# Problem Set 5

Brandon Henke<sup>1</sup>

 $^1$ Michigan State University - Department of Physics & Astronomy

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## I DEUTERIUM HYPERFINE TRANSITION

## II TIME-DEPENDENT ELECTRIC FIELD

## III SAKURAI 5.29

$$c_0^{(0)} = 1. (III.1)$$

$$c_n^{(1)} = -\frac{i}{\hbar} \int_0^t e^{i(E-E_0)t'/\hbar} \langle n|F_0 x \cos \omega t'|0\rangle dt', \qquad (III.2)$$

$$= -\frac{i}{\hbar} F_0 \sqrt{\frac{\hbar}{2m\omega_0}} \delta_{n1} \int_0^t e^{-i\omega_0 t'} \cos \omega t' \, dt'.$$
 (III.3)

$$c_1^{(1)} = -\frac{i}{\hbar} F_0 \sqrt{\frac{\hbar}{2m\omega_0}} \int_0^t e^{-i\omega_0 t'} \cos \omega t' \, dt', \qquad (III.4)$$
$$= -\frac{i}{2\hbar} F_0 \sqrt{\frac{\hbar}{2m\omega_0}}$$

$$\times \left( \frac{e^{i(\omega_0 + \omega)t} - 1}{\omega_0 + \omega} + \frac{e^{i(\omega_0 - \omega)t} - 1}{\omega_0 - \omega} \right) \equiv c_1(t). \text{ (III.5)}$$

$$\langle \hat{x} \rangle = \langle \alpha, t | \hat{x} | \alpha, t \rangle_S,$$
 (III.6)

$$= \left(e^{i\omega_0 t/2} \langle 0| - c_1^* e^{3i\omega_0 t/2} \langle 1|\right) \hat{x}$$
 (III.7)

$$\times \left(e^{-i\omega_0 t/2} |0\rangle - c_1 e^{-3i\omega_0 t/2} |1\rangle\right),$$
 (III.8)

$$= \sqrt{\frac{\hbar}{2m\omega_0}} \left( c_1 e^{-i\omega_0 t} + c_1^* e^{i\omega_0 t} \right), \tag{III.9}$$

$$= -\frac{F_0}{m} \frac{\cos \omega t - \cos \omega_0 t}{\omega_0^2 - \omega^2}.$$
 (III.10)

This does not work at resonance,  $\omega = \omega_0$ , since it tends towards infinity. It must be modified from the start for that case

## IV SAKURAI 5.33

$$F(t) = \frac{F_0 \tau / \omega}{\tau^2 + t^2}, \qquad t \in (-\infty, \infty).$$
 (IV.1)

$$\therefore V(t) = -F(t)x. \tag{IV.2}$$

$$c_1^{(1)}(\infty) = \frac{i}{\hbar} \frac{F_0 \tau}{\omega} \langle 1 | \hat{x} | 0 \rangle \int_{-\infty}^{\infty} \frac{e^{i\omega t}}{\tau^2 + t^2} dt, \qquad (IV.3)$$

$$= \frac{i}{\hbar} \frac{F_0 \tau}{\omega} \sqrt{\frac{\hbar}{2m\omega}} \frac{\pi}{\tau} e^{-\omega \tau}. \tag{IV.4}$$

This makes sense in the limit  $\tau \gg 1/\omega$ . The perturbation is essentially turned on then off very slowly, so the system stays in the ground state.