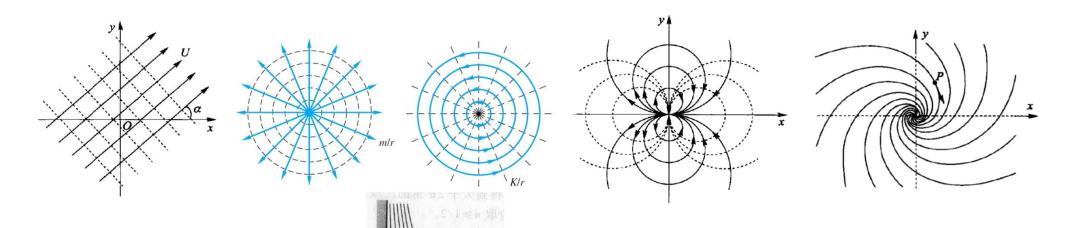
空气与气体动力学

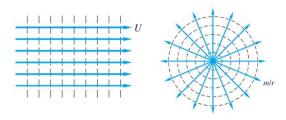
张科

回顾:

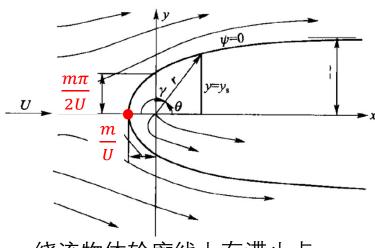
- 1.势流、势函数、流函数: $u = \frac{\partial \phi}{\partial x} = \frac{\partial \psi}{\partial y} \qquad u_r = \frac{\partial \phi}{\partial r} = \frac{1}{r} \frac{\partial \psi}{\partial \theta}$ $v = \frac{\partial \phi}{\partial y} = -\frac{\partial \psi}{\partial x} \qquad u_\theta = \frac{1}{r} \frac{\partial \phi}{\partial \theta} = -\frac{\partial \psi}{\partial r}$
- **2.基本平面势流:** 不可压势流: $\nabla^2 \phi = 0$, $\nabla^2 \psi = 0$
- 3.基本势流叠加:



3.半无穷长物体绕流(Rankine half-body,兰金半体扰流)



$$\psi = Ursin\theta$$
 $\psi = m\theta$
 $\phi = Urcos\theta$ $\phi = mlnr$



$$\psi = \psi_1 + \psi_2 = Ursin\theta + m\theta$$
 $m = \frac{Q}{2\pi}$: 点源强度 $\phi = \phi_1 + \phi_2 = Urcos\theta + mlnr$

$$u_r = \frac{\partial \phi}{\partial r} = U \cos \theta + \frac{m}{r}$$

 $u_{\theta} = \frac{1}{r} \frac{\partial \phi}{\partial \theta} = -U \sin \theta$
滞止点 $u_r = u_{\theta} = 0$

$$\begin{aligned} u_r &= U cos\theta + \frac{m}{r} = 0 \\ u_\theta &= -U sin\theta = 0 \end{aligned} \} \longrightarrow \theta = \pi, \quad r = m/U$$

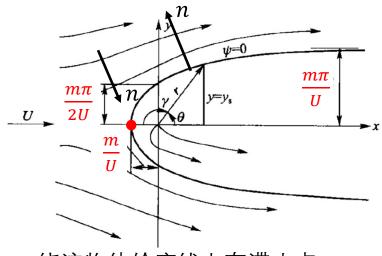
轮廓线上:
$$\psi_0 = U \frac{m}{U} sin\pi + m\pi = m\pi$$

轮廓线为:
$$\psi = Ursin\theta + m\theta = m\pi$$
 轮廓线曲线方程

$$\theta = \frac{\pi}{2}, \frac{3\pi}{2}$$
时: $Ur = m\pi/2$
$$r = \frac{m\pi}{2U}$$

$$u_r = U\cos\theta + \frac{m}{r}$$
$$u_\theta = -U\sin\theta$$

3.半无穷长物体绕流(Rankine half-body,兰金半体扰流)



轮廓线上: $\psi_0 = U \frac{m}{U} sin\pi + m\pi = m\pi$

轮廓线为: $\psi = Ursin\theta + m\theta = m\pi$ | 轮廓线曲线方程

$$\theta = \frac{\pi}{2}, \frac{3\pi}{2}$$
时: $Ur = m\pi/2$ $r = \frac{m\pi}{2U}$

$$r = \frac{m\pi}{2U}$$

 $r \rightarrow \infty$ 时, $\theta \rightarrow 0$: $Ursin\theta = m\pi$

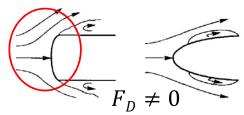
绕流物体轮廓线上有滞止点。

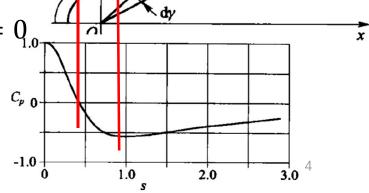
$$p_{\infty} + \frac{1}{2}\rho U^2 = p_s + \frac{1}{2}\rho U_s^2$$

$$C_{p} = \frac{p_{s} - p_{\infty}}{\frac{1}{2}\rho U^{2}} = \frac{\frac{1}{2}\rho[U^{2} - u_{r}^{2} - u_{\theta}^{2}]}{\frac{1}{2}\rho U^{2}}$$
$$= \frac{2}{\gamma}\sin\gamma\cos\gamma - \frac{1}{\gamma^{2}}\sin^{2}\gamma$$

半无限长、无粘:

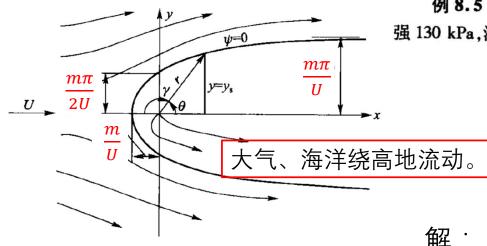
$$p_{\infty} + \frac{1}{2}\rho U^{2} = p_{S} + \frac{1}{2}\rho U_{S}^{2} \qquad F_{D} = \int_{0}^{\pi} (p_{S} - p_{\infty}) \, r cos\theta d\theta = 0$$



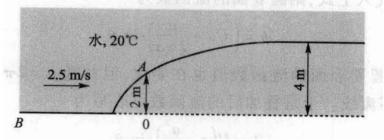


$$u_r = U\cos\theta + \frac{m}{r}$$
$$u_\theta = -U\sin\theta$$

3.半无穷长物体绕流(Rankine half-body,兰金半体扰流)



例 8.5 如图 8.21 所示,河床有一类似于半体柱的隆起,高 4 m。已知 B 点压强 130 kPa,河流速度 2.5 m/s。试利用势流理论确定高于 B 点 2 m 的 A 点的压强。



P307

图 8.21 绕河床隆起的流动

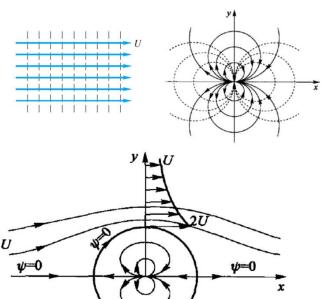
解:
$$\frac{m\pi}{U} = 4$$
m, $U = 2.5$ m/s, $p_{\infty} = 130$ Kpa

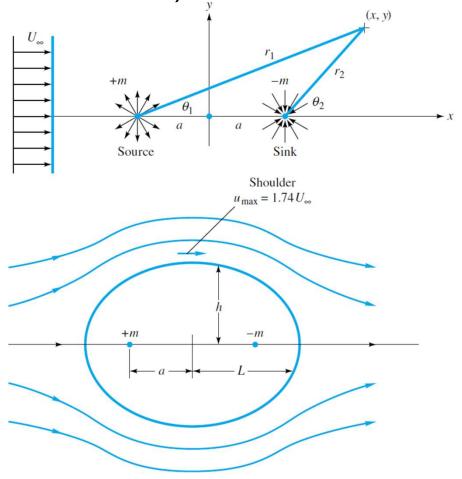
$$\psi = Ursin\theta + m\theta = m\pi$$
 A 点: $rsin\theta = 2m$
 $\theta = \pi/2, r = 2m$

 $m = 10/\pi$

$$u_{Ar} = U\cos\theta + \frac{m}{r} = 5/\pi$$
 $u_{A\theta} = -U\sin\theta = -2.5$ $p_A + \frac{1}{2}\rho U_A^2 + \rho g z_A = p_\infty + \frac{1}{2}\rho U^2 + \rho g z_B$ $p_A = 1.09 \times 10^5 Pa$

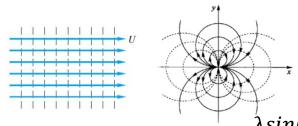
4.绕圆柱无环量流动(均匀流+偶极流):





₩ 无粘壁面条件

4.绕圆柱无环量流动(均匀流+偶极流):



$$\psi = Ursin\theta \quad \psi = -\frac{\lambda sin\theta}{r}$$

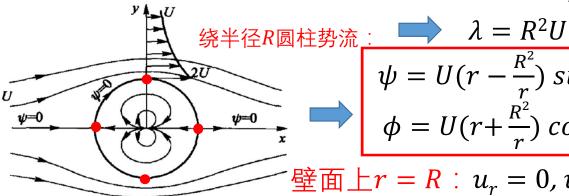
$$\phi = Urcos\theta$$
 $\phi = \frac{\lambda cos}{r}$

$$\psi = Ursin\theta - \frac{\lambda sin\theta}{r}$$
 $\phi = Urcos\theta + \frac{\lambda cos\theta}{r}$

封闭流线:
$$\psi_0 = 0$$

$$\psi = -\frac{\lambda sin\theta}{r}$$
 绕流体壁面: $\psi = Ursin\theta - \frac{\lambda sin\theta}{r} = 0$

$$\phi = \frac{\lambda cos}{r}$$



$$\lambda = R^2 U$$

$$\psi = U(r - \frac{R^2}{r}) \sin\theta$$
$$\phi = U(r + \frac{R^2}{r}) \cos\theta$$

$$\psi = U(r - \frac{R^2}{r}) \sin\theta \qquad u_r = \frac{\partial \phi}{\partial r} = U(1 - \frac{R^2}{r^2}) \cos\theta$$

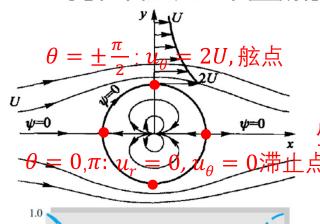
$$\phi = U(r + \frac{R^2}{r}) \cos\theta \qquad u_\theta = \frac{1}{r} \frac{\partial \phi}{\partial \theta} = -U(1 + \frac{R^2}{r^2}) \sin\theta$$

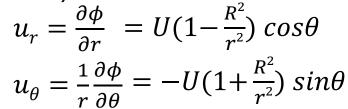
 \mathbb{E} 壁面上r=R: $u_r=0$, $u_{ heta}=-2Usin heta$

$$\theta=0$$
, π : $u_r=0$, $u_\theta=0$, 前后滯止点 $\theta=\pm\frac{\pi}{2}$: $u_\theta=2U$, 舷点

$$\theta = \pm \frac{\pi}{2}$$
: $u_{\theta} = 2U$, 舷点

4.绕圆柱无环量流动(均匀流+偶极流):





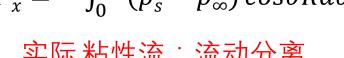
ightharpoonup 壁面上r=R: $u_r=0$, $u_{\theta}=-2Usin\theta$

$$p_{\infty} + \frac{1}{2}\rho U^2 = p_s + \frac{1}{2}\rho U_s^2$$

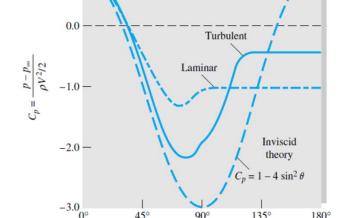
 $p_{\infty} + \frac{1}{2}\rho U^2 = p_s + \frac{1}{2}\rho U_s^2$ 圆柱上 $U_s = u_{\theta} = -2Usin\theta$

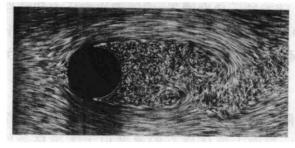
$$C_p = \frac{p_s - p_{\infty}}{\frac{1}{2}\rho U^2} = \frac{\frac{1}{2}\rho [U^2 - 4U^2 \sin^2 \theta]}{\frac{1}{2}\rho U^2} = 1 - 4\sin^2 \theta$$

$$F_x = -\int_0^{2\pi} (p_s - p_\infty) \cos\theta R d\theta = 0$$

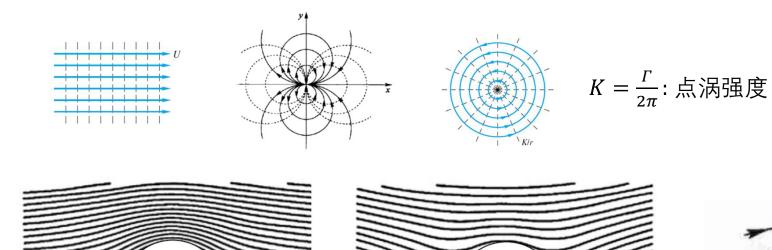


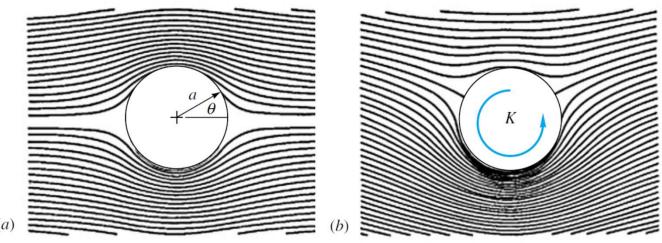
实际 粘性流:流动分离, 阻力 $F_x \neq 0$

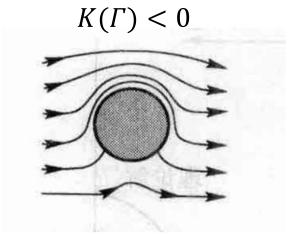




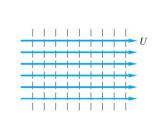
5.绕圆柱有环量流动(均匀流+偶极流+点涡):

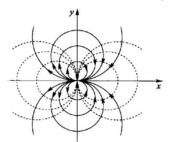


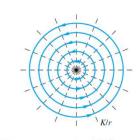




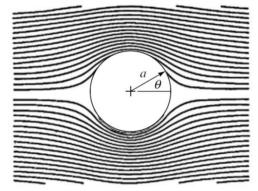
5.绕圆柱有环量流动(均匀流+偶极流+点涡):



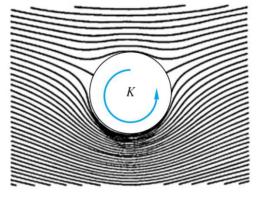




$$\psi = -K lnr$$
 $K = \frac{\Gamma}{2\pi}$: 点涡强度 $\phi = K\theta$ $(\Gamma > 0$ 逆时针, $\Gamma < 0$ 顺时针)



$$\phi = U(r + \frac{R^2}{r}) \cos\theta$$
$$\psi = U(r - \frac{R^2}{r}) \sin\theta$$



$$\psi = U(r - \frac{R^2}{r}) \sin\theta - K \ln r$$

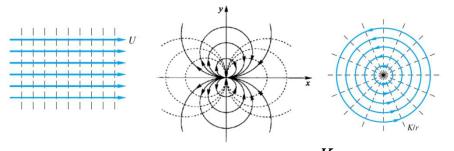
$$\phi = U(r + \frac{R^2}{r}) \cos\theta + K\theta$$

$$u_r = \frac{\partial \phi}{\partial r} = U \cos\theta (1 - \frac{R^2}{r^2})$$

$$u_\theta = \frac{1}{r} \frac{\partial \phi}{\partial \theta} = -U \sin\theta (1 + \frac{R^2}{r^2}) + \frac{K}{r}$$

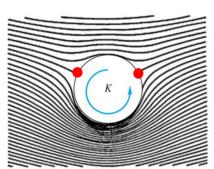
壁面上
$$r = R$$
: $u_r = 0$, $u_\theta = -2Usin\theta + \frac{K}{R}$
滞止点: $u_\theta = -2Usin\theta + \frac{K}{R} = 0 \Longrightarrow sin\theta = \frac{K}{2UR}$

5.绕圆柱有环量流动(均匀流+偶极流+点涡):

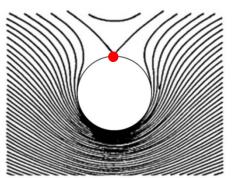


滯止点:
$$u_{\theta} = -2U\sin\theta + \frac{K}{R} = 0$$
 $\sin\theta = \frac{K}{2UR}$

$$\left| \frac{K}{2UR} \right| < 1: \quad \frac{\theta_{s1}}{\theta_{s2}} = \frac{\arcsin\frac{K}{2UR}}{\pi - \theta_{s1}} \qquad \left| \frac{K}{2UR} \right| = 1: \quad \theta_s = \pm \frac{\pi}{2} \qquad \left| \frac{K}{2UR} \right| > 1:$$
 滞止点不在圆柱上



$$\left| \frac{K}{2UR} \right| = 1$$
: $\theta_s = \pm \frac{\pi}{2}$



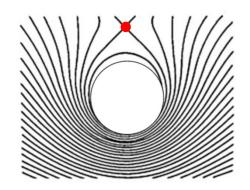
$$\psi = U(r - \frac{R^2}{r}) \sin\theta - K \ln r$$

$$\phi = U(r + \frac{R^2}{r}) \cos\theta + K\theta$$

$$u_r = \frac{\partial \phi}{\partial r} = U \cos\theta (1 - \frac{R^2}{r^2})$$

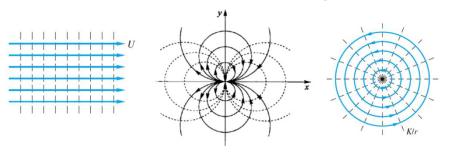
$$u_\theta = \frac{1}{r} \frac{\partial \phi}{\partial \theta} = -U \sin\theta (1 + \frac{R^2}{r^2}) + \frac{K}{r}$$

$$\left|\frac{K}{2UR}\right| > 1$$
: 滞止点不在圆柱上



 $K = \frac{\Gamma}{2\pi}$: 点涡强度 $(\Gamma > 0$ 逆时针, $\Gamma < 0$ 顺时针)

5.绕圆柱有环量流动(均匀流+偶极流+点涡): $\psi = U(r - \frac{R^2}{r}) \sin\theta - K \ln r$



 $\dot{\psi} = U(r - \frac{R^2}{r}) \sin\theta - K \ln r$ $\phi = U(r + \frac{R^2}{r}) \cos\theta + K\theta$ $u_r = \frac{\partial \phi}{\partial r} = U \cos\theta (1 - \frac{R^2}{r^2})$ $u_\theta = \frac{1}{r} \frac{\partial \phi}{\partial \theta} = -U \sin\theta (1 + \frac{R^2}{r^2}) + \frac{K}{r}$

圆柱表面
$$r = R: U_s = u_\theta = -2Usin\theta + \frac{K}{R}$$

$$p_\infty + \frac{1}{2}\rho U^2 = p_s + \frac{1}{2}\rho U_s^2$$

$$p_s = p_\infty + \frac{1}{2}\rho U^2 - \frac{1}{2}\rho(-2Usin\theta + \frac{K}{R})^2$$

$$= p_\infty + \frac{1}{2}\rho[U^2 - (2Usin\theta - \frac{K}{R})^2]$$

$$D = Fx = -\int_0^{2\pi} (p_s - p_\infty) \cos\theta R d\theta$$

$$= 0$$

$$L = Fy = -\int_0^{2\pi} (p_s - p_\infty) \sin\theta R d\theta$$

$$= -\int_0^{2\pi} 2\rho U \sin^2\theta K d\theta$$

$$= -2\rho U K\pi$$

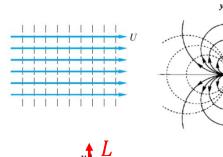
Kutta-Joukowski Lift Theorem:

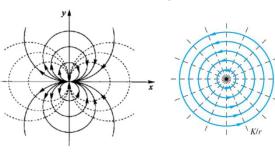
$$L = -\rho U\Gamma$$

升力理论

 $K = \frac{\Gamma}{2\pi}$: 点涡强度 $(\Gamma > 0$ 逆时针, $\Gamma < 0$ 顺时针)

5.绕圆柱有环量流动(均匀流+偶极流+点涡): $\psi = U(r - \frac{R^2}{r}) \sin \theta - K \ln r$







$$\dot{\psi} = U(r - \frac{R^2}{r}) \sin\theta - K \ln r$$

$$\phi = U(r + \frac{R^2}{r}) \cos\theta + K\theta$$

$$u_r = \frac{\partial \phi}{\partial r} = U \cos\theta (1 - \frac{R^2}{r^2})$$

$$u_\theta = \frac{1}{r} \frac{\partial \phi}{\partial \theta} = -U \sin\theta (1 + \frac{R^2}{r^2}) + \frac{K}{r}$$

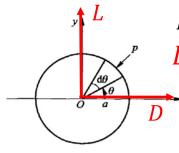
$$D=0$$

 $L=-\rho U\Gamma$ 升力理论

均匀流中任意形状柱体,单位长度受升力为 $\rho U\Gamma$,方向为流动沿 Γ 反向旋转90°

 $K = \frac{\Gamma}{2\pi}$: 点涡强度 $(\Gamma > 0$ 逆时针, $\Gamma < 0$ 顺时针)

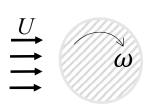
5.绕圆柱有环量流动(均匀流+偶极流+点涡):



$$D = 0$$

D=0 均匀流中任意形状柱体,单 $L=-\rho U\Gamma$ 升力理论 方向为流动沿 Γ 反向旋转90° 均匀流中任意形状柱体,单位长度受升力为 $\rho U\Gamma$,

例:圆柱D=0.5m,长L=3m,以 $\omega=180r$ /min顺时针旋转, 来流空气10m/s。求:所受升力 $F_L=?$



解:
$$\Gamma = (\omega R)(2\pi R) = 2\pi \omega R^2$$

$$F_L = \rho U \Gamma L$$

$$= \rho U 2\pi \omega R^2 L$$

$$= 266.5 N$$

作业:

复习笔记!

8.18, 8.22

大纲

流体力学基础部分

- 1. 基本概念 (2.5)
- 2. 流体静力学(3.5)
- 3. 流体运动学基础(2)
- 4. 流体动力学积分方程(6)
- 5. 流体动力学微分方程(4)
- 6. 粘性不可压流动(7)
- 7. 相似原理(3)
- 8. 无粘不可压势流理论(4)

空气动力学部分

- 1. 绕翼型不可压流动(7)
- 2. 绕机翼不可压流动(7)
- 3. 高速可压流动基础(8)
- 4. 一维定常可压管内流(3)
- 5. 绕翼型亚声速流动(3)
- 6. 绕翼型超、跨声速流动(4)

作业:

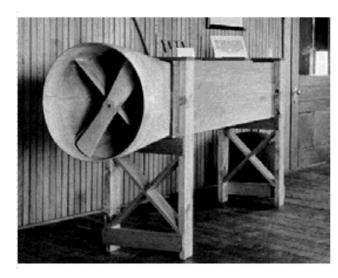
1.阅读1.1(p.1-4)"空气动力学发展",总结发展史。

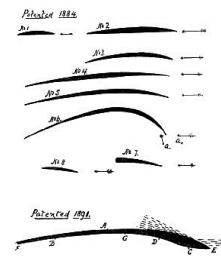
九. 不可压翼型理论(空4, Aero4)

- 9.1 标准大气
- 9.2 翼型几何、气动参数
- 9.3 翼型气动特征
- 9.4 面涡理论
- 9.5 库塔条件
- 9.6 开尔文环量定理、启动涡

- 9.7 经典薄翼理论
- 9.8 有弯度翼型扰流
- 9.9 面元法、涡板块法
- 9.10 粘性流动、翼型阻力

九. 不可压翼型理论(空4, Aero4)





1880s英国H.F.菲利普风洞测试翼型, 德国奥托利林塔尔测试曲线翼滑翔机; 1900~1902美国莱特兄弟测试200余 种翼型,1903人类第一次动力飞行。

早期局限于解释与估算, 1910s空气动力学理论精确计算低速翼型气动力。

1912~1918, Prandtle:

机翼气动分析~

翼剖面(翼型)

翼剖面修正 (三维机翼)

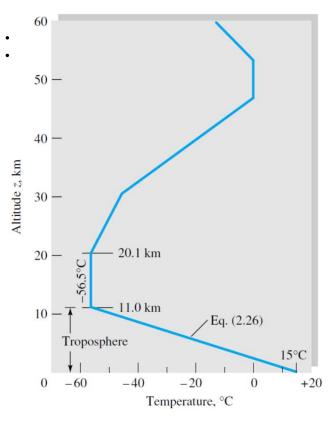
9.1标准大气(p.10~13)

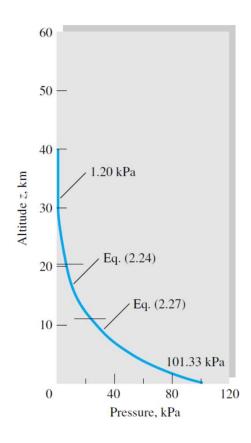
大气温度、压强分布:

$$\frac{dp}{p} = -\frac{g}{R}\frac{dz}{T}$$

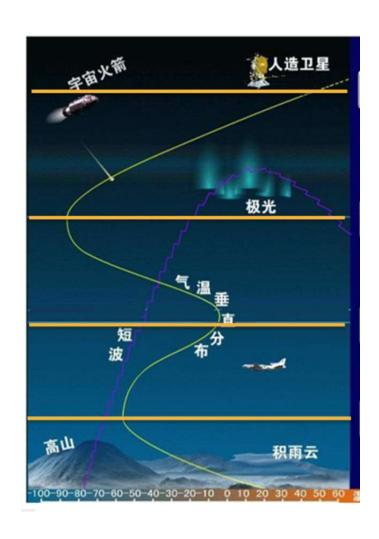
$$T = T_0 - \alpha z$$

$$p_2 = p_1 \exp\left[-\frac{g(z_2 - z_1)}{RT_0}\right]$$





9.1标准大气(p.10~13)



$$\frac{dp}{p} = -\frac{g}{R} \frac{dz}{T}$$

大气温度、成分随高度变化; 大气压强、密度不断降低。

9.1标准大气(p.10~13)

高层大气: 组分不均匀, 受紫外线辐射。

85km

低层大气:

 $N_2 78.1\%$, $O_2 21\%$

平流层底层:

能见度高,受力稳, 噪声小,安全系数高。 大飞机巡航高度

小型飞机、农用、军用飞机

