

空气与气体动力学

张科

回顾：

$$\vec{\nabla}^* \cdot \vec{V}^* = 0$$

1. Π 原理

$$\frac{D\vec{V}^*}{Dt^*} = -\vec{\nabla}^* p^* + \frac{1}{Re} \nabla^{*2} \vec{V}^* - \frac{1}{Fr^2} \vec{\nabla}^* z^*$$

2. 无量纲化N-S方程

$$\vec{V}^* = \vec{V}^* (Re) \quad p^* = p^* (Re)$$

3. 无量纲参数

$$Re = \frac{\rho U_0 L_0}{\mu} \quad Fr = \frac{U_0}{\sqrt{gL_0}} \quad Ma = \frac{U}{a} \quad Eu = \frac{\Delta p}{0.5 \rho U^2}$$

4. 相似准则及应用

$$St = \frac{fL}{U} \quad We = \frac{\rho U^2 L}{\sigma}$$

6. 流动相似原则 (similarity)

设计模型实验，分析实验数据！
例7.5~7.10

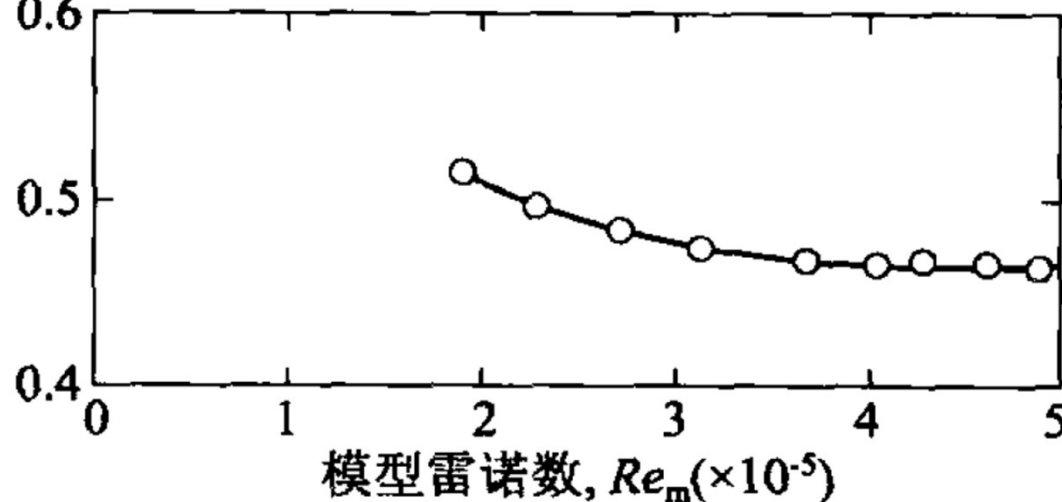
例 7.8：客车风洞实验，1/16模型(迎风面积 $A_m = 0.0305m^2$,模型宽度 $W_m = 0.152m$)，得数据如下： $\mu = 1.789 \times 10^{-5} Pa \cdot s, \rho = 1.225 kg/m^3$

速度 $U/(m/s)$	10.0	21.0	25.0	30.1	35.0	40.5	46.0	46.7
阻力 F_D/N	3.1	12.0	15.0	18.1	21.0	23.5	24.0	18.9

求：1. 绘制 C_D 图
2. 实物以100km/h速度行驶时，阻力系数

解：1. $C_D = \frac{F_D}{\frac{1}{2} \rho U^2 A_m}$

模型阻力因数, C_{Dm}



低气流速度。
阻力系数。

$$Re = \frac{\rho U_{\infty} W_m}{\mu} = \frac{1.225 \times 0.152}{1.789 \times 10^{-5}} \times U_{\infty} = 1.04 \times 10^4 U_{\infty}$$

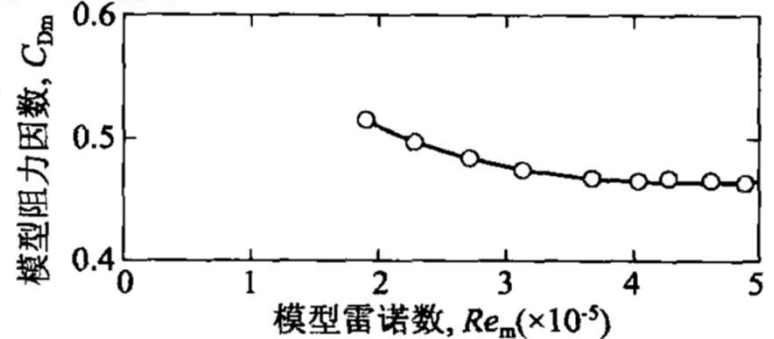
$$Re > 4 \times 10^5 \text{ 后, } C_D \approx 0.46 \rightarrow U_{\infty} > 40 m/s$$

6. 流动相似原则 (similarity)

设计模型实验，分析实验数据！
例7.5~7.10

例 7.8: 客车风洞实验，1/16模型(迎风面积 A_n
 $W_m = 0.152m$)，得数据如下： $\mu = 1.789 \times 10^{-5} I$

求：1. 绘制 $C_D(\frac{F_D}{0.5\rho U_\infty^2 A_m}) - Re(\frac{\rho U_\infty W_m}{\mu})$ 曲线， C



2. 实物以 $100km/h$ 行驶时的气动阻力，及克服其所需功率。

解：2. $U_{\infty p} = 100km/h = 100/3.6 = 27.78m/s$

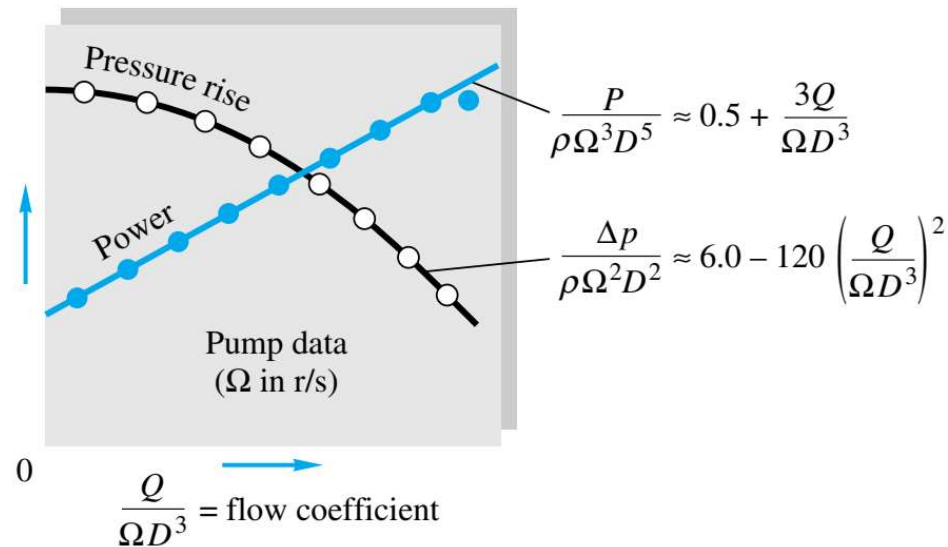
$$Re = \frac{\rho U_{\infty p} W_p}{\mu} = \frac{1.225 \times 27.78 \times 0.152 \times 16}{1.789 \times 10^{-5}} = 4.63 \times 10^6 \rightarrow C_D = 0.46$$

$$F_D = C_D \times 0.5 \rho U_{\infty p}^2 A_p = 0.46 \times 0.5 \times 1.225 \times 27.78^2 \times 0.0305 \times 16^2 = 1.698 \times 10^3 N$$

$$\rightarrow \dot{W} = F_D \times U_{\infty p} = 47.2 \times 10^3 W$$

作业:

P5.61 If viscosity is neglected, typical pump-flow results from Prob. 5.20 are shown in Fig. P5.61 for a model pump tested in water. The pressure rise decreases and the power required increases with the dimensionless flow coefficient. Curve-fit expressions are given for the data. Suppose a similar pump of 12-cm diameter is built to move gasoline at 20°C and a flow rate of 25 m³/h. If the pump rotation speed is 30 r/s, find (a) the pressure rise and (b) the power required.



P5.61

单选题 2分

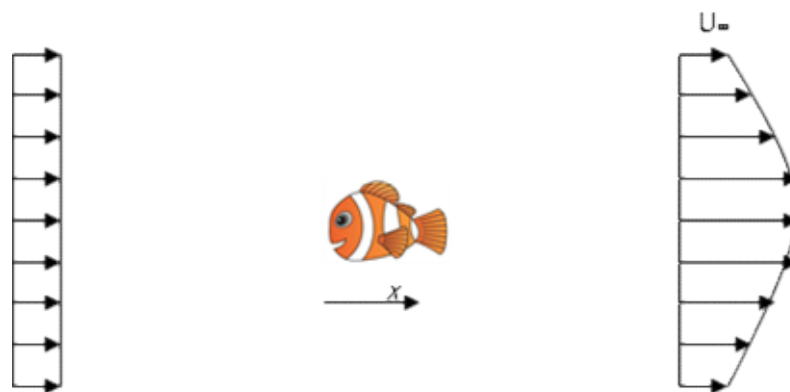
设置

此题未设置答案，请点击右侧设置按钮

鱼在水中运动过程中，前后速度分布如下图所示。

问：(1) 鱼受到水对其的作用力沿x_____？

(2) 需要给鱼施加沿x_____的力才能使鱼保持匀速运动？



☐ A 正向

☐ B 负向

提交

六. 粘性不可压流动（内流，外流）

粘性不可压内流

通道内流动一般特征（5.6、9.1）、无限大平板间（周向均匀圆管）充分发展层流（5.3、5.4）、管内流能量损失(9.2-9.4)、

管内流, \bar{V} , τ , Q , Δp , h_{LT}

粘性不可压外部扰流

边界层基本概念（10.1）、边界层动量积分方程（10.4）、边界层方程(10.2、10.3)、曲面边界层及边界层分离（10.5）、扰流物体的阻力（10.6）

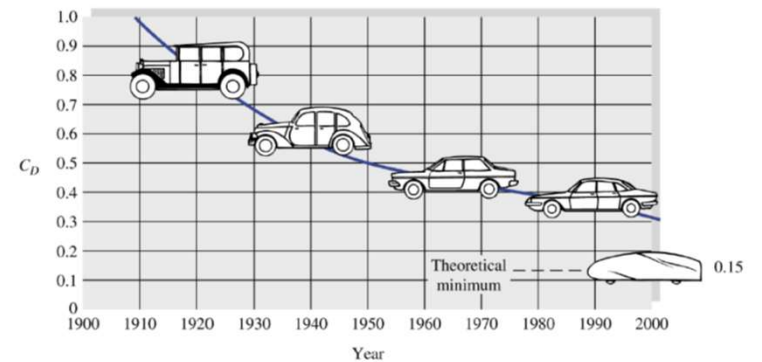
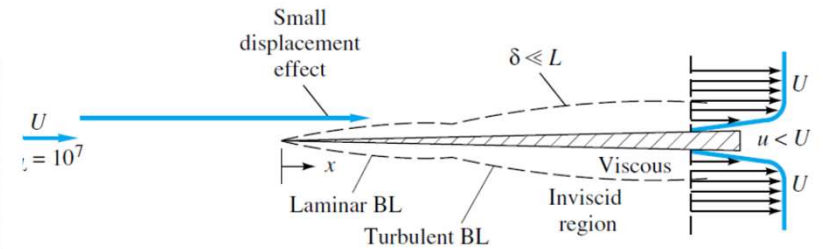
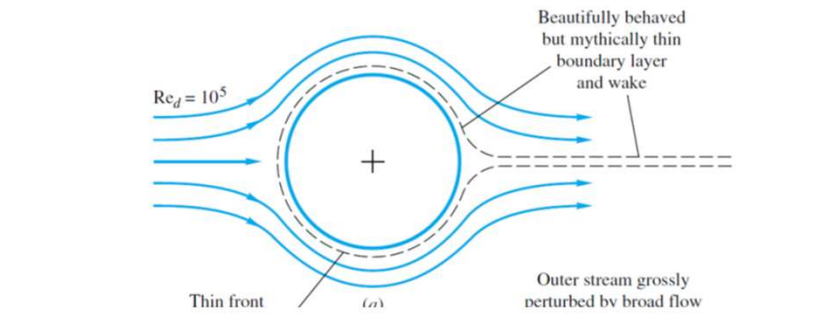
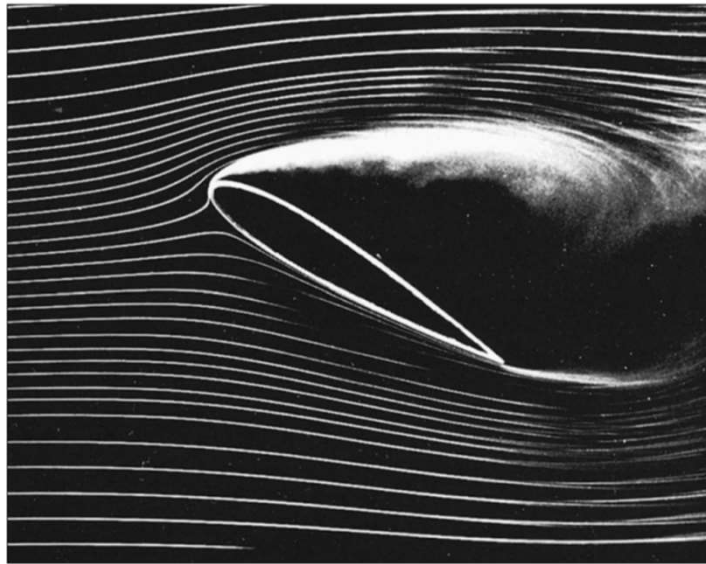
流场分布, u, D, L (飞行器, 建筑物, 桥梁等)

六. 粘性不可压流动（内流，外流）



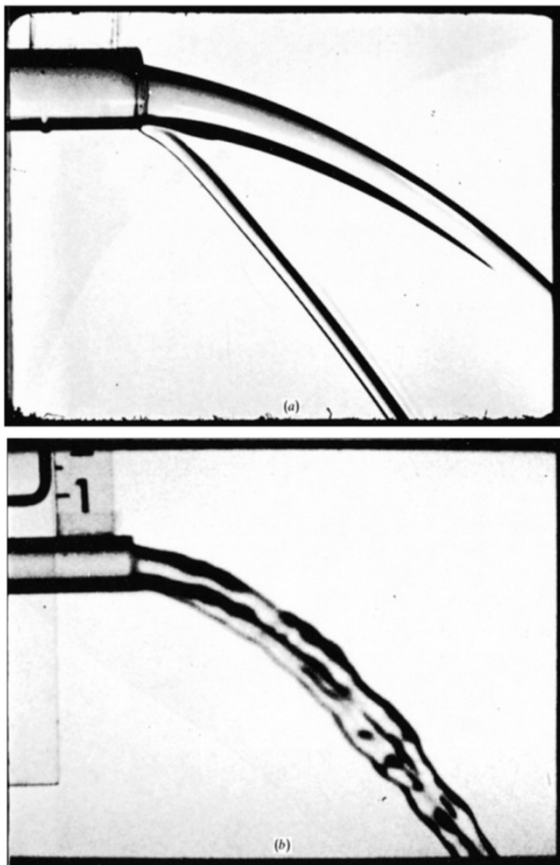
Steam pipe bridge in a geothermal power plant. Pipe flows are everywhere, often occurring in groups or networks. They are designed using the principles outlined in this chapter. (Courtesy of Dr. E. R. Degginger/Color-Pic Inc.)

- 边界层基本概念、边界层动量积分方程、边界层方程、曲面边界层及边界层分离、扰流物体的阻力



6.1 通道内流动一般特征

① 层流、湍流 (5.6)



1839年德国工程师
G.H.L. Hagen发现粘性
流存在两个区域。

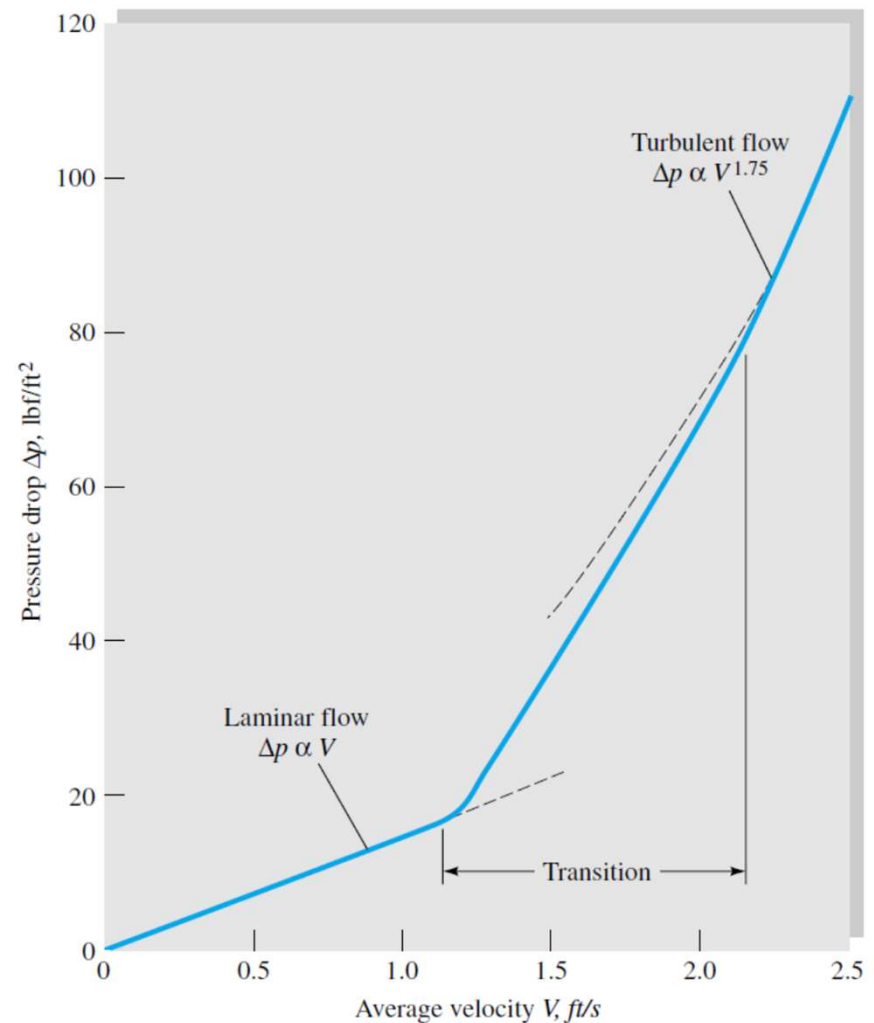
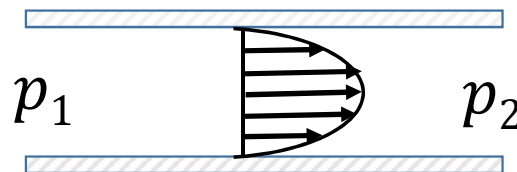
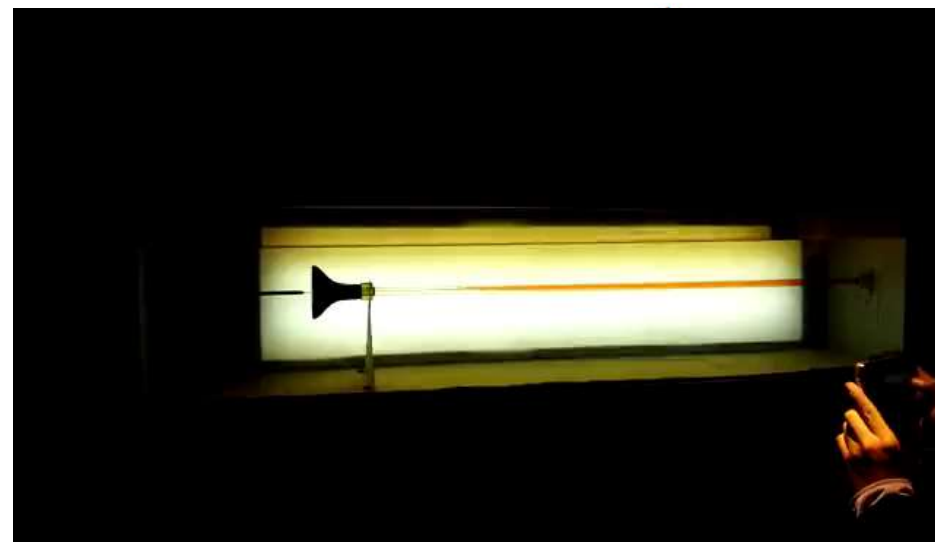
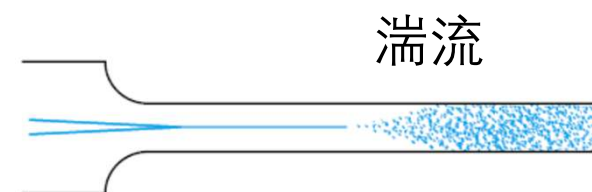
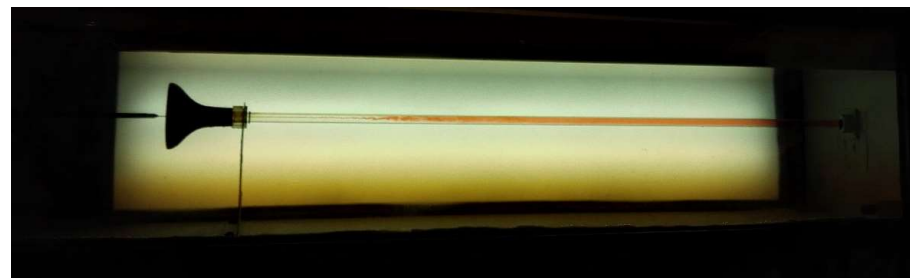


Fig. 6.4 Experimental evidence of transition for water flow in a $\frac{1}{4}$ -in smooth pipe 10 ft long.

6.1 通道内流动一般特征

① 层流、湍流 (5.6)

1883年英国学者Osborne Reynolds实验表明,
粘性流随 $\frac{\rho V d}{\mu}$ (Re 数) 变化。

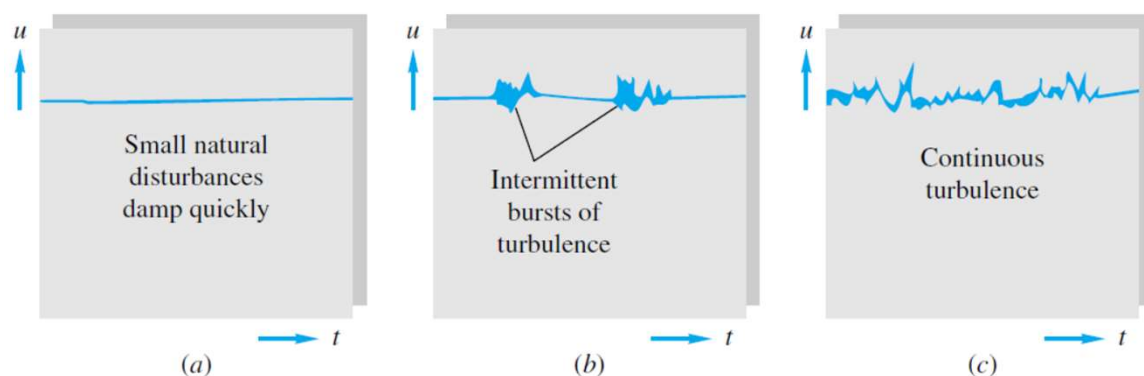


condition (b). (From Ref. 4.)

6.1 通道内流动一般特征

① 层流、湍流 (5.6)

1883年英国学者Osborne Reynolds实验表明，粘性流随 $\frac{\rho V d}{\mu}$ (Re 数) 变化。



$$\bar{V}_a < \bar{V}_b < \bar{V}_c$$

$$Re_a < Re_b < Re_c$$

$$\bar{V} = \frac{1}{A} \int u dA$$

$$Re = \frac{\rho \bar{V} D}{\mu}$$

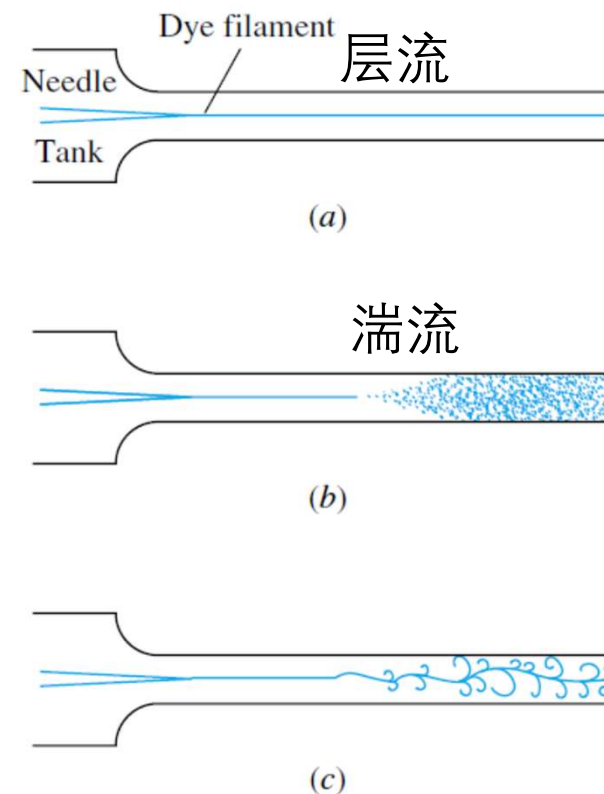


Fig. 6.5 Reynolds' sketches of pipe-flow transition: (a) low-speed, laminar flow; (b) high-speed, turbulent flow; (c) spark photograph of condition (b). (From Ref. 4.)

6.1 通道内流动一般征

① 层流、湍流 (5.6)

1883年英国学者Osborne Reynolds实验表明, 粘性流随 $\frac{\rho V d}{\mu}$ (Re 数) 变化。

$$\bar{V} = \frac{1}{A} \int u dA \quad Re = \frac{\rho \bar{V} D}{\mu}$$

圆管： $Re < 2000 \sim 2300$ 为层流。

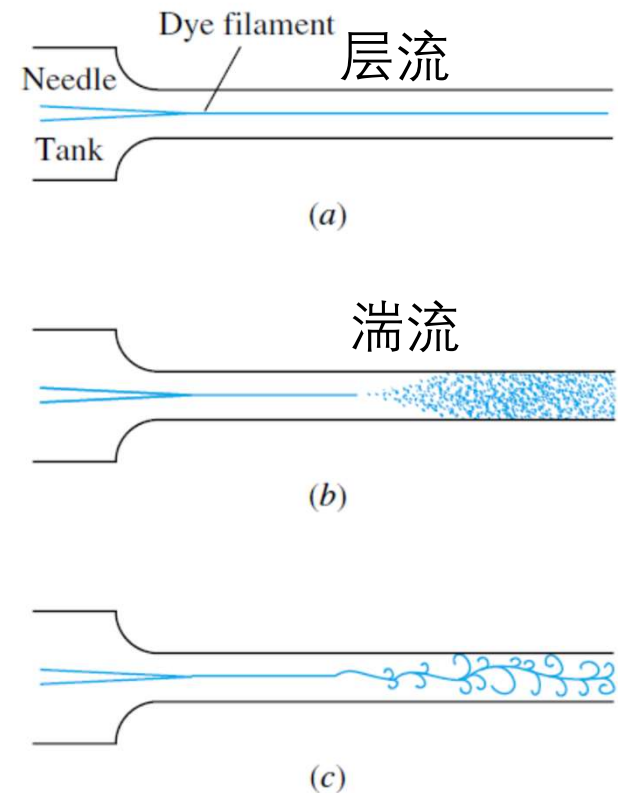
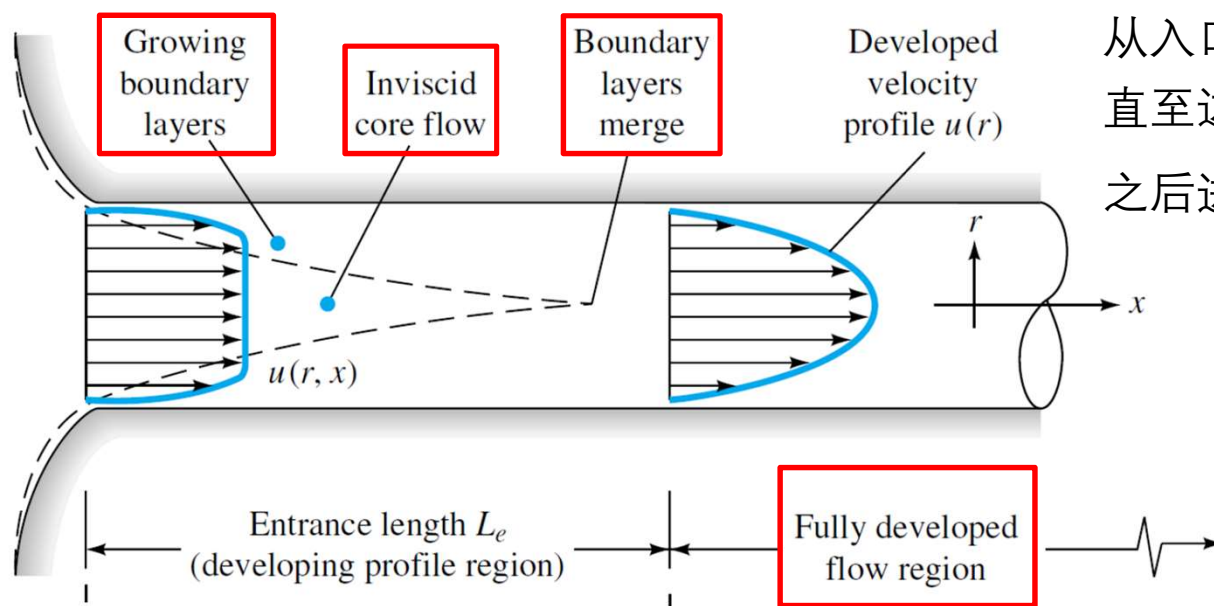


Fig. 6.5 Reynolds' sketches of pipe-flow transition: (a) low-speed, laminar flow; (b) high-speed, turbulent flow; (c) spark photograph of condition (b). (From Ref. 4.)

6.1 通道内流动一般征

② 起始段、充分发展流动 (9.1)



入口段长度 L_e

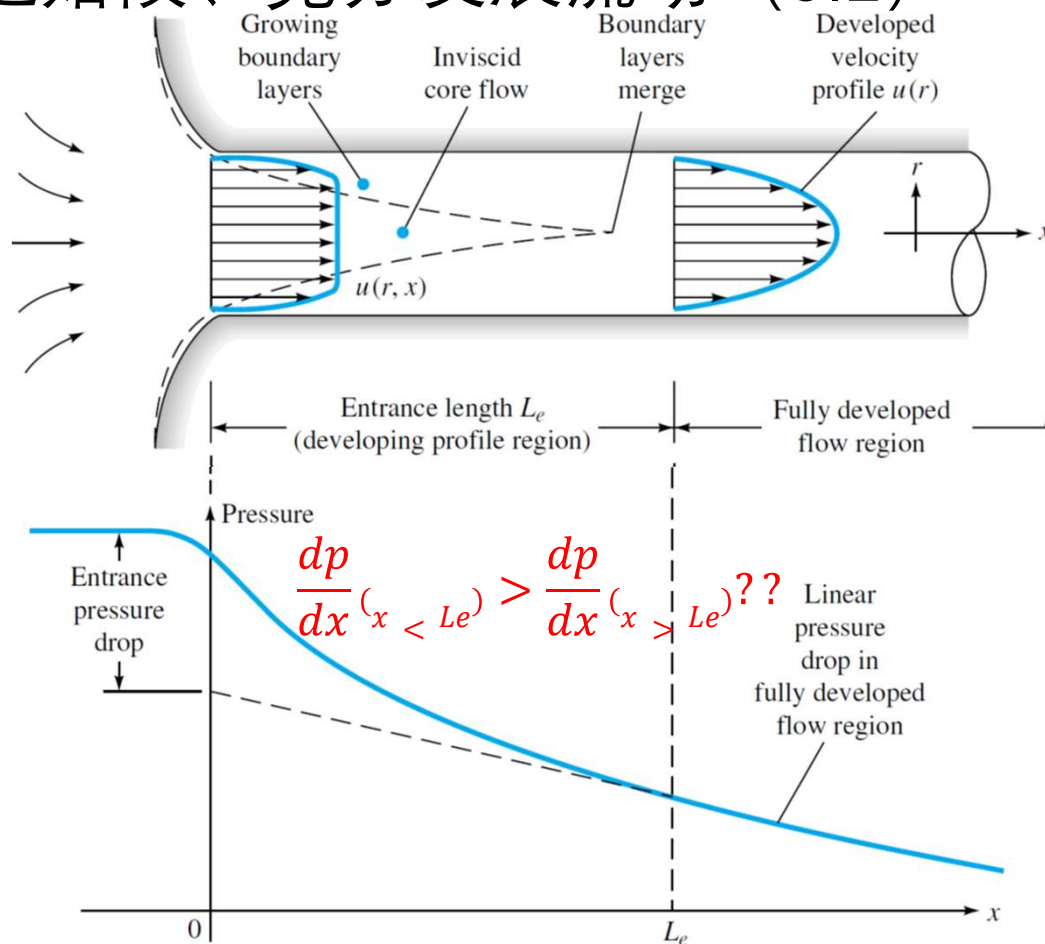
$x > L_e$ 充分发展流动

$u = u(r)$ 不再随 x 变化

内流发展受壁面限制，
壁面附近为有速度梯度的边界层区，
远离壁面为速度均匀的无粘核心区
从入口向下游壁面粘性影响范围逐渐扩大，
直至边界层在轴线处相交，
之后进入充分发展流动， $\frac{\partial u}{\partial x} = 0$

6.1 通道内流动一般征

② 起始段、充分发展流动 (9.1)



充分发展后 $\frac{\partial u}{\partial x} = 0, u = u(r)$ 不变

$$\tau = \frac{du}{dr} = \text{constant}$$

$$\Delta p \propto x$$

入口段长度 L_e

层流： $L_e/D \approx 0.06Re$

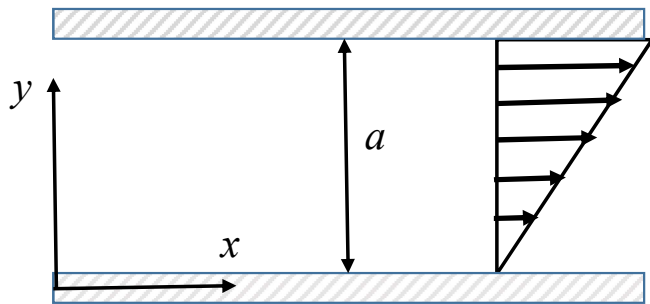
$Re = 2000, L_e \approx 120 D$

湍流： $L_e/D \approx 4.4Re^{1/6}$

$L_e \approx 25 \sim 40 D$

$L_{e\text{层}} > L_{e\text{湍}}$??

6.2 无限大平板间充分发展层流 (5.3)



求： \vec{V} , τ , Q , Δp 解微分连续性方程、N-S方程。

假设：牛顿流体、不可压， x 方向无重力、
充分发展流动($\frac{\partial \vec{V}}{\partial x} = 0$)、 z 向无限大(2D)、定常($\frac{\partial}{\partial t} = 0$)

$$\frac{L_e}{a} < \frac{x}{a}$$

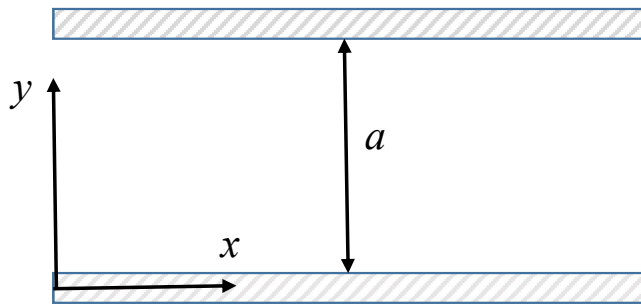
$$\frac{a}{W} \ll 1$$

解：定常、不可压连续性方程： $\vec{\nabla} \cdot \vec{V} = 0$

$$\begin{aligned} \frac{\partial \vec{V}}{\partial x} = 0 &\rightarrow \left. \begin{array}{l} \cancel{\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0} \\ \frac{\partial u}{\partial x} = 0 \quad \frac{\partial v}{\partial x} = 0 \end{array} \right\} \begin{array}{l} \frac{\partial v}{\partial y} = 0, \\ \frac{\partial v}{\partial x} = 0 \end{array} \rightarrow \left. \begin{array}{l} v = C \\ v = 0 \\ @y = 0, a \end{array} \right\} v = 0 \end{aligned}$$

$\vec{V} = (u(y), 0)$

6.2 无限大平板间充分发展层流 (5.3)



求： \vec{V} , τ , Q , Δp 解微分连续性方程、N-S方程。

假设：牛顿流体、不可压， x 方向无重力、
充分发展流动($\frac{\partial \vec{V}}{\partial x} = 0$)、 z 向无限大(2D)、定常($\frac{\partial}{\partial t} = 0$)

$$\frac{L_e}{a} < \frac{x}{a}$$

$$\frac{a}{W} \ll 1$$

解：N-S方程： $\rho \frac{D\vec{V}}{Dt} = -\vec{\nabla} p_k + \mu \nabla^2 \vec{V}$ $p_k = p + \rho g y$

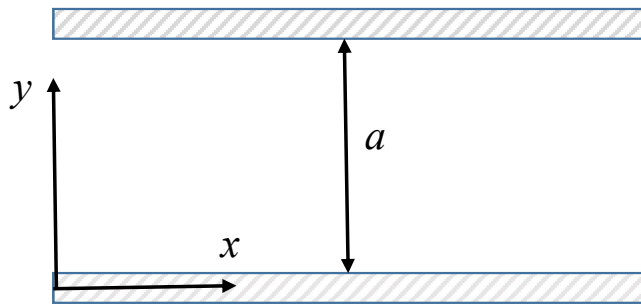
x 方向：

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = -\frac{1}{\rho} \frac{\partial p}{\partial x} + \nu \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right)$$

$\frac{\partial u}{\partial x} = 0$ $v = 0$ $\frac{\partial u}{\partial x} = 0$

$\Rightarrow \nu \frac{\partial^2 u}{\partial y^2} = \frac{1}{\rho} \frac{\partial p}{\partial x}$

6.2 无限大平板间充分发展层流 (5.3)



求： \vec{V} , τ , Q , Δp 解微分连续性方程、N-S方程。

假设：牛顿流体、不可压， x 方向无重力、
充分发展流动($\frac{\partial \vec{V}}{\partial x} = 0$)、 z 向无限大(2D)、定常($\frac{\partial}{\partial t} = 0$)

$$\frac{L_e}{a} < \frac{x}{a}$$

$$\frac{a}{W} \ll 1$$

解：N-S方程： $\rho \frac{D\vec{V}}{Dt} = -\vec{\nabla} p_k + \mu \nabla^2 \vec{V}$ $p_k = p + \rho g y$

y 方向：

$$u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} = -\frac{1}{\rho} \frac{\partial p}{\partial y} + \nu \left(\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} \right)$$

$v=0$

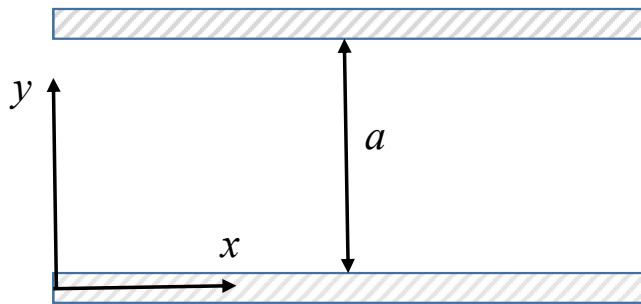
→ $\frac{\partial p}{\partial y} = 0$

z 方向： $w = 0$

$$\frac{\partial p}{\partial z} = 0$$

→ $p = p(x)$ only

6.2 无限大平板间充分发展层流 (5.3)



求： \vec{V} , τ , Q , Δp 解微分连续性方程、N-S方程。

假设：牛顿流体、不可压， x 方向无重力、
充分发展流动($\frac{\partial \vec{V}}{\partial x} = 0$)、 z 向无限大(2D)、定常($\frac{\partial}{\partial t} = 0$)

$$\frac{L_e}{a} < \frac{x}{a}$$

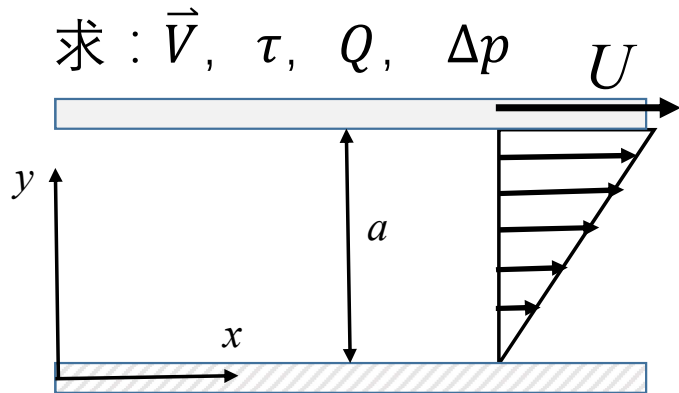
$$\frac{a}{W} \ll 1$$

解：

$$\left. \begin{aligned} \nu \frac{\partial^2 u}{\partial y^2} &= \frac{1}{\rho} \frac{\partial p}{\partial x} \\ u &= u(y), p = p(x) \end{aligned} \right\} \frac{dp}{dx} = \mu \frac{d^2 u}{dy^2} \quad ODE \text{常微分方程}$$

$$f(x) = f(y) = \text{constant}$$

6.2 无限大平板间充分发展层流 (5.3)

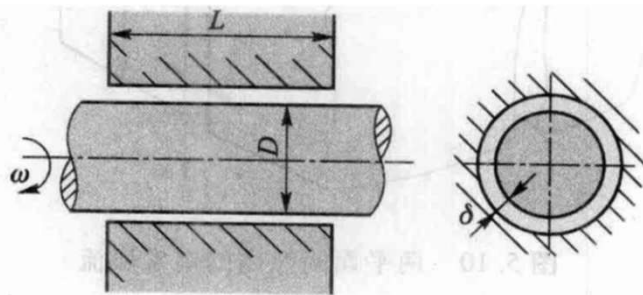
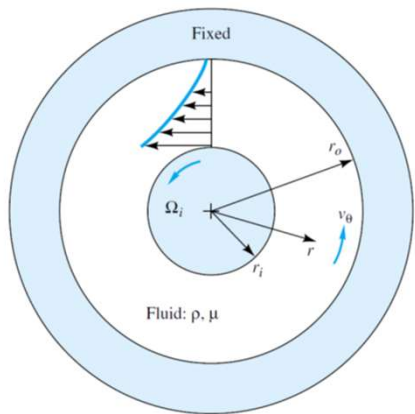


求: \vec{V} , τ , Q , Δp 解: $\frac{dp}{dx} = \mu \frac{d^2u}{dy^2} = C$

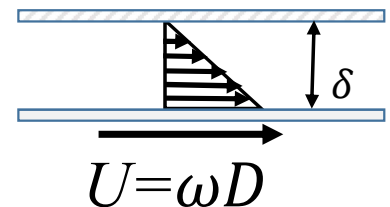
(a). 平板库埃特流动(couette flow, purely shear driven flow)

$\frac{dp}{dx} = 0$, $u(0) = 0$, $u(a) = U$ (边界条件)

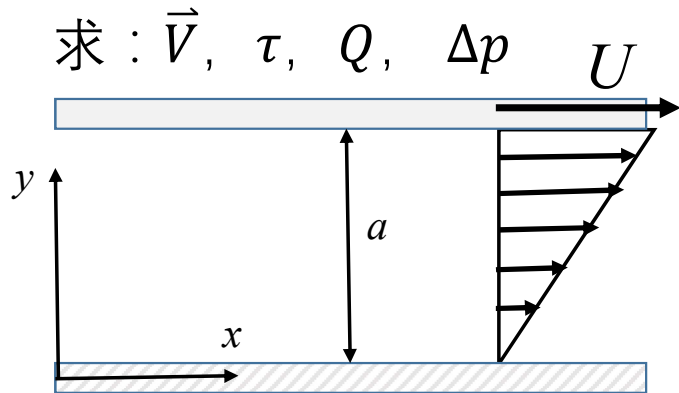
$$\frac{d^2u}{dy^2} = 0 \rightarrow \left. \begin{aligned} u(y) &= ky + B \\ u(0) &= 0 \\ u(a) &= U \end{aligned} \right\} u(y) = \frac{U}{a}y$$



$\delta \ll D$ 时可看作



6.2 无限大平板间充分发展层流 (5.3)



求: \bar{V} , τ , Q , Δp

解: $\frac{dp}{dx} = \mu \frac{d^2u}{dy^2} = C$

(a). 平板库埃特流动(couette flow, purely shear driven flow)

$$\frac{dp}{dx} = 0, \quad \Rightarrow \quad u(y) = \frac{U}{a}y$$

➤ Shear stress: $\tau_{yx} = \mu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) = \mu \frac{du}{dy} = \mu \frac{U}{a} = C$

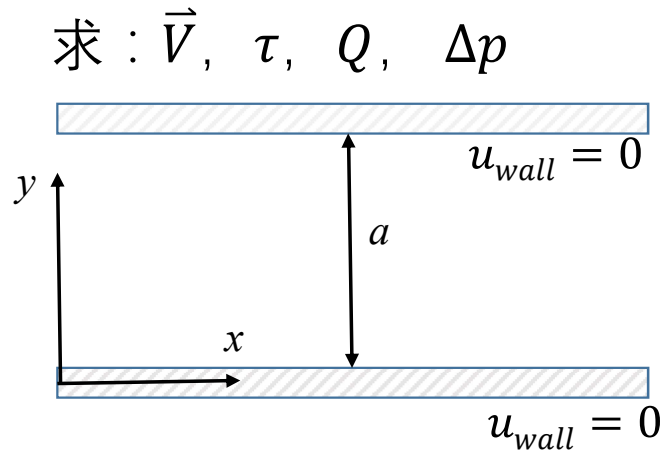
$$\tau_{wall} = \mu \frac{U}{a}$$

➤ Volume flow rate: $Q = \int u(y) dy dz = b \int \frac{U}{a} y dy$ (宽度为 b)

$$q = \frac{Q}{b} = \int_0^a \frac{U}{a} y dy = \frac{Ua}{2}$$

➤ $\bar{V} = \frac{q}{a} = \frac{U}{2}$

6.2 无限大平板间充分发展层流 (5.3)



解: $\frac{dp}{dx} = \mu \frac{d^2u}{dy^2} = C$

(b). 平板泊肃叶流动 $\frac{dp}{dx} \neq 0$ (purely pressure driven flow)

$$\frac{dp}{dx} = \mu \frac{d^2u}{dy^2} = C$$

$$\mu \frac{du}{dy} = \frac{dp}{dx} y + C_1$$

$$\mu u(y) = \frac{dp}{dx} \frac{y^2}{2} + C_1 y + C_2$$

$$u(y) = \frac{1}{2\mu} \frac{dp}{dx} y^2 + \frac{C_1}{\mu} y + \frac{C_2}{\mu}$$

$$u(y) = \frac{1}{2\mu} \frac{dp}{dx} y^2 - \frac{a}{2\mu} \frac{dp}{dx} y$$

边界条件: $u(0) = 0$
 $u(a) = 0$

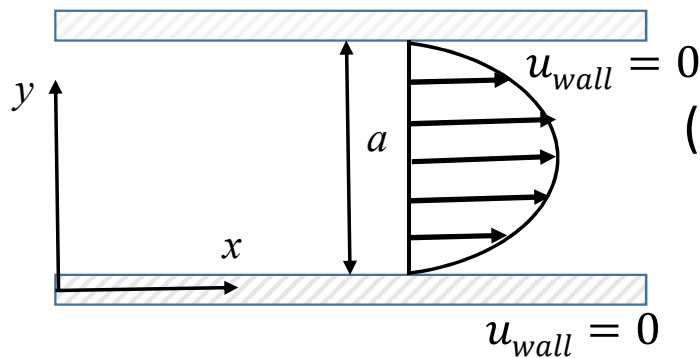
$$u(0) = \frac{C_2}{\mu} = 0 \Rightarrow C_2 = 0$$

$$u(a) = \frac{1}{2\mu} \frac{dp}{dx} a^2 + \frac{C_1}{\mu} a = 0 \Rightarrow C_1 = -\frac{a}{2} \frac{dp}{dx}$$

6.2 无限大平板间充发展层流 (5.3)

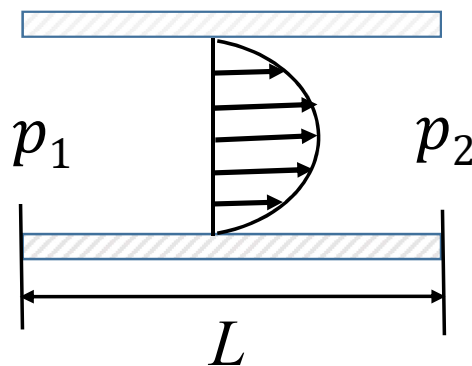
求: \vec{V} , τ , Q , Δp

解: $\frac{dp}{dx} = \mu \frac{d^2 u}{dy^2} = C$



(b). 平板泊肃叶流动 $\frac{dp}{dx} \neq 0$ (purely pressure driven flow)

$$\begin{aligned} u(y) &= \frac{1}{2\mu} \frac{dp}{dx} y^2 - \frac{a}{2\mu} \frac{dp}{dx} y \\ &= \frac{a^2}{2\mu} \frac{dp}{dx} \left[\left(\frac{y}{a}\right)^2 - \left(\frac{y}{a}\right) \right] \\ &= \frac{a^2}{2\mu} \underbrace{\left(-\frac{dp}{dx}\right)}_{>0} \underbrace{\frac{y}{a} \left(1 - \frac{y}{a}\right)}_{>0} > 0 \end{aligned}$$



$$p_1 > p_2$$

→ $u(y) > 0$, 则 $\frac{dp}{dx} < 0$ $\frac{dp}{dx} = C$ $p = kx + B$

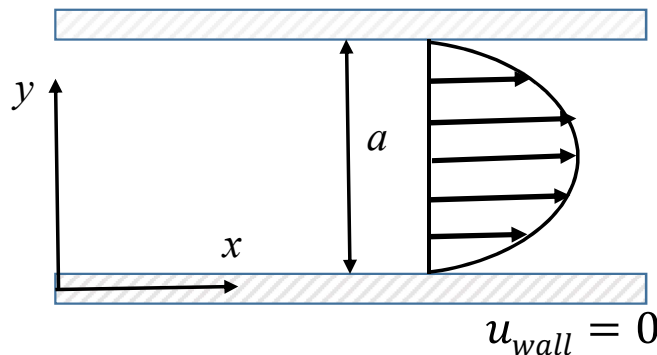
压强沿流动方向降低 → 顺压流动

$$-\frac{dp}{dx} = \frac{p_1 - p_2}{L} = \frac{\Delta p}{L} \quad u(y) = \frac{a^2}{2\mu} \left(\frac{\Delta p}{L}\right) \frac{y}{a} \left(1 - \frac{y}{a}\right)$$

6.2 无限大平板间充发展层流 (5.3)

求: \bar{V} , τ , Q , Δp

解: $\frac{dp}{dx} = \mu \frac{d^2 u}{dy^2} = C$



(b). 平板泊肃叶流动 $\frac{dp}{dx} \neq 0$ (purely pressure driven flow)

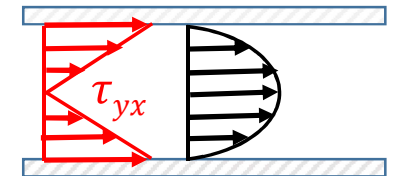
$$\text{➤ } u(y) = \frac{a^2}{2\mu} \left(-\frac{dp}{dx} \right) \frac{y}{a} \left(1 - \frac{y}{a} \right)$$

➤ Shear stress: $\tau_{yx} = \mu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) = \mu \frac{du}{dy} = a \frac{dp}{dx} \left(\frac{y}{a} - \frac{1}{2} \right)$

$$\tau_{wall} = \pm a \frac{dp}{dx}$$

➤ Volume flow rate: $q = \frac{Q}{b} = \int_0^a u(y) dy$

$$= \frac{a^3}{12\mu} \left(-\frac{dp}{dx} \right)$$

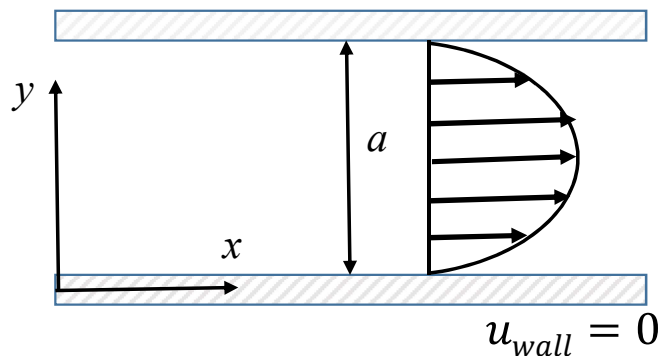


➤ $\bar{V} = \frac{q}{a} = \frac{a^2}{12\mu} \left(-\frac{dp}{dx} \right)$ ➤ @ $y = a/2$, $u_{max} = \frac{a^2}{8\mu} \left(-\frac{dp}{dx} \right) = 1.5 \bar{V}$

6.2 无限大平板间充发展层流 (5.3)

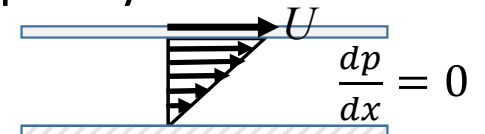
求: \vec{V} , τ , Q , Δp

解: $\frac{dp}{dx} = \mu \frac{d^2 u}{dy^2} = C$



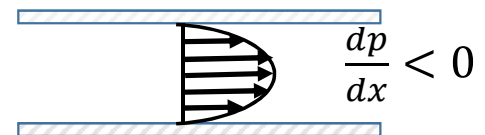
(a). 平板库埃特流动 (couette flow, purely shear driven flow)

➤ $u(y) = \frac{U}{a} y$



(b). 平板泊肃叶流动 $\frac{dp}{dx} \neq 0$ (purely pressure driven flow)

➤ $u(y) = \frac{a^2}{2\mu} \left(-\frac{dp}{dx}\right) \frac{y}{a} \left(1 - \frac{y}{a}\right)$



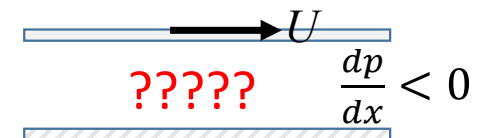
Note:

1. 仅适用层流.

$Re < 2000$ 层流 ; $Re > 7700$ 湍流

2. 入口效应, $L_e/D \approx 0.06 Re$

(c). 一般(a+b)???



作业：

复习笔记！

1. 用积分方程求 τ_w 。

5.2, 5.4, 5.14

看例5.3~5.7

