空气与气体动力学

张科

回顾:

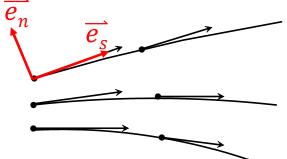
1.动量方程:
$$\rho \frac{D\vec{V}}{Dt} = \rho \vec{g} - \vec{\nabla} p + \vec{\nabla} \cdot \tau_{ij}$$

2.N-S方程:
$$\rho \frac{D\vec{V}}{Dt} = \rho \vec{g} - \vec{\nabla} p + \mu \nabla^2 \vec{V}$$

3.欧拉方程:
$$\rho \frac{D\vec{V}}{Dt} = \rho \vec{g} - \vec{\nabla} p$$

$$\rho \frac{D\vec{V}}{Dt} = \rho \vec{g} - \vec{\nabla} p + \mu \nabla^2 \vec{V}$$

欧拉方程: $\rho \frac{D\vec{V}}{Dt} = -\vec{\nabla} p_k$



欧拉 方程, 定常,沿 流线和流线法向积分!

欧拉方程:
$$\rho \frac{D\vec{V}}{Dt} = -\vec{\nabla} p_k$$

$$\rho\left[\frac{\partial \vec{V}}{\partial t} + (\vec{V} \cdot \vec{\nabla}) \vec{V}\right] = -\vec{\nabla} p_k$$

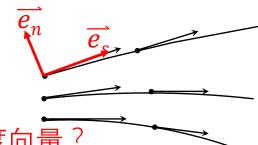
$$\rho(\vec{V}\cdot\vec{\nabla})\,\vec{V}\,=-\vec{\nabla}p_k$$

$$\rho V \frac{\partial (V \overrightarrow{e_s})}{\partial s} = -\frac{\partial p_k}{\partial s} \overrightarrow{e_s} + \frac{\partial p_k}{\partial n} \overrightarrow{e_n}$$

$$\frac{\partial (V\overrightarrow{e_s})}{\partial s} \times \frac{\partial V}{\partial s} \overrightarrow{e_s} ? ?$$

$$\frac{\frac{\partial(V\overrightarrow{e_s})}{\partial s}}{\frac{\partial s}{\partial s}} = \frac{\partial V}{\partial s}\overrightarrow{e_s} + \frac{\partial \overrightarrow{e_s}}{\partial s}V \longrightarrow \frac{\partial(V\overrightarrow{e_s})}{\partial s} = \frac{\partial V}{\partial s}\overrightarrow{e_s} - V\frac{\overrightarrow{e_n}}{R} \qquad \rho(\overrightarrow{V} \cdot \overrightarrow{\nabla})\overrightarrow{V} = \rho V\frac{\partial(V\overrightarrow{e_s})}{\partial s}$$

$$\frac{\partial \overrightarrow{e_s}}{\partial s} = -\frac{\overrightarrow{e_n}}{R}$$



沿流向和法向梯度向量?

定常:
$$\frac{\partial}{\partial t} = 0$$

$$\vec{\nabla} = \frac{\partial}{\partial s} \vec{e_s} + \frac{\partial}{\partial n} \vec{e_n}$$

$$\vec{V} = u_s \vec{e_s} + y_n \vec{e_n} = V \vec{e_s}$$

$$\vec{V} \cdot \vec{\nabla} = V \frac{\partial}{\partial s}$$

$$\rho(\vec{V} \cdot \vec{\nabla}) \vec{V} = \rho V \frac{\partial (V \vec{e_s})}{\partial s}$$

欧拉方程:
$$\rho \frac{D\vec{V}}{Dt} = -\vec{\nabla} p_k$$

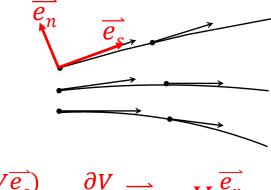
$$\rho \left[\frac{\partial \vec{V}}{\partial t} + (\vec{V} \cdot \vec{\nabla}) \vec{V} \right] = -\vec{\nabla} p_k$$

$$\rho (\vec{V} \cdot \vec{\nabla}) \vec{V} = -\vec{\nabla} p_k$$

$$\rho V \frac{\partial (V\vec{e_s})}{\partial s} = -\frac{\partial p_k}{\partial s} \vec{e_s} - \frac{\partial p_k}{\partial n} \vec{e_n}$$

$$\rho V \left(\frac{\partial V}{\partial s} \vec{e_s} - V \frac{\vec{e_n}}{R} \right) = -\frac{\partial p_k}{\partial s} \vec{e_s} - \frac{\partial p_k}{\partial n} \vec{e_n}$$

$$\rho V \frac{\partial V}{\partial s} \vec{e_s} - \frac{\rho V^2}{R} \vec{e_n} = -\frac{\partial p_k}{\partial s} \vec{e_s} - \frac{\partial p_k}{\partial n} \vec{e_n}$$



$$\frac{\partial (V\overrightarrow{e_s})}{\partial s} = \frac{\partial V}{\partial s} \overrightarrow{e_s} - V \frac{\overrightarrow{e_n}}{R}$$

$$s = -\frac{\partial p_k}{\partial s} = -\frac{\partial p_k}{\partial s}$$

$$\mathbf{n} = \frac{\partial p_k}{\partial n} = \frac{\rho V^2}{R}$$

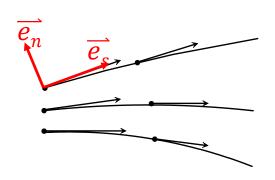
欧拉方程:
$$\rho \frac{D\vec{V}}{Dt} = -\vec{\nabla} p_k$$

定常: s向:
$$\rho V \frac{\partial V}{\partial s} = -\frac{\partial p_k}{\partial s}$$

不可压:
$$\frac{\partial}{\partial s} \left[\frac{\rho V^2}{2} + p_k \right] = 0$$

沿流线:
$$\frac{\rho V^2}{2} + p_k = C$$

$$\frac{\rho V^2}{2} + p + \rho gz = C$$



$$n = \frac{\partial p_k}{\partial n} = \frac{\rho V^2}{R} > 0$$

$$\frac{\vec{n}}{p_3} \qquad \frac{\partial p_k}{\partial n} > 0$$

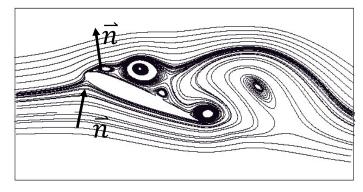
$$p_2 \qquad p_1$$

$$p_3 > p_2 > p_1$$

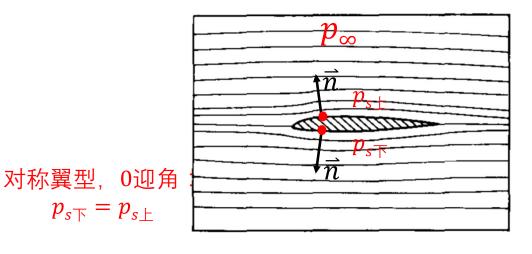
伯努利方程(定常、不可压、沿流线、无粘)

欧拉方程: $\rho \frac{D\vec{V}}{Dt} = -\vec{\nabla} p_k$

 $n = \frac{\partial p_k}{\partial n} = \frac{\rho V^2}{R} > 0$ $\frac{\partial p_k}{\partial n} > 0$

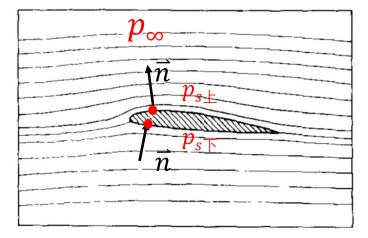


流动分离区粘性不可忽略, 压力接近于0;



$$p_{s\perp} < p_{\infty}$$
 $p_{s\top} < p_{\infty}$

 $p_{s\top} = p_{s\perp}$

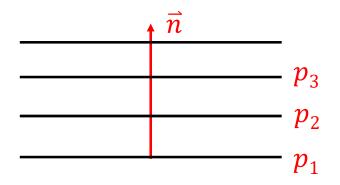


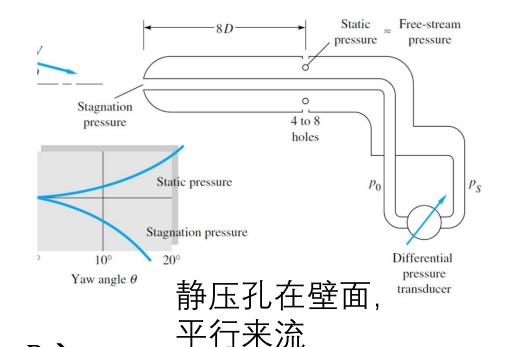
$$p_{s+} < p_{\infty}$$
 $p_{s\mp} > p_{\infty}$

 $p_{s \top} > p_{s \perp}$ 升力!

欧拉方程:
$$\rho \frac{D\vec{V}}{Dt} = -\vec{\nabla} p_k$$

$$n = \frac{\partial p_k}{\partial n} = \frac{\rho V^2}{R} > 0 \quad \frac{\partial p_k}{\partial n} > 0$$





$$R \rightarrow \infty$$

$$\frac{\partial p_k}{\partial n} = \frac{\rho V^2}{R} \approx 0$$

$$p_3 = p_2 = p_1$$

垂直平行流线各处静压不变。

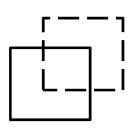
$$\frac{dudt}{dy}$$
角变形率 $\frac{d\alpha}{dt} = \frac{du}{dy}$

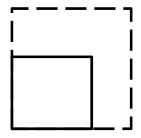
$$\rho \frac{D\vec{V}}{Dt} = \rho \vec{g} - \vec{\nabla} p + \vec{\nabla} \cdot \tau_{ij} \qquad \tau_{ij} \propto \frac{\partial u_i}{\partial x_i}$$

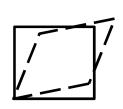
$$\tau_{ij} \propto \frac{\partial u_i}{\partial x_i}$$

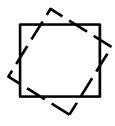
$$\rho \frac{D\vec{V}}{Dt} = \rho \vec{g} - \vec{\nabla} p + \mu \nabla^2 \vec{V}$$

流体运动和变形:





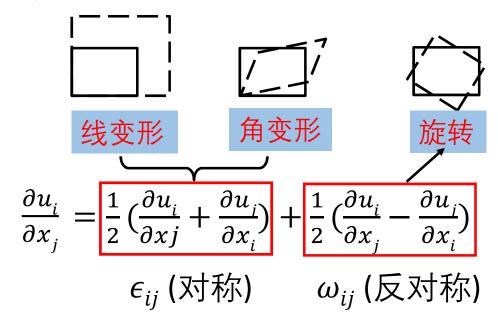




线变形

角变形

① 流体运动和变形:



$$\epsilon_{11}$$
、 ϵ_{22} 、 ϵ_{33} **>**线性变形率 ϵ_{12} 、 ϵ_{23} 、 ϵ_{13} **>**角变形率

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = \vec{\nabla} \cdot \vec{V}$$

流体微团的相对体膨胀率

$$\vec{\nabla} \cdot \vec{V} = 0$$

流体微团体积不变, 即不可压!

1 流体运动和变形:



$$\omega_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} - \frac{\partial u_j}{\partial x_i} \right)$$
 旋转

旋转角速度: $\vec{\omega} = \frac{1}{2}\vec{\nabla} \times \vec{V}$

$$= \frac{1}{2} \begin{vmatrix} i & j & k \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ u & v & w \end{vmatrix} \qquad \omega_y = \frac{1}{2} \left(\frac{\partial u}{\partial z} - \frac{\partial u}{\partial x} \right)$$
$$\omega_z = \frac{1}{2} \left(\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \right)$$

$$\omega_{x} = \frac{1}{2} \left(\frac{\partial w}{\partial y} - \frac{\partial v}{\partial z} \right)$$

$$\omega_y = \frac{1}{2} \left(\frac{\partial u}{\partial z} - \frac{\partial w}{\partial x} \right)$$

$$\omega_z = \frac{1}{2} \left(\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \right)$$



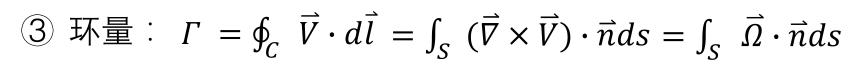
② 涡量:
$$\vec{\Omega} = \vec{\nabla} \times \vec{V} = 2\vec{\omega}$$
 $\Omega_x = (\frac{\partial w}{\partial y} - \frac{\partial v}{\partial z}) \Omega_y = (\frac{\partial u}{\partial z} - \frac{\partial w}{\partial x})$

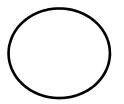
反应流体质点旋转角速度。 $\Omega_z = (\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y})$

$$\Omega_z = \left(\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}\right)$$

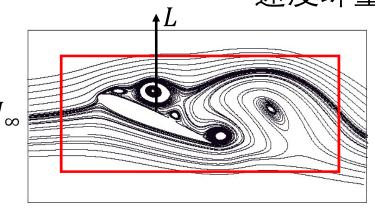
$$\vec{\Omega} = 0$$
, 无旋流动; $\vec{\Omega} \neq 0$, 有旋流动

$$\vec{\Omega} \neq 0$$
, 有旋流动





速度环量 涡通量



$$L = \rho U_{\infty} \Gamma$$

$$\vec{\Omega} = 0 \longrightarrow \Gamma = 0$$

$$\Gamma = 0 \Longrightarrow \vec{\Omega} = 0$$
 ? ?

作业:

复习笔记!

P130. 4.21, 4.8, 4.3

P99. 3.18, 3.20

七. 量纲分析与相似原理(7.1-7.5)

1. 简介

1 提供实验相似准则





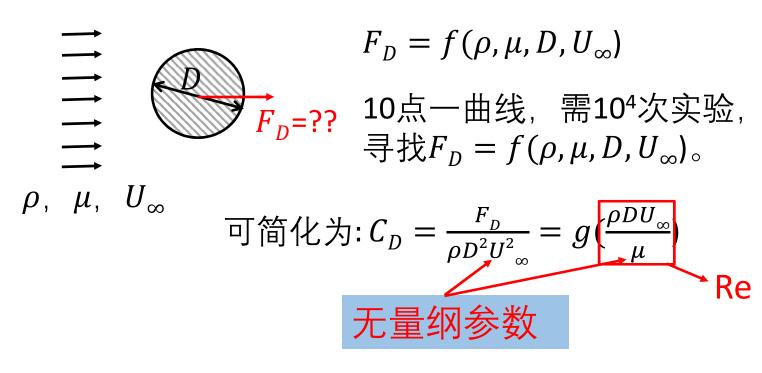




实验参数如何选择?

1.简介 量纲分析!

② 多变量组合 →减少变量→分析数据、简化实验



 $C_D = g(\text{Re})$,仅需10次实验!

1.简介

③ 分析各项相对重要性

可简化为:
$$CD = \frac{F_D}{\rho D^2 U^2_{\infty}} = g(\frac{\rho D U_{\infty}}{\mu})$$

$$F^* \sim \mu^*$$

$$C_D' = \frac{\rho F_D}{\mu^2} = f(\frac{\rho D U_{\infty}}{\mu})$$

$$F^* \sim U_{\infty}^* (L^*)$$

2.量纲分析

量纲一致性原则:
$$F_D = f(\rho, \mu, D, U_{\infty}) \quad \dim(F_D) = MLT^{-2}$$

$$\frac{F_D}{\rho D^2 U^2} = g(\frac{\rho D U_{\infty}}{\mu}) \quad \dim\left(\frac{F_D}{\rho D^2 U^2}\right) = 1$$

量纲分析:组合物理量,使其成为量纲为1的组合变量,减少变量数。量纲为一化,或无量纲化!

2.量纲分析

$$\frac{p}{\rho} + \frac{V^2}{2} + gz = C_1$$
 $p + \frac{\rho V^2}{2} + \rho gz = C_2$

p,V,z——有量纲变量,无量纲化目标

 ρ , g——有量纲常数

无量纲化方法: ① 未知表达式——Π原理

- ② 已知表达式——方程无量纲化
- ◆ 所有变量都用到;
- ◆ 不同选择,不同含义。

3. □ □ E. Buckingham 1914年

- ① 物理过程必须可表达为无量纲(量纲为一)参数的关系式;
- ② 无量纲变量为Ⅱ(有量纲变量的乘积形式)。

Ⅱ的个数=变量数-问题基本量纲数

$$k = n - m$$

$$g(\phi_1,\phi_2,\ldots,\phi_n)=0$$
 n 个变量,m个基本量纲

$$G(\Pi_1, \Pi_2, \dots, \Pi_n \quad m)=0$$
 $k=n-m$

∏不唯一!

$$g(\phi_1, \phi_2, \dots, \phi_n)=0$$

 $G(\Pi_1, \Pi_2, \dots, \Pi_{n-m})=0$

- ③ 无量纲化步骤:
- 1) 确定问题的影响参数n, (独立变量);
- 2) 确定问题基本变量数*m*;
- 3) 确定 Π 个数k = n m;
- 4) 选m个变量(不能互相组成 Π 、包含所有量纲、含单一基本量纲);
- 5) m个变量与剩余n-m个组成 Π 。

例: $F_D = f(\rho, \mu, D, U_\infty)$

解: 1) 确定问题的影响参数n=5

F_D	ρ	μ	D	U_{∞}
MLT^{-2}	ML^{-3}	$ML^{-1}T^{-1}$	L	LT^{-1}

- 2) 确定基本量纲数*m* =3;
- 3) 确定 Π 个数k = n m = 2; 2个 Π ;
- 4) 选D, U_{∞} , ρ , 来无量纲化 F_D , μ ;

F_{D}	ρ	μ	D	U_{∞}
MLT^{-2}	ML^{-3}	$ML^{-1}T^{-1}$	L	LT^{-1}

例:
$$F_D = f(\rho, \mu, D, U_\infty)$$

- 解: 4) 选D, U_{∞} , ρ ,来无量纲化 F_D , μ ;
 - 5) m个变量与剩余n-m个组成 Π 。

$$\Pi_1 = F_D D^a U^b_{\infty} \rho^c = M^0 L^0 T^0$$

$$(MLT^{-2})L^a(LT^{-1})^b(ML^{-3})^c = M^0L^0T^0$$

$$\begin{array}{l}
M: 1+c=0 \\
L: 1+a+b-3c=0 \\
T: -2-b=0
\end{array}
\qquad
\begin{array}{l}
a=-2 \\
b=-2 \longrightarrow \Pi_1 = F_D D^{-2} U^{-2} {}_{\infty} \rho^{-1} \\
c=-1 = F_D /(\rho U^2 {}_{\infty} D^2)
\end{array}$$

$$F_D$$
 ρ μ D U_∞ ML^{-2} ML^{-3} $ML^{-1}T^{-1}$ L LT^{-1}

例:
$$F_D = f(\rho, \mu, D, U_\infty)$$

解:
$$\Pi_1 = F_D / (\rho U^2_{\infty} D^2)$$

$$\Pi_2 = \mu D^a U^b_{\infty} \rho^c = M^0 L^0 T^0$$

$$(ML^{-1}T^{-1})L^a(LT^{-1})^b(ML^{-3})^c = M^0L^0T^0$$

$$M: 1 + c = 0
L: -1 + a + b - 3c = 0$$

$$T: -1 - b = 0$$

$$a = -1
b = -1
c = -1$$

$$c = -1$$

$$\Pi_1 = g(\Pi_2) \implies \frac{F_D}{\rho U_{\infty}^2 D^2} = g(\frac{\mu}{\rho U_{\infty} D})$$

作业:

复习笔记!

例7.1用∏原理做, 7.1, 7.7

回顾:

- **1.**欧拉方程: 沿流线: $\frac{\rho V^2}{2} + p_k = C$ n向: $\frac{\partial p_k}{\partial n} = \frac{\rho V^2}{R}$
- 2.流体变形、涡量、环量
- 3. ∏原理