

# Aerospace Propulsion

Lecture 12

Airbreathing Propulsion II

# Airbreathing Propulsion: Part II

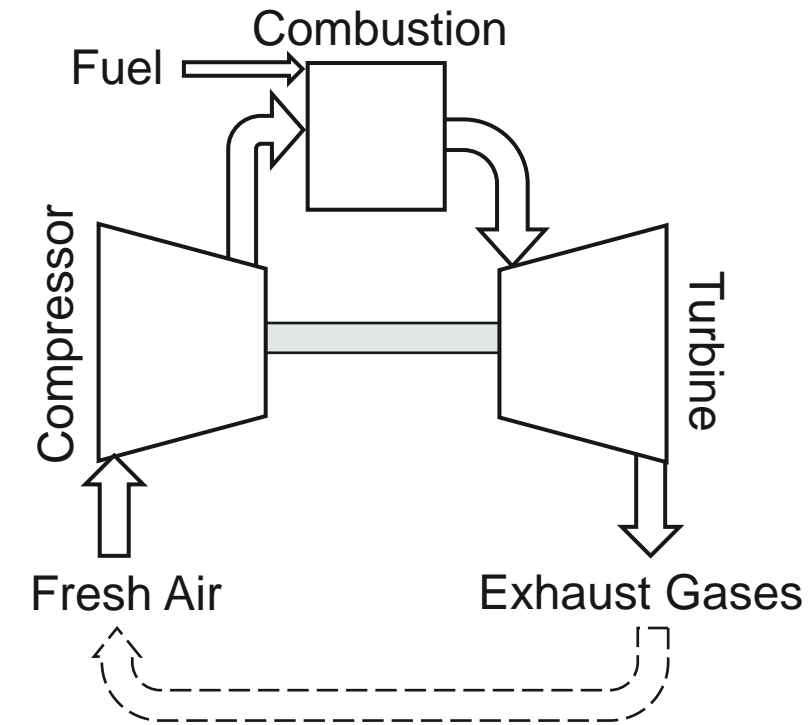
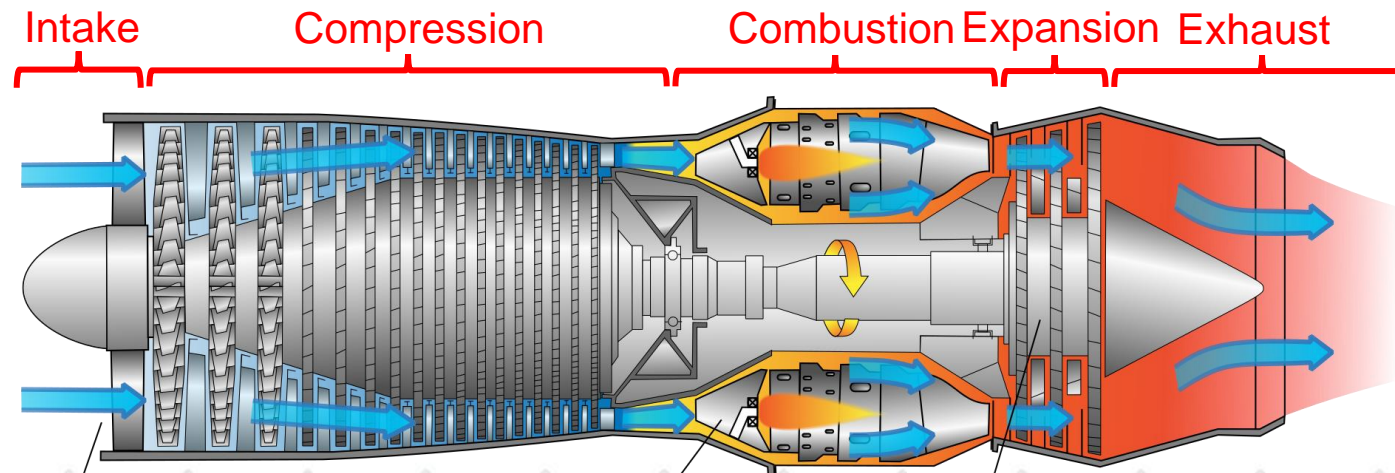
- First law for open cycles
- Ideal Brayton Cycle
- Turbojet Analysis

# First law for open cycles

- So far, we generally considered closed systems
  - First Law:  $Q_{in} - W_{out} = m\Delta h_t$
- Now, we consider open systems
  - First Law:  $\dot{Q}_{in} - \dot{W}_{out} = \dot{m}\Delta h_t$
- Keep in mind, gas turbines are constantly flowing (at steady state)
  - We care about power, heat release rate and mass flow rate

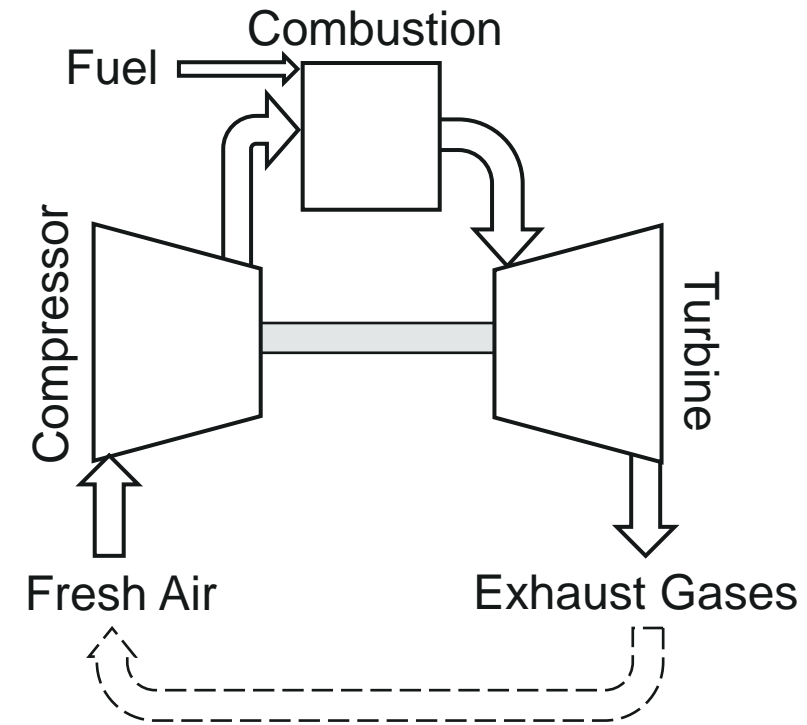
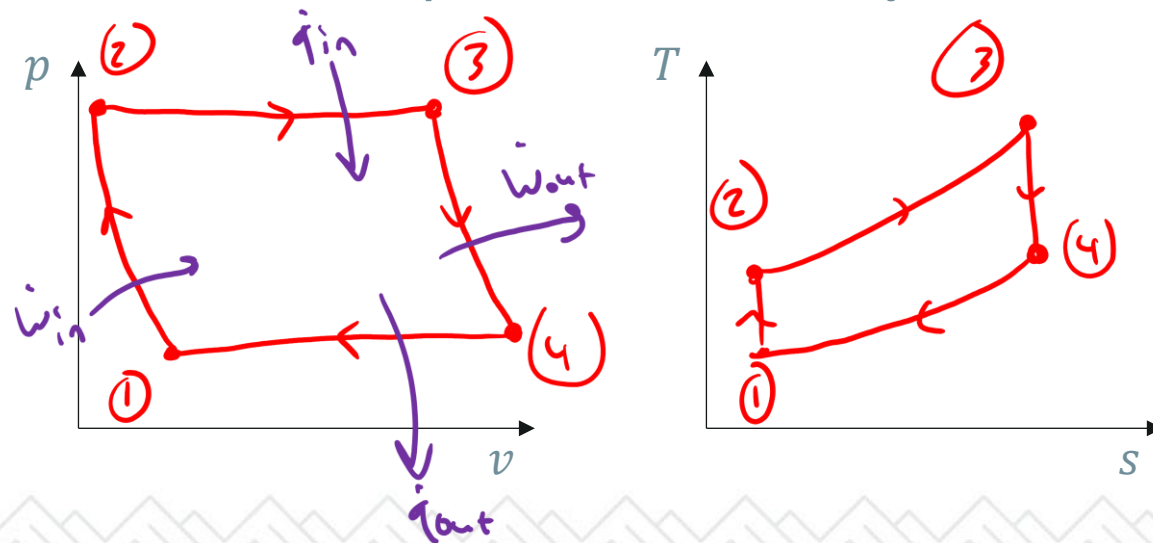
# Ideal Brayton Cycle

- Basic Turbojet represented by ideal Brayton cycle
  - 1-2: Isentropic compression
  - 2-3: Constant pressure heat addition
  - 3-4: Isentropic expansion
  - 4-1: Constant pressure heat rejection



# Ideal Brayton Cycle

- Basic Turbojet represented by ideal Brayton cycle
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Assume ideal gas

Assume constant  $c_p$

# Ideal Brayton Cycle

- Basic Turbojet represented by ideal Brayton cycle

- 1-2: Isentropic compression
- 2-3: Constant pressure heat addition
- 3-4: Isentropic expansion
- 4-1: Constant pressure heat rejection

1 → 2: isentropic compression

$$\frac{T_2}{T_1} = \left( \frac{p_2}{p_1} \right)^{\frac{\gamma-1}{\gamma}} = r_p^{\frac{\gamma-1}{\gamma}}$$

2 → 3: isobaric heat addition

$$Q_{in} = m \Delta h$$

$$m_f(LHV) = m c_p (T_3 - T_2)$$

3 → 4: isentropic expansion

$$\frac{T_3}{T_4} = \left( \frac{p_3}{p_4} \right)^{\frac{\gamma-1}{\gamma}} = r_p^{\frac{\gamma-1}{\gamma}}$$

4 → 1: isobaric heat release

$$Q_{out} = m c_p (T_4 - T_1)$$

- Introduce new quantities

measure of compression in compressor

$$\left\{ \begin{array}{l} \bullet r = p_2/p_1 = v_1/v_2 \\ \bullet r_p = p_2/p_1 \end{array} \right.$$

# Ideal Brayton Cycle

- Ideal Brayton cycle efficiency

- $$\eta = \frac{W_{out} - W_{in}}{Q_{in}}$$

$$= 1 - \frac{Q_{out}}{Q_{in}}$$

$$= 1 - \frac{m c_p (T_4 - T_1)}{m c_p (T_3 - T_2)}$$

$$= 1 - \frac{T_4 - T_1}{T_3 - T_2}$$

- $$\eta = 1 - r_p^{\frac{1-\gamma}{\gamma}} = 1 - r^{1-\gamma}$$

Can show  $r_p = r^\gamma$

$$\eta = 1 - \frac{T_1 (T_4/T_1 - 1)}{T_2 (T_3/T_2 - 1)}$$

$$\eta = 1 - \frac{T_1}{T_2}$$

$$\eta = 1 - r_p^{\frac{1-\gamma}{\gamma}}$$

$$P_1 = P_4 \Rightarrow \frac{P_4}{P_3} = \frac{P_1}{P_2}$$

$$\left( \frac{T_4}{T_3} \right)^{\frac{\gamma}{\gamma-1}} = \left( \frac{T_1}{T_2} \right)^{\frac{\gamma}{\gamma-1}}$$

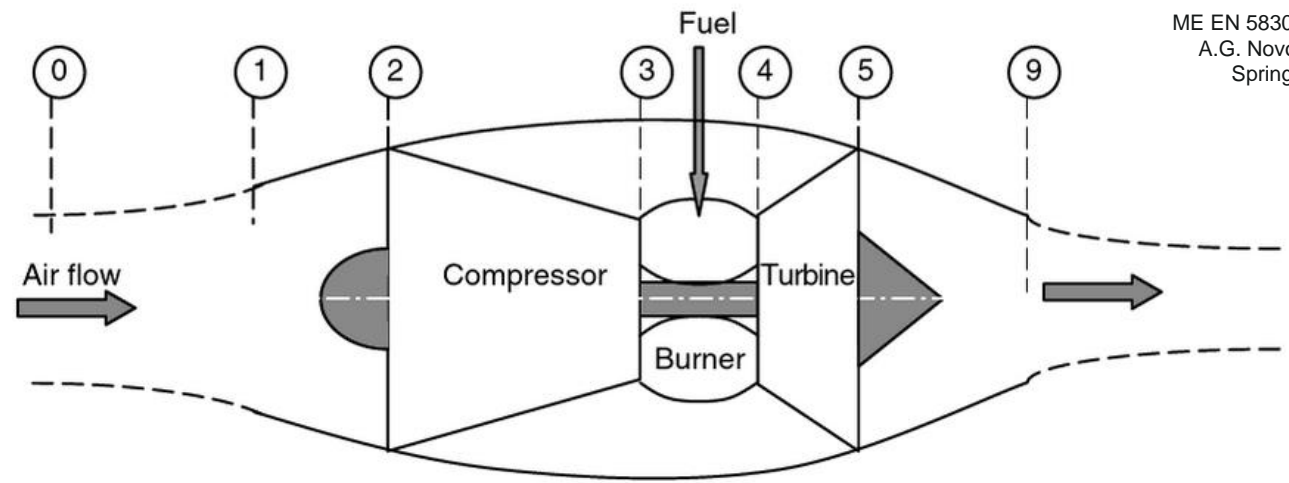
$$\frac{T_4}{T_3} = \frac{T_1}{T_2}$$

$$\frac{T_4}{T_1} = \frac{T_3}{T_2}$$

# Turbojet Analysis

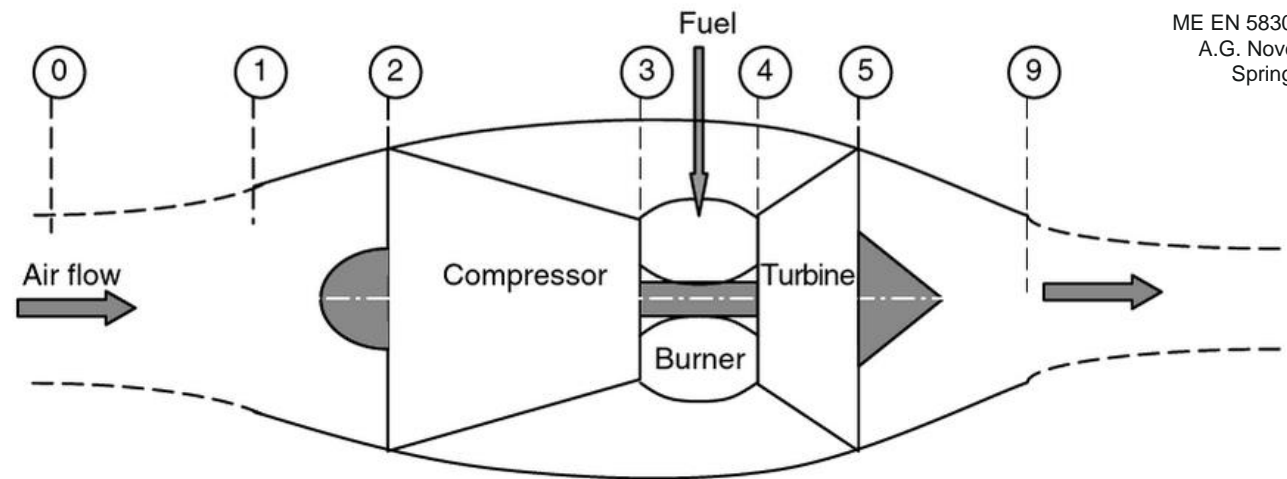
- Station Numbering

- 0: Unperturbed flow
- 1: Diffuser inlet
- 2: Diffuser outlet/Compressor inlet
- 3: Compressor outlet/Combustor inlet
- 4: Combustor outlet/Turbine inlet
- 5: Turbine outlet/Nozzle inlet
- 6-8: Space for additional processes (later slides)
- 9: Nozzle outlet



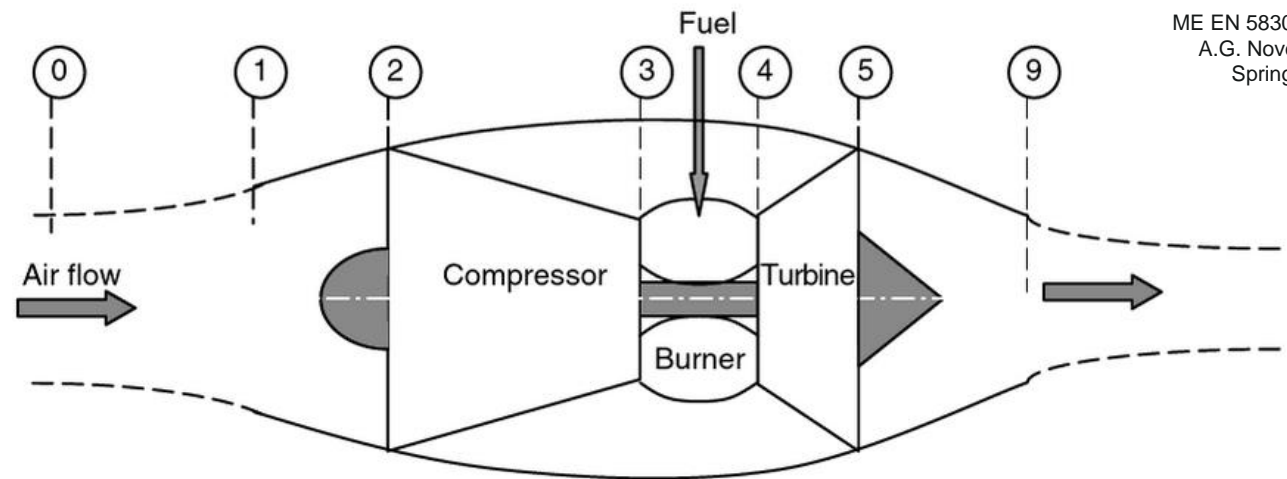


# Turbojet Analysis



- Inlet (Diffuser) (0-2)
  - Delivers air to the compressor at the necessary Mach number
  - Commercial/Military aircraft fly at Mach numbers  $>0.7$  ( $M_0$ )
  - Compressors are usually designed for inlet Mach numbers  $<0.6$  ( $M_2$ )
    - Higher Mach numbers increase the risk of shocks
- Decelerates flow
- Increases pressure
- Try to avoid flow separation in the boundary layer

# Turbojet Analysis



- Inlet (Diffuser) (0-2)

- Ambient conditions (0) are not the same as inlet conditions (1)
- We will lump all outside/inside acceleration/deceleration together
  - Consider 0-2 as a single process

- No work input/output

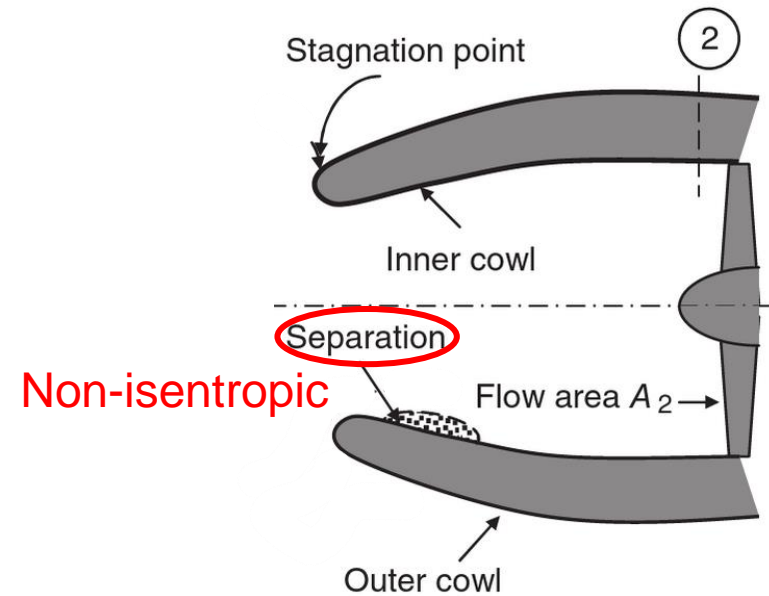
- Diffuser is adiabatic

- $\Delta h_t = 0$

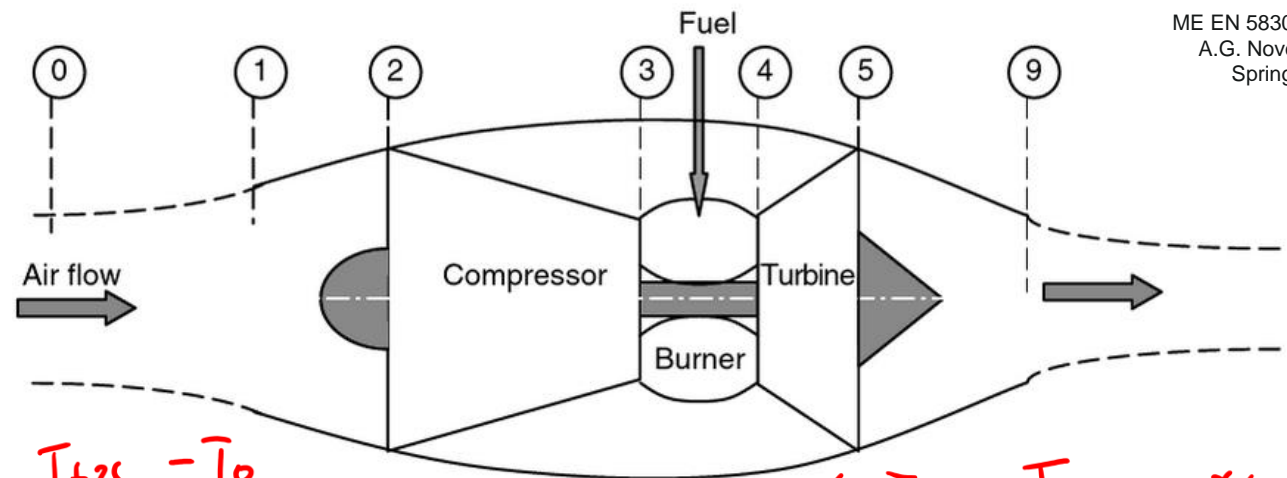
- $T_{t2} = T_{t0} = T_0 \left( 1 + \frac{\gamma-1}{2} M_0^2 \right)$

- Diffuser is not necessarily isentropic

- $\eta_d = \frac{h_{t2s} - h_0}{h_{t2} - h_0} = \frac{T_{t2s} - T_0}{T_{t2} - T_0}$



# Turbojet Analysis



- Inlet (Diffuser) (0-2)
  - Diffuser outlet pressure

$$\eta_d = \frac{T_{+2s} - T_0}{T_{+2} - T_0} = \frac{T_{+2s}/T_0 - 1}{T_{+2}/T_0 - 1}$$

$$\frac{T_{+2}}{T_0} = \frac{T_{+0}}{T_0} \left(1 + \frac{\gamma-1}{2} M_0^2\right)$$

isov. flow eqn.

$$\frac{T_{+2s}}{T_0} = 1 + \eta_d \left( \frac{T_{+2}}{T_0} - 1 \right)$$

$$\left( \frac{p_{+2}}{p_0} \right)^{\frac{\gamma-1}{\gamma}} = 1 + \eta_d \left( \frac{T_{+2}}{T_0} - 1 \right)$$

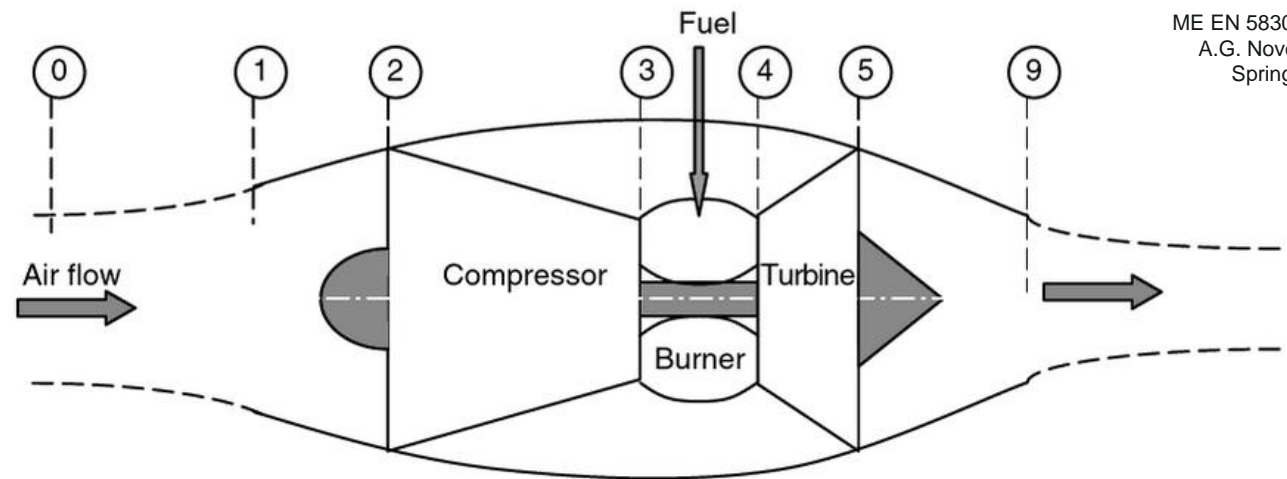
$$\left( \frac{p_{+2}}{p_0} \right) = \left[ 1 + \eta_d \left( \frac{\gamma-1}{2} M_0^2 \right) \right]^{\frac{\gamma}{\gamma-1}}$$

$$\frac{p_{t2}}{p_0} = \left( 1 + \eta_d \frac{\gamma-1}{2} M_0^2 \right)^{\frac{\gamma}{\gamma-1}}$$

$$\frac{p_{t2}}{p_{t0}} = \left( \frac{p_{t2}}{p_0} \right) \left( \frac{p_0}{p_{t0}} \right) = \frac{\left( 1 + \eta_d \frac{\gamma-1}{2} M_0^2 \right)^{\frac{\gamma}{\gamma-1}}}{\left( 1 + \frac{\gamma-1}{2} M_0^2 \right)^{\frac{\gamma}{\gamma-1}}}$$

↑  
Stagnation  
pressure definition

# Turbojet Analysis



## • Compressor (2-3)

- Increases the pressure and energy in the gas before the combustor
  - Power comes from turbine

• Work input

• Compressor is adiabatic

•  $\dot{W}_{in} = \dot{m}\Delta h_t$

• Compressor is not necessarily isentropic

•  $\eta_c = \frac{h_{t3s} - h_{t2}}{h_{t3} - h_{t2}} = \frac{T_{t3s} - T_{t2}}{T_{t3} - T_{t2}}$

•  $\frac{T_{t3}}{T_{t2}} = 1 + \frac{1}{\eta_c} \left( r_p^{\frac{\gamma-1}{\gamma}} - 1 \right)$

Pressure change  
across compressor

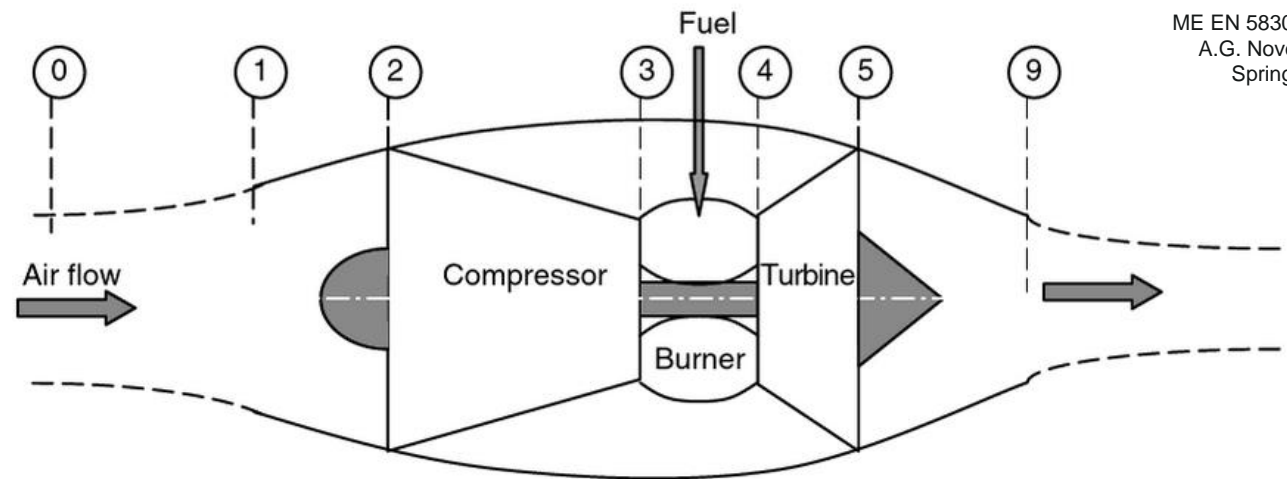
$$r_p = \frac{p_{t3}}{p_{t2}}$$

$$\eta_c = \frac{T_{t3s} - T_{t2}}{T_{t3} - T_{t2}} = \frac{T_{t3s}/T_{t2} - 1}{T_{t3}/T_{t2} - 1}$$

$$\frac{T_{t3}}{T_{t2}} = 1 + \frac{1}{\eta_c} \left( \frac{T_{t3s}}{T_{t2}} - 1 \right)$$

$$\frac{T_{t3}}{T_{t2}} = 1 + \frac{1}{\eta_c} \left[ \left( \frac{p_{t3}}{p_{t2}} \right)^{\frac{\gamma-1}{\gamma}} - 1 \right]$$

# Turbojet Analysis



## • Burner (3-4)

- Adds tremendous energy to the flow through combustion
- Assume isentropic (in practice ~3% friction losses are typical)
- Assume pressure is approximately constant between 3 and 4

- $p_{t4} \approx p_{t3}$

- No work input/output
- Heat input from combustion

- $\dot{Q}_{in} = \dot{m}\Delta h_t$

- $T_{t4} = T_{t3} + \frac{\phi \left(\frac{F}{A}\right)_{st} LHV}{c_p} = T_{t3} + \frac{\left(\frac{F}{A}\right)_{st} LHV}{c_p}$

Combustion:

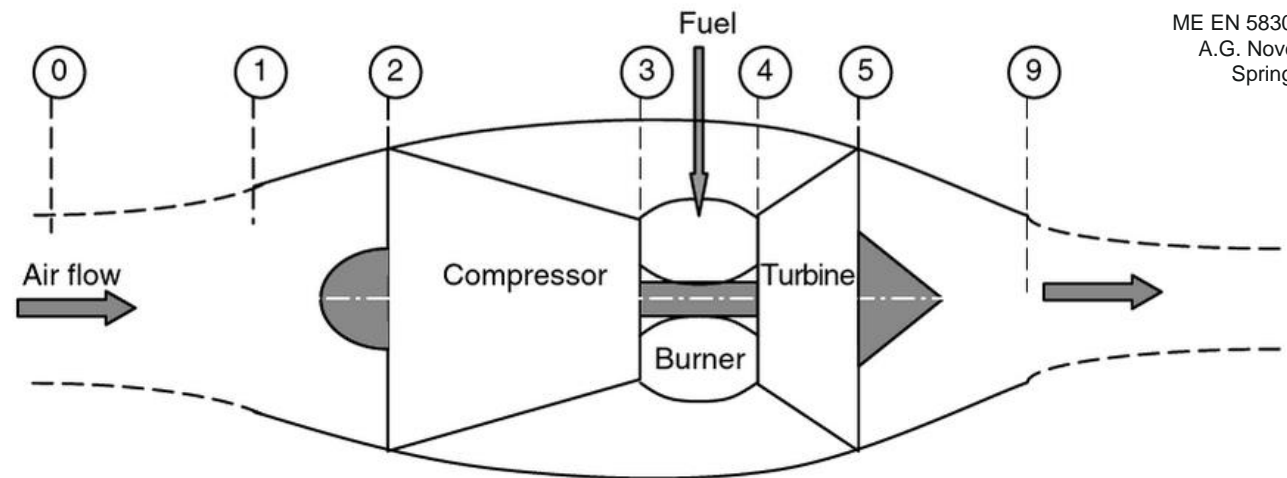
$$\begin{aligned}\dot{Q}_{in} &= \dot{m}_f (LHV) \\ &= \dot{m}_a \left(\frac{F}{A}\right) (LHV) \\ &= \dot{m}_a \left(\frac{F}{A}\right)_{st} \phi (LHV)\end{aligned}$$

$$\approx \dot{m} \left(\frac{F}{A}\right)_{st} \phi (LHV)$$

First Law:

$$\begin{aligned}\dot{Q}_{in} &= \dot{m} c_p (T_{t4} - T_{t3}) \\ \cancel{\dot{m} \left(\frac{F}{A}\right)_{st} \phi (LHV)} &= \cancel{\dot{m} c_p (T_{t4} - T_{t3})}\end{aligned}$$

# Turbojet Analysis



## • Turbine (4-5)

- Lowers pressure and extracts energy from flow to power compressor
- Work output
- “Ideal” turbine is adiabatic
- $-\dot{W}_{out} = \dot{m}\Delta h_t$
- Turbine is not necessarily isentropic

$$\eta_t = \frac{h_{t4} - h_{t5}}{h_{t4} - h_{t5s}} = \frac{T_{t4} - T_{t5}}{T_{t4} - T_{t5s}}$$

① + ② → •  $T_{t5} = T_{t4} - \frac{T_{t2}}{\eta_c} \left( r_p^{\frac{\gamma-1}{\gamma}} - 1 \right)$

① + ② + ③ → •  $p_{t5} = p_{t4} \left[ 1 - \frac{1}{\eta_c \eta_t} \frac{T_{t2}}{T_{t4}} \left( r_p^{\frac{\gamma-1}{\gamma}} - 1 \right) \right]^{\frac{\gamma}{\gamma-1}}$

$$\dot{W}_{in,c} = \dot{W}_{out,T} = \dot{m} c_p (T_{t3} - T_{t2})$$

Eqn. on Slide 12

①  $\dot{W}_{in,c} = \dot{W}_{out,T} = \frac{\dot{m} c_p T_{t2}}{\eta_c} \left( r_p^{\frac{\gamma-1}{\gamma}} - 1 \right)$

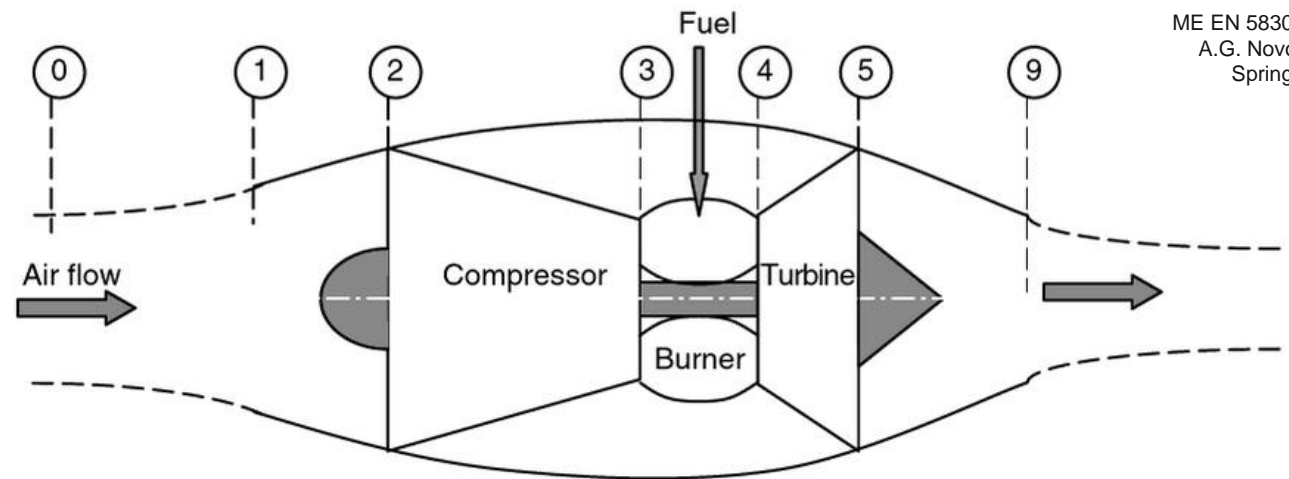
②  $\dot{W}_{out,T} = \dot{m} c_p (T_{t4} - T_{t5})$

$$\eta_t = \frac{T_{t4} - T_{t5}}{T_{t4} - T_{t5s}} \quad \dots \quad \frac{T_{t5}}{T_{t4}} = 1 - \eta_t \left( 1 - \left( \frac{T_{t5s}}{T_{t4}} \right) \right)$$

③  $\frac{T_{t5}}{T_{t4}} = 1 - \eta_t \left[ 1 - \left( \frac{p_{t5}}{p_{t4}} \right)^{\frac{\gamma-1}{\gamma}} \right]$

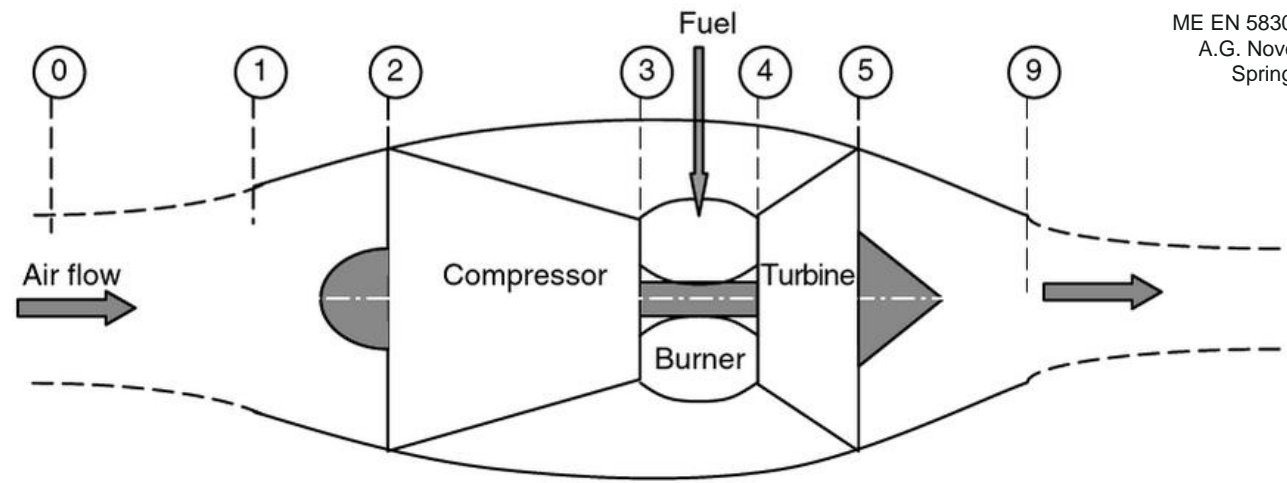


# Turbojet Analysis



- Exhaust (Nozzle) (5-9)
  - Efficiently accelerates exhaust gas to increase thrust
  - Pressure at exit (9) is not necessarily the same as ambient
    - Recall what we learned about supersonic nozzles as an example
  - We will not consider nozzle shocks for ideal turbojets
  - We will lump acceleration in both ambient and nozzle together (i.e., assume  $p_9 = p_a$ )

# Turbojet Analysis



- Exhaust (Nozzle) (5-9)

- No work input/output
- Nozzle is adiabatic
- $\Delta h_t = 0$ 
  - $T_{t9} = T_{t5}$

- Compute exhaust velocity from stagnation temperature ( $T_t = T + \frac{1}{2c_p} V^2$ ):

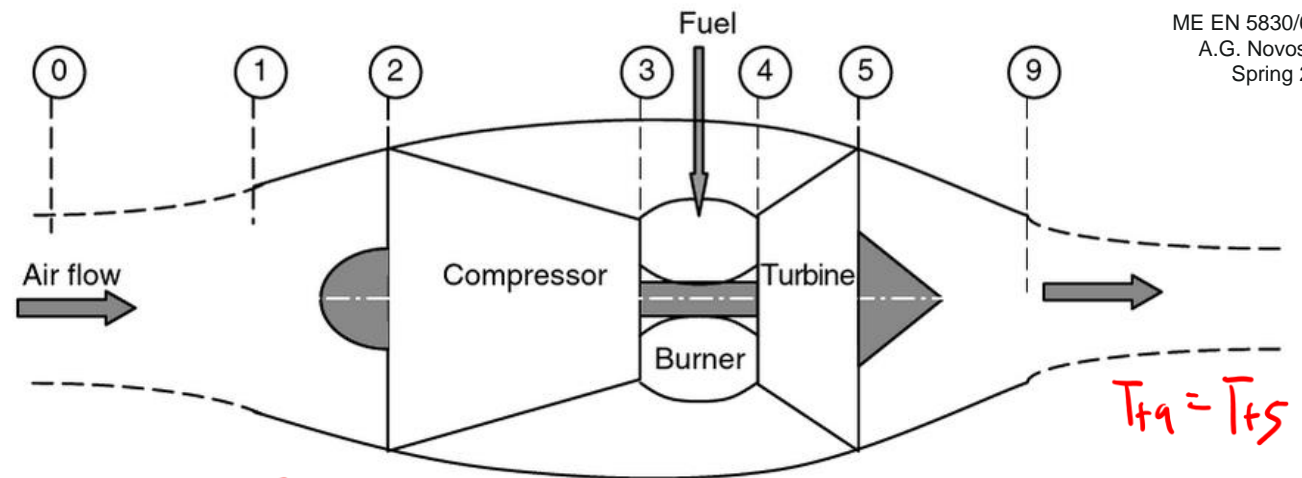
- $V_e = \sqrt{2c_p(T_{t9} - T_9)}$

- Nozzle is not necessarily isentropic

- $\eta_n = \frac{h_{t5} - h_9}{h_{t5} - h_{9s}} = \frac{T_{t5} - T_9}{T_{t5} - T_{9s}}$



# Turbojet Analysis



## • Exhaust (Nozzle) (5-9)

### • Exhaust Velocity

$$\eta_n = \frac{T_{t5} - T_a}{T_{t5} - T_{a5}} = \frac{1 - T_a/T_{t5}}{1 - T_{a5}/T_{t5}}$$

$$\Rightarrow \frac{T_a}{T_{t5}} = 1 - \eta_n \left(1 - \frac{T_{a5}}{T_{t5}}\right)$$

$$\frac{T_a}{T_{t5}} = 1 - \eta_n \left[1 - \left(\frac{p_a}{p_{t5}}\right)^{\frac{\gamma-1}{\gamma}}\right]$$

$$\bullet V_e = \sqrt{2 \frac{\gamma}{\gamma-1} \eta_n R T_{t5} \left[1 - \left(\frac{p_a}{p_{t5}}\right)^{\frac{\gamma-1}{\gamma}}\right]}$$

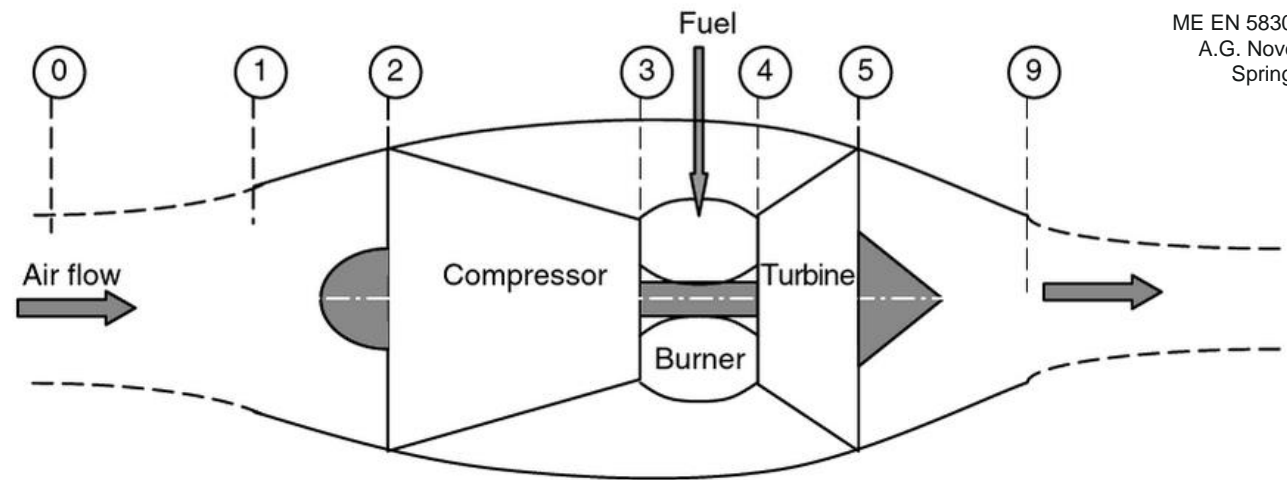
$$T_q = T_{t5} - T_{t5} \eta_n \left[1 - \left(\frac{p_a}{p_{t5}}\right)^{\frac{\gamma-1}{\gamma}}\right]$$

$$V_e = \sqrt{2 c_p (T_{t9} - T_a)} = \sqrt{2 c_p (T_{t5} - T_a)}$$

$$= \sqrt{2 c_p T_{t5} \eta_n \left[1 - \left(\frac{p_a}{p_{t5}}\right)^{\frac{\gamma-1}{\gamma}}\right]}$$

$$c_p = \frac{\gamma}{\gamma-1} R$$

# Turbojet Analysis



- Exhaust (Nozzle) (5-9)

- $$T = \dot{m}(V_e - V) = \dot{m} \left( V_e - M \sqrt{\gamma R T_a} \right)$$

$\uparrow$   $n_o$        $\uparrow$   $T_o$

- Propulsive Efficiency

- $$\eta_p = \frac{TV}{\dot{m} \left[ \frac{V_e^2}{2} - \frac{V^2}{2} \right]} = \frac{2V}{V_e + V}$$