

空气与气体动力学

张科

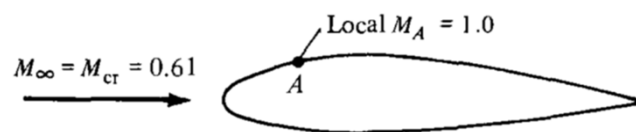
回顾：

1. Prandtl-Glauert压缩性修正: $C_p = \frac{C_{p,0}}{\sqrt{1-Ma_\infty^2}}$ $C_l = \frac{C_{l,0}}{\sqrt{1-Ma_\infty^2}}$ $C_m = \frac{C_{m,0}}{\sqrt{1-Ma_\infty^2}}$

相同翼型：(x,y)空间可压流动 → (ξ,η)空间不可压流动！

$$\beta^2 \frac{\partial^2 \hat{\phi}}{\partial x^2} + \frac{\partial^2 \hat{\phi}}{\partial y^2} = 0 \quad \frac{\partial^2 \bar{\phi}}{\partial \xi^2} + \frac{\partial^2 \bar{\phi}}{\partial \eta^2} = 0$$

2. 临界马赫数：



3. 阻力发散，声障

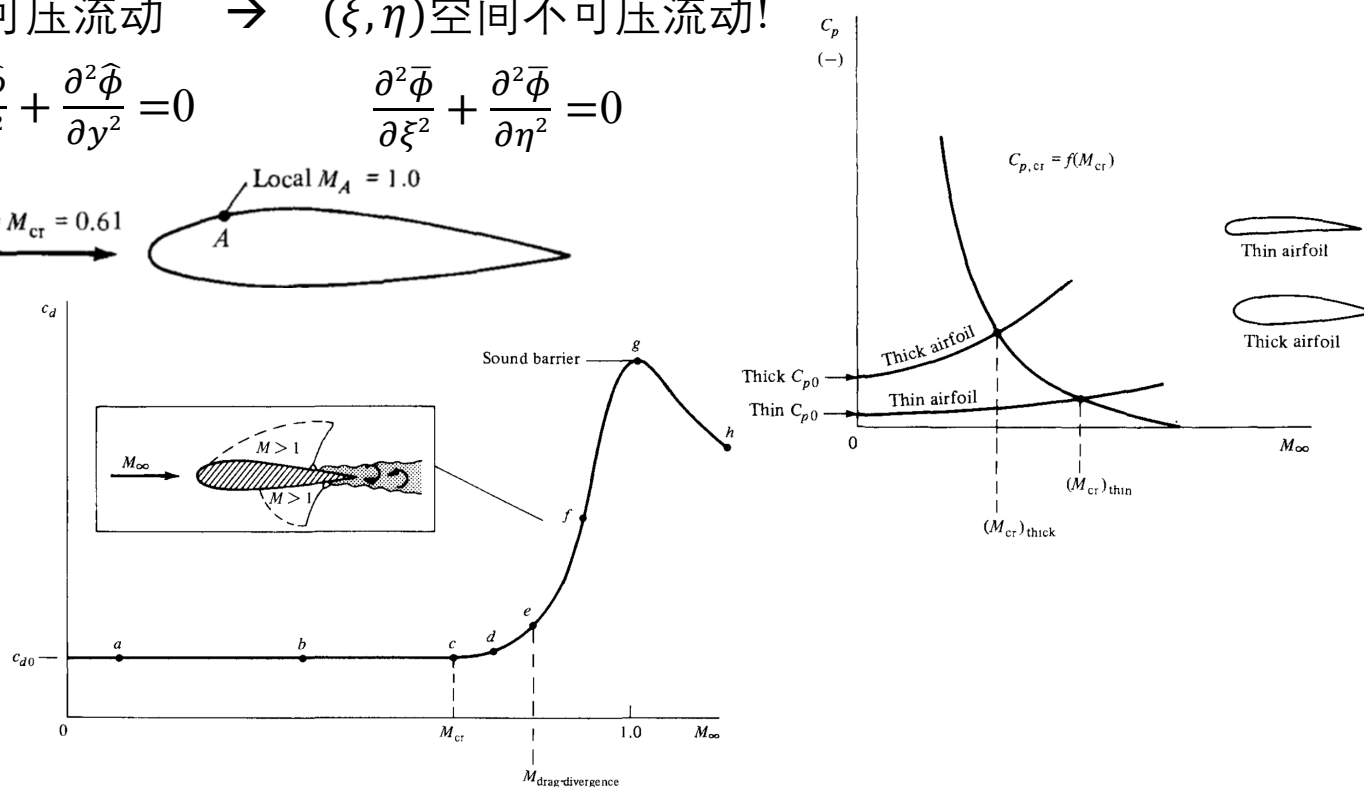
减 C_d , C_{dmax} ：

薄翼型；后掠翼。

超临界翼型；面积律。

4. 超声速翼型特点

尖头尖尾、对称、薄；激波，膨胀波。



13.7超声速线化理论(10.9)

$$(1-Ma_\infty^2)\frac{\partial^2\hat{\phi}}{\partial x^2} + \frac{\partial^2\hat{\phi}}{\partial y^2} = 0 \quad Ma_\infty < 1 \text{ 椭圆型方程} \quad Ma_\infty > 1 \text{ 双曲型方程}$$

1. 线化扰动速度势方程:

$$Ma_\infty > 1 : \quad \lambda^2 \frac{\partial^2\hat{\phi}}{\partial x^2} - \frac{\partial^2\hat{\phi}}{\partial y^2} = 0 \quad \lambda = \sqrt{Ma_\infty^2 - 1}$$

解为: $\hat{\phi} = f(x - \lambda y)$

$$\begin{array}{ll} \frac{\partial\hat{\phi}}{\partial x} = f'(x - \lambda y) & \frac{\partial^2\hat{\phi}}{\partial x^2} = f'' \\ \frac{\partial\hat{\phi}}{\partial y} = (-\lambda)f'(x - \lambda y) & \frac{\partial^2\hat{\phi}}{\partial y^2} = \lambda^2 f'' \end{array} \quad \longrightarrow \quad \lambda^2 \frac{\partial^2\hat{\phi}}{\partial x^2} - \frac{\partial^2\hat{\phi}}{\partial y^2} = \lambda^2 f'' - \lambda^2 f'' = 0$$

$\longrightarrow \hat{\phi} = f(x - \lambda y)$ 为 $\lambda^2 \frac{\partial^2\hat{\phi}}{\partial x^2} - \frac{\partial^2\hat{\phi}}{\partial y^2} = 0$ 的解。

13.7 超声速线化理论(10.9)

1. 线化扰动速度势方程:

$$Ma_\infty > 1: \quad \lambda^2 \frac{\partial^2 \hat{\phi}}{\partial x^2} - \frac{\partial^2 \hat{\phi}}{\partial y^2} = 0 \quad \lambda = \sqrt{Ma_\infty^2 - 1}$$

$\hat{\phi} = f(x - \lambda y)$ 为解!

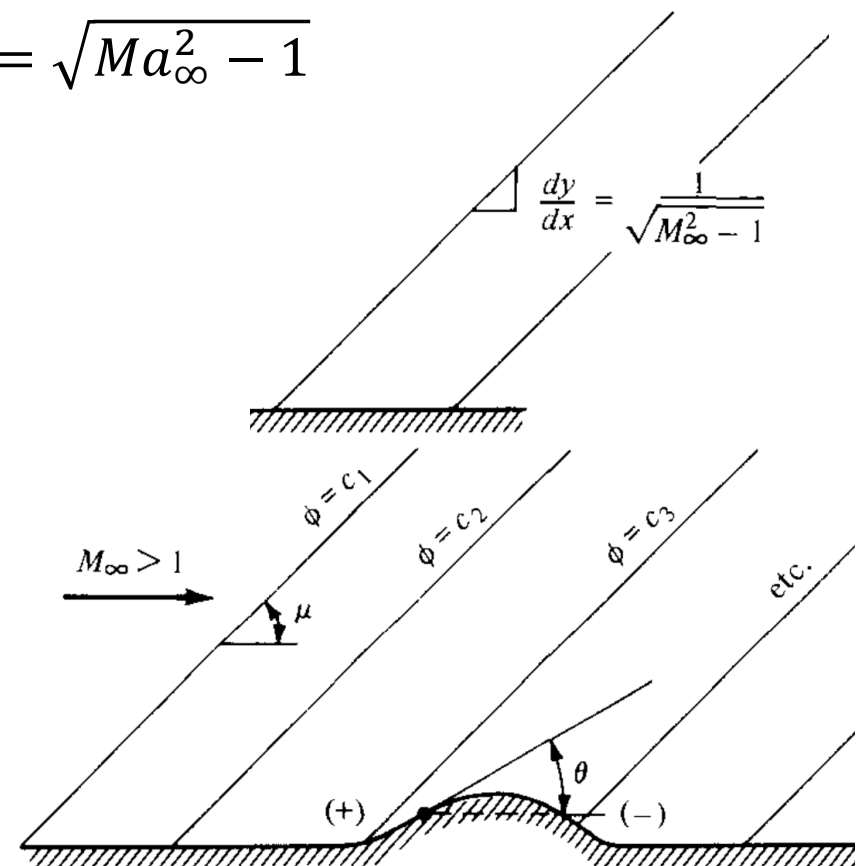
$\hat{\phi}$ 沿 $x - \lambda y = C$ 为常数。

$x - \lambda y = C$ 是斜率 $\frac{dy}{dx} = \frac{1}{\lambda} = \frac{1}{\sqrt{Ma_\infty^2 - 1}}$ 的直线。

$\tan \mu = \frac{1}{\sqrt{Ma_\infty^2 - 1}}$, μ 为马赫角。

$\hat{\phi}$ 沿马赫线为常数, 小扰动沿 Ma 线传播!

适用: 线化超声流 (小扰动), 不适用激波、膨胀波。



$Ma_\infty > 1$ 扰动无法向上游传播。

13.7 超声速线化理论(10.9)

2. 线化 C_p : $C_p = -\frac{2\hat{u}}{V_\infty}$

势流场中:

$$\left. \begin{aligned} \hat{u} &= \frac{\partial \hat{\phi}}{\partial x} = f' \\ \hat{v} &= \frac{\partial \hat{\phi}}{\partial y} = -\lambda f' \end{aligned} \right\} \hat{u} = -\frac{\hat{v}}{\lambda}$$

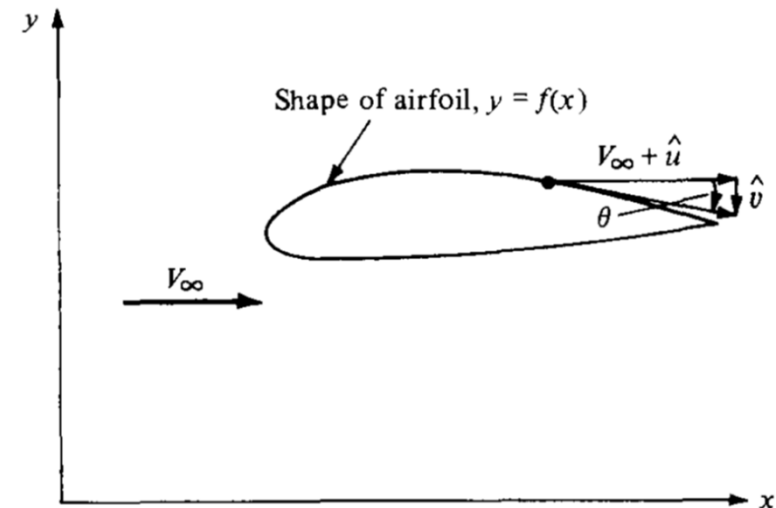
翼面上: $\hat{v} \approx V_\infty \tan \theta \approx V_\infty \theta$

翼面上: $\hat{u} = -\frac{V_\infty \theta}{\lambda}$

$C_p = -\frac{2\hat{u}}{V_\infty} = \frac{2\theta}{\lambda} = \frac{2\theta}{\sqrt{Ma_\infty^2 - 1}}$ 超声速线化压强系数!

$$C_p = \frac{2\theta}{\sqrt{Ma_\infty^2 - 1}} \quad C_p \text{ 仅与 } \theta、Ma_\infty \text{ 有关! 一级近似。}$$

$$\lambda^2 \frac{\partial^2 \hat{\phi}}{\partial x^2} - \frac{\partial^2 \hat{\phi}}{\partial y^2} = 0 \quad \hat{\phi} = f(x - \lambda y)$$

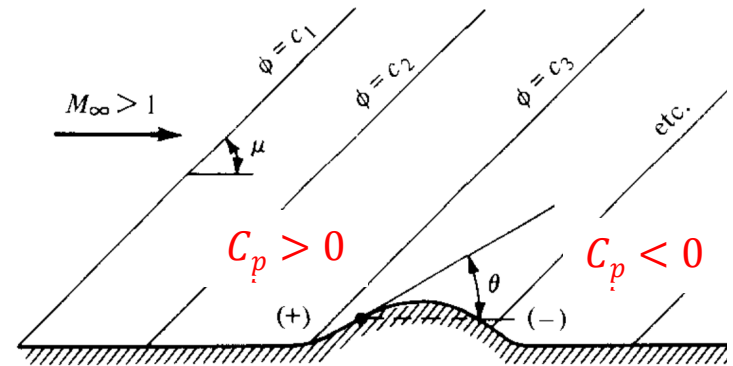


13.7 超声速线化理论(10.9)

2. 线化 C_p :

$$C_p = \frac{2\theta}{\sqrt{Ma_\infty^2 - 1}}$$

$$Ma_\infty > 1$$

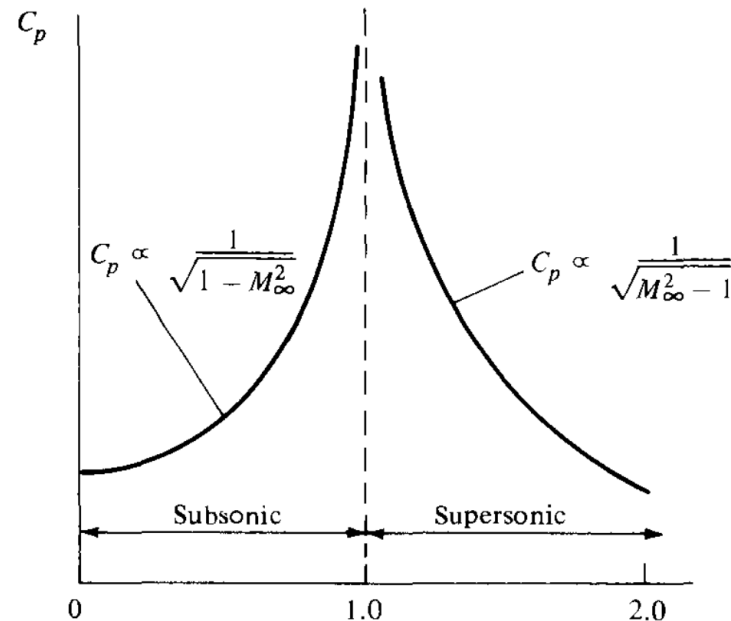


(1) $C_p \propto \theta$: 突起前高压, 后低压; \rightarrow 波阻

(2) $C_p \propto \frac{1}{\sqrt{Ma_\infty^2 - 1}}$: $Ma_\infty \uparrow C_p \downarrow$;

$$\text{亚声速: } C_p = \frac{C_{p,0}}{\sqrt{1 - Ma_\infty^2}}$$

超声速 C_p 与 θ 、 Ma_∞ 有关。



13.7超声速线化理论(10.9)

$$C_p = \frac{2\theta}{\sqrt{Ma_\infty^2 - 1}} \quad Ma_\infty > 1$$

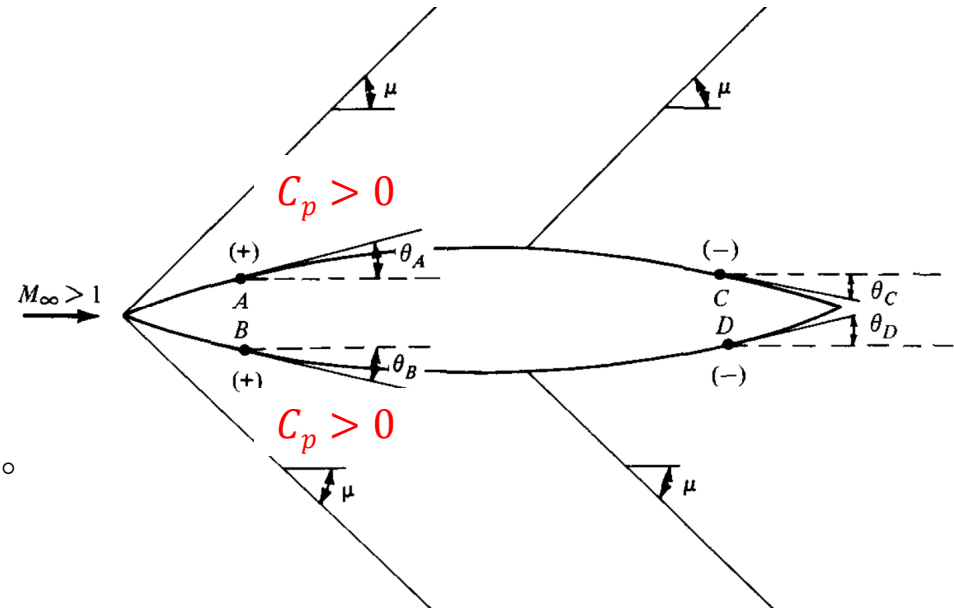
3. 超声速翼型理论:

$\theta = \frac{dy}{dx}$: 上表面 $\frac{dy}{dx} > 0$, 下表面 $\frac{dy}{dx} < 0$ 。

$$C_p = \frac{\pm 2\theta}{\sqrt{Ma_\infty^2 - 1}}$$

$+$: 壁面向气流内折; $-$: 壁面远离气流外折。

$$C_{p,u} = \frac{2 \frac{dy}{dx}}{\sqrt{Ma_\infty^2 - 1}} \quad C_{p,l} = \frac{-2 \frac{dy}{dx}}{\sqrt{Ma_\infty^2 - 1}}$$



13.7 超声速线化理论(10.9)

3. 超声速翼型理论: $C_p = \frac{\pm 2\theta}{\sqrt{Ma_\infty^2 - 1}}$

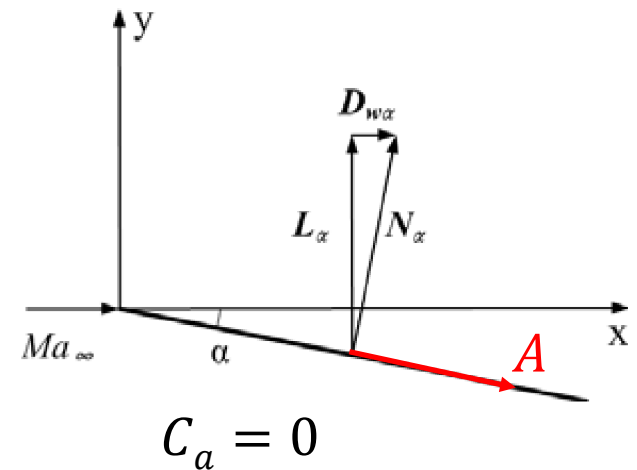
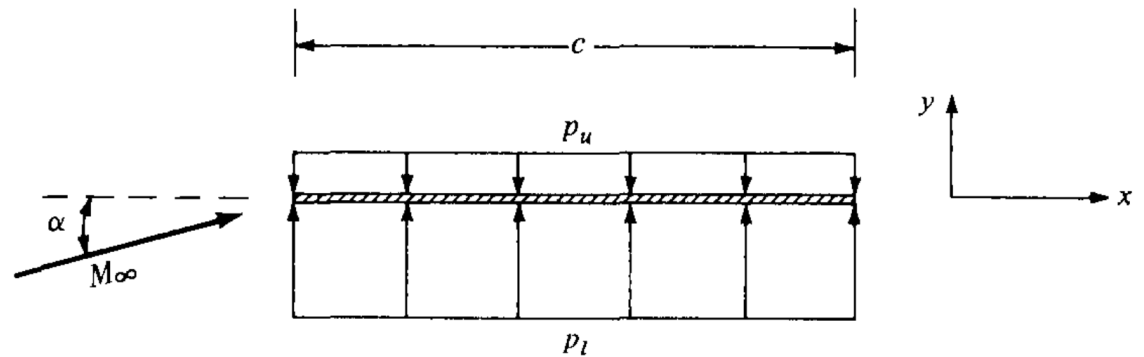
1) 平板: $C_{p,u} = \frac{-2\alpha}{\sqrt{Ma_\infty^2 - 1}}$

$$C_{p,l} = \frac{2\alpha}{\sqrt{Ma_\infty^2 - 1}}$$

$$\Delta C_{p,\alpha} = C_{p,l} - C_{p,u} = \frac{4\alpha}{\sqrt{Ma_\infty^2 - 1}} = \frac{4\alpha}{B}$$

$$C_n = \frac{1}{c} \int_0^c (C_{p,l} - C_{p,u}) dx$$

$$= \frac{4\alpha}{\sqrt{Ma_\infty^2 - 1}} \frac{1}{c} \int_0^c dx = \frac{4\alpha}{\sqrt{Ma_\infty^2 - 1}}$$



13.7超声速线化理论(10.9)

3. 超声速翼型理论:

1) 平板: $C_n = \frac{4\alpha}{\sqrt{Ma_\infty^2 - 1}} \quad C_a = 0$

$$C_l = C_n \cos \alpha - C_a \sin \alpha$$

$$C_d = C_n \sin \alpha + C_a \cos \alpha$$

$$\sin \alpha \approx \alpha, \cos \alpha \approx 1$$

→ $C_l = C_n \quad C_d = C_n \alpha$

$Ma_\infty > 1$, 平板:

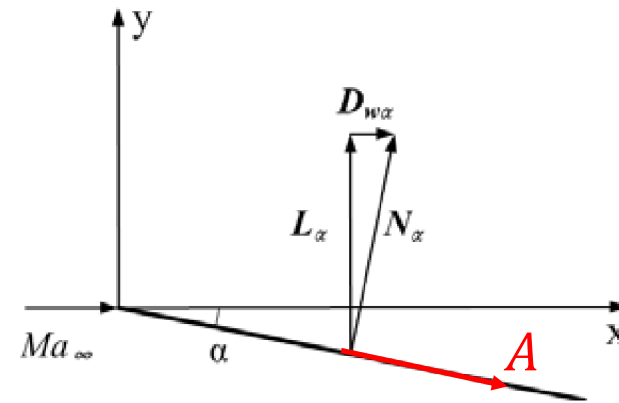
$$C_l = \frac{4\alpha}{\sqrt{Ma_\infty^2 - 1}}$$

$$C_d = \frac{4\alpha^2}{\sqrt{Ma_\infty^2 - 1}}$$

无粘：波阻!

$$C_d = C_l \alpha \quad \text{升致波阻!}$$

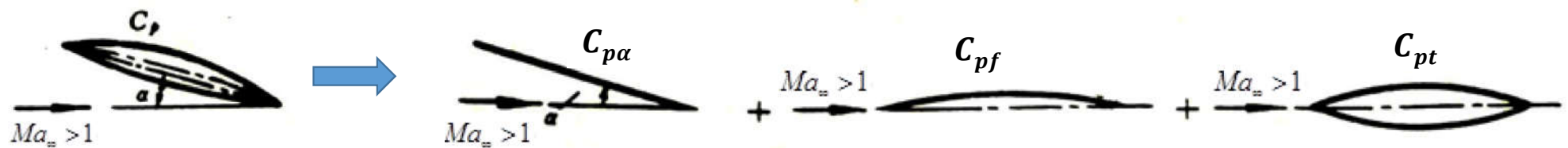
$Ma_\infty < 1$ 无粘： $C_d = 0!!$



13.7 超声速线化理论(10.9)

$$C_{p,u} = \frac{2\theta}{B} \quad C_{p,l} = \frac{-2\theta}{B} \quad B = \sqrt{Ma_\infty^2 - 1}$$

4. 线化压强系数叠加（迎角、弯度、厚度）：



$$C_p = C_{p\alpha} + C_{pf} + C_{pt}$$

α :迎角为 α 的平板绕流；

f :迎角为0，中弧线弯度为 f 的弯板绕流；

t :迎角、弯度均为0，厚度为 t 的翼型绕流。

$$C_{p,u} = C_{p\alpha,u} + C_{pf,u} + C_{pt,u} = \frac{2}{B} \left[\left(\frac{dy_\alpha}{dx} \right)_u + \left(\frac{dy_f}{dx} \right)_u + \left(\frac{dy_t}{dx} \right)_u \right]$$

$$C_{p,l} = C_{p\alpha,l} + C_{pf,l} + C_{pt,l} = \frac{-2}{B} \left[\left(\frac{dy_\alpha}{dx} \right)_l + \left(\frac{dy_f}{dx} \right)_l + \left(\frac{dy_t}{dx} \right)_l \right]$$

13.7 超声速线化理论(10.9)

$$C_{p,u} = \frac{2\theta}{B} \quad C_{p,l} = \frac{-2\theta}{B} \quad B = \sqrt{Ma_\infty^2 - 1}$$

4. 线化压强系数叠加（迎角、弯度、厚度）：

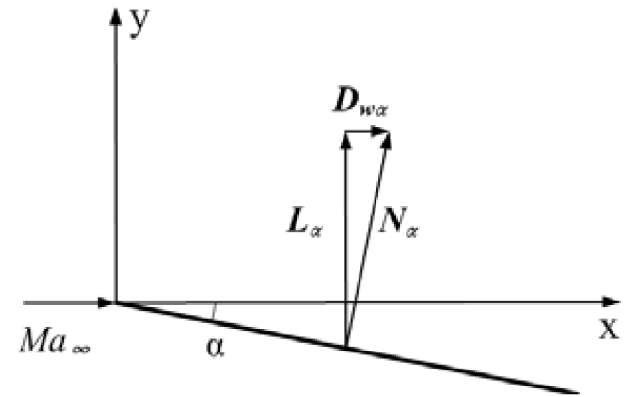
$$C_p = C_{p\alpha} + C_{pf} + C_{pt}$$

1) 平板绕流($C_{p\alpha}$)：

$$C_{p\alpha,u} = \frac{2}{B} \left(\frac{dy_\alpha}{dx} \right)_u = \frac{-2\alpha}{B}$$

$$C_{p\alpha,l} = \frac{-2}{B} \left(\frac{dy_\alpha}{dx} \right)_u = \frac{2\alpha}{B}$$

$$\Delta C_{p\alpha} = \frac{4\alpha}{B} \quad B = \sqrt{Ma_\infty^2 - 1}$$



13.7 超声速线化理论(10.9)

$$C_{p,u} = \frac{2\theta}{B} \quad C_{p,l} = \frac{-2\theta}{B} \quad B = \sqrt{Ma_\infty^2 - 1}$$

4. 线化压强系数叠加（迎角、弯度、厚度）：

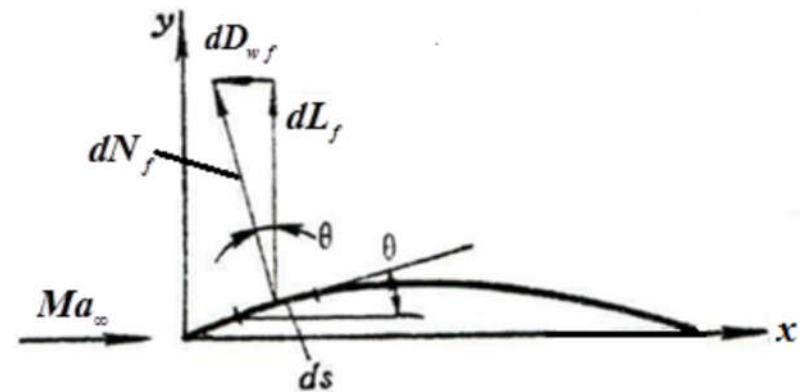
$$C_p = C_{p\alpha} + C_{pf} + C_{pt}$$

2) 弯板绕流(C_{pf})：

$$C_{pf,u} = \frac{2}{B} \left(\frac{dy_f}{dx} \right)_u$$

$$C_{pf,l} = \frac{-2}{B} \left(\frac{dy_f}{dx} \right)_l \quad \text{上下面} \left(\frac{dy_f}{dx} \right)_u = \left(\frac{dy_f}{dx} \right)_l$$

$$\Delta C_{pf} = \frac{-4}{B} \left(\frac{dy_f}{dx} \right) \quad B = \sqrt{Ma_\infty^2 - 1}$$



13.7 超声速线化理论(10.9)

$$C_{p,u} = \frac{2\theta}{B} \quad C_{p,l} = \frac{-2\theta}{B} \quad B = \sqrt{Ma_\infty^2 - 1}$$

4. 线化压强系数叠加（迎角、弯度、厚度）：

$$C_p = C_{p\alpha} + C_{pf} + C_{pt}$$

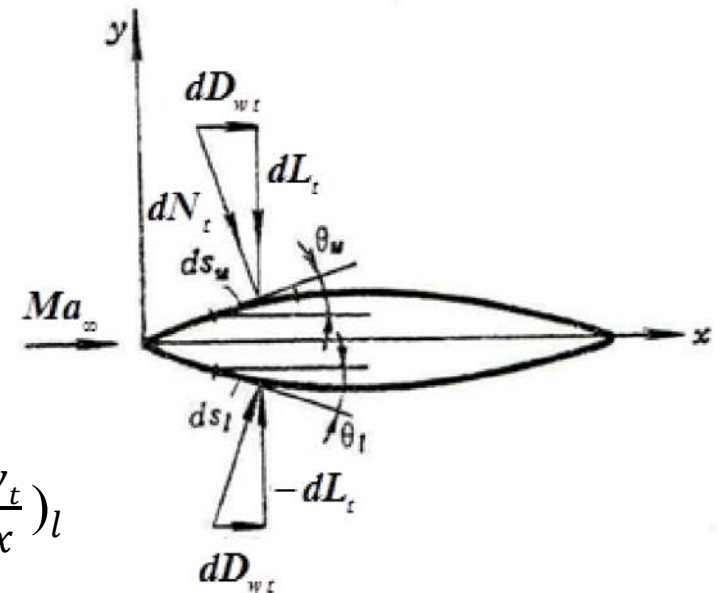
3) 厚度影响(C_{pt})：

$$C_{pt,u} = \frac{2}{B} \left(\frac{dy_t}{dx} \right)_u$$

$$C_{pt,l} = \frac{-2}{B} \left(\frac{dy_t}{dx} \right)_l \quad \text{上下面对称} \left(\frac{dy_t}{dx} \right)_u = - \left(\frac{dy_t}{dx} \right)_l$$

$$\Delta C_{pt} = 0$$

厚度引起的压差载荷为0!!



13.7超声速线化理论(10.9)

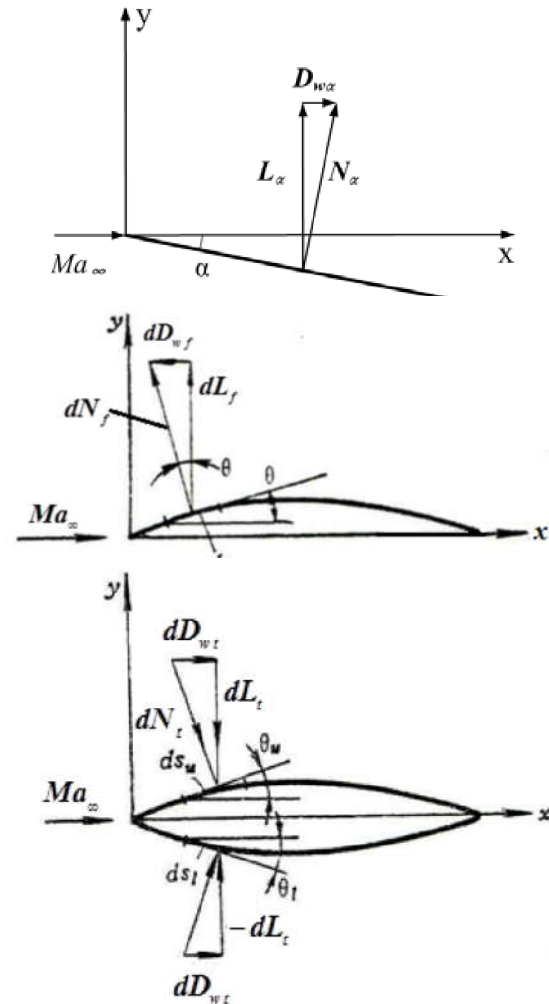
$$C_{p,u} = \frac{2\theta}{B} \quad C_{p,l} = \frac{-2\theta}{B} \quad B = \sqrt{Ma_\infty^2 - 1}$$

4. 线化压强系数叠加（迎角、弯度、厚度）：

$$C_p = C_{p\alpha} + C_{pf} + C_{pt}$$

$$\Delta C_{p\alpha} = \frac{4\alpha}{B} \quad \Delta C_{pf} = \frac{-4}{B} \left(\frac{dy_f}{dx} \right) \quad \Delta C_{pt} = 0$$

$$\begin{aligned} \Delta C_p &= \Delta C_{p\alpha} + \Delta C_{pf} + \Delta C_{pt} \\ &= \frac{4\alpha}{B} - \frac{4}{B} \left(\frac{dy_f}{dx} \right) \end{aligned}$$



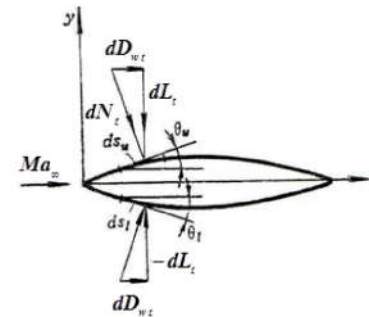
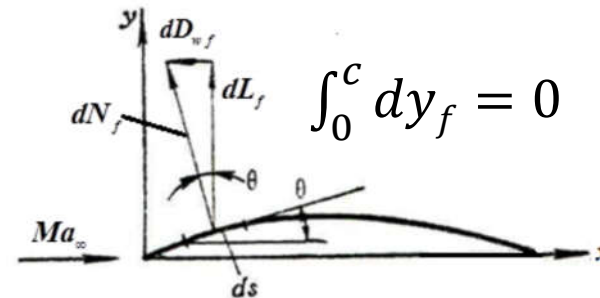
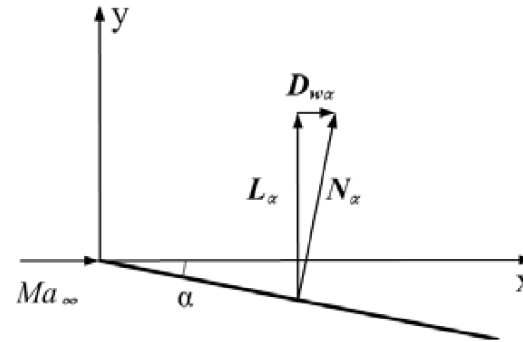
13.7超声速线化理论(10.9) $\Delta C_p = \Delta C_{p\alpha} + \Delta C_{pf} + \Delta C_{pt} = \frac{4\alpha}{B} - \frac{4}{B} \left(\frac{dy_f}{dx} \right)$

5. C_l : $C_l = C_{l\alpha} + C_{lf} + C_{lt}$

$$\begin{aligned} C_{l\alpha} &= \frac{1}{c} \int_0^c \Delta C_{p\alpha} dx / \cos \alpha \cdot \cos \alpha \\ &= \frac{1}{c} \int_0^c \Delta C_{p\alpha} dx \\ &= \frac{4\alpha}{B} \end{aligned}$$

$$\begin{aligned} C_{lf} &= \frac{1}{c} \int_0^c \Delta C_{pf} ds \cos \theta \\ &= \frac{1}{c} \int_0^c -\frac{4}{B} \left(\frac{dy_f}{dx} \right) dx \\ &= \frac{1}{c} \int_0^c -\frac{4}{B} dy_f \\ &= -\frac{4}{Bc} \int_0^c dy_f = 0 \end{aligned}$$

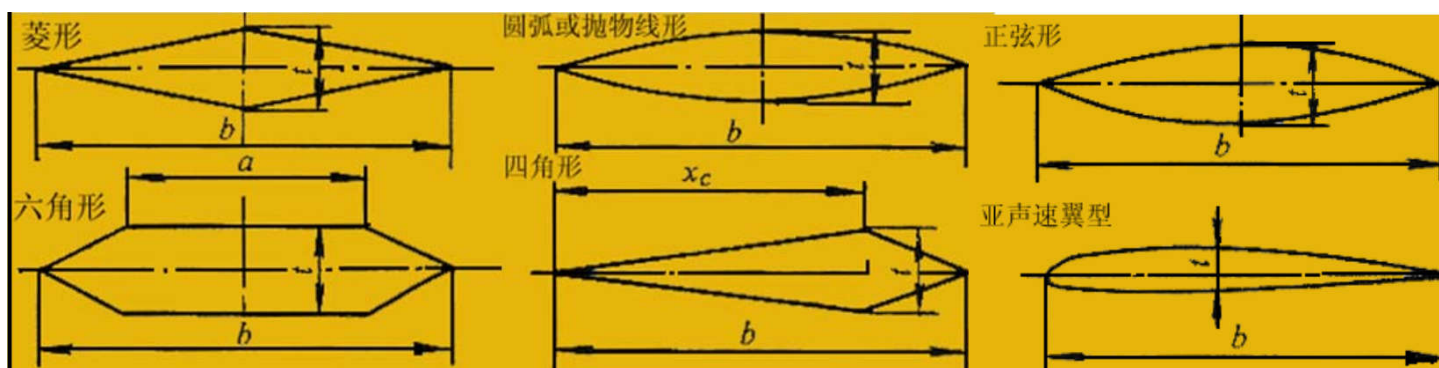
$$\Delta C_{pt} = 0 \rightarrow C_{lt} = 0$$



13.7超声速线化理论(10.9)

5. C_l : $C_l = C_{l\alpha} + C_{lf} + C_{lt}$ $C_{l\alpha} = \frac{4\alpha}{B}$ $C_{lf} = 0$ $C_{lt} = 0$ 弯度、厚度对升力无贡献！

→ $C_l = \frac{4\alpha}{B}$ 超声速翼型：升力由 α 决定，与 f, t 无关！



超声速翼型：无弯度对称、厚度小！

13.7 超声速线化理论(10.9)

$$C_l = \frac{4\alpha}{B}$$

超声速翼型：升力由 α 决定，与 f, t 无关！
弯度、厚度对升力无贡献！

6. C_{dw} : $C_{dw} = C_{dw\alpha} + C_{dwf} + C_{dwt}$
(波阻)

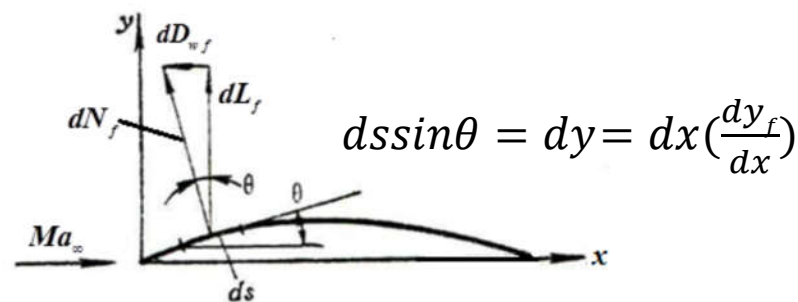
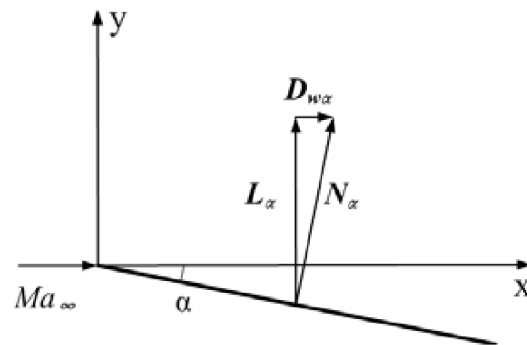
$$C_{dw\alpha} = C_{l\alpha}\alpha = \frac{4\alpha^2}{B} \text{ 升致波阻!}$$

$$C_{dwf} = \frac{1}{c} \int_0^c -\Delta C_{pf} dssin\theta$$

$$= \frac{1}{c} \int_0^c \frac{4}{B} \left(\frac{dy_f}{dx}\right) \left(\frac{dy_f}{dx}\right) dx$$

$$= \frac{4}{cB} \int_0^c \left(\frac{dy_f}{dx}\right)^2 dx$$

弯度产生波阻!



13.7超声速线化理论(10.9)

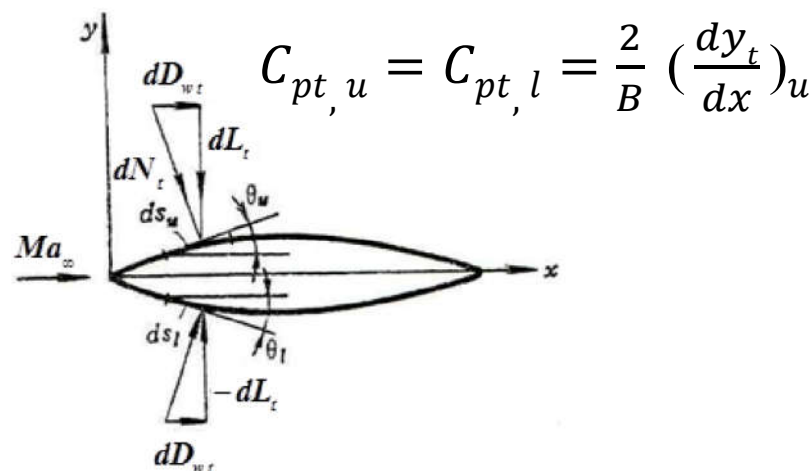
$$C_{dw\alpha} = \frac{4\alpha^2}{B} \quad C_{dwf} = \frac{4}{cB} \int_0^c \left(\frac{dy_f}{dx}\right)^2 dx$$

6. C_{dw} : $C_{dw} = C_{dw\alpha} + C_{dwf} + C_{dwt}$

$$C_{dwt} = \frac{1}{c} \int_0^c \frac{4}{B} \left(\frac{dy_t}{dx}\right) ds \sin\theta$$

$$= \frac{1}{c} \int_0^c \frac{4}{B} \left(\frac{dy_t}{dx}\right) \left(\frac{dy_t}{dx}\right) dx$$

$$= \frac{4}{cB} \int_0^c \left(\frac{dy_t}{dx}\right)^2 dx \quad \text{弯度产生波阻!}$$



$$C_{dw} = C_{dw\alpha} + C_{dwf} + C_{dwt}$$

$$C_{dw} = \frac{4\alpha^2}{B} + \frac{4}{cB} \int_0^c \left[\left(\frac{dy_f}{dx}\right)^2 + \left(\frac{dy_t}{dx}\right)^2 \right] dx = \frac{4}{B} (\alpha^2 + g_f^2 + g_t^2)$$

升致波阻

(α 引起, 与 f, t 无关)!

零升波阻

(与 f, t 有关)!

f, t 仅产生波阻, 不产生升力!

13.7超声速线化理论(10.9)

$$C_l = \frac{4\alpha}{B}$$

$$C_{dw} = \frac{4\alpha^2}{B} + \frac{4}{cB} \int_0^c \left[\left(\frac{dy_f}{dx} \right)^2 + \left(\frac{dy_t}{dx} \right)^2 \right] dx$$

$$C_l \propto \alpha, C_{dw} \propto \alpha^2 \quad \alpha \downarrow C_{dw} \downarrow \quad \text{超声速：小}\alpha\text{飞行！}$$

可见：薄翼型的波阻系数由两部分组成，一部分与升力有关，另一部分与弯度和厚度有关。

升致波阻：由迎角引起的波阻，与翼型的形状无关

零升波阻：与升力无关，只与翼剖面有关，特别是与翼型厚度有关。

13.7 超声速线化理论(10.9)

7. C_m : ① $C_{m\alpha}$: $\Delta C_{p\alpha}$ 分布均匀, 作用点在 $\frac{c}{2}$ 处

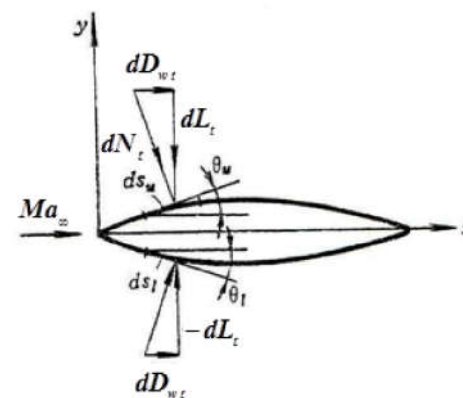
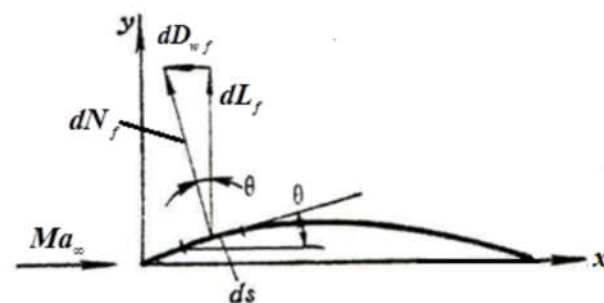
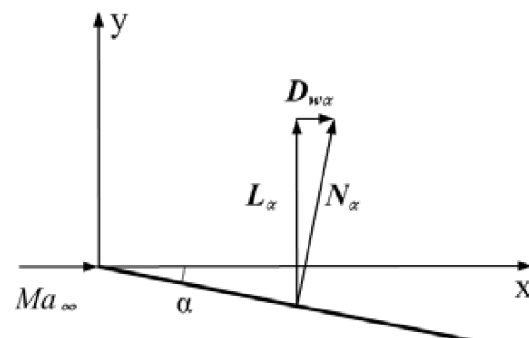
$$C_{m\alpha, 0} = -\frac{C_{l\alpha}}{2}$$

$$\begin{aligned} \text{② } C_{mf}: C_{mf, 0} &= \frac{1}{c^2} \int_0^c \frac{4}{B} \left(\frac{dy_f}{dx} \right) dx \cdot x \\ &= \frac{-4}{Bc^2} \left(\int_0^c y_f dx - y_f x \Big|_0^c \right) \\ &= \frac{-4}{Bc^2} \int_0^c y_f dx \end{aligned}$$

$$\text{③ } C_{mt}: C_{mt} = 0$$

$$C_{m, 0} = -\frac{C_l}{2} - \frac{4}{Bc^2} \int_0^c y_f dx$$

$$C_{m, 0} + \frac{C_l}{2} = -\frac{4}{Bc^2} \int_0^c y_f dx \quad \text{与 } \alpha \text{ 无关!}$$



13.7 超声速线化理论(10.9)

7. C_m : $C_{m,0} = -\frac{C_l}{2} - \frac{4}{Bc^2} \int_0^c y_f dx$

$$C_{m,0} + \frac{C_l}{2} = -\frac{4}{Bc^2} \int_0^c y_f dx \quad \text{与}\alpha\text{无关!}$$

$$C_{m,1/2} = C_{m,0} + \frac{C_l}{2} = -\frac{4}{Bc^2} \int_0^c y_f dx \quad \text{与}\alpha\text{无关!}$$

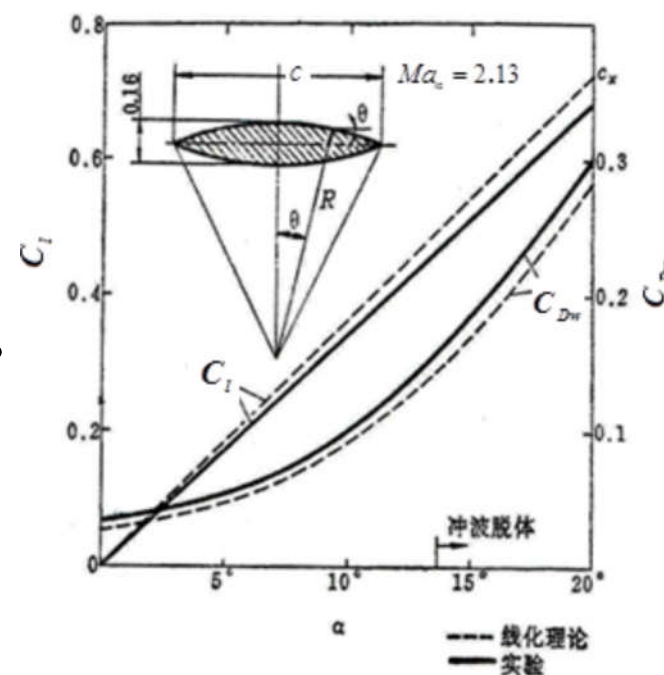
➡ $\frac{c}{2}$ 为气动中心! $Ma_\infty > 1$ 气动中心后移!

$Ma_\infty < 1$, $\frac{c}{4}$ 为气动中心!

13.7 超声速线化理论(10.9)

8. 一级线化理论的适用性:

- 升力线斜率较实验较高：
线化理论未考虑边界层及其与激波相互干扰。
- 波阻较实验较低：
线化理论为无粘阻力，未考虑粘性压差、摩擦阻力。



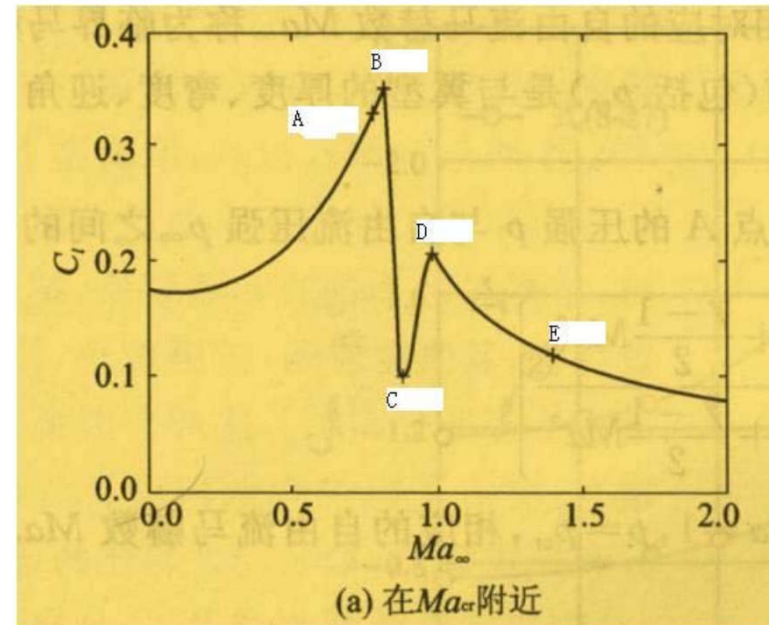
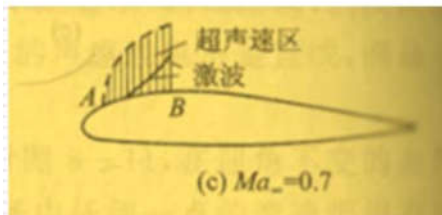
13.8 Ma 数对机翼气动特性影响

1. $C_l \sim Ma_\infty$

➤ $Ma_\infty < Ma_{cr}$: $Ma_\infty \uparrow \rightarrow C_l \uparrow$

$$\text{亚声速: } C_l = \frac{C_{l0}}{\sqrt{1 - Ma_\infty^2}}$$

➤ A - B: 刚进入跨声速区,
上翼面有超音速区, $C_p \downarrow \rightarrow C_l \uparrow$

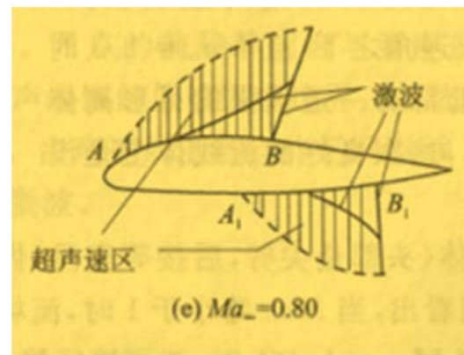
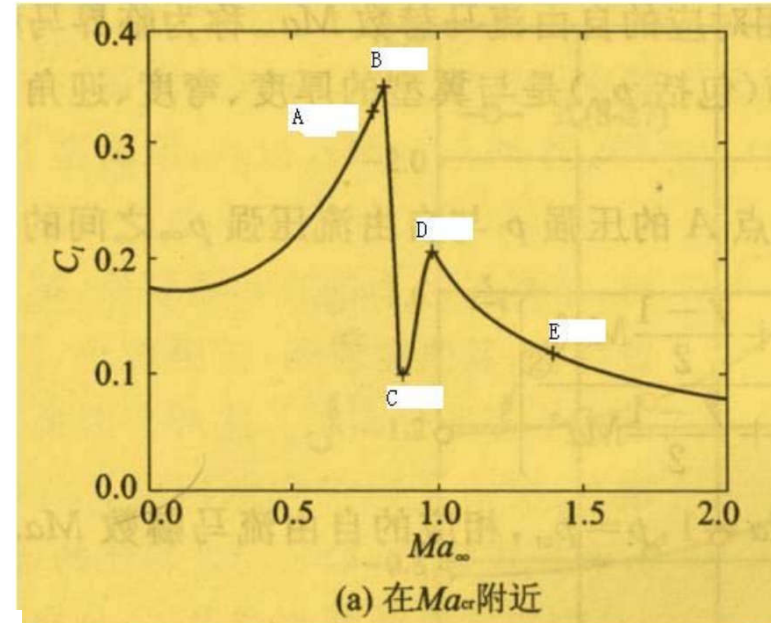


13.8 Ma 数对机翼气动特性影响

1. $C_l \sim Ma_\infty$

➤ B 后: $Ma_\infty \uparrow$, 上翼面超声速区扩大,
激波后移, 强度增大,
波后边界层内逆压梯度剧增
→ 流动分离, 升力 $C_l \uparrow$ 骤降! (激波失速)

➤ C 前: $Ma_\infty \uparrow$, 下翼面也出现超声速区,
激波快速移至后缘,
→ 下翼面压强降低,
→ 升力 $C_l \downarrow$!



13.8 Ma 数对机翼气动特性影响

1. $C_l \sim Ma_\infty$

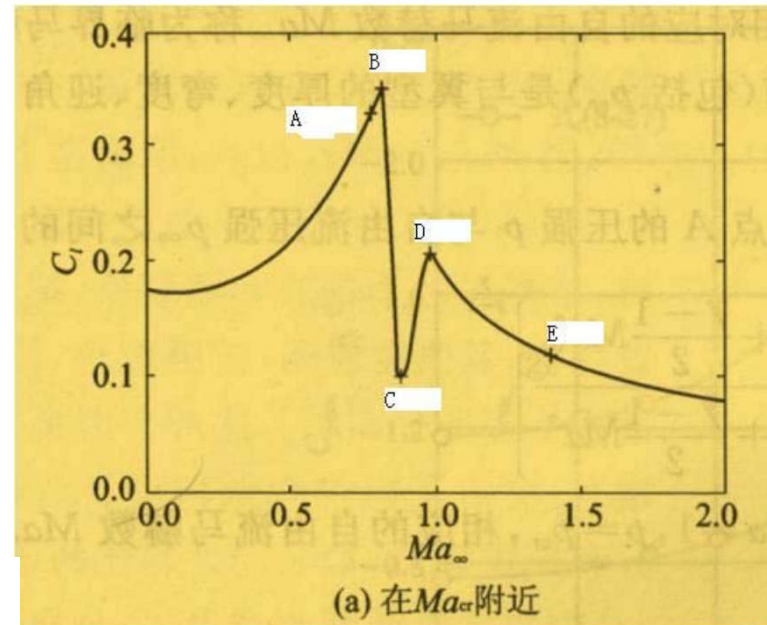
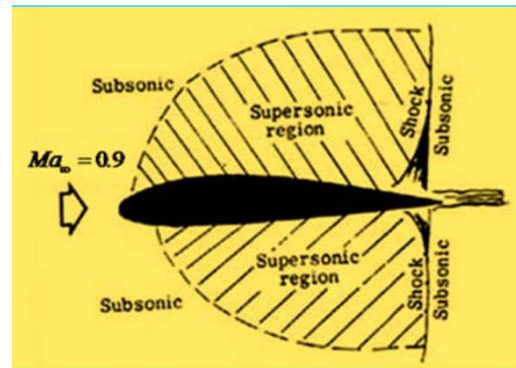
➤ $C \sim D$:

$Ma_\infty \uparrow$, 下翼面超声速区不变,

上翼面激波移至后缘,

上翼面边界层分离点后移

→ 上翼面 $C_p \downarrow \rightarrow C_l \uparrow$



13.8 Ma 数对机翼气动特性影响

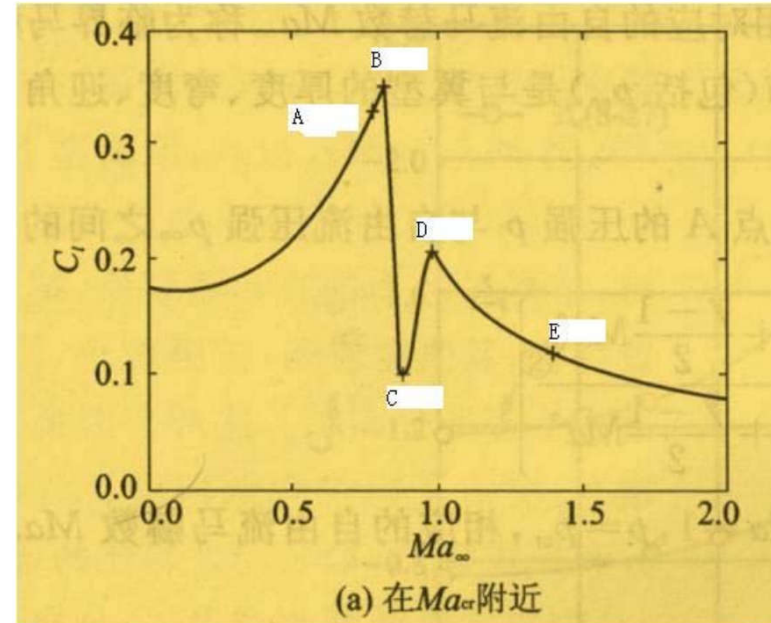
1. $C_l \sim Ma_\infty$

➤ $D \sim E$: $Ma_\infty > 1$, 翼型超声速绕流,
翼型前脱体激波, 斜激波

→ $Ma_\infty \uparrow$, $C_l \downarrow$

$$C_l = \frac{4\alpha}{\sqrt{Ma_\infty^2 - 1}}$$

跨声速飞行中, C_l 几上几下, 波动大。



13.8 Ma 数对机翼气动特性影响

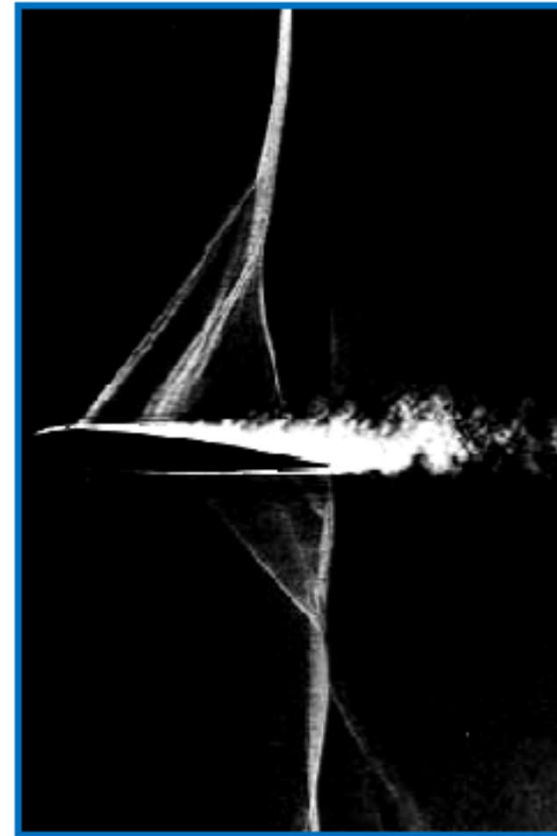
2. λ 形激波系：

激波 \rightarrow 逆压梯度 \uparrow

\rightarrow 边界层厚度 \uparrow

\rightarrow 增厚的边界层使外流形成系列压缩波

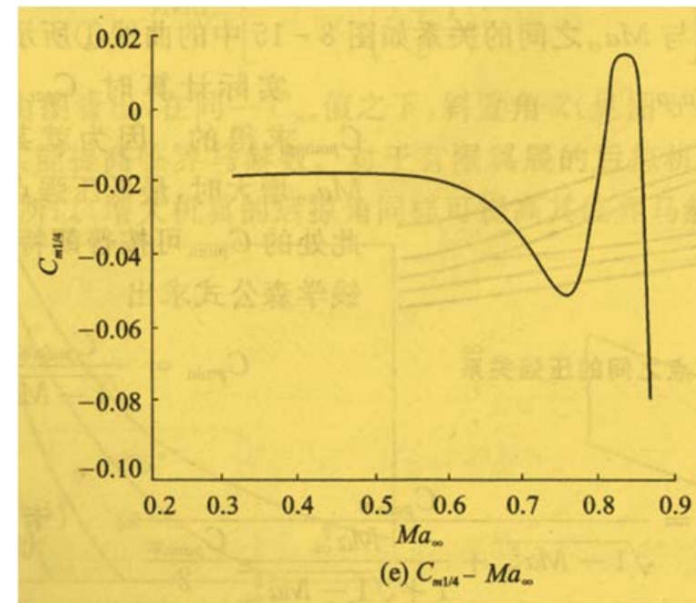
\rightarrow λ 形激波系



13.8 Ma 数对机翼气动特性影响

3. $C_m \sim Ma_\infty$:

跨声速飞行中，纵向力矩变化剧烈，
飞行器操作困难，跨越声障易出事故。



例：F104战斗机平面面积 $S=18.21\text{m}^2$ ，在11km高空（ $\rho_\infty=0.3648\text{kg/m}^3$ ， $T_\infty=216.78\text{K}$ ， $R=287\text{J/kgK}$ ）以 $Ma_\infty=2$ 飞行。飞机重量为9400kgf。假设飞机重量全由机翼承担，机翼升力系数近似等于翼型升力系数。求飞机机翼的迎角 α 。

解： 11km高空的当地声速为：

$$c = \sqrt{\lambda R T_\infty} = \sqrt{1.4 \times 287 \times 216.78} = 295 \text{ m/s}$$

$$C_l = \frac{4\alpha}{\sqrt{Ma_\infty^2 - 1}}$$

飞机飞行速度为：

$$V_\infty = Ma_\infty \cdot c = 2 \times 295 = 590 \text{ m/s}$$

飞机升力为：

$$L = 9400 \times 9.8 = 9.212 \times 10^4 \text{ N}$$

机翼升力系数为：

$$C_L = \frac{L}{\frac{1}{2} \rho_\infty V_\infty^2 \cdot S} = \frac{9.212 \times 10^4}{\frac{1}{2} \times 0.3648 \times 590^2 \times 18.21} = 0.08$$

例：F104战斗机平面面积 $S=18.21\text{m}^2$ ，在11km高空（ $\rho_\infty=0.3648\text{kg/m}^3$ ， $T_\infty=216.78\text{K}$ ， $R=287\text{J/kgK}$ ）以 $Ma_\infty=2$ 飞行。飞机重量为9400kgf。假设飞机重量全由机翼承担，机翼升力系数近似等于翼型升力系数。求飞机机翼的迎角 α 。

解：翼型升力系数为： $C_l = C_L = 0.08$

超声速翼型升力系数与迎角、来流马赫数的关系为：

$$C_l = \frac{4\alpha}{\sqrt{Ma_\infty^2 - 1}}$$

则，机翼的迎角为：

迎角很小，符合小扰动假设。

$$\alpha = \frac{C_l}{4} \sqrt{Ma_\infty^2 - 1} = \frac{0.08}{4} \sqrt{2^2 - 1} = 0.035\text{rad} = 1.98^\circ$$

作业：

复习笔记！

空气动力学书10.1, 10.2, 10.3