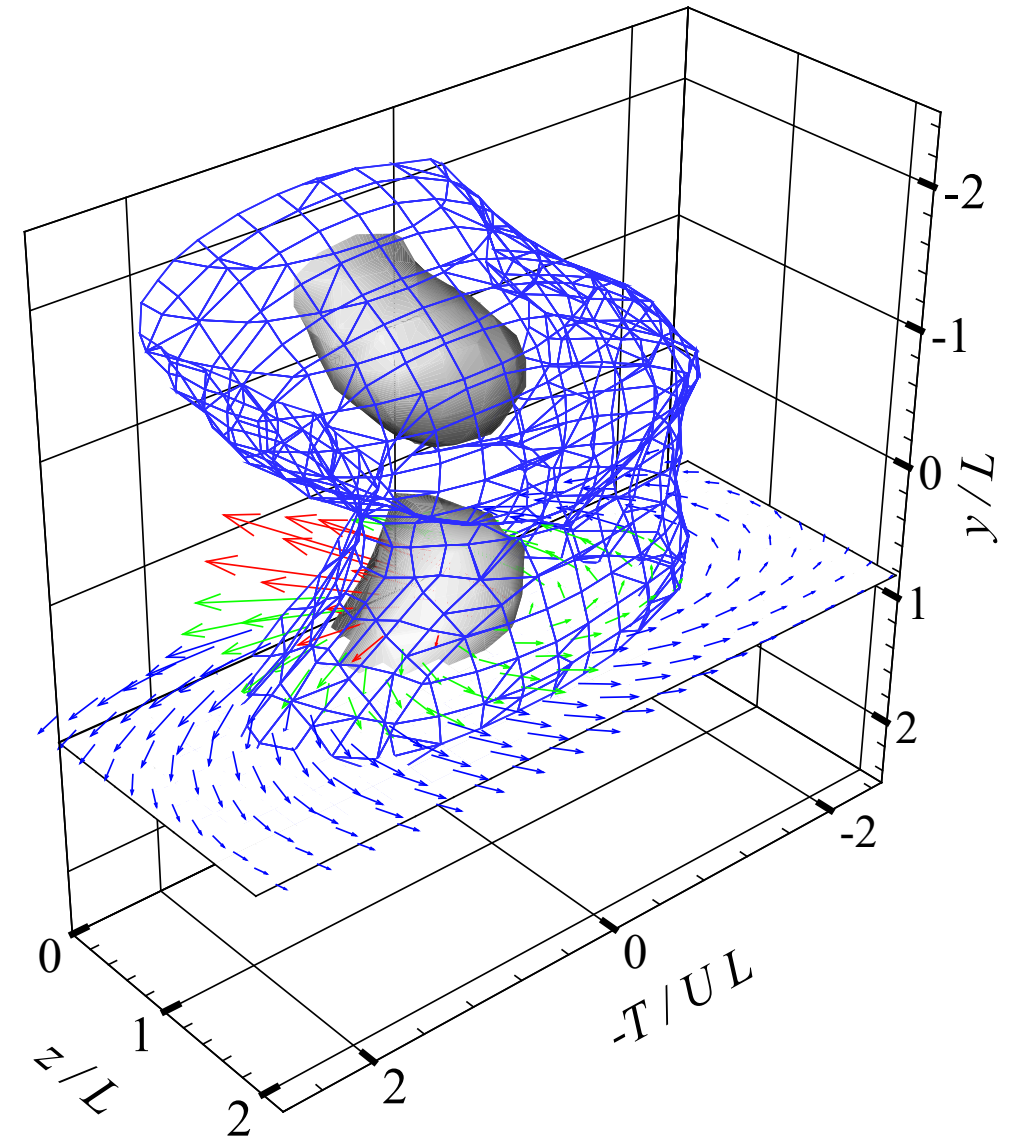


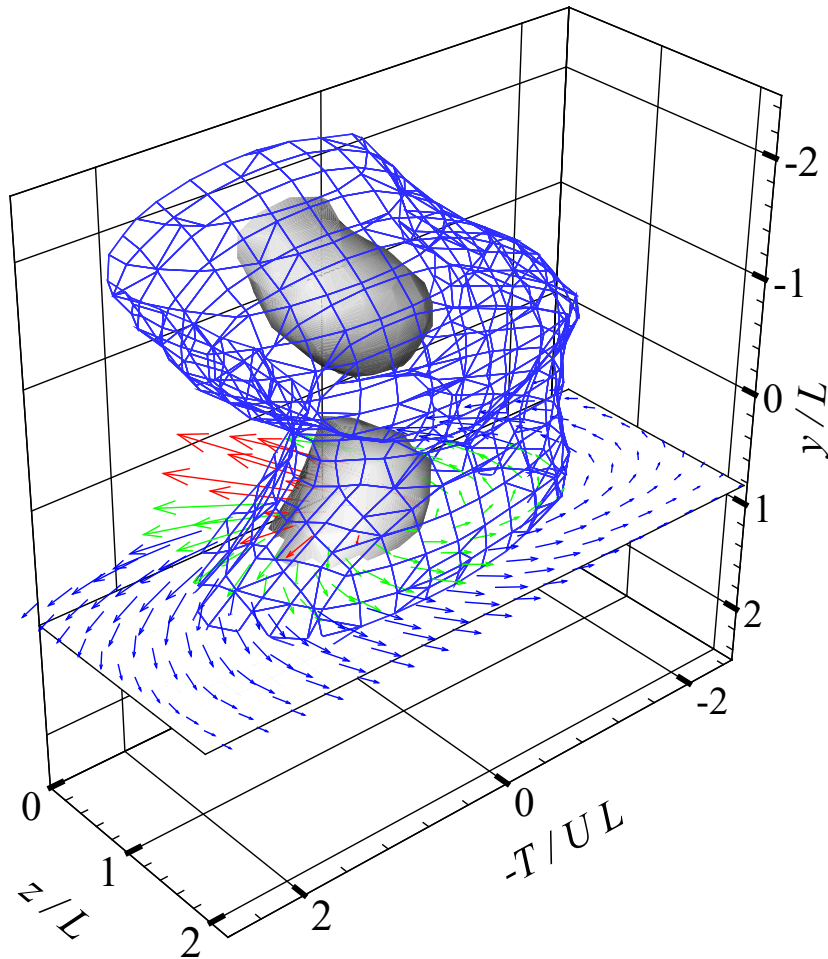
3. Vector Calculus

Fluid particle: Differentially
Small Piece of the Fluid
Material

- Concept of differential change in a vector, vector field
- Changes in unit vectors
- Calculus w.r.t. time
- Integral calculus w.r.t. space
- Differential calculus w.r.t. space
- Integral theorems, second order operators



Concept of Differential Change In a Vector. The Vector Field.



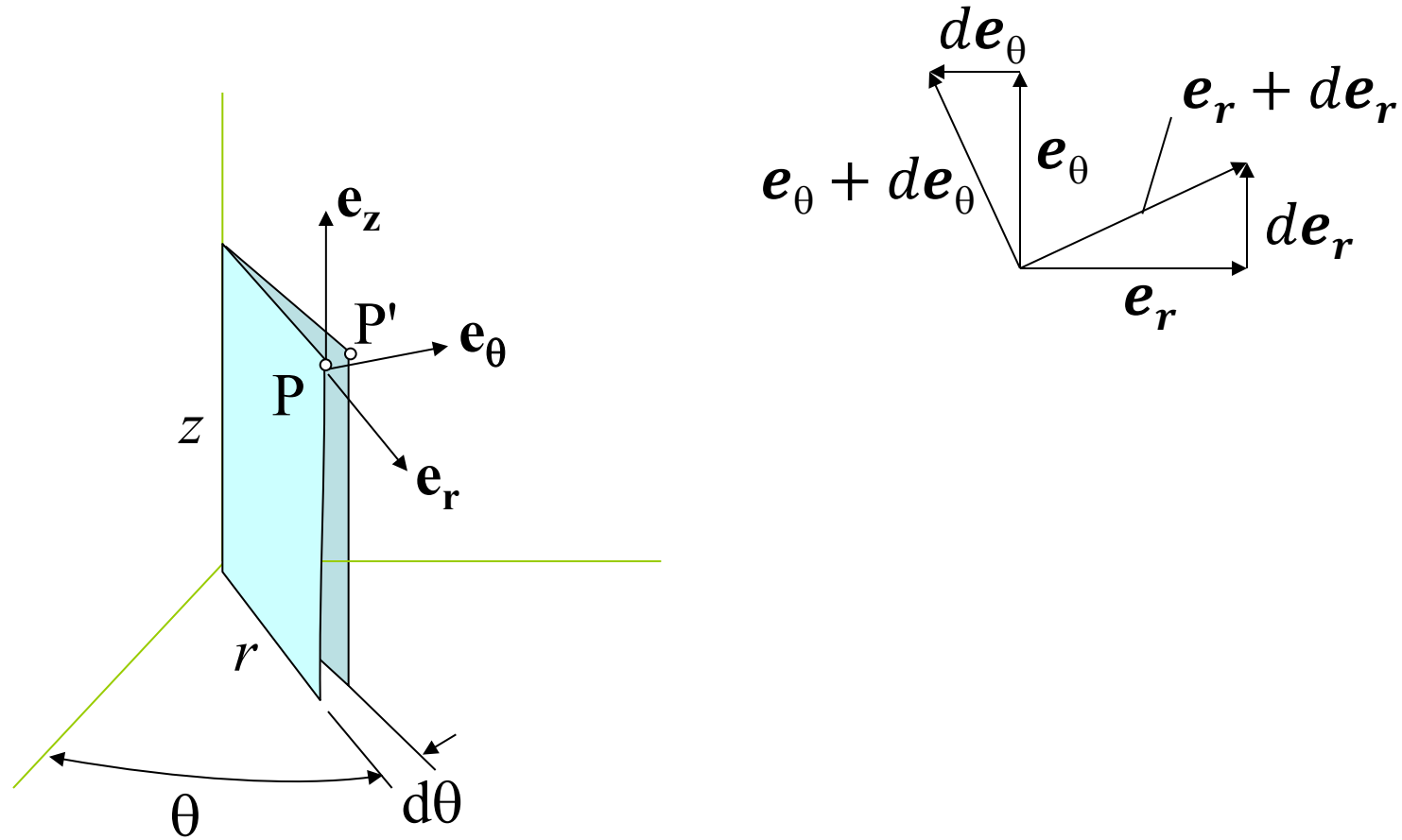
Scalar field

$$\phi = \phi(\mathbf{r}, t)$$

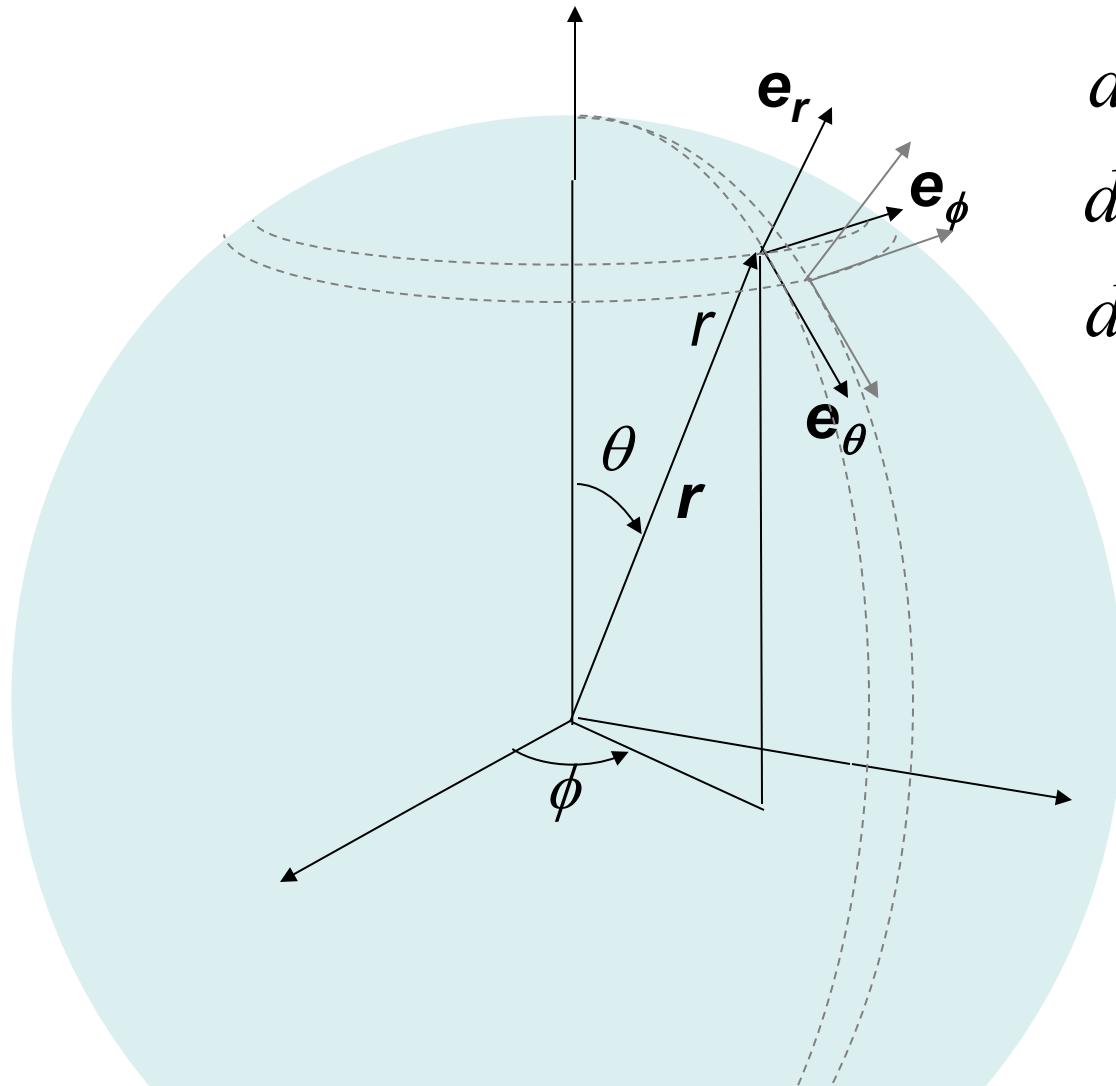
Vector field

$$\mathbf{V} = \mathbf{V}(\mathbf{r}, t)$$

Change in Unit Vectors – Cylindrical System



Change in Unit Vectors – Spherical System



$$d\mathbf{e}_r = d\theta\mathbf{e}_\theta + d\phi\sin\theta\mathbf{e}_\phi$$

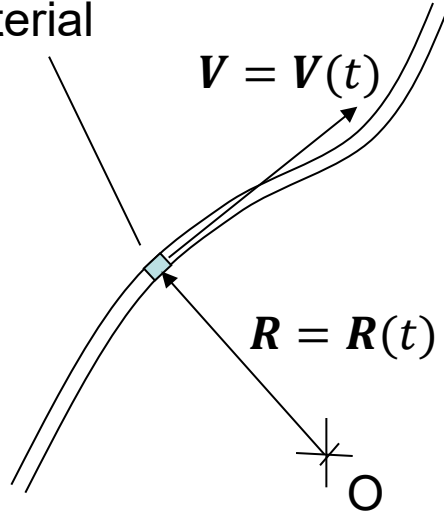
$$d\mathbf{e}_\theta = -d\theta\mathbf{e}_r + d\phi\cos\theta\mathbf{e}_\phi$$

$$d\mathbf{e}_\phi = -d\phi\sin\theta\mathbf{e}_r - d\phi\cos\theta\mathbf{e}_\theta$$

See “Formulae for Vector
Algebra and Calculus”

Fluid particle

Differentially small
piece of the fluid
material

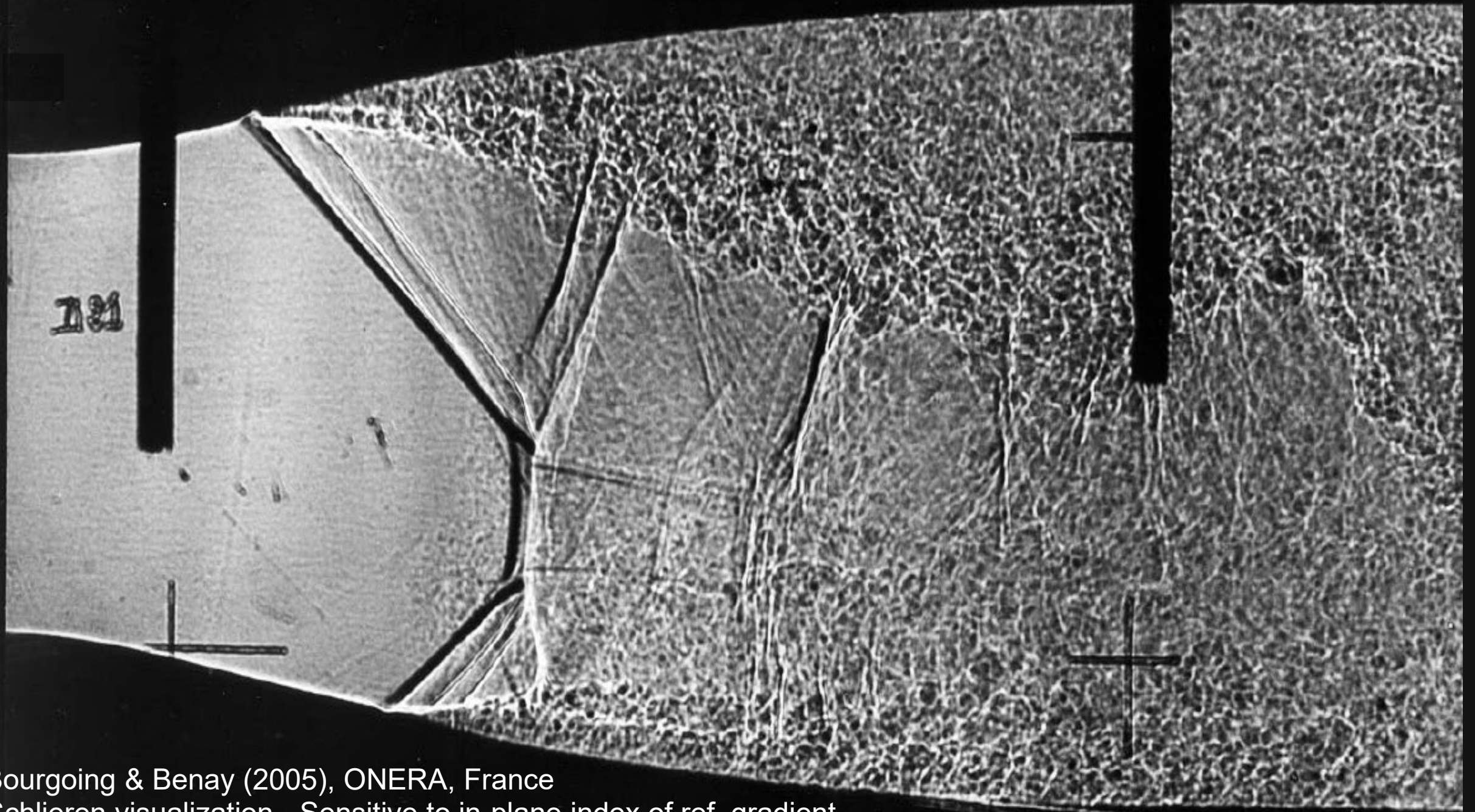


Example

The position of fluid particle moving in a flow varies with time. Working in different coordinate systems write down expressions for the position and, by differentiation, the velocity vectors.

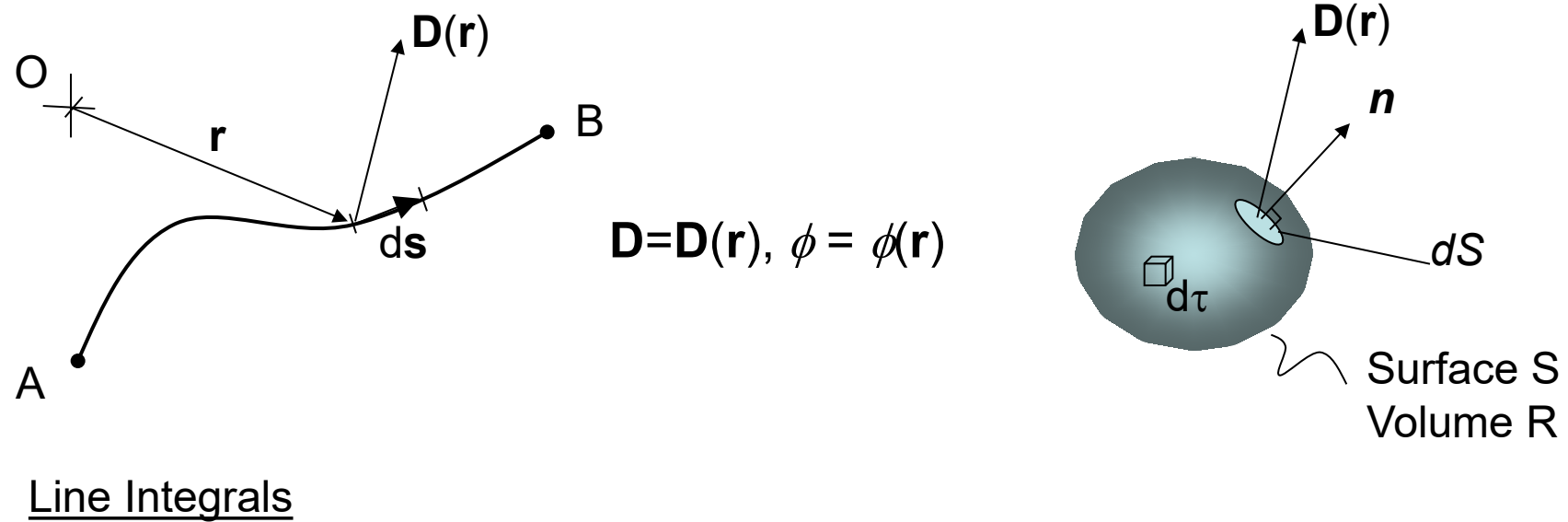
Vector Calculus w.r.t. Time

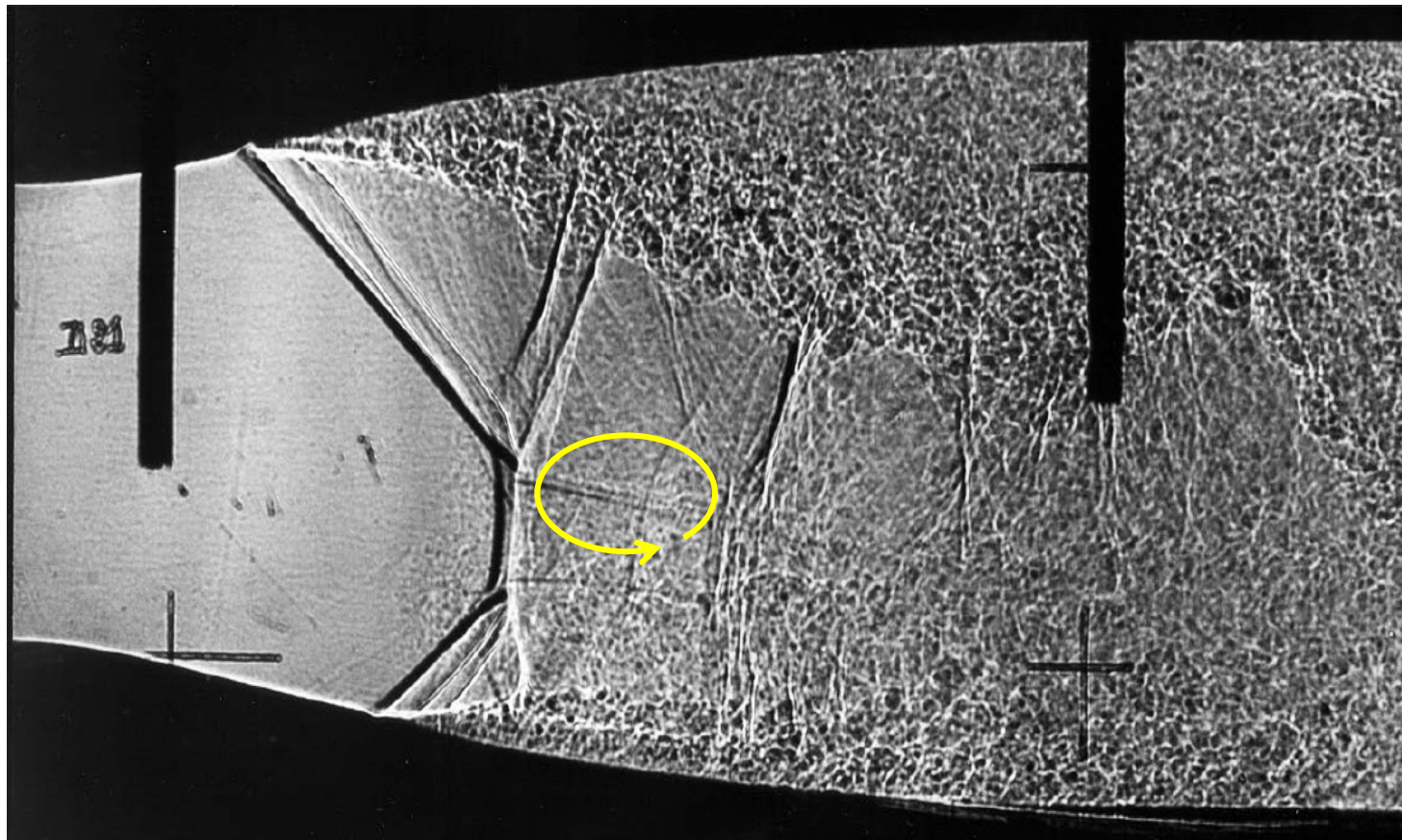
- Since *any* vector may be decomposed into scalar components, calculus w.r.t. time, only involves *scalar* calculus of the components



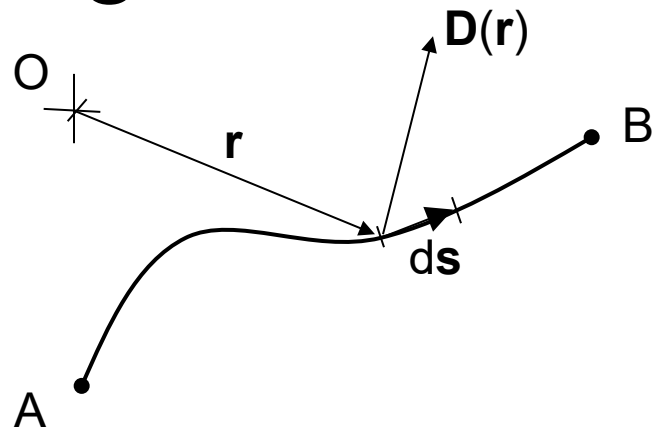
Bourgoing & Benay (2005), ONERA, France
Schlieren visualization - Sensitive to in-plane index of ref. gradient

Integral Calculus With Respect to Space

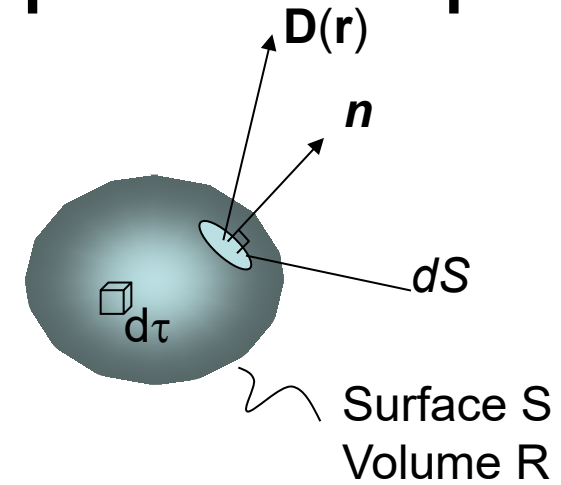




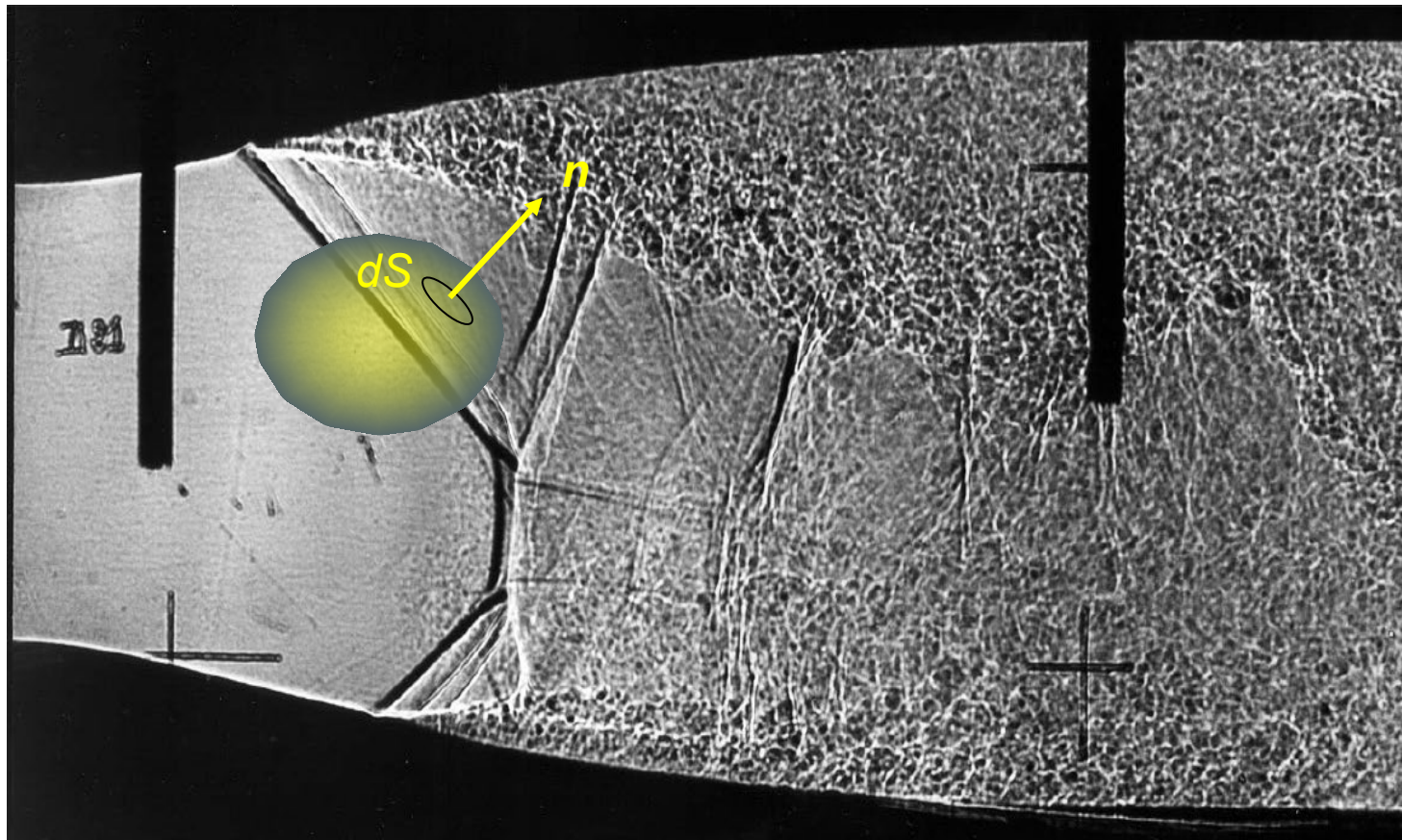
Integral Calculus With Respect to Space



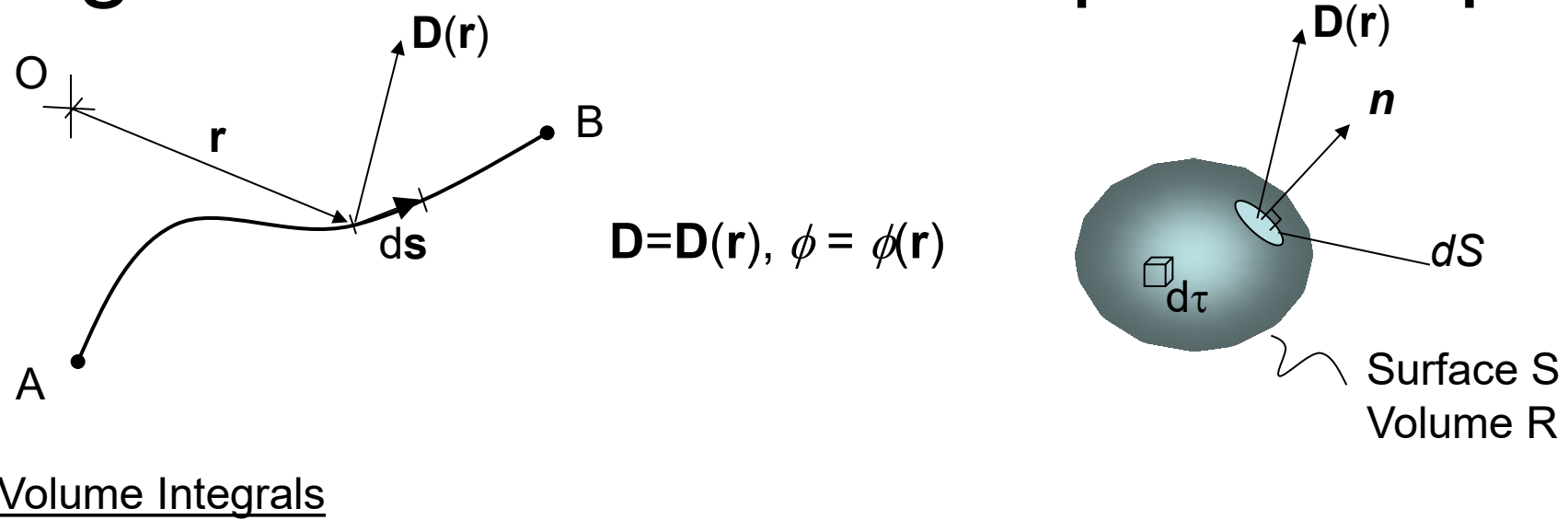
$$\mathbf{D} = \mathbf{D}(\mathbf{r}), \phi = \phi(\mathbf{r})$$



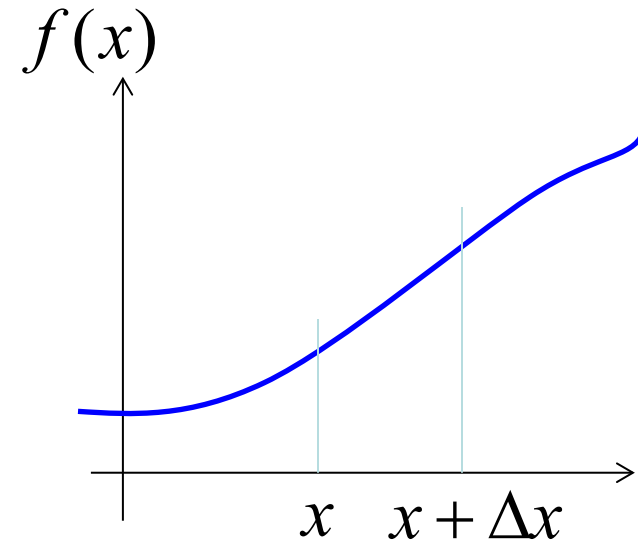
Surface Integrals



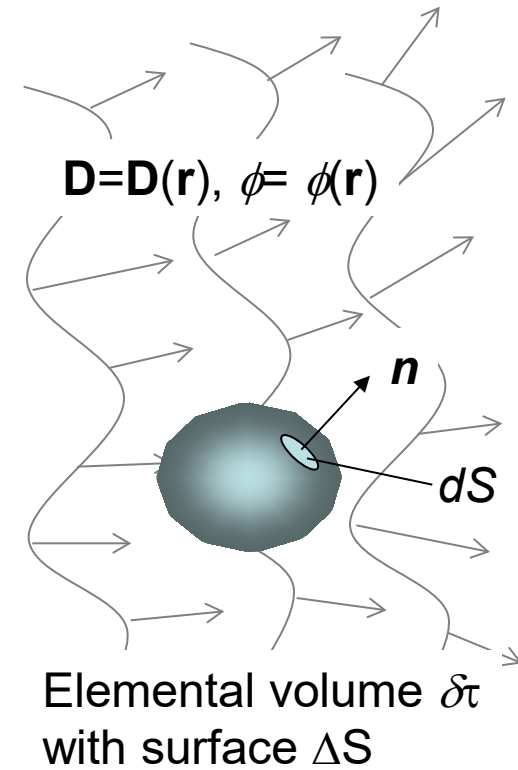
Integral Calculus With Respect to Space



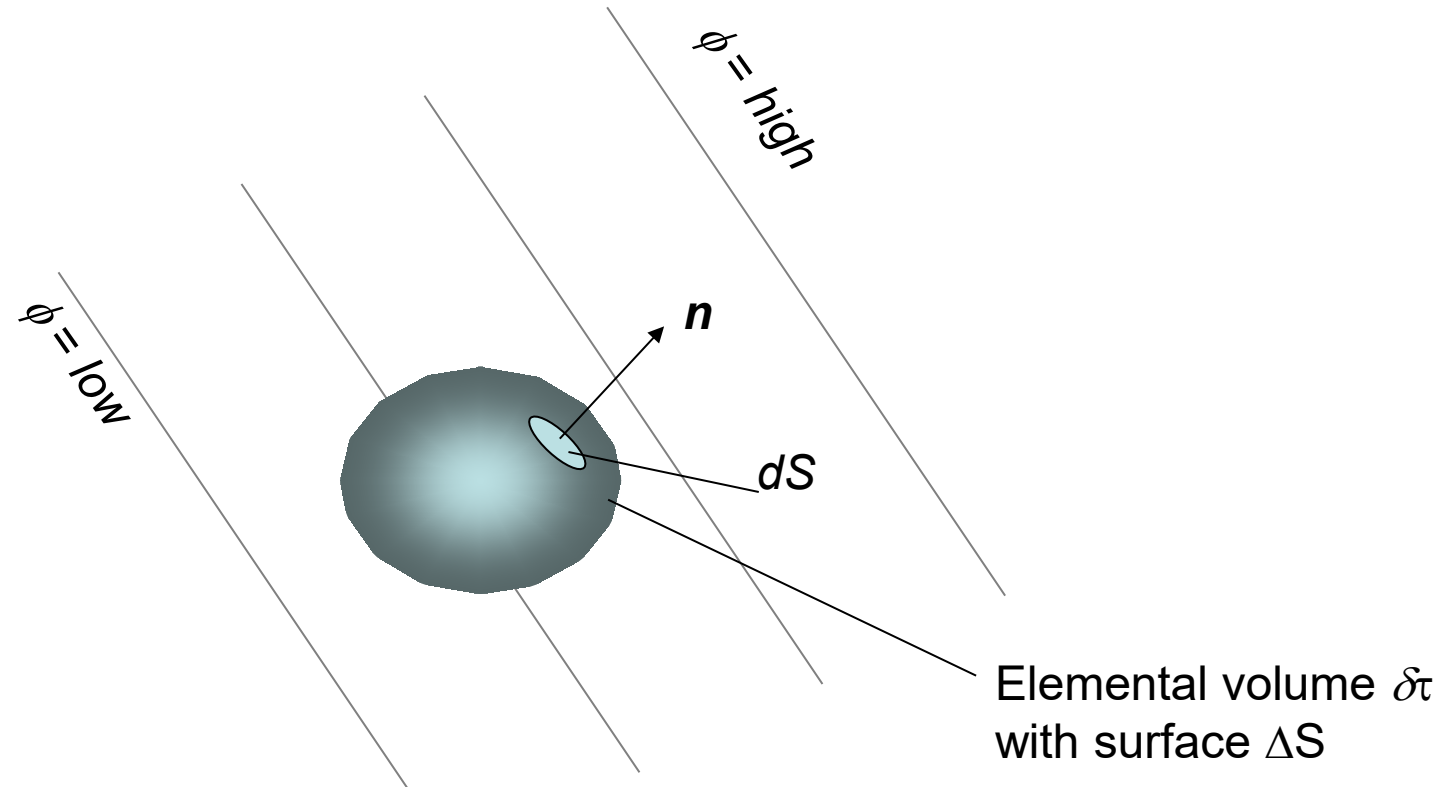
Differential Calculus w.r.t. Space in 1D



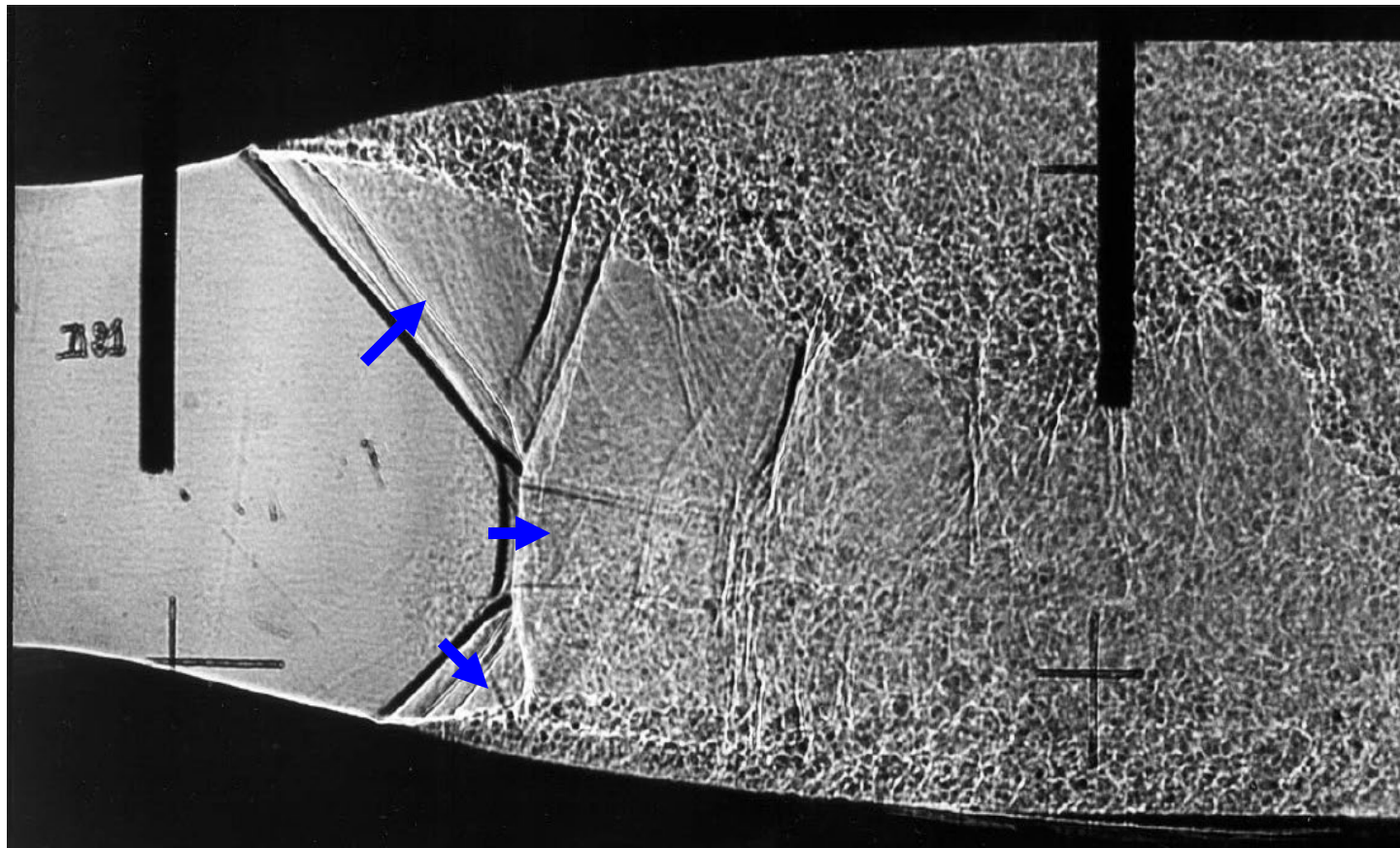
Differential Calculus w.r.t. Space in 3D



Gradient



$$\text{grad } \phi \equiv \lim_{\delta\tau \rightarrow 0} \frac{1}{\delta\tau} \oint_{\Delta S} \phi \mathbf{n} dS$$



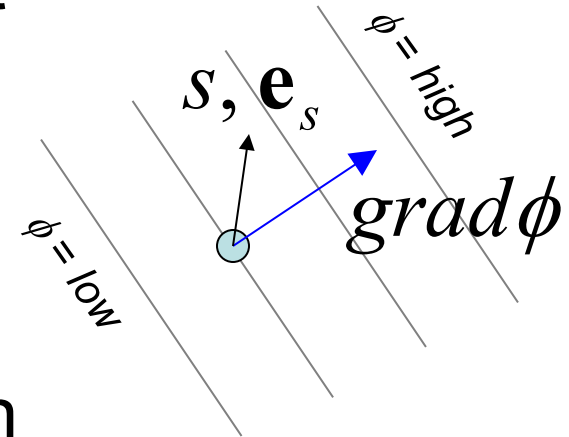
Gradient

$$\text{grad } \phi \equiv \lim_{\Delta S \rightarrow 0} \frac{1}{\Delta S} \oint \phi \, \mathbf{n} \, dS$$

Magnitude and direction
of the slope in the scalar
field at a point

Gradient

- Component of gradient is the partial derivative in the direction of that component



- Fourier's Law of Heat Conduction

Differential form of the Gradient

Cartesian system

Evaluate integral by expanding the variation in ϕ about a point P at the center of an elemental Cartesian volume. Consider the two x faces:

$$\int_{\text{Face 1}} \phi \mathbf{n} dS \approx \left(\phi - \frac{\partial \phi}{\partial x} \frac{dx}{2} \right) (-\mathbf{i}) dy dz$$

$$\int_{\text{Face 2}} \phi \mathbf{n} dS \approx \left(\phi + \frac{\partial \phi}{\partial x} \frac{dx}{2} \right) (+\mathbf{i}) dy dz$$

adding these gives $\mathbf{i} \frac{\partial \phi}{\partial x} dx dy dz$

Proceeding in the same way for y and z

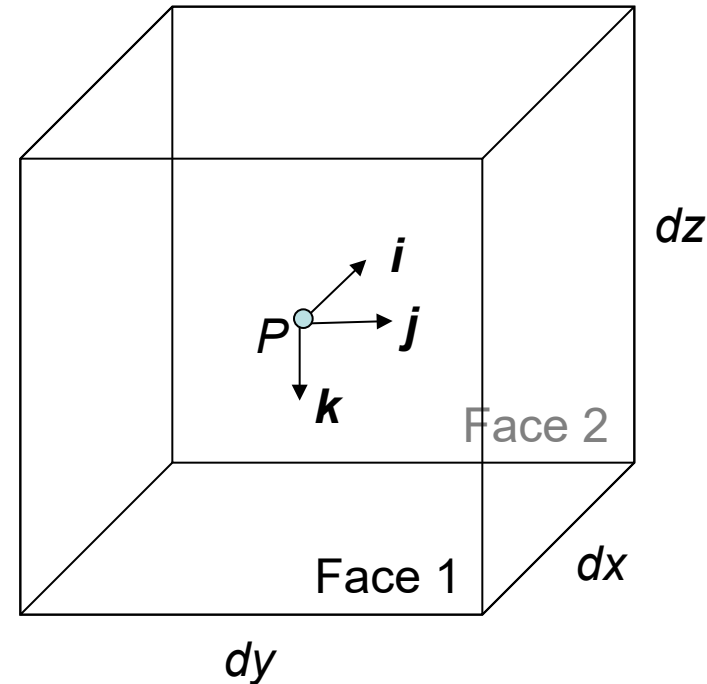
we get $\mathbf{j} \frac{\partial \phi}{\partial y} dx dy dz$ and $\mathbf{k} \frac{\partial \phi}{\partial z} dx dy dz$, so

$$\text{grad } \phi \equiv \text{Lim}_{\delta\tau \rightarrow 0} \frac{1}{\delta\tau} \oint_{\Delta S} \phi \mathbf{n} dS$$

$$= \text{Lim}_{\delta\tau \rightarrow 0} \frac{1}{\delta\tau} \left(\mathbf{i} \frac{\partial \phi}{\partial x} dx dy dz + \mathbf{j} \frac{\partial \phi}{\partial y} dx dy dz + \mathbf{k} \frac{\partial \phi}{\partial z} dx dy dz \right) = \underline{\underline{\mathbf{i} \frac{\partial \phi}{\partial x} + \mathbf{j} \frac{\partial \phi}{\partial y} + \mathbf{k} \frac{\partial \phi}{\partial z}}}$$

$$\text{grad } \phi \equiv \text{Lim}_{\delta\tau \rightarrow 0} \frac{1}{\delta\tau} \oint_{\Delta S} \phi \mathbf{n} dS$$

$$\phi = \phi(x, y, z)$$



Differential Forms of the Gradient

$$\text{grad}\Phi = \nabla\Phi$$

Cartesian

$$\vec{i} \frac{\partial\Phi}{\partial x} + \vec{j} \frac{\partial\Phi}{\partial y} + \vec{k} \frac{\partial\Phi}{\partial z} = \left(\vec{i} \frac{\partial}{\partial x} + \vec{j} \frac{\partial}{\partial y} + \vec{k} \frac{\partial}{\partial z} \right) \Phi$$

Cylindrical

$$\vec{e}_r \frac{\partial\Phi}{\partial r} + \frac{\vec{e}_\theta}{r} \frac{\partial\Phi}{\partial \theta} + \vec{e}_z \frac{\partial\Phi}{\partial z} = \left(\vec{e}_r \frac{\partial}{\partial r} + \frac{\vec{e}_\theta}{r} \frac{\partial}{\partial \theta} + \vec{e}_z \frac{\partial}{\partial z} \right) \Phi$$

Spherical

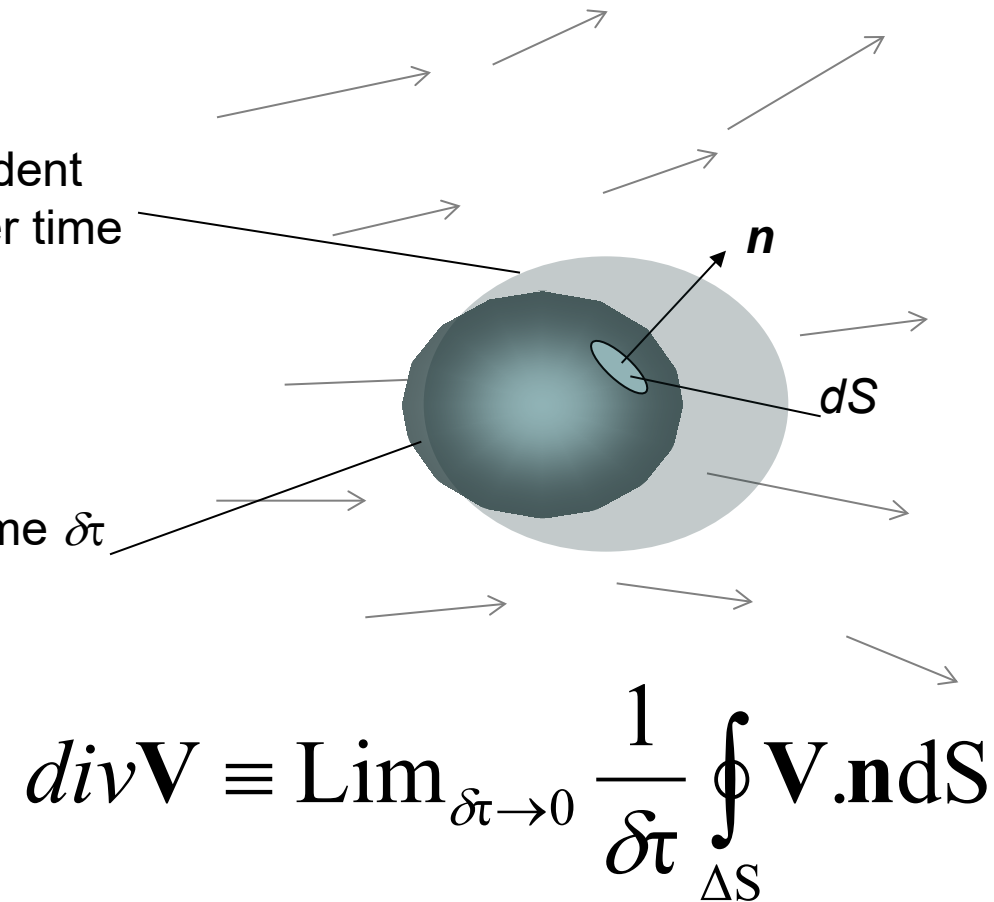
$$\vec{e}_r \frac{\partial\Phi}{\partial r} + \frac{\vec{e}_\theta}{r} \frac{\partial\Phi}{\partial \theta} + \frac{\vec{e}_\phi}{r \sin \theta} \frac{\partial\Phi}{\partial \phi} = \left(\vec{e}_r \frac{\partial}{\partial r} + \frac{\vec{e}_\theta}{r} \frac{\partial}{\partial \theta} + \frac{\vec{e}_\phi}{r \sin \theta} \frac{\partial}{\partial \phi} \right) \Phi$$

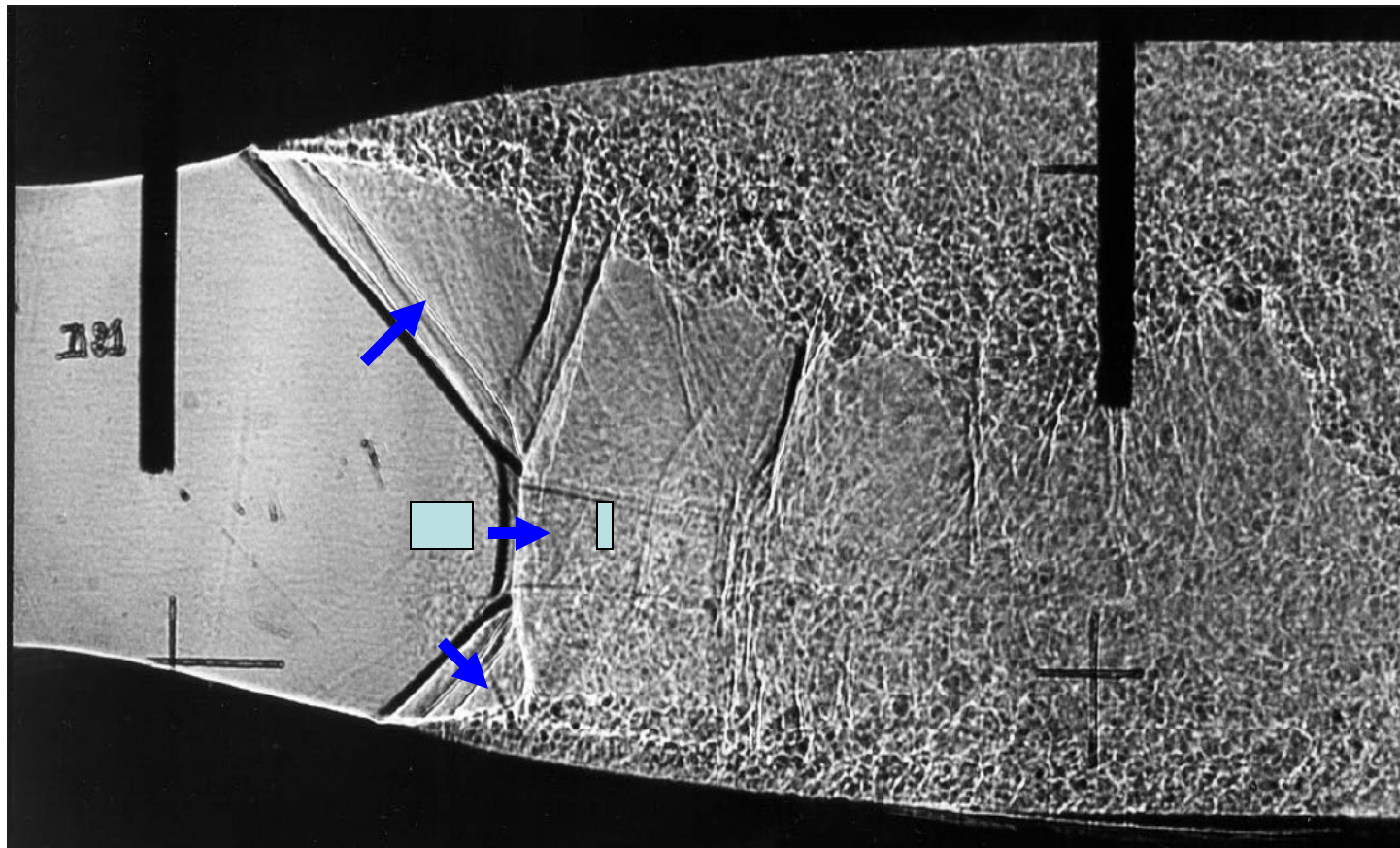
These differential forms define the vector operator ∇

Divergence

Fluid particle, coincident with $\delta\tau$ at time t , after time δt has elapsed.

Elemental volume $\delta\tau$ with surface ΔS





Gradient

$$\text{grad } \phi \equiv \lim_{\delta\tau \rightarrow 0} \frac{1}{\delta\tau} \oint_{\Delta S} \phi \mathbf{n} dS$$

Magnitude and direction
of the slope in the scalar
field at a point

Divergence

$$\text{div } \mathbf{V} \equiv \lim_{\delta\tau \rightarrow 0} \frac{1}{\delta\tau} \oint_{\Delta S} \mathbf{V} \cdot \mathbf{n} dS$$

For velocity: proportionate
rate of change of volume
of a fluid particle

Differential Forms of the Divergence

$$\operatorname{div} \vec{A} = \nabla \cdot \vec{A}$$

Cartesian

$$\frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z} = \left(\vec{i} \frac{\partial}{\partial x} + \vec{j} \frac{\partial}{\partial y} + \vec{k} \frac{\partial}{\partial z} \right) \cdot \vec{A}$$

Cylindrical

$$\frac{1}{r} \frac{\partial}{\partial r} (r A_r) + \frac{1}{r} \frac{\partial A_\theta}{\partial \theta} + \frac{\partial A_z}{\partial z} = \left(\vec{e}_r \frac{\partial}{\partial r} + \frac{\vec{e}_\theta}{r} \frac{\partial}{\partial \theta} + \vec{e}_z \frac{\partial}{\partial z} \right) \cdot \vec{A}$$

Spherical

$$\frac{1}{r^2} \frac{\partial}{\partial r} (r^2 A_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (A_\theta \sin \theta) + \frac{1}{r \sin \theta} \frac{\partial A_\phi}{\partial \phi} = \left(\vec{e}_r \frac{\partial}{\partial r} + \frac{\vec{e}_\theta}{r} \frac{\partial}{\partial \theta} + \frac{\vec{e}_\phi}{r \sin \theta} \frac{\partial}{\partial \phi} \right) \cdot \vec{A}$$

Differential Forms of the Curl

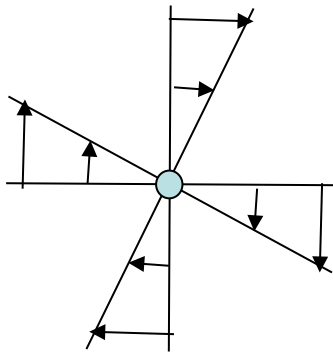
$$\text{curl} \vec{A} \equiv -\lim_{\delta\tau \rightarrow 0} \frac{1}{\delta\tau} \oint_{\Delta S} \vec{A} \times \vec{n} dS$$

$$\text{curl} \vec{A} = \nabla \times A = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ A_x & A_y & A_z \end{vmatrix} = \frac{1}{r} \begin{vmatrix} \vec{e}_r & r \vec{e}_\theta & \vec{e}_z \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial z} \\ A_r & r A_\theta & A_z \end{vmatrix} = \frac{1}{r^2 \sin \theta} \begin{vmatrix} \vec{e}_r & r \vec{e}_\theta & r \sin \theta \vec{e}_\phi \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial \phi} \\ A_r & r A_\theta & r \sin \theta A_\phi \end{vmatrix}$$

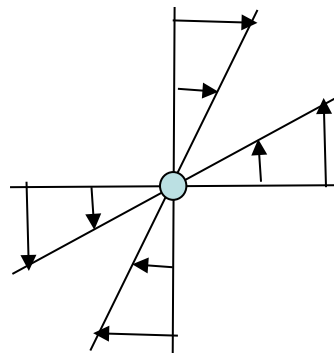
Cartesian
Cylindrical
Spherical

Physical Interpretation of the Curl

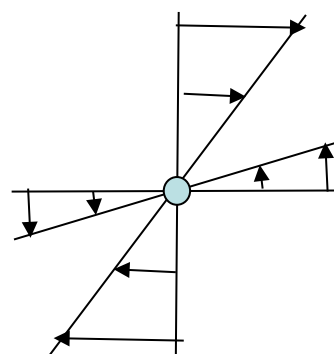
$$\nabla \times \mathbf{V} =$$



Pure rotation



No rotation



Rotation

Curl

$$\text{curl} \mathbf{V} \equiv -\lim_{\delta\tau \rightarrow 0} \frac{1}{\delta\tau} \oint_{\Delta S} \mathbf{V} \times \mathbf{n} dS$$

$$\mathbf{e} \cdot \text{curl} \mathbf{V} \equiv -\lim_{\delta\tau \rightarrow 0} \frac{1}{\delta\tau} \oint_{\Delta S} \mathbf{e} \cdot \mathbf{V} \times \mathbf{n} dS$$

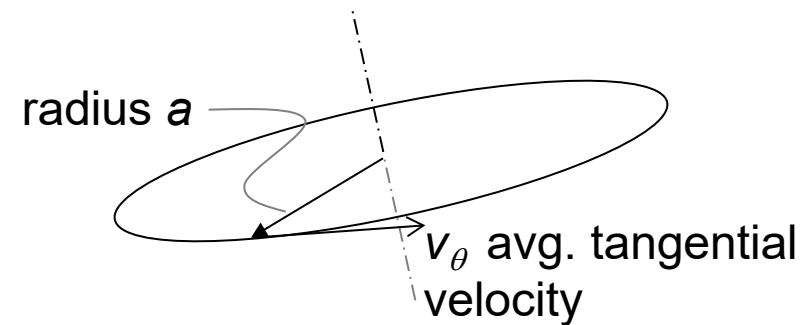
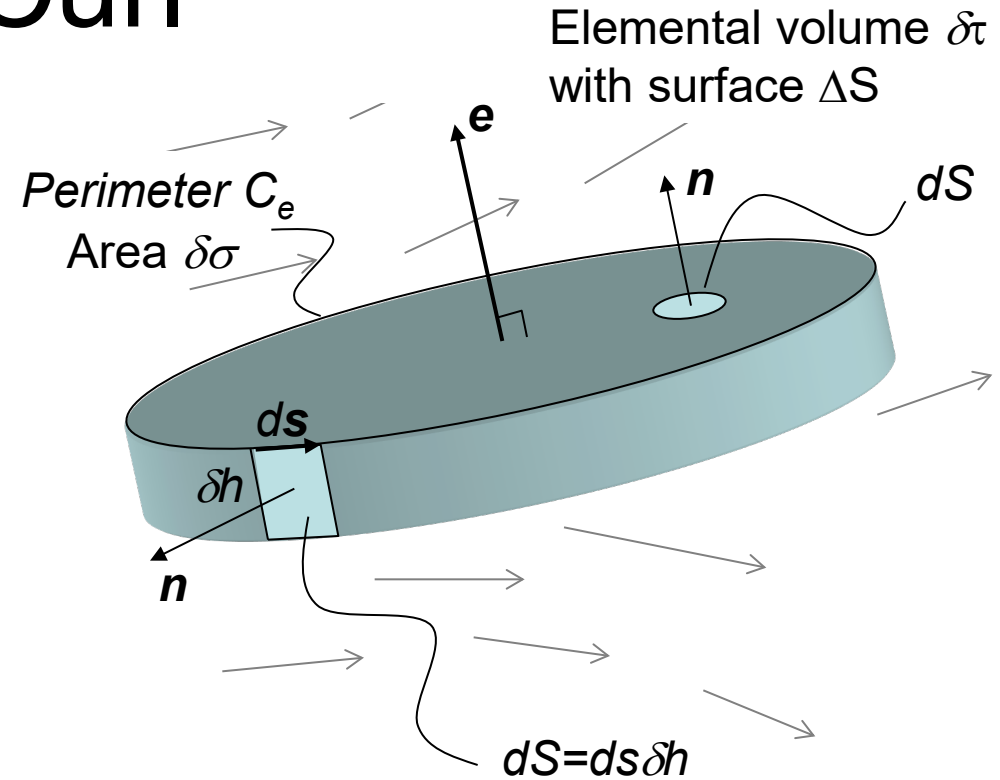
$$\mathbf{e} \cdot \text{curl} \mathbf{V} \equiv \lim_{\delta\tau \rightarrow 0} \frac{1}{\delta\sigma \delta h} \oint_{\Delta S} \mathbf{V} \cdot \mathbf{e} \times \mathbf{n} dS$$

$$\mathbf{e} \cdot \text{curl} \mathbf{V} \equiv \lim_{\delta\tau \rightarrow 0} \frac{1}{\delta\sigma \delta h} \oint_{\Delta S} \mathbf{V} \cdot \mathbf{e} \times \mathbf{n} ds \delta h$$

$$\mathbf{e} \cdot \text{curl} \mathbf{V} \equiv \lim_{\delta\sigma \rightarrow 0} \frac{1}{\delta\sigma} \oint_{C_e} \mathbf{V} \cdot d\mathbf{s} = \lim_{\delta\sigma \rightarrow 0} \frac{\Gamma_{C_e}}{\delta\sigma}$$

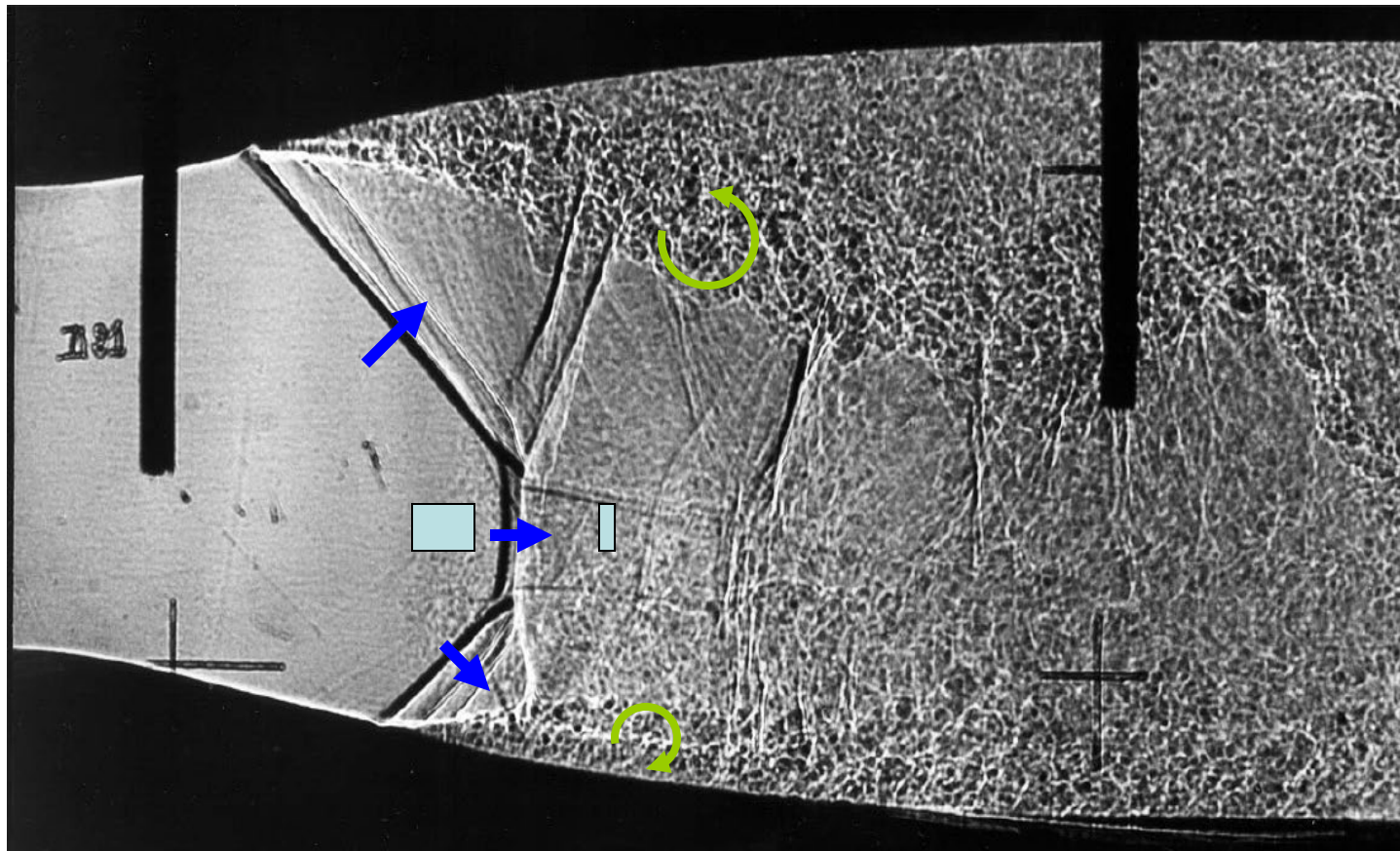
$$\mathbf{e} \cdot \text{curl} \mathbf{V} \equiv \lim_{a \rightarrow 0} \frac{1}{\pi a^2} v_\theta 2\pi a = 2 \lim_{a \rightarrow 0} \frac{v_\theta}{a}$$

= twice the avg. angular velocity
about \mathbf{e}



Curl

$$\mathbf{e}.curl\mathbf{V} \equiv \lim_{\delta\sigma \rightarrow 0} \frac{1}{\delta\sigma} \oint_{C_e} \mathbf{V} \cdot d\mathbf{s} = \lim_{\delta\sigma \rightarrow 0} \frac{\Gamma_{Ce}}{\delta\sigma}$$



Gradient

$$\text{grad} \phi \equiv \lim_{\delta\tau \rightarrow 0} \frac{1}{\delta\tau} \oint_{\Delta S} \phi \mathbf{n} dS$$

Magnitude and direction of the slope in the scalar field at a point

Divergence

$$\text{div} \mathbf{V} \equiv \lim_{\delta\tau \rightarrow 0} \frac{1}{\delta\tau} \oint_{\Delta S} \mathbf{V} \cdot \mathbf{n} dS$$

For velocity: proportionate rate of change of volume of a fluid particle

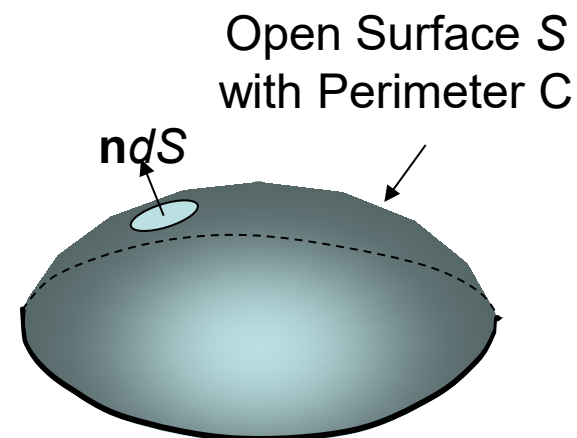
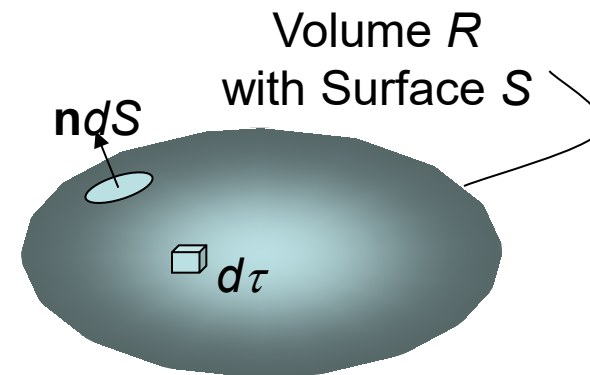
Curl

$$\text{curl} \mathbf{V} \equiv -\lim_{\delta\tau \rightarrow 0} \frac{1}{\delta\tau} \oint_{\Delta S} \mathbf{V} \times \mathbf{n} dS$$

For velocity: twice the circumferentially averaged angular velocity of a fluid particle = Vorticity $\boldsymbol{\Omega}$

1st Order Integral Theorems

- Gradient theorem
- Divergence theorem
- Curl theorem
- Stokes' theorem

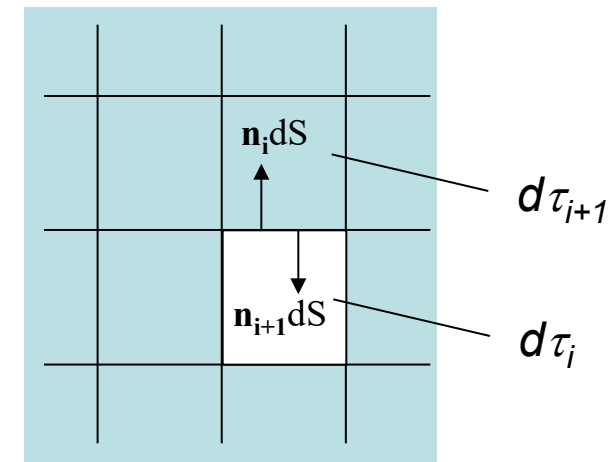
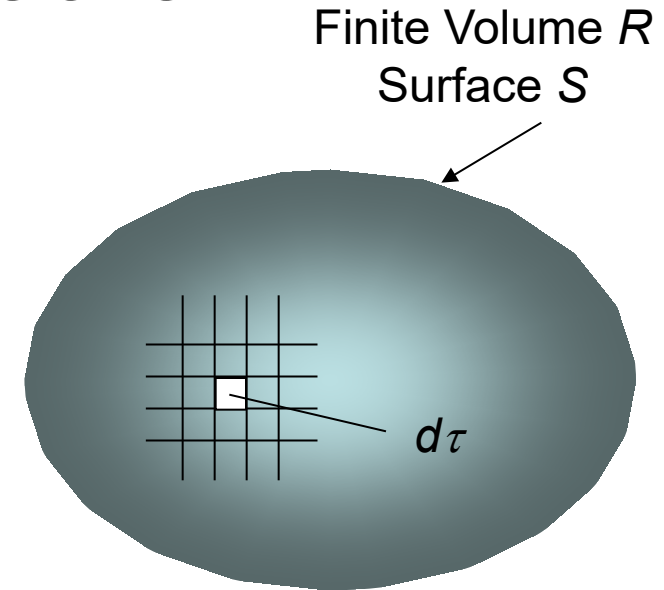


The Gradient Theorem

Begin with the definition of grad:

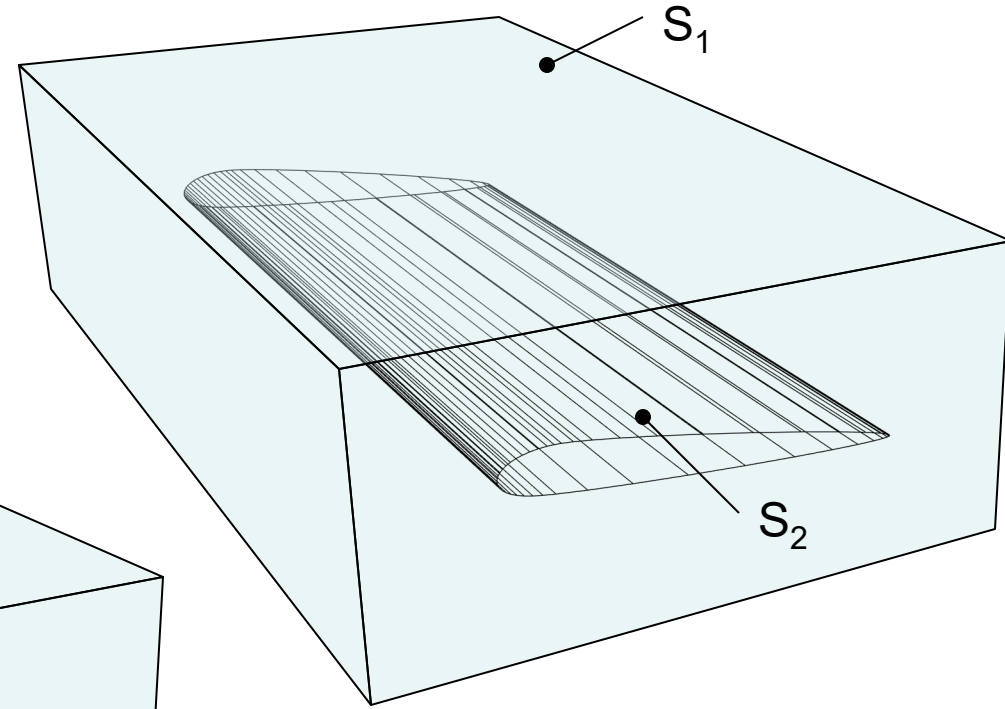
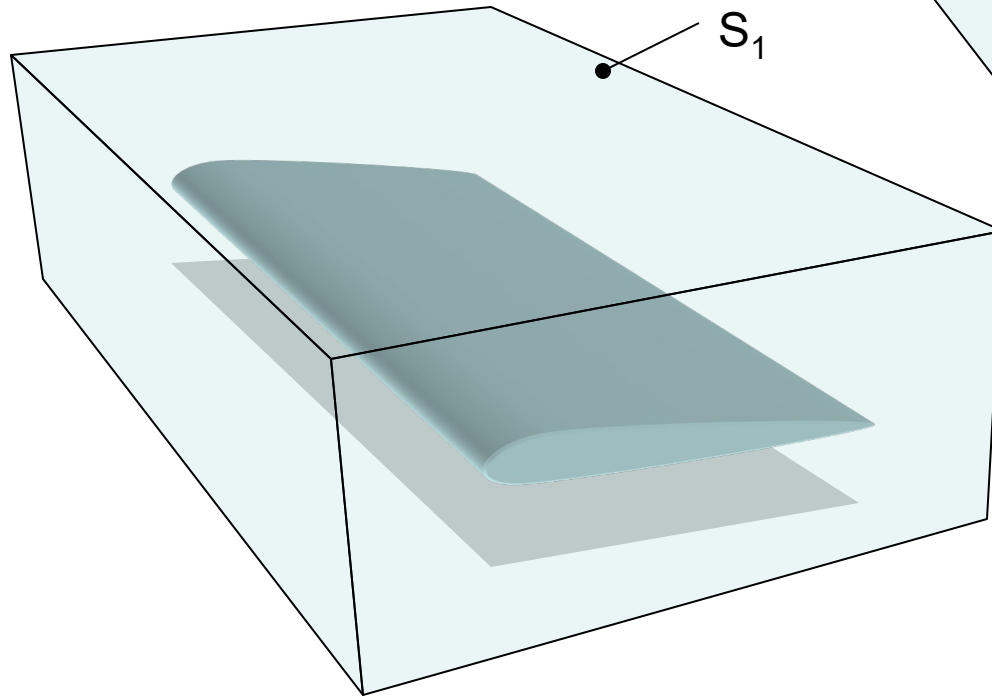
$$\text{grad} \phi \equiv \lim_{\delta \tau \rightarrow 0} \frac{1}{\delta \tau} \oint_{\Delta S} \phi \mathbf{n} dS$$

Sum over all the $d\tau$ in R:



Assumptions in Gradient Theorem

Flow over a finite wing



$$S = S_1 + S_2$$

R is the volume of fluid enclosed between S_1 and S_2

$$\int_R \nabla p d\tau = \int_S p \mathbf{n} dS$$

p is not defined inside the wing so the wing itself must be excluded from the integral

1st Order Integral Theorems

- Gradient theorem

$$\int_R \nabla \phi d\tau = \oint_S \phi \mathbf{n} dS$$

- Divergence theorem

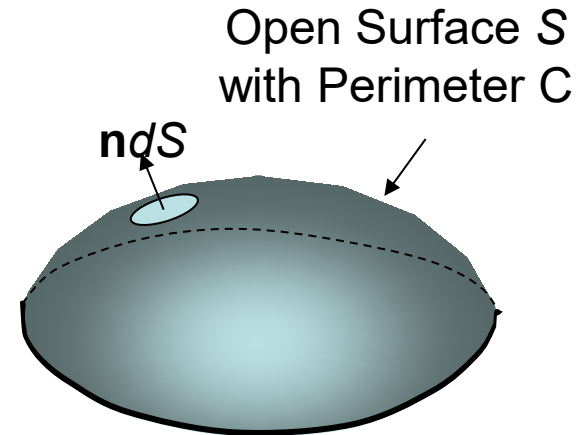
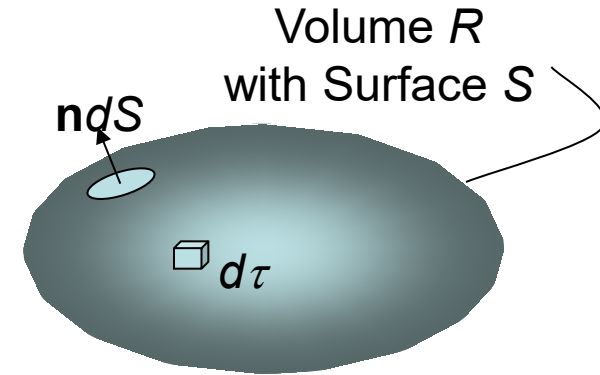
$$\int_R \nabla \cdot \mathbf{A} d\tau = \oint_S \mathbf{A} \cdot \mathbf{n} dS$$

- Curl theorem

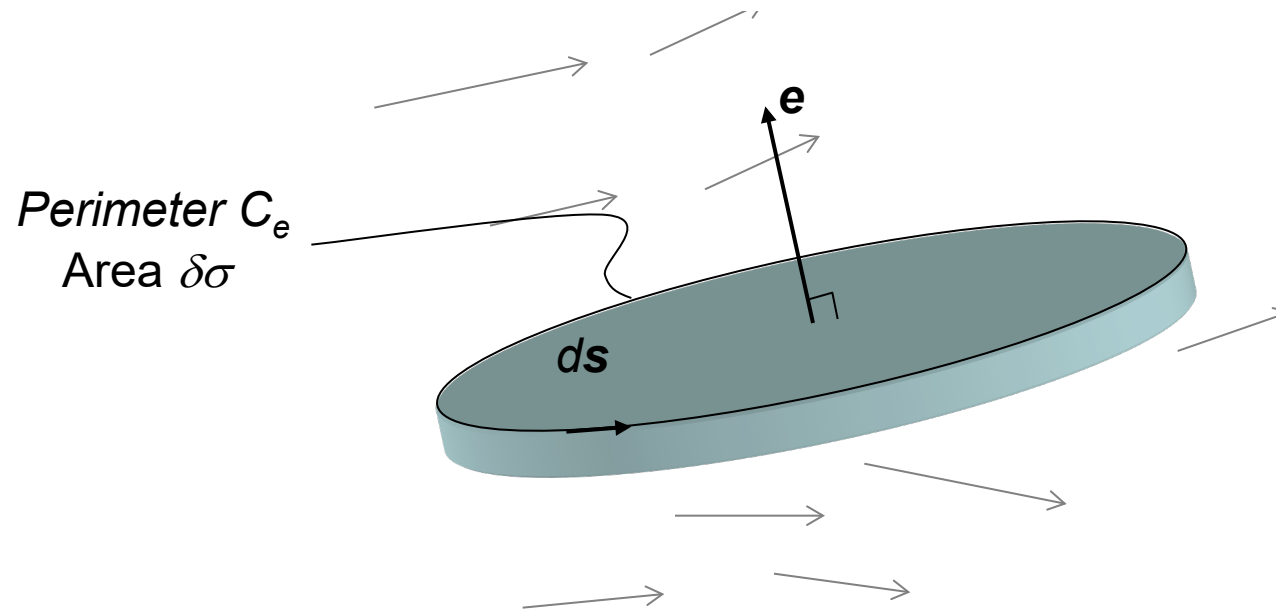
$$\int_R \nabla \times \mathbf{A} d\tau = - \oint_S \mathbf{A} \times \mathbf{n} dS$$

- Stokes' theorem

$$\int_S \nabla \times \mathbf{A} \cdot \mathbf{n} dS = \oint_C \mathbf{A} \cdot d\mathbf{s}$$



Alternative Definition of the Curl



$$\mathbf{e} \cdot \text{curl} \mathbf{A} \equiv \lim_{\delta\sigma \rightarrow 0} \frac{1}{\delta\sigma} \oint_{C_e} \mathbf{A} \cdot d\mathbf{s} = \lim_{\delta\sigma \rightarrow 0} \frac{\Gamma_{C_e}}{\delta\sigma}$$

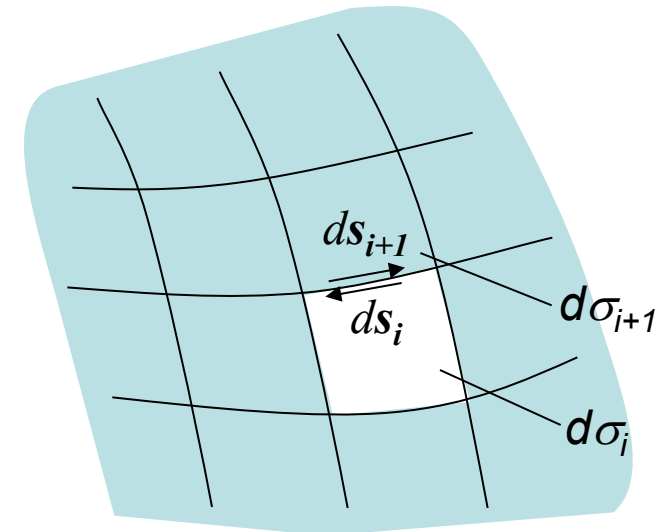
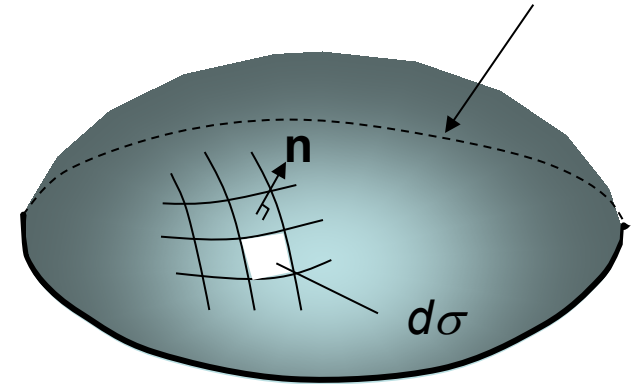
Stokes' Theorem

Begin with the alternative definition of curl, choosing the direction \mathbf{e} to be the outward normal to the surface \mathbf{n} :

$$\mathbf{n} \cdot \nabla \times \mathbf{A} \equiv \lim_{\delta\sigma \rightarrow 0} \frac{1}{\delta\sigma} \oint_{C_e} \mathbf{A} \cdot d\mathbf{s}$$

Sum over all the $d\sigma$ in S :

Finite Surface S
With Perimeter C



Stokes' Theorem and Velocity

- Apply Stokes' Theorem to a velocity field

$$\int_S \nabla \times \mathbf{V} \cdot \mathbf{n} dS = \oint_C \mathbf{V} \cdot d\mathbf{s}$$

- Or, in terms of vorticity and circulation

$$\int_S \boldsymbol{\Omega} \cdot \mathbf{n} dS = \oint_C \mathbf{V} \cdot d\mathbf{s} = \Gamma_C$$

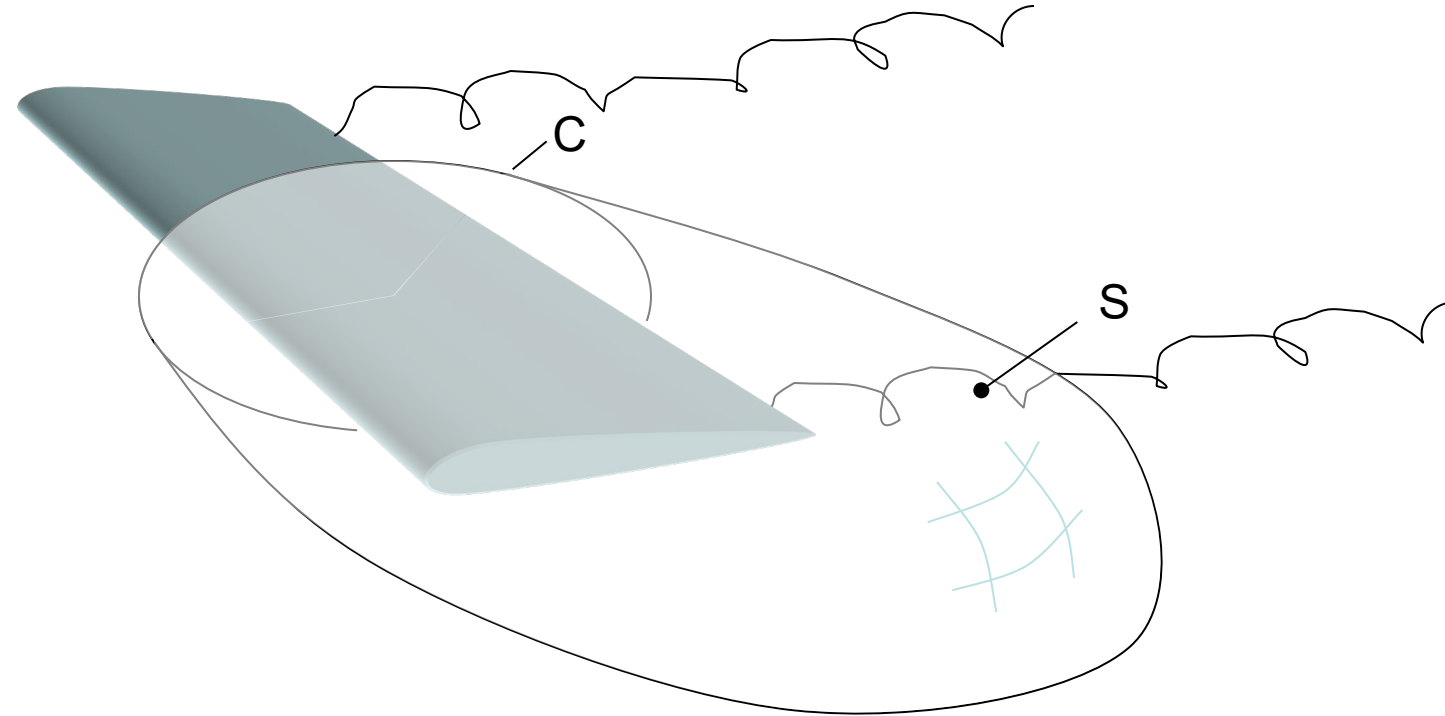
- What about a closed surface?

$$\oint_S \boldsymbol{\Omega} \cdot \mathbf{n} dS = 0$$

Assumptions of Stokes' Theorem

Flow over a finite wing

$$\int_S \nabla \times \mathbf{V} \cdot \mathbf{n} dS = \oint_C \mathbf{V} \cdot d\mathbf{s}$$



Wing with circulation must trail vorticity. *Always.*

Vector Operators of Vector Products

$$\nabla(\psi\Phi) = \psi\nabla\Phi + \Phi\nabla\psi$$

$$\nabla.(\Phi\vec{A}) = \Phi\nabla.\vec{A} + \nabla\Phi.\vec{A}$$

$$\nabla\times(\Phi\vec{A}) = \Phi\nabla\times\vec{A} + \nabla\Phi\times\vec{A}$$

$$\nabla(\vec{A}.\vec{B}) = (\vec{A}.\nabla)\vec{B} + (\vec{B}.\nabla)\vec{A} + \vec{A}\times(\nabla\times\vec{B}) + \vec{B}\times(\nabla\times\vec{A})$$

$$\nabla.(\vec{A}\times\vec{B}) = \vec{B}.\nabla\times\vec{A} - \vec{A}.\nabla\times\vec{B}$$

$$\nabla\times(\vec{A}\times\vec{B}) = \vec{A}(\nabla.\vec{B}) + (\vec{B}.\nabla)\vec{A} - \vec{B}(\nabla.\vec{A}) - (\vec{A}.\nabla)\vec{B}$$

Convective Operator

$$\begin{aligned}
 (\vec{A} \cdot \nabla) \Phi &= \left(A_x \frac{\partial}{\partial x} + A_y \frac{\partial}{\partial y} + A_z \frac{\partial}{\partial z} \right) \phi \\
 &= \vec{A} \cdot (\nabla \Phi)
 \end{aligned}$$

$\mathbf{V} \cdot \nabla \rho$ = change in density in direction of \mathbf{V} , multiplied by magnitude of \mathbf{V}

$$\begin{aligned}
 (\vec{A} \cdot \nabla) \vec{B} &= \left(A_x \frac{\partial}{\partial x} + A_y \frac{\partial}{\partial y} + A_z \frac{\partial}{\partial z} \right) \vec{B} \\
 &= \frac{1}{2} [\nabla(\vec{A} \cdot \vec{B}) - \vec{A} \times (\nabla \times \vec{B}) - \vec{B} \times (\nabla \times \vec{A}) \\
 &\quad - \nabla \times (\vec{A} \times \vec{B}) + \vec{A}(\nabla \cdot \vec{B}) - \vec{B}(\nabla \cdot \vec{A})]
 \end{aligned}$$

Second Order Operators

$$\nabla \cdot \nabla \phi = \nabla^2 \phi = \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} + \frac{\partial^2 \phi}{\partial z^2}$$

The Laplacian, may also be applied to a vector field.

$$\nabla(\nabla \cdot \mathbf{A})$$

$$\nabla \times \nabla \times \mathbf{A} = \nabla(\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A}$$

$$\nabla \times \nabla \phi \equiv 0$$

- So, any vector differential equation of the form $\nabla \times \mathbf{B} = 0$ can be solved identically by writing $\mathbf{B} = \nabla \phi$.
- We say \mathbf{B} is ***irrotational***.
- We refer to ϕ as the ***scalar potential***.

$$\nabla \cdot \nabla \times \mathbf{A} \equiv 0$$

- So, any vector differential equation of the form $\nabla \cdot \mathbf{B} = 0$ can be solved identically by writing $\mathbf{B} = \nabla \times \mathbf{A}$.
- We say \mathbf{B} is ***solenoidal*** or ***incompressible***.
- We refer to \mathbf{A} as the ***vector potential***.

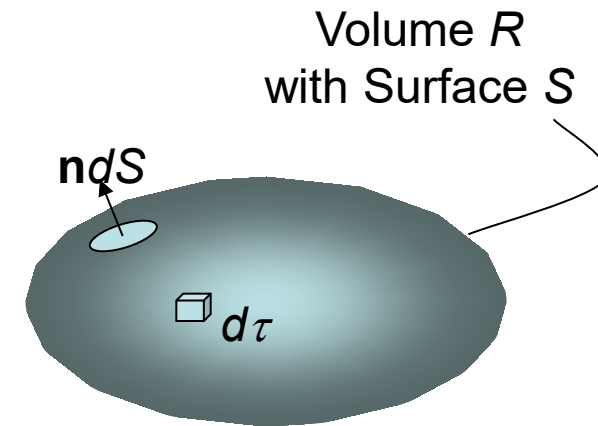
2nd Order Integral Theorems

- Green's theorem (1st form)

$$\int_R \psi \nabla^2 \phi + \nabla \psi \nabla \phi d\tau = \oint_S \psi \frac{\partial \phi}{\partial n} dS$$

- Green's theorem (2nd form)

$$\int_R \psi \nabla^2 \phi - \phi \nabla^2 \psi d\tau = \oint_S \psi \frac{\partial \phi}{\partial n} - \phi \frac{\partial \psi}{\partial n} dS$$



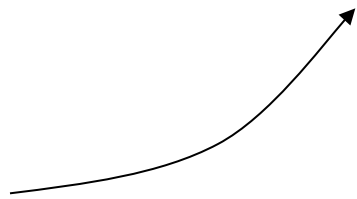
These are both re-expressions of the divergence theorem.

Helmholz Decomposition Theorem

- Any vector field may be expressed as the sum of a gradient vector field and a curl vector field.

$$\mathbf{B} = \nabla \times \mathbf{A} + \nabla \phi$$

Vector Potential



Scalar Potential

