

空气与气体动力学

张科

回顾：准一维可压内流

$$d(\rho u A) = 0 \quad dp = -\rho u du \quad dh + u du = 0$$

1. 速度面积关系式: $\frac{dA}{A} = (Ma^2 - 1) \frac{du}{u}$

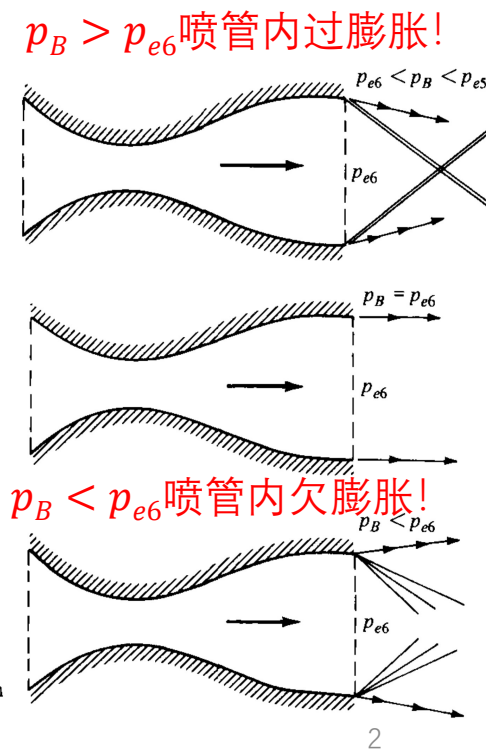
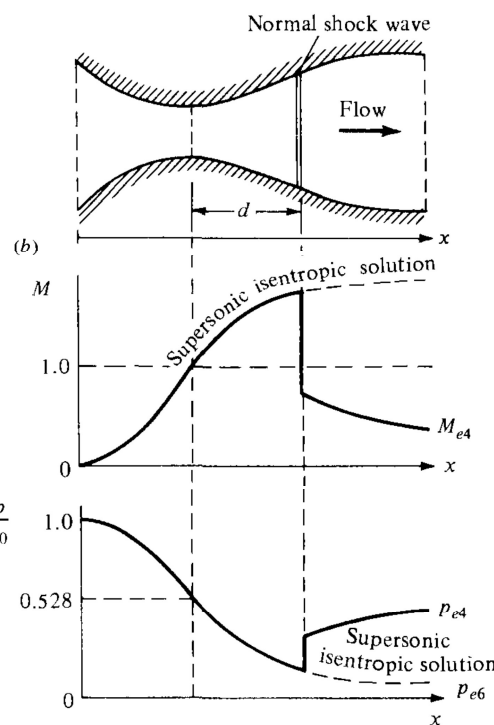
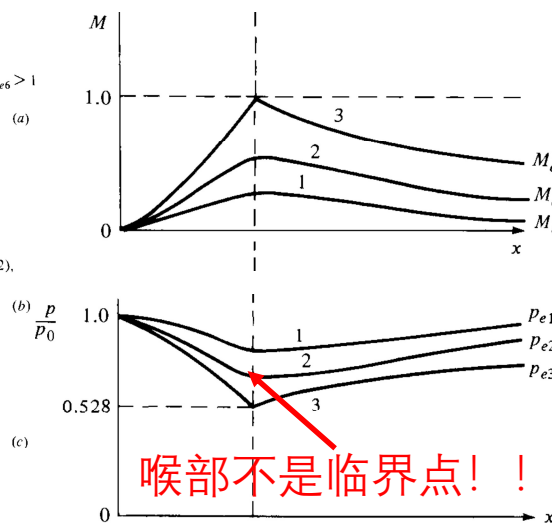
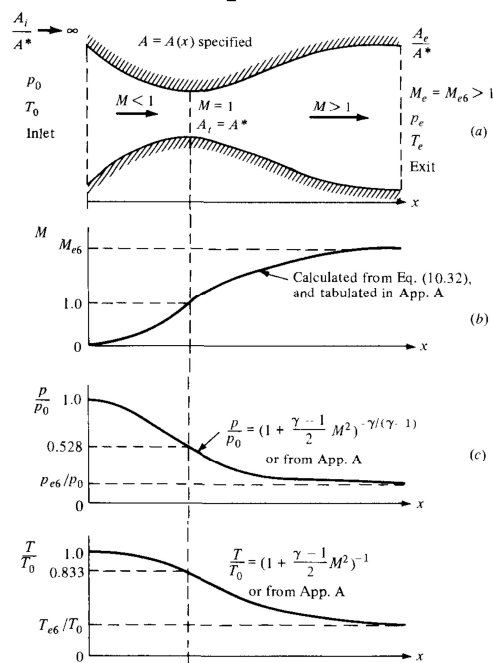
$0 < Ma < 1$ 时,

$dA > 0 \quad du < 0, A \uparrow \quad u \downarrow$

$Ma > 1$ 时,

$dA > 0 \quad du > 0, A \uparrow \quad u \uparrow$

2. 喷管(减压、增速):



12.3 扩压器(9.4)

扩压器(diffuser): **减速、增压**。

等熵: $s_1 = s_2, p_{02} = p_{01}$, 可用能不变。

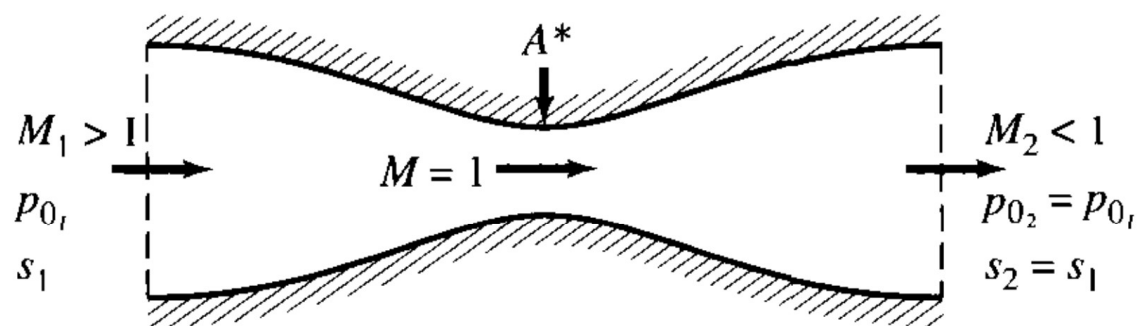
超声速气流等熵减速——极不可能!

壁面内折——斜激波;

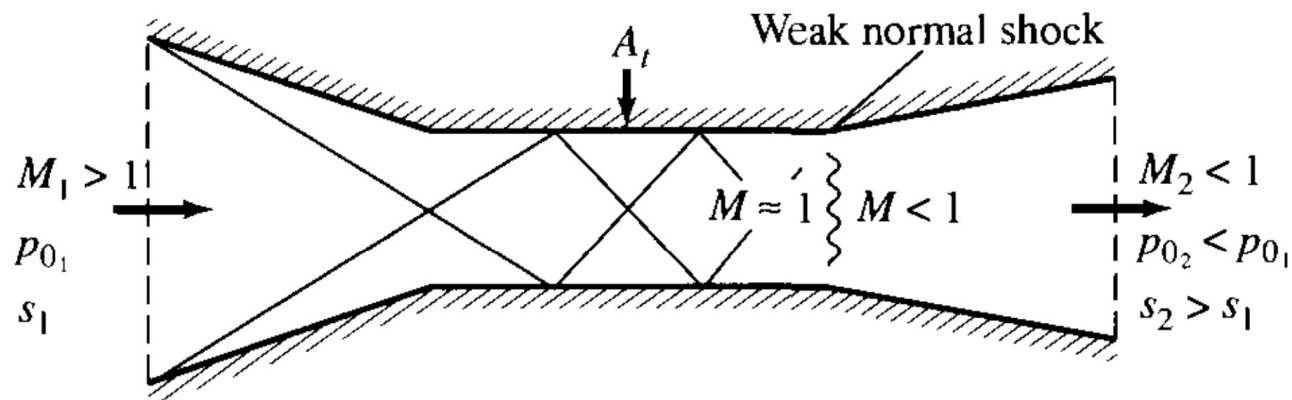
边界层——粘性损失。

$p_0 \downarrow$, 可用功减小!

扩压器设计: 提高 $\frac{p_{02}}{p_{01}}$ 。

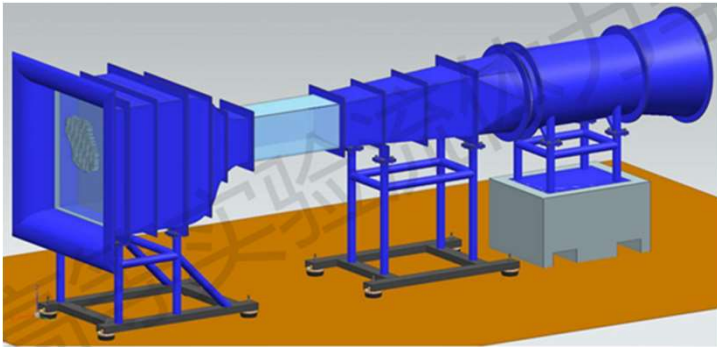
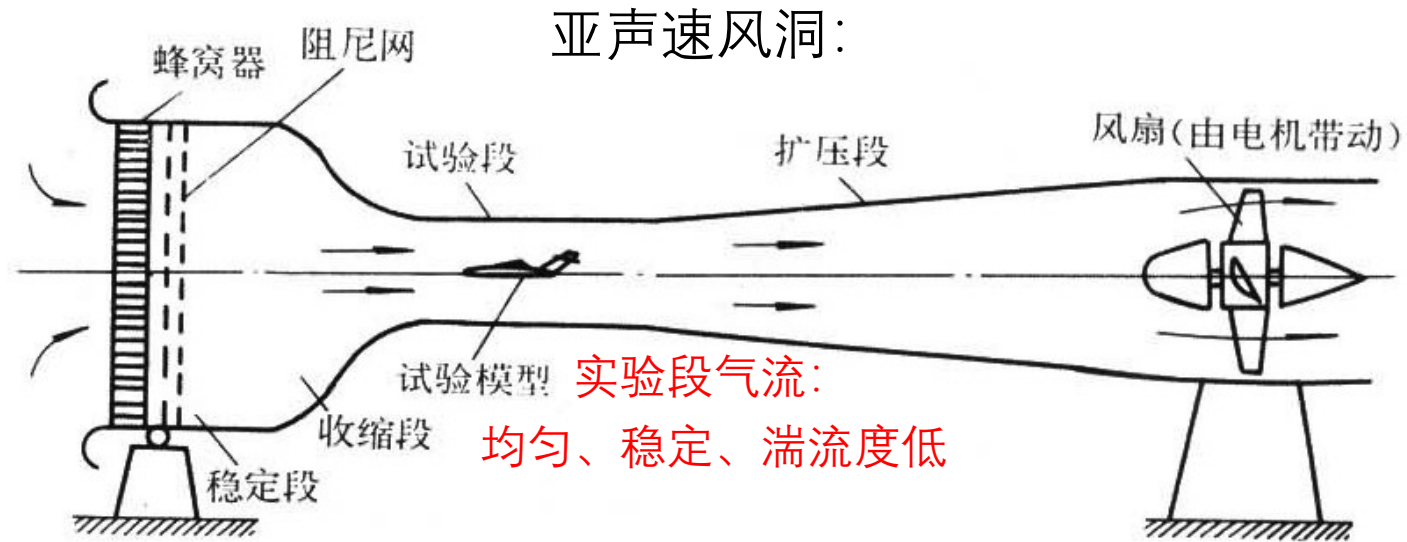


(a) Ideal (isentropic) supersonic diffuser

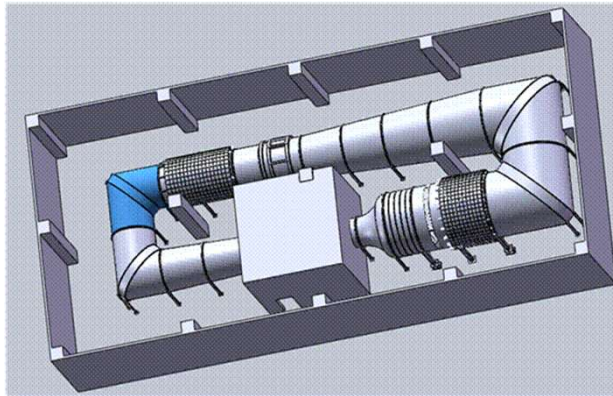


(b) Actual supersonic diffuser

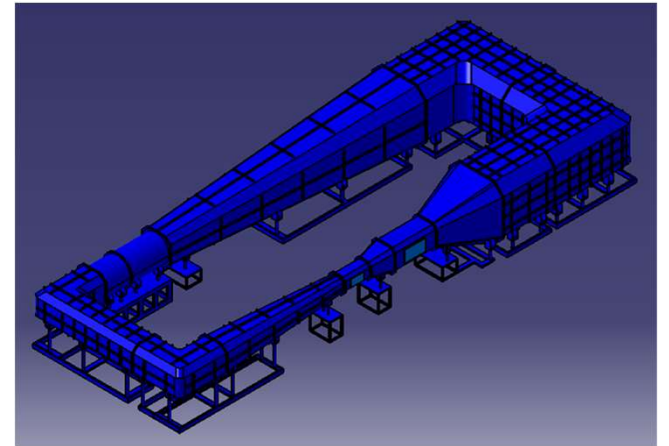
12.4超声速风洞(9.5)



直流式

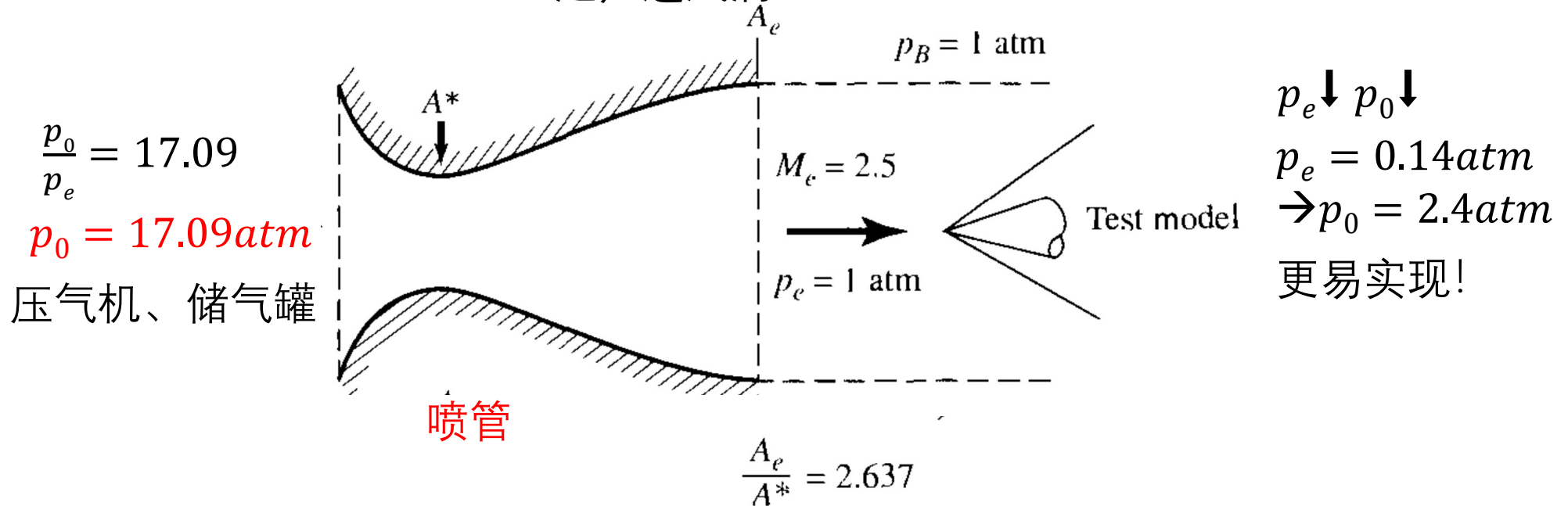


回流式

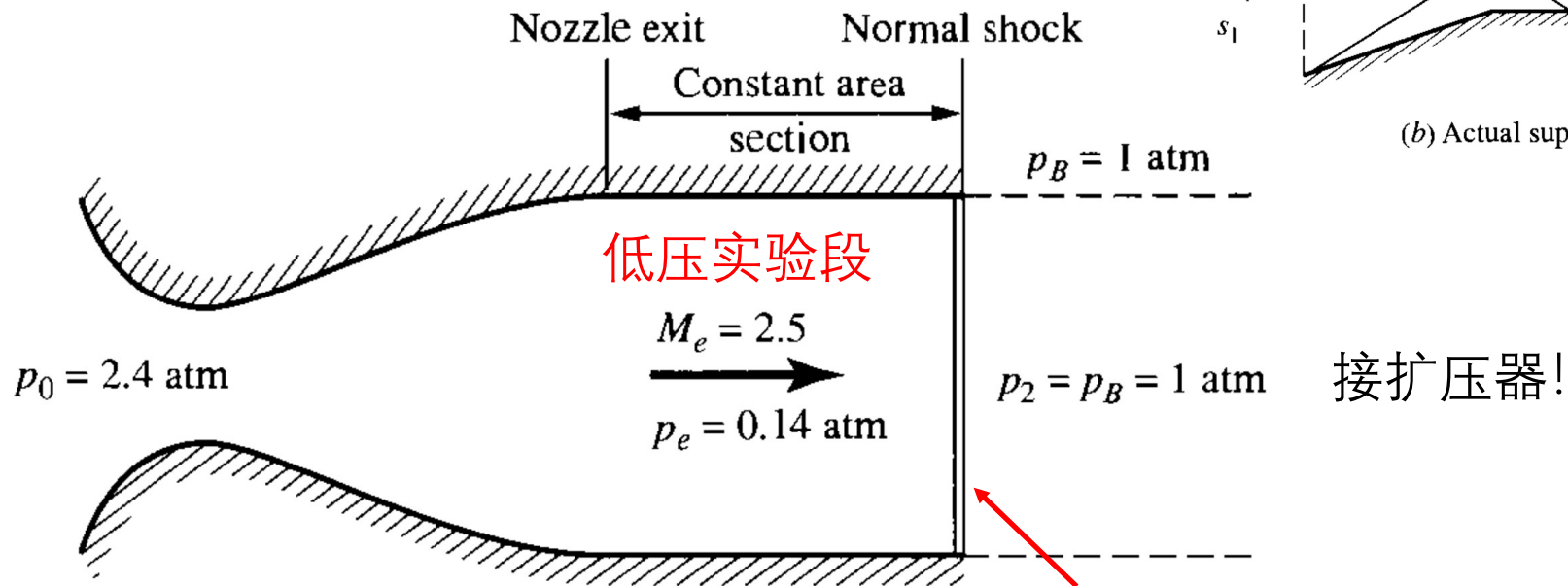


12.4超声速风洞(9.5)

超声速风洞： 低压实验段！

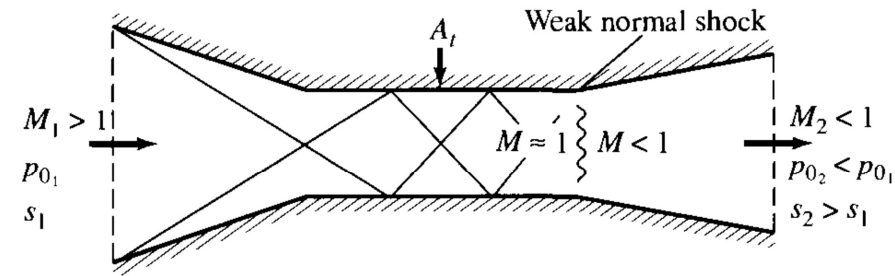


12.4 超声速风洞(9.5)



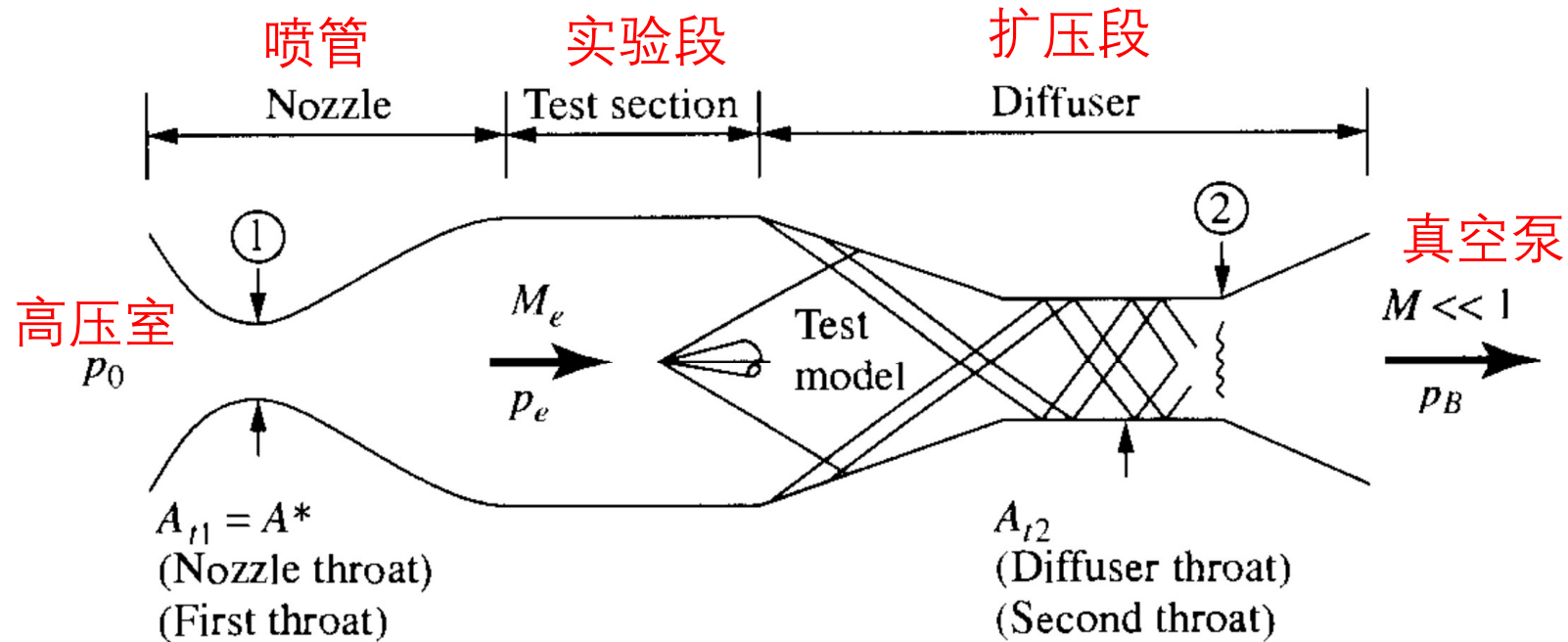
正激波，增压、减速（扩压器）！

正激波耗散大，不稳定。



(b) Actual supersonic diffuser

12.4超声速风洞(9.5)



损失：模型表面、扩压段——激波；
边界层粘性流动。

7.1 超声速风洞的储气罐内的温度为288K, 实验段速度为450m/s。假定风洞中气流为绝热流动, 计算实验段Ma。

解: $a = \sqrt{\gamma R T} = \sqrt{1.4 \times 287 \times 288} = 340.17 \text{ m/s}$
 $\therefore Ma = \frac{V}{a} = \frac{450}{340.17} = 1.32$

储气罐内为滞止参数!!
总温

解: $T_0 = 288 \text{ K}, V = 450 \text{ m/s}$

能量方程: $c_p T_0 = c_p T + \frac{V^2}{2}$

$$T = T_0 - \frac{V^2}{2c_p} = 188.2 \text{ K}$$

实验段 $a = \sqrt{\gamma R T} = 274.2 \text{ m/s}$

$$Ma = \frac{V}{a} = 1.64$$

7.2 - 给定空气流温度为 300 K, 压力为 1.2 atm, 速度为 250 m/s.
 计算该点处的总压、总温、临界压力、临界温度和特征马赫数
 解: $a = \sqrt{\gamma R T} = \sqrt{1.4 \times 287 \times 300} = 347.2 \text{ m/s}$.

$$Ma = \frac{V}{a} = \frac{250}{347.2} = 0.72$$

$$\frac{T_0}{T} = 1 + \frac{\gamma-1}{2} Ma^2 \Rightarrow T_0 = (1 + \frac{\gamma-1}{2} Ma^2) \cdot T = (1 + \frac{0.4}{2} \cdot 0.72^2) 300 = 268.9 \text{ K}$$

$$\frac{P_0}{P} = (1 + \frac{\gamma-1}{2} Ma^2)^{\frac{\gamma}{\gamma-1}} \Rightarrow P_0 = (1 + \frac{\gamma-1}{2} Ma^2)^{\frac{\gamma}{\gamma-1}} P = (1 + \frac{0.4}{2} \cdot 0.72^2)^{\frac{1.4}{0.4}} 1.2 = 0.818 \text{ atm}$$

$$\frac{T^*}{T_0} = 0.833 \Rightarrow T^* = 0.833 \times 268.9 = 223.99 \text{ K}$$

$$\frac{P^*}{P_0} = 0.528 \Rightarrow P^* = 0.528 \times 0.818 = 0.432 \text{ atm}$$

$$\cancel{Ma^* = \frac{V}{a^*} = \frac{250}{\sqrt{1.4 \times 287 \times 223.99}} = 0.833} \quad a^* = \sqrt{\gamma R T^*} = \frac{V}{Ma^*} = \frac{250}{0.833} = 300$$

滞止压力、温度求错!

7.2. 解: $a = \sqrt{\gamma R T} = \sqrt{1.4 \times 287 \times 300} = 347.2 \text{ (m/s)}$

$$Ma = \frac{V}{a} = 0.72$$

$$Ma = 0.72 \text{ 附表 A, } \frac{P_0}{P} = 1.412 \quad \frac{T_0}{T} = 1.104$$

$$P_0 = 1.694 \text{ atm} \quad T_0 = 331.2 \text{ K}$$

$$\frac{T^*}{T_0} = 0.833 \quad T^* = 275.9 \text{ K} \quad a^* = \sqrt{\gamma R T^*} = 332.95 \text{ m/s}$$

$$\frac{P^*}{P_0} = 0.528 \quad P^* = 0.894 \text{ atm}$$

$$Ma^* = \frac{V}{a^*} = \frac{250}{332.95} = 0.75$$

7.6 如果通过激波的熵增为 $199.5 \text{ J/kg}\cdot\text{K}$, 问来流马赫数为多大?

$$\Delta S = S_2 - S_1 = C_p \ln \left\{ \left[\frac{2 + (\gamma - 1) Ma_1^2}{(\gamma + 1) Ma_1^2} \right]^\gamma \left[1 + \frac{2\gamma}{\gamma + 1} (Ma_1^2 - 1) \right] \right\}$$

$$Ma = 1.09$$

计算过程??

7.6 解: $S_2 - S_1 = -R \ln \frac{P_{02}}{P_{01}} = 199.5 \text{ J/kg}\cdot\text{K}$

$$\ln \left(\frac{P_{02}}{P_{01}} \right) = \left(\frac{-199.5}{287} \right) = -0.6951$$

$$\frac{P_{02}}{P_{01}} = 0.499 \xRightarrow{\text{附表B}} Ma_1 = 2.5$$

7.8. 根据不可压缩流体假设:

$$\frac{P_0}{\rho_0} + \frac{V_0^2}{2} = \frac{P}{\rho} + \frac{V^2}{2} \Rightarrow \frac{1.555 \text{ atm}}{\rho_0} = \frac{1 \text{ atm}}{\rho} + \frac{V^2}{2}$$

$$\frac{\rho_0}{\rho} = 1.371 \quad \rho = \frac{P}{RT} = \frac{1 \times 1.01325 \times 10^5}{287 \times 288} = 1.226 \text{ kg/m}^3$$

$$\therefore \rho_0 = 1.371 \rho = 1.681 \text{ kg/m}^3$$

$$\therefore V = 148.884 \text{ m/s}$$

\therefore 误差为 46.62%

不可压??

7.7. $T_\infty = 288 \text{ K}$, $P_\infty = 1 \text{ atm}$

若 $Ma_\infty = 1$, 则 $P_0/P_\infty = 1.893$

$P_{\text{pitot}} = 1.555 \text{ atm} < 1.893 \text{ atm} \Rightarrow$ 亚声速流

$P_{\text{pitot}} = P_0$, $\frac{P_0}{P_\infty} = 1.555 \Rightarrow Ma_\infty = 0.82$

$a_\infty = \sqrt{\gamma R T} = \sqrt{1.4 \times 286.6 \times 288} = 340 \text{ m/s}$

$V_\infty = Ma_\infty \cdot a_\infty = 278.8 \text{ m/s}$

7.8. 解. 不可压. $P_0 = 1.555 \text{ atm}$, $P_\infty = 1 \text{ atm}$

$$V_\infty' = \sqrt{\frac{2 \Delta P}{\rho}} = \sqrt{\frac{2 \times 0.555 \times 1.013 \times 10^5}{1.22}}$$

$$= 303.6 \text{ (m/s)}$$

$$\frac{|V_\infty' - V_\infty|}{V_\infty} = 8.9\%$$

十三. 绕翼型可压缩流动 (空10)



b: 79.8m S: 845m² $\chi=33.5^\circ$

Ma_{巡航}=0.85 巡航时气动特性如何? ?

$0.3 < Ma < 1$:

C_l, C_d ??

$Ma > 1$ 激波产生??

$Ma < 0.3$:

不可压势流

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = 0$$

可压势流 ϕ ? ?

13.1速度势方程(10.2)

无粘可压无旋流——存在势函数 ϕ 。 $\vec{\Omega} = \vec{\nabla} \times \vec{V} = 0$ 若: $\vec{\nabla} \times \vec{V} = 0$, 则 $\vec{V} = \vec{\nabla} \phi$

定常、无旋、等熵、2D: $\vec{V} = \vec{\nabla} \phi$, $u = \frac{\partial \phi}{\partial x}$, $v = \frac{\partial \phi}{\partial y}$

连续性方程: $\frac{\partial(\rho u)}{\partial x} + \frac{\partial(\rho v)}{\partial y} = 0$

不可压: $\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$

$$\rho \frac{\partial u}{\partial x} + u \frac{\partial \rho}{\partial x} + \rho \frac{\partial v}{\partial y} + v \frac{\partial \rho}{\partial y} = 0$$

$$\rho \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial \phi}{\partial x} \frac{\partial \rho}{\partial x} + \rho \frac{\partial^2 \phi}{\partial y^2} + \frac{\partial \phi}{\partial y} \frac{\partial \rho}{\partial y} = 0 \quad (1)$$

欧拉方程: $dp = -\rho V dV = -\frac{\rho}{2} dV^2 = -\frac{\rho}{2} d(u^2 + v^2)$

等熵: $\frac{dp}{d\rho} = a^2 \rightarrow dp = a^2 d\rho$

$$d\rho = -\frac{\rho}{2a^2} d(u^2 + v^2) = -\frac{\rho}{2a^2} d\left[\left(\frac{\partial \phi}{\partial x}\right)^2 + \left(\frac{\partial \phi}{\partial y}\right)^2\right] \quad (2)$$

13.1速度势方程(10.2)

$$d\rho = -\frac{\rho}{2a^2} d[(\frac{\partial\phi}{\partial x})^2 + (\frac{\partial\phi}{\partial y})^2] \quad (2)$$

$$\frac{\partial(2)}{\partial x}: \quad \frac{\partial\rho}{\partial x} = -\frac{\rho}{2a^2} \frac{\partial}{\partial x} [(\frac{\partial\phi}{\partial x})^2 + (\frac{\partial\phi}{\partial y})^2] = -\frac{\rho}{a^2} (\frac{\partial\phi}{\partial x} \frac{\partial^2\phi}{\partial x^2} + \frac{\partial\phi}{\partial y} \frac{\partial^2\phi}{\partial y\partial x}) \quad (3)$$

$$\frac{\partial(2)}{\partial y}: \quad \frac{\partial\rho}{\partial y} = -\frac{\rho}{2a^2} \frac{\partial}{\partial y} [(\frac{\partial\phi}{\partial x})^2 + (\frac{\partial\phi}{\partial y})^2] = -\frac{\rho}{a^2} (\frac{\partial\phi}{\partial x} \frac{\partial^2\phi}{\partial x\partial y} + \frac{\partial\phi}{\partial y} \frac{\partial^2\phi}{\partial y^2}) \quad (4)$$

$$\frac{\partial(\rho u)}{\partial x} + \frac{\partial(\rho v)}{\partial y} = 0 \quad \rightarrow \quad \rho \frac{\partial^2\phi}{\partial x^2} + \frac{\partial\phi}{\partial x} \frac{\partial\rho}{\partial x} + \rho \frac{\partial^2\phi}{\partial y^2} + \frac{\partial\phi}{\partial y} \frac{\partial\rho}{\partial y} = 0 \quad (1)$$

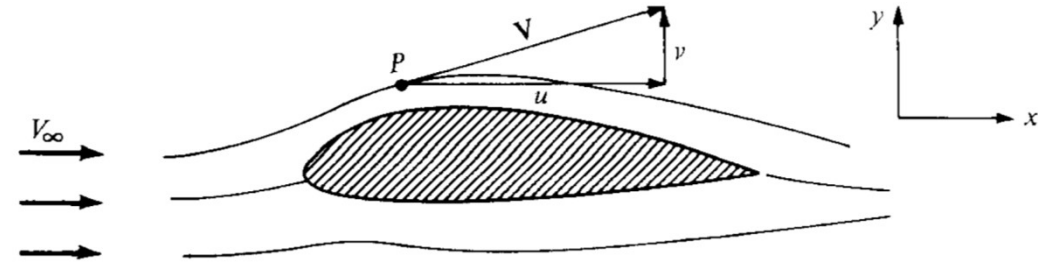
③+④→①:

$$\left\{ \begin{aligned} & \left[1 - \frac{1}{a^2} \left(\frac{\partial\phi}{\partial x} \right)^2 \right] \frac{\partial^2\phi}{\partial x^2} + \left[1 - \frac{1}{a^2} \left(\frac{\partial\phi}{\partial y} \right)^2 \right] \frac{\partial^2\phi}{\partial y^2} - \frac{2}{a^2} \left(\frac{\partial\phi}{\partial x} \right) \left(\frac{\partial\phi}{\partial y} \right) \frac{\partial^2\phi}{\partial x\partial y} = 0 \\ & a^2 = a_0^2 - \frac{\gamma-1}{2} V^2 = a_0^2 - \frac{\gamma-1}{2} [(\frac{\partial\phi}{\partial x})^2 + (\frac{\partial\phi}{\partial y})^2] \end{aligned} \right\} \quad (5)$$

关于 ϕ 的非线性偏微分方程!! ($Ma > 0.3$)

$$Ma < 0.3 : \frac{\partial^2\phi}{\partial x^2} + \frac{\partial^2\phi}{\partial y^2} = 0$$

13.2线化速度势方程(10.3)



t 、 f 、 α 较小时，翼型周围=来流+小扰动

$$u = V_\infty + \hat{u}, v = \hat{v} \quad \hat{u}, \hat{v} \text{ 扰动速度}$$

$$\phi = V_\infty x + \hat{\phi}, \quad \hat{\phi} \text{ 扰动速度势} \quad \hat{u} = \frac{\partial \hat{\phi}}{\partial x}, \hat{v} = \frac{\partial \hat{\phi}}{\partial y}$$

$$\begin{aligned} \frac{\partial \phi}{\partial x} &= V_\infty + \frac{\partial \hat{\phi}}{\partial x} \\ \frac{\partial \phi}{\partial y} &= \frac{\partial \hat{\phi}}{\partial y} \end{aligned} \quad \longrightarrow \quad \begin{aligned} \frac{\partial^2 \phi}{\partial x^2} &= \frac{\partial^2 \hat{\phi}}{\partial x^2} \\ \frac{\partial^2 \phi}{\partial y^2} &= \frac{\partial^2 \hat{\phi}}{\partial y^2}, \quad \frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \hat{\phi}}{\partial x \partial y} \end{aligned}$$

代入速度势方程⑤ \rightarrow

$$\left[a^2 - \left(V_\infty + \frac{\partial \hat{\phi}}{\partial x} \right)^2 \right] \frac{\partial^2 \hat{\phi}}{\partial x^2} + \left[a^2 - \left(\frac{\partial \hat{\phi}}{\partial y} \right)^2 \right] \frac{\partial^2 \hat{\phi}}{\partial y^2} - 2 \left(V_\infty + \frac{\partial \hat{\phi}}{\partial x} \right) \left(\frac{\partial \hat{\phi}}{\partial y} \right) \frac{\partial^2 \hat{\phi}}{\partial x \partial y} = 0$$

$$\left[a^2 - (V_\infty + \hat{u})^2 \right] \frac{\partial \hat{u}}{\partial x} + (a^2 - \hat{v}^2) \frac{\partial \hat{v}}{\partial y} - 2(V_\infty + \hat{u})\hat{v} \frac{\partial \hat{u}}{\partial y} = 0$$

$$\text{能量: } \frac{a_\infty^2}{\gamma-1} + \frac{V_\infty^2}{2} = \frac{a^2}{\gamma-1} + \frac{(V_\infty + \hat{u})^2 + \hat{v}^2}{2}$$



13.2线化速度势方程(10.3)

$$\begin{aligned}
 (1 - M_\infty^2) \frac{\partial \hat{u}}{\partial x} + \frac{\partial \hat{v}}{\partial y} = & \boxed{M_\infty^2 \left[(\gamma + 1) \frac{\hat{u}}{V_\infty} + \frac{\gamma + 1}{2} \frac{\hat{u}^2}{V_\infty^2} + \frac{\gamma - 1}{2} \frac{\hat{v}^2}{V_\infty^2} \right] \frac{\partial \hat{u}}{\partial x}} \quad \text{A} \\
 & + \boxed{M_\infty^2 \left[(\gamma - 1) \frac{\hat{u}}{V_\infty} + \frac{\gamma + 1}{2} \frac{\hat{v}^2}{V_\infty^2} + \frac{\gamma - 1}{2} \frac{\hat{u}^2}{V_\infty^2} \right] \frac{\partial \hat{v}}{\partial y}} \quad \text{B} \\
 & + \boxed{M_\infty^2 \left[\frac{\hat{v}}{V_\infty} \left(1 + \frac{\hat{u}}{V_\infty} \right) \left(\frac{\partial \hat{u}}{\partial y} + \frac{\partial \hat{v}}{\partial x} \right) \right]} \quad \text{C 非线性偏微分方程!!}
 \end{aligned} \quad \text{⑥}$$

小扰动(细长体、小攻角): $\frac{\hat{u}}{V_\infty}, \frac{\hat{v}}{V_\infty} \ll 1 \quad \frac{\hat{u}^2}{V_\infty^2}, \frac{\hat{v}^2}{V_\infty^2} \ll 1$

1. $Ma_\infty \leq 0.8$ 或 $Ma_\infty > 1.2$ 时: $A \ll (1 - Ma_\infty^2) \frac{\partial \hat{u}}{\partial x}$

2. $Ma_\infty < 5$: $B \ll \frac{\partial \hat{v}}{\partial y} \quad C \approx 0$

→ $(1 - Ma_\infty^2) \frac{\partial \hat{u}}{\partial x} + \frac{\partial \hat{v}}{\partial y} = 0$ 线性偏微分方程!!

$$\boxed{(1 - Ma_\infty^2) \frac{\partial^2 \hat{\phi}}{\partial x^2} + \frac{\partial^2 \hat{\phi}}{\partial y^2} = 0}$$

13.2线化速度势方程(10.3)

$$(1-Ma_{\infty}^2)\frac{\partial^2\hat{\phi}}{\partial x^2} + \frac{\partial^2\hat{\phi}}{\partial y^2} = 0$$

适用条件：①小扰动(薄体、小攻角)

②亚、超声速： $Ma_{\infty} \leq 0.8$, $Ma_{\infty} > 1.2$, $Ma_{\infty} < 5$

不适用：①厚体、跨声速($0.8 < Ma_{\infty} < 1.2$)

②高超声速($Ma_{\infty} > 5$)

13.2线化速度势方程(10.3)

$$C_p \sim \hat{u}, \hat{v}, \hat{\phi} \quad ??$$

➤ 压力系数线化: $C_p = \frac{p-p_\infty}{\frac{1}{2}\rho_\infty V_\infty^2}$ $\frac{1}{2}\rho_\infty V_\infty^2 = \frac{1}{2}\frac{\gamma p_\infty}{\gamma p_\infty}\rho_\infty V_\infty^2$

$$= \frac{\gamma}{2}p_\infty \frac{\rho_\infty}{\gamma p_\infty} V_\infty^2 \quad a_\infty^2 = \frac{\gamma p_\infty}{\rho_\infty}$$

$$= \frac{\gamma}{2}p_\infty Ma_\infty^2$$

➡ $C_p = (p - p_\infty) \frac{2}{\gamma p_\infty Ma_\infty^2} = \left(\frac{p}{p_\infty} - 1\right) \frac{2}{\gamma Ma_\infty^2} \quad (1)$

能量: $T + \frac{V^2}{2C_p} = T_\infty + \frac{V_\infty^2}{2C_p}$

$$T - T_\infty = \frac{V_\infty^2 - V^2}{\frac{2\gamma}{(\gamma-1)}R}$$

$$\frac{T}{T_\infty} - 1 = \frac{\gamma-1}{2} \frac{V_\infty^2 - V^2}{\gamma R T_\infty} = \frac{\gamma-1}{2} \frac{V_\infty^2 - V^2}{a_\infty^2} = -\frac{\gamma-1}{2} \frac{(\hat{u}^2 + \hat{v}^2 + 2\hat{u}V_\infty)}{a_\infty^2}$$

$$V^2 = (V_\infty + \hat{u})^2 + \hat{v}^2$$

13.2线化速度势方程(10.3) $C_p = (\frac{p}{p_\infty} - 1) \frac{2}{\gamma Ma_\infty^2}$ ①

$$\frac{T}{T_\infty} = 1 - \frac{\gamma-1}{2} \frac{(\hat{u}^2 + \hat{v}^2 + 2\hat{u}V_\infty)}{a_\infty^2}$$

等熵: $\frac{p}{p_\infty} = \left(\frac{T}{T_\infty}\right)^{\frac{\gamma}{\gamma-1}} = \left[1 - \frac{\gamma-1}{2} \frac{(\hat{u}^2 + \hat{v}^2 + 2\hat{u}V_\infty)}{a_\infty^2}\right]^{\frac{\gamma}{\gamma-1}}$

$$= \left[1 - \frac{\gamma-1}{2} Ma_\infty^2 \left(\frac{2\hat{u}}{V_\infty} + \frac{\hat{u}^2 + \hat{v}^2}{V_\infty^2}\right)\right]^{\frac{\gamma}{\gamma-1}}$$

$$\varepsilon = \frac{\gamma-1}{2} Ma_\infty^2 \left(\frac{2\hat{u}}{V_\infty} + \frac{\hat{u}^2 + \hat{v}^2}{V_\infty^2}\right)$$

小扰动: $\frac{\hat{u}}{V_\infty} \ll 1, \quad \frac{\hat{u}^2}{V_\infty^2}, \frac{\hat{v}^2}{V_\infty^2} \ll 1$

$$\frac{p}{p_\infty} = (1 - \varepsilon)^{\frac{\gamma}{\gamma-1}} \approx 1 - \frac{\gamma}{\gamma-1} \varepsilon + \dots = 1 - \frac{\gamma}{2} Ma_\infty^2 \left(\frac{2\hat{u}}{V_\infty} + \frac{\hat{u}^2 + \hat{v}^2}{V_\infty^2}\right) \quad \text{②}$$

①+② $\rightarrow C_p = -\frac{2\hat{u}}{V_\infty} - \frac{\hat{u}^2 + \hat{v}^2}{V_\infty^2} \approx -\frac{2\hat{u}}{V_\infty}$ 线化压力系数!

13.2线化速度势方程(10.3)

$$\begin{cases} (1-Ma_\infty^2)\frac{\partial^2\hat{\phi}}{\partial x^2} + \frac{\partial^2\hat{\phi}}{\partial y^2} = 0 & \rightarrow \hat{\phi}, \hat{u} \rightarrow C_p \\ C_p \approx -\frac{2\hat{u}}{V_\infty} \end{cases}$$

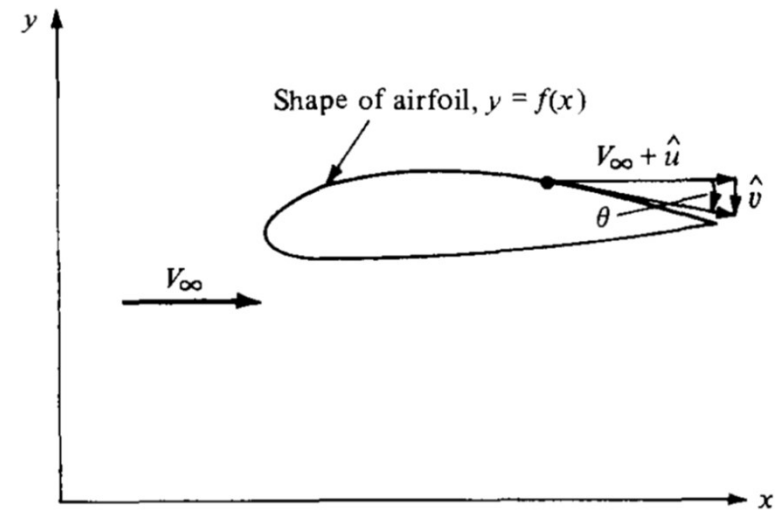
边界条件: $@\infty \hat{\phi} = \text{constant}, \hat{u} = \hat{v} = 0$

$@\text{物面 } v_\perp = 0$

$$\tan\theta = \frac{v}{u} = \frac{\hat{v}}{V_\infty + \hat{u}}$$

$$\hat{u} \ll V_\infty \rightarrow \hat{v} = V_\infty \tan\theta = V_\infty \left(\frac{dy}{dx}\right)_s$$

$$\frac{\partial\hat{\phi}}{\partial y} = V_\infty \tan\theta \quad \text{物面为流线近似方程}$$



13.3 Prandtl-Glauert 压缩性修正(10.4)

1904-1940, 低 Ma 飞行, 不可压空气动力学, 低速翼型理论、实验;
二战期间, $Ma \uparrow$, 高速翼型理论, 可压流动, 压缩性修正。

1. Prandtl-Glauert 方法

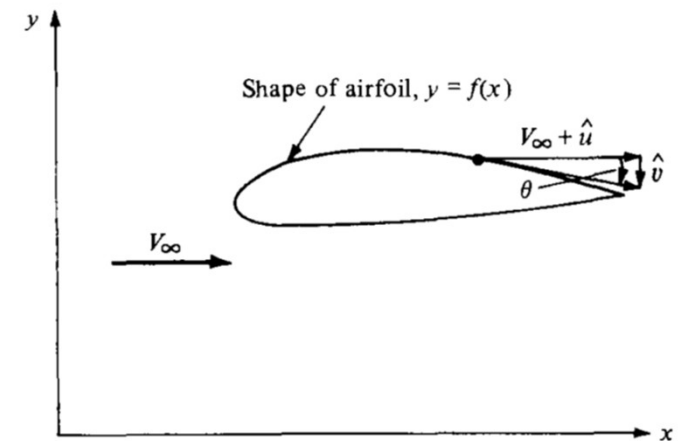
(基于线化扰动速度势方程, 薄翼、小攻角、 $Ma < 0.7$)

亚声速、可压、无粘流:

$$(1 - Ma_\infty^2) \frac{\partial^2 \hat{\phi}}{\partial x^2} + \frac{\partial^2 \hat{\phi}}{\partial y^2} = 0 \quad Ma < 1 \text{ 椭圆型方程}$$

$$\boxed{\beta^2 \frac{\partial^2 \hat{\phi}}{\partial x^2} + \frac{\partial^2 \hat{\phi}}{\partial y^2} = 0} \quad \beta^2 = 1 - Ma_\infty^2$$

$$Ma < 1: \boxed{\frac{\partial^2 \hat{\phi}}{\partial x^2} + \frac{\partial^2 \hat{\phi}}{\partial y^2} = 0} \quad \text{关联? ?}$$



13.3 Prandtl-Glauert 压缩性修正(10.4)

$$\beta^2 \frac{\partial^2 \hat{\phi}}{\partial x^2} + \frac{\partial^2 \hat{\phi}}{\partial y^2} = 0$$

仿射变换: $\xi = x, \eta = \beta y$ 引入: $\bar{\phi}(\xi, \eta) = \beta \hat{\phi}(x, y)$

$$\begin{aligned} \frac{\partial \xi}{\partial x} &= 1, \frac{\partial \xi}{\partial y} = 0 \\ \frac{\partial \eta}{\partial x} &= 0, \frac{\partial \eta}{\partial y} = \beta \end{aligned}$$

$$\begin{aligned} \bar{\phi}(\xi, \eta) = \beta \hat{\phi}(x, y) &\longrightarrow \frac{\partial \hat{\phi}}{\partial x} = \frac{1}{\beta} \frac{\partial \bar{\phi}}{\partial \xi} \frac{\partial \xi}{\partial x} = \frac{1}{\beta} \frac{\partial \bar{\phi}}{\partial \xi} \longrightarrow \frac{\partial^2 \hat{\phi}}{\partial x^2} = \frac{1}{\beta} \frac{\partial^2 \bar{\phi}}{\partial \xi^2} \frac{\partial \xi}{\partial x} = \frac{1}{\beta} \frac{\partial^2 \bar{\phi}}{\partial \xi^2} \\ \frac{\partial \hat{\phi}}{\partial y} &= \frac{1}{\beta} \frac{\partial \bar{\phi}}{\partial \eta} \frac{\partial \eta}{\partial y} = \frac{\partial \bar{\phi}}{\partial \eta} \longrightarrow \frac{\partial^2 \hat{\phi}}{\partial y^2} = \frac{\partial^2 \bar{\phi}}{\partial \eta^2} \frac{\partial \eta}{\partial y} = \beta \frac{\partial^2 \bar{\phi}}{\partial \eta^2} \end{aligned}$$

$$\beta^2 \frac{\partial^2 \hat{\phi}}{\partial x^2} + \frac{\partial^2 \hat{\phi}}{\partial y^2} = 0 \longrightarrow \beta^2 \frac{1}{\beta} \frac{\partial^2 \bar{\phi}}{\partial \xi^2} + \beta \frac{\partial^2 \bar{\phi}}{\partial \eta^2} = 0$$

$$\frac{\partial^2 \bar{\phi}}{\partial \xi^2} + \frac{\partial^2 \bar{\phi}}{\partial \eta^2} = 0 \quad \nabla^2 \bar{\phi} = 0 \quad \text{不可压势方程! !}$$