$$\sigma_{\theta\theta} = 2c_1 \lambda_\theta^2 - p \quad (1.39)$$

$$\sigma_{zz} = 2c_1 \lambda_z^2 - p \quad (1.40)$$

with

$$\lambda_\theta = \frac{r}{R} \quad (1.41)$$

$$r = \sqrt{\lambda_{\theta,i}^2 R_i^2 + \frac{R^2 - R_i^2}{\lambda_z}} \quad (1.42)$$

### 1.4.3 The Finite Element Method

## 1.5 Other biomechanical applications