

Problem Set 1

Brandon Henke¹

¹Michigan State University - Department of Physics & Astronomy

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I SAKURAI 4.4

A

$$y_0^{1/2,1/2} = \begin{pmatrix} Y_0^0 \\ 0 \end{pmatrix}, \quad (\text{I.1})$$

$$= \frac{1}{\sqrt{4\pi}} \begin{pmatrix} 1 \\ 0 \end{pmatrix}. \quad (\text{I.2})$$

B

$$(\boldsymbol{\sigma} \cdot \mathbf{x}) y_0^{1/2,1/2} = \frac{r}{\sqrt{4\pi}} \begin{pmatrix} \cos \theta \\ e^{i\phi} \sin \theta \end{pmatrix}, \quad (\text{I.3})$$

$$= \frac{r}{\sqrt{3}} \begin{pmatrix} Y_1^0 \\ -\sqrt{2} Y_1^1 \end{pmatrix}, \quad (\text{I.4})$$

$$= -r y_1^{1/2,1/2}. \quad (\text{I.5})$$

Here, j and m remained the same, but l increased from 0 to 1.

C

Since $\mathbf{S} \cdot \mathbf{x}$ is a scalar under rotations, it can only connect states with $j' = j$ and $m' = m$ (Wigner-Eckart Theorem). However, the operator connects states of opposite parity. This is done by changing l by one unit.

II SAKURAI 4.5

The symmetry of V is determined by $\mathbf{S} \cdot \mathbf{p}$, which is a pseudoscalar. Therefore the matrix element $C_{n'l'j'm'}$ is zero unless $j' = j$, $m' = m$, and $l' = l \pm 1$. The radial wave functions go as r^l , so, since $V(\mathbf{x}) \propto \delta^3(\mathbf{x})$, the matrix element will only be nonzero for states with $l = 0$. Consequently, this interaction only connects $S_{1/2}$ and $P_{1/2}$ states.

III SAKURAI 4.7

A

$$\psi_p(x, t) = e^{i(p \cdot x - Et)/\hbar}. \quad (\text{III.1})$$

$$\psi_p^*(x, -t) = e^{-i(p \cdot x + Et)/\hbar}, \quad (\text{III.2})$$

$$= e^{i(-p \cdot x - Et)/\hbar}, \quad (\text{III.3})$$

$$= \psi_{-p}(x, t). \quad \checkmark \quad (\text{III.4})$$

B

Option 1

$$\chi_+(n) = \begin{pmatrix} e^{-i\phi/2} \cos \frac{\theta}{2} \\ e^{i\phi/2} \sin \frac{\theta}{2} \end{pmatrix}. \quad (\text{III.5})$$

$$i\sigma_2 \chi_+^*(n) = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} e^{i\phi/2} \cos \frac{\theta}{2} \\ e^{-i\phi/2} \sin \frac{\theta}{2} \end{pmatrix}, \quad (\text{III.6})$$

$$= \begin{pmatrix} -e^{-i\phi/2} \sin \frac{\theta}{2} \\ e^{i\phi/2} \cos \frac{\theta}{2} \end{pmatrix}. \quad \checkmark \quad (\text{III.7})$$

Option 2

$$\chi(n) = e^{-i\sigma_3 \phi/2} e^{-i\sigma_2 \theta/2}. \quad (\text{III.8})$$

$$-i\sigma_2 \chi^*(n) = -i\sigma_2 e^{i\sigma_2 \theta/2} e^{i\sigma_3 \phi/2}, \quad (\text{III.9})$$

$$= e^{-i\sigma_2 \pi/2} e^{i\sigma_2 \theta/2} e^{i\sigma_3 \phi/2}. \quad \checkmark \quad (\text{III.10})$$

IV SAKURAI 4.11

$$\langle \hat{\mathbf{L}} \rangle = \langle E | \hat{\mathbf{L}} | E \rangle, \quad (\text{IV.1})$$

$$= \langle \tilde{E} | \hat{\Theta} \hat{\mathbf{L}} \hat{\Theta}^{-1} | \tilde{E} \rangle, \quad (\text{IV.2})$$

$$= -\langle \tilde{E} | \hat{\mathbf{L}} | \tilde{E} \rangle, \quad (\text{IV.3})$$

$$= -\langle E | \hat{\mathbf{L}} | E \rangle. \quad (\text{IV.4})$$

$$\therefore \langle \hat{\mathbf{L}} \rangle = 0. \quad (\text{IV.5})$$

$$\langle \mathbf{x}, \tilde{E} | \mathbf{x}, \tilde{E} \rangle = e^{-i\delta} \langle \mathbf{x}, E | \mathbf{x}, E \rangle, \quad (\text{IV.6})$$

$$= e^{-i\delta} \sum_{lm} F_{lm}(r) Y_l^m(\theta, \phi). \quad (\text{IV.7})$$

$$\langle \mathbf{x}, \tilde{E} | \mathbf{x}, \tilde{E} \rangle = \langle E | \mathbf{x} \rangle, \quad (\text{IV.8})$$

$$= \sum_{lm} F_{lm}^*(r) (Y_l^m)^*(\theta, \phi), \quad (\text{IV.9})$$

$$= \sum_{lm} F_{l,-m}^*(r) (-1)^{-m} Y_l^{-m}(\theta, \phi), \quad (\text{IV.10})$$

$$\therefore F_{lm}(r) = (-1)^m e^{-i\delta} F_{l,-m}^*(r). \quad (\text{IV.11})$$