

3D holographic head mounted display using holographic optical elements with astigmatism aberration compensation

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Abstract: We propose a bar-type three-dimensional holographic head mounted display using two holographic optical elements. Conventional stereoscopic head mounted displays may suffer from eye fatigue because the images presented to each eye are two-dimensional ones, which causes mismatch between the accommodation and vergence responses of the eye. The proposed holographic head mounted display delivers three-dimensional holographic images to each eye, removing the eye fatigue problem. In this paper, we discuss the configuration of the bar-type waveguide head mounted displays and analyze the aberration caused by the non-symmetric diffraction angle of the holographic optical elements which are used as input and output couplers. Pre-distortion of the hologram is also proposed in the paper to compensate the aberration. The experimental results show that proposed head mounted display can present three-dimensional see-through holographic images to each eye with correct focus cues.

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1. Introduction

Head mounted displays (HMDs) have been attracting growing attention recently. Two micro displays dedicated to the eyes of the user enable separated presentation of two parallax images to corresponding eyes, giving strong depth sensation by stereopsis. See-through optics enable the presentation of the stereoscopic three-dimensional (3D) images on top of the real environment, which is crucial in augmented reality applications.

In usual stereoscopic HMDs, however, the image presented to each eye is two-dimensional (2D) one formed at the conjugate plane of the micro display [1–9]. Thus each eye of the user focuses on the conjugate plane. On the contrary, the optical axes of two eyes converge at a distance which is given by the disparity between two parallax images presented to two eyes. Therefore, the focusing distance and the converging distance are different in the stereoscopic HMDs, which is called accommodation-vergence mismatch. The accommodation-vergence mismatch has been considered as a fundamental cause of the eye fatigue in 3D viewing [10, 11]. Besides the accommodation-vergence mismatch, the fixed distance of the conjugate plane also limits the presentation of the interactive objects in AR that has dynamic depth range. These drawbacks make users uncomfortable when using conventional stereoscopic HMDs.

In this paper, we propose a waveguide-type holographic HMD using two holographic optical elements (HOEs) and a spatial light modulator (SLM) for each eye. The proposed system delivers a holographic 3D image to each eye of the user instead of a 2D parallax image, removing the accommodation-vergence mismatch problem. The holographic 3D image is not confined in the conjugate plane of the SLM but it can be located at arbitrary distance from the eye, enabling realistic registration of the holographic 3D images to real scene in AR applications. The limitation of narrow viewing angle of usual holographic displays is relaxed in the proposed system since the relative position of the eye to the SLM is fixed in HMD configuration. The waveguide structure with two HOEs as input and output couplers enables the implementation of the proposed system in a thin and compact form factor and the presentation of the holographic 3D images in a see-through fashion. Note that the waveguide-type HMD configuration using two HOE couplers for each eye has already been reported [3–8]. What is unique in the system proposed in this paper is that the image delivered by the waveguide and presented to the eye is not a 2D image but a holographic 3D image, which removes the accommodation-vergence mismatch.

Main contribution of this paper is the analysis and the holographic compensation of the aberration in the waveguide-type holographic HMD system using HOE couplers. We find that the non-symmetric diffraction characteristics of the HOEs cause astigmatism of the holographic 3D images. This aberration is analyzed using ray tracing based on the coupled wave theory [12], and compensated by applying pre-distortion in the hologram generation for 3D images. The experimental results show that the proposed pre-distortion compensates the aberration successfully, displaying clear 3D holographic images in dynamic depth range. Although the holographic 3D HMD has been reported [13], to the authors' best knowledge, the implementation with see-through waveguide-type configuration has not been proposed yet. The aberration analysis and the holographic compensation in the waveguide-type HMD with two HOE couplers are also original contribution of this paper.

In the following sections, we explain the system configuration, aberration analysis, holographic compensation, and the experimental verifications.

2. System configuration

Figure 1 shows the conceptual diagram of the proposed method. The system consists of a SLM, a bar-type waveguide, and two reflection-type HOEs. Linear grating is recorded in the HOEs so that they work as input and output couplers. After the laser light is diffracted by a SLM, it enters the waveguide and is incident on the in-coupling HOE in normal direction. The in-coupling HOE diffracts the incident normal light rays to an angle larger than the critical angle θ_c for the total internal reflection inside the waveguide. The light rays diffracted by the in-coupling HOE travel the waveguide through several internal

reflections to reach the out-coupling HOE. At last, the light rays are diffracted by the out-coupling HOE to be directed to the normal direction and arrive at the user's eye. A virtual image of the SLM is located at the distance of the optical path length that the light rays experience. The initial holographic 3D images formed by the SLM is also delivered to the space around the virtual image of the SLM. Therefore, one can see the holographic 3D images. Moreover, as the waveguide and the out-coupling HOE are transparent, the holographic 3D images are observed on top of the real environment, enabling see-through 3D presentation.

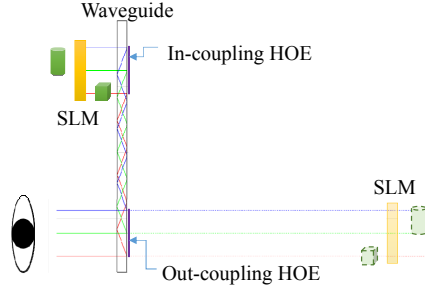


Fig. 1. Conceptual diagram of the proposed waveguide-type holographic HMD with two HOEs.

3. Aberration analysis and holographic compensation

In order to find the ray propagation characteristics in the waveguide, a ray tracing based on the H. Kogelnik's coupled wave theory is performed. Figure 2 shows the vector diagrams of the in-coupling and out-coupling HOEs. With the coordinate system shown in Fig. 2, the grating vectors of two HOEs, i.e. \mathbf{K}_{in} and \mathbf{K}_{out} lie in the yz plane. The reference, signal and the grating vectors of the out-coupling HOE, i.e. $\mathbf{k}_{\text{ro,out}}$, $\mathbf{k}_{\text{so,out}}$, and \mathbf{K}_{out} are the mirror reflections of the signal, reference, and the grating vectors of the in-coupling HOE, i.e. $\mathbf{k}_{\text{so,in}}$, $\mathbf{k}_{\text{ro,in}}$, and \mathbf{K}_{in} with respect to the xy plane. In both HOEs, the reference and the signal beam vectors are related to the grating vector by

$$\mathbf{k}_{\text{so}} = \mathbf{k}_{\text{ro}} - \mathbf{K}. \quad (1)$$

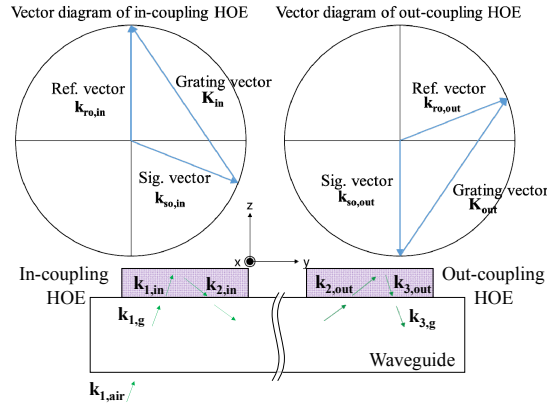


Fig. 2. Vector diagram of the in-coupling and out-coupling HOEs.

When a beam $\mathbf{k}_{1,\text{air}}$ from the SLM in the air is incident on the in-coupling HOE through the waveguide, its wave vector $\mathbf{k}_{1,\text{in}}$ in the in-coupling HOE region is represented by

$$\mathbf{k}_{1,\text{in}} = \left[k_{1,\text{air},x}, k_{1,\text{air},y}, \sqrt{k_{o,h}^2 - k_{1,\text{air},x}^2 - k_{1,\text{air},y}^2} \right], \quad (2)$$

where $k_{o,h} = 2\pi n_h/\lambda_{air}$ is the wave number in the HOE region with the mean refractive index of the HOE n_h and the wavelength in the air λ_{air} . This beam is diffracted by the in-coupling HOE to have the corresponding wave vector $\mathbf{k}_{2,in}$ given by

$$k_{2,in,x} = k_{1,in,x} - K_{in,x} = k_{1,in,x}, \quad (3)$$

$$k_{2,in,y} = k_{1,in,y} - K_{in,y}, \quad (4)$$

$$k_{2,in,z} = -\sqrt{k_{o,h}^2 - k_{2,in,x}^2 - k_{2,in,y}^2}. \quad (5)$$

Note that the z component of the $\mathbf{k}_{2,in}$ is not given by the grating equation of Eq. (1) but given by Eq. (5) in order to meet the wave number condition, considering that the HOE is thin in z -direction so that the grating vector \mathbf{K}_{in} is spread in the same direction as shown in Fig. 3.

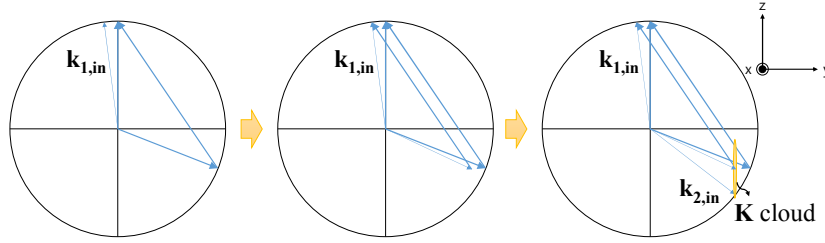


Fig. 3. Diffraction of arbitrary ray by grating vector in in-coupling HOE.

The diffracted beam $\mathbf{k}_{2,in}$ is refracted into the waveguide of the refractive index n_g , guided inside the waveguide through total internal reflections, and refracted into the out-coupling HOE region. The wave vector $\mathbf{k}_{2,out}$ of the incident beam on the out-coupling HOE is the mirror reflection of the $\mathbf{k}_{2,in}$ with respect to the xy plane, i.e.

$$\mathbf{k}_{2,out} = [k_{2,in,x}, k_{2,in,y}, -k_{2,in,z}], \quad (6)$$

and thus the diffracted beam by the out-coupling HOE has the wave vector $\mathbf{k}_{3,out}$ given by

$$k_{3,out,x} = k_{2,out,x} - K_{out,x} = k_{2,out,x}, \quad (7)$$

$$k_{3,out,y} = k_{2,out,y} - K_{out,y}, \quad (8)$$

$$k_{3,out,z} = -\sqrt{k_{o,h}^2 - k_{3,out,x}^2 - k_{3,out,y}^2}. \quad (9)$$

From Eqs. (2)-(9) and the fact that two grating vectors are in the mirror symmetry with respect to the xy plane, i.e. $\mathbf{K}_{in,xy} = \mathbf{K}_{out,xy}$, $K_{in,z} = -K_{out,z}$, it can be easily seen that the diffracted wave from the out-coupling HOE, $\mathbf{k}_{3,out}$ is also the mirror reflection of the incident wave on the in-coupling HOE, $\mathbf{k}_{1,in}$. Therefore, the light ray from the SLM enters the user's eye, maintaining its angle.

However, the optical path length that the light ray experiences inside the waveguide is dependent on the orientation of the plane of the incidence at the same angular deviation from the normal direction. Suppose that the light ray from the SLM is incident on the in-coupling HOE with the polar angle θ_{in} and the azimuthal angle φ_{in} in the in-coupling HOE region. As illustrated in Fig. 4, the diffracted wave by the in-coupling HOE, $\mathbf{k}_{2,in}$ is given by Eqs. (3)-(5) with the transverse vector component

$$\mathbf{k}_{2,in,xy} = [k_{o,h} \sin \theta_{in} \cos \varphi_{in}, k_{o,h} \sin \theta_{in} \sin \varphi_{in} - K_{in,y}]. \quad (10)$$

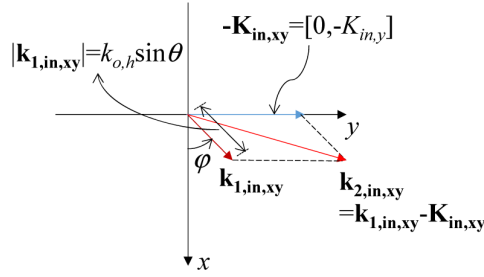


Fig. 4. Diffracted wave from the in-coupling HOE

Then the angle between $\mathbf{k}_{2,\text{in},\theta_{\text{in}}=0} = \mathbf{k}_{\text{so},\text{in}}$ and $\mathbf{k}_{2,\text{in},\theta_{\text{in}}=\theta_{\text{in}}}$, i.e. $\Delta\theta'$ is given by

$$\Delta\theta' = \cos^{-1} \left(\frac{K_{m,y}^2 - K_{m,y} k_{o,h} \sin \theta \sin \varphi_m + \sqrt{(k_{o,h}^2 - K_{m,y}^2)(k_{o,h}^2 - K_{m,y}^2 - k_{o,h}^2 \sin^2 \theta \sin^2 \varphi_m + 2K_{m,y} K_{m,y} \sin \theta \sin \varphi_m)}}{k_{o,h}^2} \right). \quad (11)$$

Equation (11) indicates that the angular deviation $\Delta\theta'$ of the diffracted light ray $\mathbf{k}_{2,\text{in},\theta_{\text{in}}=\theta_{\text{in}}}$ with respect to $\mathbf{k}_{2,\text{in},\theta_{\text{in}}=0} = \mathbf{k}_{\text{so},\text{in}}$ is dependent on the azimuthal angle φ_m of the incident ray $\mathbf{k}_{1,\text{in}}$. Therefore the incident beam cone around the normal direction (z -axis) is distorted around the $\mathbf{k}_{2,\text{in},\theta_{\text{in}}=0} = \mathbf{k}_{\text{so},\text{in}}$ after it is diffracted from the in-coupling HOE. This distortion makes the optical path length that the light rays experience in the waveguide be different for different φ_m , causing astigmatism in the final images.

More specifically, the ray path of a light ray emanating from a point $(0, 0, -z_o)$ in the SLM space with a polar angle θ and azimuthal angle φ can be calculated. For example, the transverse position of the ray on the bottom side of the waveguide when it exits the waveguide toward the user's eye $(x_{b,\text{exit}}, y_{b,\text{exit}})$ is given by

$$x_{b,\text{exit}} = \sin \theta \cos \varphi \left(\frac{z_o}{\cos \theta} + \frac{2t}{\sqrt{n_g^2 - \sin^2 \theta}} + \frac{(m+1)tk_{o,\text{air}}}{|k_{2,g,z}|} \right), \quad (12)$$

$$y_{b,\text{exit}} = \sin \theta \sin \varphi \left(\frac{z_o}{\cos \theta} + \frac{2t}{\sqrt{n_g^2 - \sin^2 \theta}} + \frac{(m+1)tk_{o,\text{air}}}{|k_{2,g,z}|} \right) - \frac{(m+1)tK_{in,y}}{|k_{2,g,z}|}, \quad (13)$$

where

$$|k_{2,g,z}| = \sqrt{k_{o,\text{air}}^2 (n_g^2 - \sin^2 \theta) - K_{in,y}^2 + 2K_{in,y} k_{o,\text{air}} \sin \theta \sin \varphi}, \quad (14)$$

t is the thickness of the waveguide, and m is the number of the reflections inside the waveguide which is always odd.

Figures 5–7 show ray-tracing results for the verification of the analysis. In all simulations shown in Figs. 5–7, a glass waveguide of the thickness $t = 3\text{mm}$ is assumed. The number of the internal reflection between the in-coupling and out-coupling HOEs is 5, i.e. $m = 5$, and the angle of the ray guided inside the waveguide is 70° with respect to the normal direction of the HOE. Figure 5 shows the footprint of the rays on the bottom side of the waveguide when they enter from the SLM and exit toward the eye. In Fig. 5, it can be found that the spread of the footprint of the exiting rays is larger in y -axis than in x -axis due to φ dependency of the angular deviation of Eqs. (11)–(14). Figures 6(a) and 6(b) show the rays coming to the eye. Note that the rays are extended to the virtual image space in Figs. 6(a) and 6(b) to identify the virtual image point clearly. From Fig. 6, it is obvious that the image point is focused at different distances in x and y directions, showing the astigmatism as expected.

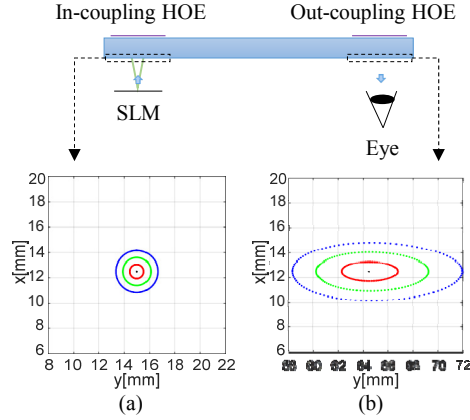


Fig. 5. Footprint of rays on bottom side of the waveguide; (a) when the rays enter the waveguide from the SLM and, (b) when the rays exit the waveguide toward the eye.

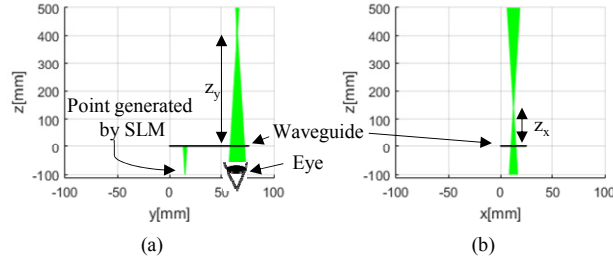


Fig. 6. Simulation results; (a) ray tracing in yz-plane, (b) ray tracing in zx-plane.

The rays exiting from the waveguide toward the eye are mirror reflection of original rays from the point source as discussed before. Therefore the virtual image distances from the bottom side of the waveguide in x and y directions, i.e. z_x , z_y can be approximately estimated by

$$z_x = \frac{|x_{b,exit,\varphi=0} - x_{b,exit,\varphi=\pi}|}{2 \tan \theta}, \quad z_y = \frac{|y_{b,exit,\varphi=\pi/2} - y_{b,exit,\varphi=3\pi/2}|}{2 \tan \theta}. \quad (15)$$

Figure 7(a) shows the virtual image distances for different polar angles. As can be seen from Fig. 7(a), the astigmatism is dominant while spherical aberration is also identified. One notable thing is that the difference between directional image distances, i.e. $\Delta z = z_y - z_x$ do not change with the point source distance z_o as shown in Fig. 7(b). This property can also be verified from Eqs. (12)-(15).

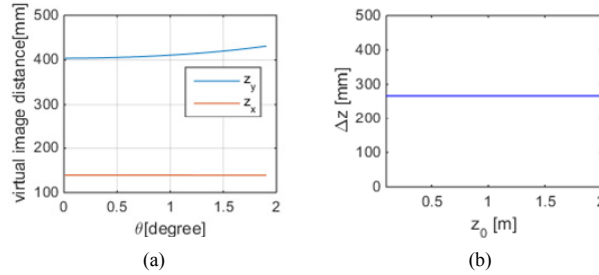


Fig. 7. (a) Virtual image distances for different polar angles and (b) difference between virtual image distances in x and y directions, i.e. $\Delta z = z_y - z_x$ for different object point source distances z_o .

In the proposed method, the constant difference Δz is simply compensated in the hologram generation process for the SLM by introducing the distance offset to y direction. For a 3D image which is composed of a set of 3D image points, the hologram is generated in the proposed method by

$$H(u, v) = \sum_i A_i \exp \left\{ \frac{jk_{o,air}}{2} \left(\frac{(x_i - u)^2}{z_i} + \frac{(y_i - v)^2}{z_i - \Delta z} \right) \right\}, \quad (16)$$

where (x_i, y_i, z_i) is the position of the 3D image point and the A_i is its amplitude. Note that the distance of the 3D image point is adjusted by Δz in y direction in Eq. (16), compensating the astigmatism of the waveguide with two HOEs structure.

4. Experimental verification

Holographic aberration compensation for the waveguide type HMD with HOE input/output couplers and the presentation of the see-through holographic 3D images with correct focus cues were verified experimentally. In experiment, the HOEs were implemented by recording the linear gratings on the photopolymer sheet with $\lambda_{air} = 532\text{nm}$ laser. The experimental setup to record the HOE is shown in Fig. 8. A polarizing beam splitter (PBS) divides a plane wave to the reference and signal beams which are incident on the photopolymer in normal direction and an angle larger than θ_c , respectively. Note that a right angle prism was attached on the photopolymer in the signal beam path in order to ensure that the signal beam angle inside the photopolymer is 70° which is larger than the critical angle θ_c of the waveguide glass – air boundary. In the configuration shown in Fig. 8, the half wave plate (HWP) 1 controls the power between reference beam and signal beam and the HWP 2 matches the polarization state of two beams to s-polarization. The power of the reference and the signal beams were adjusted to be the same as $0.5\text{mW}/\text{cm}^2$. The diffraction efficiency of the each of the recorded HOEs was measured to be around 34%.

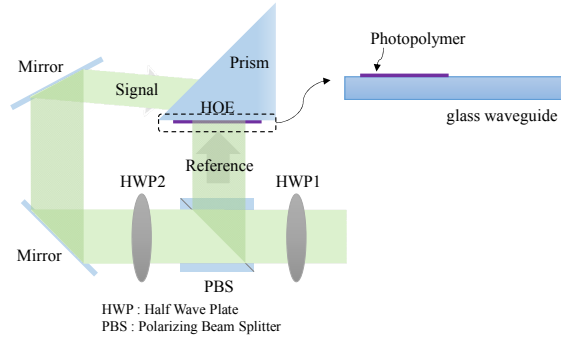


Fig. 8. Experimental setup for recording HOE.

Using the recorded HOEs, the experimental setup for the waveguide type holographic HMD was implemented. Figure 9 shows the implemented experimental setup. The waveguide is a glass slab with a thickness 3mm. The separation between two HOEs is about 49.5mm which corresponds to 5 times internal reflections of the guided light rays. The SLM used in the experiment is a reflection type with 8um pixel pitch. The experimental setup was configured such that the plane beam from the laser source passes through the in-coupling HOE and the waveguide as shown in Fig. 9 to be reflected by the SLM. The reflected and diffracted beams by the SLM are coupled into the waveguide by the in-coupling HOE. The coupled light is guided inside the waveguide and diffracted toward the camera or the user's eye by the out-coupling HOE.

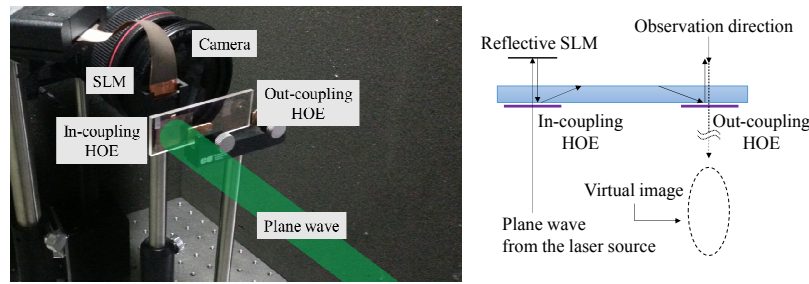


Fig. 9. Experiment setup.

The astigmatism aberration and its holographic compensation were first verified experimentally. A hologram of a mesh grid pattern at 60cm distance behind the SLM was generated without aberration compensation and loaded in the SLM. Figure 10(a) and 10(b) show the images captured at different focal distances of the camera. As can be observed in Figs. 10(a) and 10(b), the horizontal and vertical lines of the reconstructed mesh grid pattern are focused at different focal distances of the camera, revealing the astigmatism caused by the HOEs and the waveguide structure clearly.

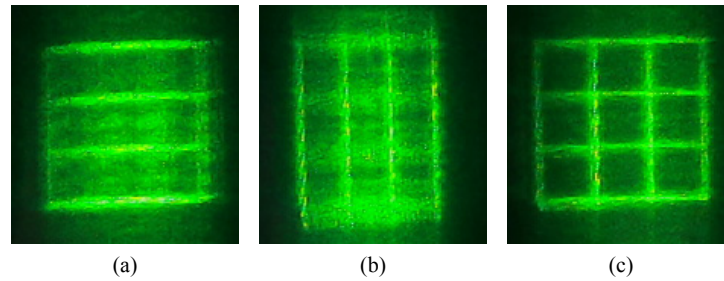


Fig. 10. Reconstructed mesh grid pattern (a) without compensation, focused on horizontal lines (b) without compensation, focused on vertical lines, (c) with proposed compensation.

Figure 10(c) shows the reconstructed image when the holographic compensation is applied. In Fig. 10(c), the camera was focused at 65cm distance. In the holographic compensation, the distance difference Δz was set to 17.7cm. It is obvious from Fig. 10(c) that the reconstructed image becomes much clearer by the proposed holographic compensation. Figure 11 shows another example of the effect of the aberration and its compensation by the proposed method, proving its feasibility.

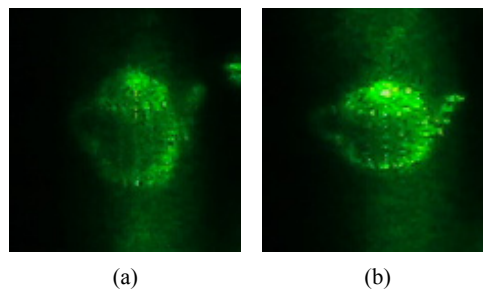


Fig. 11. Reconstructed teapot image (a) without compensation (b) with proposed compensation.

Figure 12 shows the spreads of a single point reconstructed with and without proposed compensation. The intensity profiles in Fig. 12 were obtained by summing columns or rows of the captured images and normalizing them. The point reconstructed with the proposed compensation shows narrow spread both in horizontal and vertical directions as shown in Fig. 12(a). When the proposed compensation is not applied, the reconstructed

point shows asymmetric spread according to the focusing distance of the camera as shown in Figs. 12(b)-12(d), failing to be focused sharply both in the horizontal and vertical directions simultaneously. Comparing Figs. 12(a) and 12(b), it can be observed that the full-width-at-half-maximum (FWHM) of the reconstructed point is reduced from $19(x) \times 26(y)$ pixels to $6(x) \times 6(y)$ pixels by the proposed compensation, showing enhancement over 3 times. In our experimental setup, the 6 pixel FWHM of Fig. 12(a) corresponds approximately to 522 μ m width at 56.5cm distance from the camera, which indicates the spatial resolution of the proposed system roughly.

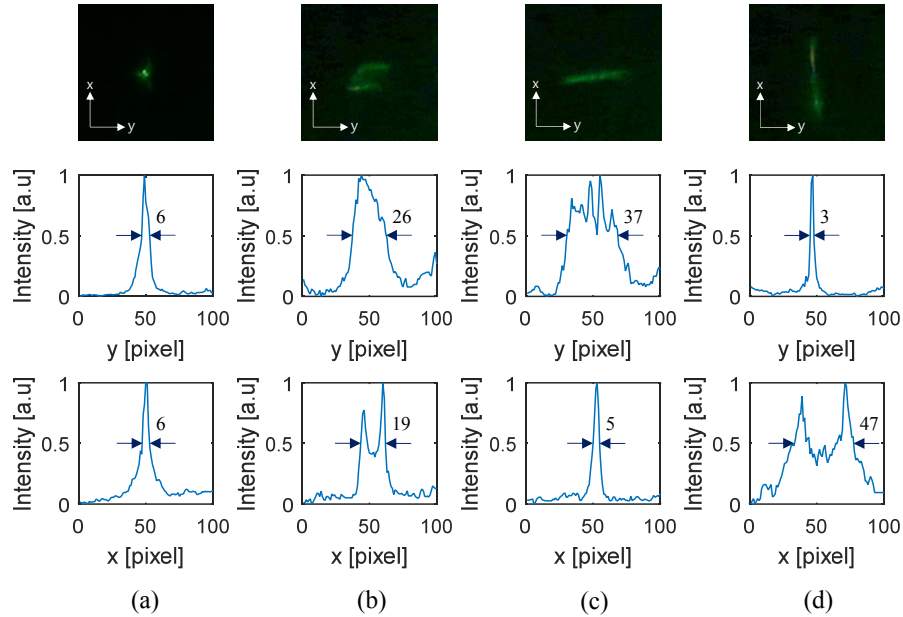


Fig. 12. Point spread of the reconstruction (a) with proposed compensation, (b) without compensation, focused in the middle of the vertical and horizontal focal lengths (c) without compensation, focused at vertical focal length, (d) without compensation, focused at horizontal focal length.

Figure 13 shows the reconstructed images at different distances from the camera. In all images, the proposed holographic compensation was applied with constant $\Delta z = 17.7$ cm. The red object in Fig. 13 is a real object at 32cm distance which is seen through the waveguide and the out-coupling HOE. From Fig. 13, it can be seen that the proposed waveguide-type holographic HMD can present a see-through holographic 3D images in a wide depth range.

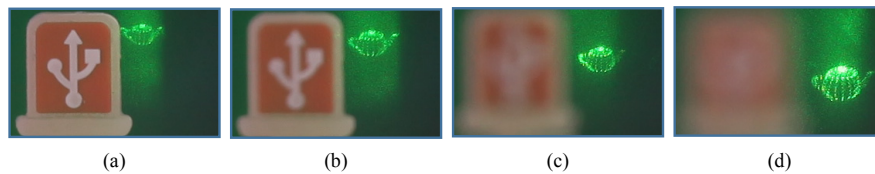


Fig. 13. Holographic images at different depths: (a) 32cm, (b) 37cm, (c) 57cm, (d) 107cm.

Figure 14 shows reconstructed images at different distances at the same time. The holographic cube and the teapot images are displayed at the same distances as the real red and black objects, respectively. The depth difference between the teapot and the cube images is 25cm. From Fig. 14, it can be seen that the reconstructed holographic images are focused and blurred in the same way as the real objects at the same distances, showing the 3D nature of the reconstructed holographic images successfully.

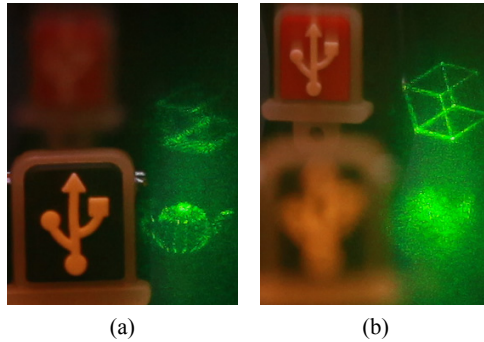


Fig. 14. Real objects and holographic images with 25cm depth difference.

5. Conclusion

In this paper, we proposed a waveguide type 3D holographic HMD using HOE input-output couplers with aberration compensation. The astigmatism aberration caused by the asymmetric diffraction of the HOEs for a given angular deviation of the light from the SLM is analyzed and its simple compensation by introducing distance offset for x and y directions is proposed. The experimental results show that the proposed system can present holographic 3D images in a see-through fashion clearly in a wide depth range.

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