# Problem Set 1

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### I SAKURAI 4.4

Α

$$\mathcal{Y}_0^{1/2,1/2} = \begin{pmatrix} Y_0^0 \\ 0 \end{pmatrix},$$
 (I.1)

$$= \frac{1}{\sqrt{4\pi}} \begin{pmatrix} 1\\0 \end{pmatrix}. \tag{I.2}$$

В

$$(\boldsymbol{\sigma} \cdot \mathbf{x}) \mathcal{Y}_0^{1/2,1/2} = \frac{r}{\sqrt{4\pi}} \begin{pmatrix} \cos \theta \\ e^{i\phi} \sin \theta \end{pmatrix}, \tag{I.3}$$

$$= \frac{r}{\sqrt{3}} \begin{pmatrix} Y_1^0 \\ -\sqrt{2}Y_1^1 \end{pmatrix}, \tag{I.4}$$

$$= -r\mathcal{V}_1^{1/2,1/2}.\tag{I.5}$$

Here, j and m remained the same, but l increased from 0 to 1.

 $\mathbf{C}$ 

Since  $\mathbf{S} \cdot \mathbf{x}$  is a scalar under rotations, it can only connect states with j' = j and m' = m (Wigner-Ekhart Theorem). However, the operator connects states of opposite parity. This is done by changing l by one unit.

## II SAKURAI 4.5

The symmetry of V is determined by  $\mathbf{S} \cdot \mathbf{p}$ , which is a pseudoscalar. Therefore the matrix element  $C_{n'l'j'm'}$  is zero unless  $j'=j,\ m'=m,$  and  $l'=l\pm 1.$  The radial wave functions go as  $r^l$ , so, since  $V(\mathbf{x}) \propto \delta^3(\mathbf{x})$ , the matrix element will only be nonzero for states with l=0. Consequently, this interaction only connects  $S_{1/2}$  and  $P_{1/2}$  states.

## III SAKURAI 4.7

A

$$\psi_p(x,t) = e^{i(p \cdot x - Et)/\hbar}.$$
 (III.1)

$$\psi_p^*(x, -t) = e^{-i(p \cdot x + Et)/\hbar}, \qquad (III.2)$$

$$= e^{i(-p \cdot x - Et)/\hbar}, \qquad (III.3)$$

$$=\psi_{-n}(x,t). \qquad \checkmark \qquad (III.4)$$

 $\mathbf{B}$ 

Option 1

$$\chi_{+}(n) = \begin{pmatrix} e^{-i\phi/2} \cos \frac{\theta}{2} \\ e^{i\phi/2} \sin \frac{\theta}{2} \end{pmatrix}.$$
 (III.5)

$$i\sigma_2 \chi_+^*(n) = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} e^{i\phi/2} \cos\frac{\theta}{2} \\ e^{-i\phi/2} \sin\frac{\theta}{2} \end{pmatrix},$$
 (III.6)

$$= \begin{pmatrix} -e^{-i\phi/2} \sin\frac{\theta}{2} \\ e^{i\phi/2} \cos\frac{\theta}{2} \end{pmatrix}. \qquad \checkmark$$
 (III.7)

Option 2

$$\chi(n) = e^{-i\sigma_3\phi/2}e^{-i\sigma_2\theta/2}.$$
 (III.8)

$$-i\sigma_2\chi^*(n) = -i\sigma_2 e^{i\sigma_2\theta/2} e^{i\sigma_3\phi/2}, \qquad (III.9)$$

$$= e^{-i\sigma_2\pi/2} e^{i\sigma_2\theta/2} e^{i\sigma_3\phi/2}. \qquad \checkmark \qquad \text{(III.10)}$$

### IV SAKURAI 4.11

$$\langle \widehat{\mathbf{L}} \rangle = \langle E | \widehat{\mathbf{L}} | E \rangle \,, \tag{IV.1}$$

$$= \langle \tilde{E} | \hat{\Theta} \hat{\mathbf{L}} \hat{\Theta}^{-1} | \tilde{E} \rangle, \qquad (IV.2)$$

$$= -\langle \tilde{E} | \hat{\mathbf{L}} | \tilde{E} \rangle, \qquad (IV.3)$$

$$= -\langle E|\widehat{\mathbf{L}}|E\rangle. \tag{IV.4}$$

$$\therefore \langle \widehat{\mathbf{L}} \rangle = 0. \tag{IV.5}$$

$$\langle \mathbf{x}|\tilde{E}\rangle = e^{-i\delta} \langle \mathbf{x}|E\rangle,$$
 (IV.6)

$$= e^{-i\delta} \sum_{lm} F_{lm}(r) Y_l^m(\theta, \phi). \tag{IV.7}$$

$$\langle \mathbf{x} | \tilde{E} \rangle = \langle E | \mathbf{x} \rangle,$$
 (IV.8)

$$= \sum_{lm} F_{lm}^{*}(r)(Y_{l}^{m})^{*}(\theta, \phi), \qquad (IV.9)$$

$$= \sum_{l,m} F_{l,-m}^*(r)(-1)^{-m} Y_l^{-m}(\theta,\phi), \qquad \text{(IV.10)}$$

$$\therefore F_{lm}(r) = (-1)^m e^{-i\delta} F_{l,-m}^*(r).$$
 (IV.11)