Solutions

1. Charged kaon decay:

(a) From the particle data group listing we find the mean lifetime of the K^+ to be

$$\tau = 1.238 \cdot 10^{-8} \text{ s.} \tag{1}$$

Note that this is the life expectancy of the kaon in its rest frame. In the laboratory frame this time interval appears to be longer:

$$\tau_{\text{lab}} = \gamma \cdot \tau = 8.78 \cdot 10^{-8} \text{ s} \tag{2}$$

since $\gamma = 1/\sqrt{1-\beta^2} \approx 7.1$ for $\beta = 0.99$. The distance traveled during this time is

$$x = \beta c \tau_{\text{lab}} = 26 \text{ m.} \tag{3}$$

2. Alarm clock:

(a) Let the position of your alarm clock be the origin of the coordinate system, with the x axis pointing toward U of M. Also let the time for your alarm clock ringing be t=0. Then event 1 (your alarm clock ringing) would occur at $(ct_1,x_1)=(0,0)$ and event 2 (UM alarm clock ringing) would occur at $(ct_2,x_2)=(c\cdot 10^{-5}\text{ s},10^5\text{ m})=(3\cdot 10^3,10^5)\text{ m}$. (We suppress the irrelevant coordinates y and z here.) In a Lorentz frame moving with speed v in the x direction relative to yours, the two events would occur at

$$(ct'_1, x'_1) = (0, 0),$$
 (4)

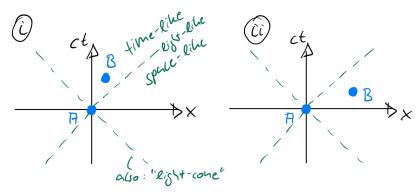
$$(ct'_2, x'_2) = \left(\gamma(ct_2 - \beta x_2), \gamma(x_2 - \beta ct_2)\right),$$
 (5)

where $\beta = v/c$ and $\gamma = (1 - \beta^2)^{-1/2}$. In order to have $t_2' < t_1'$, we need

$$ct_2 - \beta x_2 < 0 \tag{6}$$

or $v > (ct_2/x_2)c = 0.03c$. Thus, the UM alarm clock would ring first in any frame moving in the direction from East Lansing towards Ann Arbor at a speed v > 0.3c.

(b) The two situations can be displayed as follows.



(c) From the arguments in part (a), it is only possible to find a frame in which $t'_2 < t'_1$ if $ct_2/x_2 < 1$ or

$$t_2 < x_2/c \approx 3.34 \times 10^{-4} \text{s.}$$
 (7)

In other words, the time difference must be less than the time for light to travel the 100 km. Equivalently, the interval between the two events must be space-like or $s_{12}^2 = c^2(t_2 - t_1)^2 - (x_1 - x_2)^2 < 0$.

3. Colliding rockets:

(a) We can obtain the coordinates of B in A's frame (primed) by a Lorentz transformation from the Earth frame (unprimed):

$$\Delta x' = \gamma (\Delta x - v_A \Delta t) \tag{8}$$

$$\Delta t' = \gamma \left(\Delta t - \frac{v_A}{c^2} \Delta x \right) \tag{9}$$

Using these equations, the correct signs for the different velocities and $|v_B| = |\Delta x/\Delta t|$ in the Earth frame we obtain

$$v_B' = \frac{\Delta x'}{\Delta t'} = \frac{-|v_B| - |v_A|}{1 + |v_A||v_B|/c^2} = -0.95c.$$
(10)

The same result but with opposite sign holds for the velocity of A in B's frame.

(b) In the Earth frame it takes

$$\Delta t = \frac{l}{v_A + v_B} = 1 s \tag{11}$$

until the collision (if Δx_A is the distance from a to the collision point we have $\Delta x_A = v_A t$, $l - \Delta x_A = v_B t$, and thus (??)). The clock in frame A goes slower and that time interval is observed as

$$\Delta t_A = \frac{\Delta t}{\gamma_A} = \sqrt{1 - v_A^2/c^2} = 0.6 \ s. \tag{12}$$

Similarly, the time interval seen in B's frame is

$$\Delta t_B = \frac{\Delta t}{\gamma_B} = \sqrt{1 - v_B^2/c^2} = 0.8 \ s. \tag{13}$$

Note that also the initial distance to the other rocket will look different in the frames of A and B, since the starting point is defined with respect to a measurement in the Earth frame (which breaks the symmetry).

4. Explicit boost:

(a) The problem could be solved by rotating the coordinate system to align the x axis with the direction of the boost. A simpler method is as follows. It is only the component of the position vector \vec{x} in the direction of the boost, which requires a non-trivial transformation. We therefore decompose the position vector as follows:

$$\vec{x} = \vec{x}_{\parallel} + \vec{x}_{\perp} \tag{14}$$

Starting from \vec{x} and \vec{v} , the components are determined by

$$\vec{x}_{\parallel} = x_{\parallel} \hat{v}$$
 with $x_{\parallel} = \vec{x} \cdot \hat{v}$ and $\hat{v} = \vec{v}/|\vec{v}|$, (15)

$$\vec{x}_{\perp} = \vec{x} - \vec{x}_{\parallel}.\tag{16}$$

Plugging in numbers we have

$$x_{\parallel} \approx (2.23, 2.79, 3.35) m,$$
 (17)

$$x_{\perp} \approx (1.77, 0.21, -1.35) m$$
 (18)

We can obtain the vector in the boosted frame by transformation both terms in (??) separately, since Lorentz transformations are linear:

$$\vec{x}' = \vec{x}_{\parallel}' + \vec{x}_{\perp}'. \tag{19}$$

We use the Lorentz transformation for a boost with $\beta = |\vec{v}|/c = 0.88$ in x direction and replace x with x_{\parallel} and the y and z components by \vec{x}_{\perp} :

$$ct' = \gamma(ct - \beta x_{\parallel})$$
 $\approx 2.09(5 - 0.88 \cdot 4.90) \ m$ $\approx 1.46 \ m,$ (20)

$$ct' = \gamma(ct - \beta x_{\parallel})$$
 $\approx 2.09(5 - 0.88 \cdot 4.90) m$ $\approx 1.46 m,$ (20)
 $x'_{\parallel} = \gamma(x_{\parallel} - \beta ct)$ $\approx 2.09(4.90 - 0.88 \cdot 5) m$ $\approx 1.07 m,$ (21)

$$\vec{x}'_{\perp} = \vec{x}_{\perp} \tag{22}$$

For the 3-vectors we have

$$\vec{x}'_{\parallel} = x'_{\parallel} \hat{v}$$
 $\approx (0.49, 0.61, 0.73) m,$ (23)

$$\vec{x}' = \vec{x}'_{\parallel} + \vec{x}_{\perp} \qquad \approx (2.25, 0.82, -0.62) \ m$$
 (24)

and thus for the four-vector in the boosted frame:

$$x'^{\mu} \approx (2.50, 1.24, -0.20, -1.64) m.$$
 (25)

(b) The rapidity ζ is given by

$$\zeta = \operatorname{arctanh} |\beta| \approx 1.36$$
. (26)

(c) We consider the Lorentz invariant quantity

$$s^{2} = x_{\mu}x^{\mu} = (x^{0})^{2} - (x^{1})^{2} - (x^{2})^{2} - (x^{3})^{2} = -4 \text{ m}^{2} < 0.$$
 (27)

With our metric convention, a positive value means that the four-vector x is spacelike. From the fact that s^2 is Lorentz invariant it is immediately clear that also x'^{μ} is space-like. Events which differ by a space-like vector are not causally connected.