

Physics 115A – Homework 1

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Question 1 (Discrete statistics at the Tour de France). *Suppose we're lucky enough to gather all 61 official winners of the Tour de France in one room, alive and dead. Four people have won it 5 times apiece, one person has won it 4 times, three people have won it 3 times apiece, twelve people have won it 2 times apiece, and forty-one people have won it once. If we let $N(j)$ represent the number of winners with j victories in the Tour de France, then $N(5) = 4$, $N(4) = 1$, $N(3) = 3$, $N(2) = 12$, $N(1) = 41$, and all other $N(j)$ for $j \neq 5, 4, 3, 2, 1$ are zero.*

(a) Compute $\langle j^2 \rangle$ and $\langle j \rangle^2$.

(b) Determine $\Delta j \equiv j - \langle j \rangle$ for each j , and compute the standard deviation using Griffiths Eq. 1.11 (i.e., $\sigma^2 = \langle (\Delta j)^2 \rangle$).

(c) Use your results from (a) and (b) to verify Griffiths Eq. 1.12 (i.e., $\sigma = \sqrt{\langle j^2 \rangle - \langle j \rangle^2}$).

Proof.

- (a) First, consider the total number of winners, we get $N_{\text{tot}} = \sum_{i=1}^5 N(i) = 61$ (in particular, since for $j > 5$ we have $N(j) = 0$, one can exclude all cases for $j > 5$). So, the probability for $j = 1, 2, 3, 4, 5$ is given as $\mathbb{P}(j) = \frac{N(j)}{N_{\text{tot}}}$. Then, the expectation value of j and j^2 is given as follow, respectively:

$$\langle j \rangle = \sum_{i=1}^5 j \cdot \mathbb{P}(j) = 1 \cdot \frac{41}{61} + 2 \cdot \frac{12}{61} + 3 \cdot \frac{3}{61} + 4 \cdot \frac{1}{61} + 5 \cdot \frac{4}{61} = \frac{98}{61} \quad (1)$$

$$\langle j^2 \rangle = \sum_{j=1}^5 j^2 \cdot \mathbb{P}(j) = 1 \cdot \frac{41}{61} + 4 \cdot \frac{12}{61} + 9 \cdot \frac{3}{61} + 16 \cdot \frac{1}{61} + 25 \cdot \frac{4}{61} = \frac{232}{61} \quad (2)$$

So, we get $\langle j^2 \rangle = \frac{232}{61}$, and $\langle j \rangle^2 = \left(\frac{98}{61}\right)^2$.

- (b) For each j , the difference from the expectation value $\langle j \rangle$ is given as:

$$\Delta 1 = 1 - \frac{98}{61} = \frac{-37}{61}, \quad \Delta 2 = 2 - \frac{98}{61} = \frac{24}{61}, \quad \Delta 3 = 3 - \frac{98}{61} = \frac{85}{61} \quad (3)$$

$$\Delta 4 = 4 - \frac{98}{61} = \frac{146}{61}, \quad \Delta 5 = 5 - \frac{98}{61} = \frac{207}{61} \quad (4)$$

So, the variance is given as:

$$\sigma^2 = \langle (\Delta j)^2 \rangle = \sum_{j=1}^5 (\Delta j)^2 \cdot \mathbb{P}(j) = \frac{4548}{3721} \quad (5)$$

Hence, the standard deviation $\sigma = \sqrt{\frac{4548}{3721}} \approx 1.106$.

- (c) If we consider $\langle j^2 \rangle - \langle j \rangle^2 = \frac{232}{61} - \left(\frac{98}{61}\right)^2 = \frac{4548}{3721} = \sigma^2$ found in part (b), hence $\sigma = \sqrt{\langle j^2 \rangle - \langle j \rangle^2}$ in this case.

□

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Question 2 (Fun with gaussians). We'll encounter physical situations involving gaussian probability densities repeatedly throughout the course, so now's a good time to get comfortable with them. If you're finding yourself stuck, please take a look at the accompanying note on gaussian integrals on Gauchospace.

(a) Consider the probability density $\rho(x) = Ae^{-\lambda(x-\mu)^2}$, where A, μ, λ are positive, real constants. Determine A using Griffiths Eq. 1.16 (i.e.,

$$\int_{-\infty}^{\infty} \rho(x) dx = 1).$$

Find $\langle x \rangle$, $\langle x^2 \rangle$, and σ . Sketch $\rho(x)$.

(b) Consider now the probability density $\rho(x) = Axe^{-\lambda(x-\mu)^2}$ for positive, real A, μ, λ . Determine A , $\langle x \rangle$, $\langle x^2 \rangle$, and σ . Sketch $\rho(x)$. Although it is normalizable, can $\rho(x)$ always be considered a probability density?

(c) Now consider the probability density $\rho(x) = Ax^2e^{-\lambda(x-\mu)^2}$ for positive, real A, μ, λ . Determine A , $\langle x \rangle$, $\langle x^2 \rangle$, and σ . Sketch $\rho(x)$.

(d) Shifting back to QM, consider the gaussian wavefunction, $\Psi(x) = Ae^{-x^2/b^2}$, find A from the normalization condition (remember how the wavefunction relates to the probability density). For this wavefunction, calculate $\langle x \rangle$, $\langle x^2 \rangle$, and σ_x , as well as $\langle p \rangle$, $\langle p^2 \rangle$, and σ_p . Test the uncertainty principle on σ_p and σ_x for this distribution.

Proof.

- (a) First, $\rho(x) = Ae^{-\lambda(x-\mu)^2} = Ae^{-(\sqrt{\lambda}(x-\mu))^2}$. Hence, if perform the integration, doing substitution $u = \sqrt{\lambda}(x - \mu)$ and $du = \sqrt{\lambda}dx$, we get:

$$\int_{-\infty}^{\infty} \rho(x) dx = \int_{-\infty}^{\infty} Ae^{-(\sqrt{\lambda}(x-\mu))^2} dx = \frac{A}{\sqrt{\lambda}} \int_{-\infty}^{\infty} e^{-u^2} du = A\sqrt{\frac{\pi}{\lambda}} \quad (6)$$

Hence, for the normalization to happen (or the above integral being 1), we need $A = \sqrt{\frac{\lambda}{\pi}}$.

To find $\langle x \rangle$ and $\langle x^2 \rangle$, consider the following integrals (using similar substitutions as above):

$$\langle x \rangle = \int_{-\infty}^{\infty} x\rho(x) dx = \int_{-\infty}^{\infty} x \cdot Ae^{-(\sqrt{\lambda}(x-\mu))^2} dx \quad (7)$$

$$= \frac{1}{\sqrt{\lambda}} \int_{-\infty}^{\infty} \sqrt{\lambda}(x - \mu) Ae^{-(\sqrt{\lambda}(x-\mu))^2} dx + \mu \int_{-\infty}^{\infty} Ae^{-(\sqrt{\lambda}(x-\mu))^2} dx \quad (8)$$

$$= \frac{A}{\lambda} \int_{-\infty}^{\infty} ue^{-u^2} du + \mu = \mu \quad (9)$$

$$\langle x^2 \rangle = \int_{-\infty}^{\infty} x^2 \rho(x) dx = \int_{-\infty}^{\infty} x^2 \cdot A e^{-(\sqrt{\lambda}(x-\mu))^2} dx \quad (10)$$

$$= \frac{1}{\lambda} \int_{-\infty}^{\infty} (\sqrt{\lambda}(x-\mu))^2 A e^{-(\sqrt{\lambda}(x-\mu))^2} dx + \int_{-\infty}^{\infty} 2\mu x A e^{-(\sqrt{\lambda}(x-\mu))^2} dx - \int_{-\infty}^{\infty} \mu^2 A e^{-(\sqrt{\lambda}(x-\mu))^2} dx \quad (11)$$

$$= \frac{A}{\lambda \sqrt{\lambda}} \int_{-\infty}^{\infty} u^2 e^{-u^2} du + 2\mu \langle x \rangle - \mu^2 \cdot \int_{-\infty}^{\infty} \rho(x) dx \quad (12)$$

$$= \frac{A}{\lambda^{\frac{3}{2}}} \left(-\frac{1}{2} u e^{-u^2} \Big|_{-\infty}^{\infty} + \frac{1}{2} \int_{-\infty}^{\infty} e^{-u^2} du \right) + 2\mu^2 - \mu^2 = \sqrt{\frac{\lambda}{\pi}} \cdot \frac{1}{\lambda^{\frac{3}{2}}} \cdot \frac{\sqrt{\pi}}{2} + \mu^2 = \mu^2 + \frac{1}{2\lambda} \quad (13)$$

With these values in mind, we get $\sigma = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \sqrt{\mu^2 + \frac{1}{2\lambda} - \mu^2} = \frac{1}{\sqrt{2\lambda}}$.

And, the following is a sketch of $\rho(x)$:

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- (b) If now consider the function $\rho(x) = A x e^{-\lambda(x-\mu)^2}$, where $\rho(x)$ is normalized. To get the constant A , consider the following integral (with the same substitution as part (a)):

$$\int_{-\infty}^{\infty} \rho(x) dx = \int_{-\infty}^{\infty} A x e^{-(\sqrt{\lambda}(x-\mu))^2} dx = \frac{A}{\sqrt{\lambda}} \int_{-\infty}^{\infty} \sqrt{\lambda}(x-\mu) e^{-(\sqrt{\lambda}(x-\mu))^2} dx + A\mu \int_{-\infty}^{\infty} e^{-(\sqrt{\lambda}(x-\mu))^2} dx \quad (14)$$

$$= \frac{A}{\lambda} \int_{-\infty}^{\infty} u e^{-u^2} du + \frac{A\mu}{\sqrt{\lambda}} \int_{-\infty}^{\infty} e^{-u^2} du = A\mu \sqrt{\frac{\pi}{\lambda}} \quad (15)$$

Hence, for normalization to happen (or the above integral is 1), we need $A = \sqrt{\frac{\lambda}{\pi\mu^2}}$.

To find $\langle x \rangle$ and $\langle x^2 \rangle$, consider the following integrals:

$$\langle x \rangle = \int_{-\infty}^{\infty} x \rho(x) dx = \int_{-\infty}^{\infty} x \cdot A x e^{-(\sqrt{\lambda}(x-\mu))^2} dx = \frac{A}{\lambda^{\frac{3}{2}}} \left(-\frac{1}{2} u e^{-u^2} \Big|_{-\infty}^{\infty} + \frac{1}{2} \int_{-\infty}^{\infty} e^{-u^2} du \right) + \mu^2 \quad (16)$$

$$= \sqrt{\frac{\lambda}{\pi\mu^2}} \cdot \frac{1}{\lambda^{\frac{3}{2}}} \cdot \frac{\sqrt{\pi}}{2} + \mu^2 = \frac{1}{2\mu\lambda} + \mu^2 \quad (17)$$

(**Note:** The above integral is the same form of $\langle x^2 \rangle$ in part (a), except A is different, so we can continue from Equation (13)).

$$\langle x^2 \rangle = \int_{-\infty}^{\infty} x^2 \rho(x) dx = \int_{-\infty}^{\infty} x^2 \cdot A x e^{-(\sqrt{\lambda}(x-\mu))^2} dx \quad (18)$$

$$= \frac{1}{\sqrt{\lambda}^3} \int_{-\infty}^{\infty} (\sqrt{\lambda}(x-\mu))^3 \cdot A e^{-(\sqrt{\lambda}(x-\mu))^2} dx + \int_{-\infty}^{\infty} 3\mu x^2 \cdot A e^{-(\sqrt{\lambda}(x-\mu))^2} dx \quad (19)$$

$$- \int_{-\infty}^{\infty} 3\mu^2 x \cdot A e^{-(\sqrt{\lambda}(x-\mu))^2} dx + \int_{-\infty}^{\infty} \mu^3 \cdot A e^{-(\sqrt{\lambda}(x-\mu))^2} dx \quad (20)$$

$$= \frac{A}{\lambda^2} \int_{-\infty}^{\infty} u^3 e^{-u^2} du + 3\mu \langle x \rangle - 3\mu^2 \int_{-\infty}^{\infty} \rho(x) dx + \frac{\mu^3 A}{\sqrt{\lambda}} \int_{-\infty}^{\infty} e^{-u^2} du \quad (21)$$

$$= \frac{3}{2\lambda} + 3\mu^3 - 3\mu^2 + \mu^3 A \sqrt{\frac{\pi}{\lambda}} = \frac{3}{2\lambda} + 3\mu^3 - 2\mu^2 \quad (22)$$

With the above equations, we get $\sigma = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \sqrt{\frac{3}{2\lambda} + 3\mu^3 - 2\mu^2 - \frac{1}{4\mu^2\lambda^2} - \frac{\mu}{\lambda} - \mu^4}$. And, here is a sketch of $\rho(x)$:

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Overall, even though $\rho(x)$ is normalizable, the reason why $\rho(x)$ is not a probability density function, is because it yields negative values for $x < 0$, which is not allowed for general probability density functions.

(c) If consider $\rho(x) = Ax^2e^{-\lambda(x-\mu)^2}$, for normalization, consider the following:

$$\int_{-\infty}^{\infty} \rho(x) dx = \int_{-\infty}^{\infty} Ax^2 e^{-(\sqrt{\lambda}(x-\mu))^2} dx \quad (23)$$

(d) Given $\Psi(x) = Ae^{-(x/b)^2}$, for normalization, we get the following integral:

$$1 = \int_{-\infty}^{\infty} |\Psi|^2 dx = \int_{-\infty}^{\infty} A^2 e^{-(\sqrt{2}x/b)^2} dx = \frac{bA^2}{\sqrt{2}} \int_{-\infty}^{\infty} e^{-u^2} du = A^2 \cdot \sqrt{\frac{b^2\pi}{2}} \quad (24)$$

Hence, $A = \left(\frac{2}{b^2\pi}\right)^{\frac{1}{4}}$.

Given the normalization, $\langle x \rangle$ and $\langle x^2 \rangle$ are given as:

$$\langle x \rangle = \int_{-\infty}^{\infty} x |\Psi|^2 dx = \int_{-\infty}^{\infty} x A^2 e^{-2(x/b)^2} dx = 0 \quad (25)$$

(The above is the property of odd functions).

$$\langle x^2 \rangle = \int_{-\infty}^{\infty} x^2 |\Psi|^2 dx = \frac{b^2}{2} \int_{-\infty}^{\infty} A^2 \left(\frac{\sqrt{2}x}{b}\right)^2 e^{-(\sqrt{2}x/b)^2} dx \quad (26)$$

$$= \frac{b^3}{2\sqrt{2}} \int_{-\infty}^{\infty} A^2 u^2 e^{-u^2} du \quad (27)$$

$$= \frac{b^3}{2\sqrt{2}} A^2 \cdot \left(-\frac{1}{2}\right) u \cdot e^{-u^2} \Big|_{-\infty}^{\infty} + \frac{b^3}{4\sqrt{2}} A^2 \int_{-\infty}^{\infty} e^{-u^2} du \quad (28)$$

$$= \frac{b^3}{4} A^2 \cdot \sqrt{\frac{\pi}{2}} = \frac{b^3}{4} \cdot \sqrt{\frac{2}{b^2\pi}} \cdot \sqrt{\frac{\pi}{2}} = \frac{b^2}{4} \quad (29)$$

□

Question 3 (Fun with probability densities). Consider the function

$$\rho(r) = \begin{cases} Ar^2 & r \leq R, \\ 0 & r \geq R \end{cases}$$

such that $dP = 4\pi\rho(r)r^2dr$ represents the probability of finding a particle in a spherical shell of thickness dr located at the radial coordinate r in three-dimensional spherical coordinates.

(a) Determine A by requiring that the probability of finding the particle anywhere is unity. Don't forget that we're in three spatial dimensions.

(b) What is the probability of finding the particle between $r = R/2$ and $r = R$?

(c) Determine $\langle r \rangle$, the average radial position of the particle.

Proof.

- (a) Since $dP = 4\pi\rho(r)r^2dr$ represents the probability of finding a particle in a concentric spherical shell of thickness dr and radius r , then the total integral of the probability is given by the following (which runs through all radius $r \geq 0$):

$$1 = \int_0^\infty 4\pi\rho(r)r^2dr = \int_0^R 4\pi Ar^4dr = \frac{4\pi A}{5} r^5 \Big|_0^R = \frac{4\pi AR^5}{5} \quad (30)$$

Hence, for normalization to happen, $A = \frac{5}{4\pi R^5}$.

- (b) For the probability to find the particle in $\frac{R}{2} \leq r \leq R$, we get:

$$\mathbb{P}\left(\frac{R}{2} \leq r \leq R\right) = \int_{\frac{R}{2}}^R 4\pi\rho(r)r^2dr = \frac{4\pi A}{5} r^5 \Big|_{\frac{R}{2}}^R = \frac{4\pi}{5} \cdot \frac{5}{4\pi R^5} \left(1 - \frac{1}{2^5}\right) R^5 = \frac{31}{32} \quad (31)$$

- (c) To calculate $\langle r \rangle$, it's given by the following integral:

$$\langle r \rangle = \int_0^\infty r \cdot 4\pi\rho(r)r^2dr = \int_0^R 4\pi Ar^5dr = \frac{4\pi A}{6} r^6 \Big|_0^R = \frac{2\pi}{3} \cdot \frac{5}{4\pi R^5} R^6 = \frac{5R}{6} \quad (32)$$

□

Question 4 (Probability currents, Griffiths problem 1.9 and 1.14 modified). Let $P_{ab}(t)$ be the probability of finding a particle in the range $a < x < b$ at time t . Show that

$$\frac{dP_{ab}}{dt} = J(a, t) - J(b, t)$$

where

$$J(x, t) \equiv \frac{i\hbar}{2m} \left(\Psi \frac{\partial \Psi^*}{\partial x} - \Psi^* \frac{\partial \Psi}{\partial x} \right).$$

- (a) What are the units of $J(x, t)$? J is called the probability current because it tells you the rate at which probability is “flowing” past the point x . If $P_{ab}(t)$ is increasing, then more probability is flowing into the region at one end than flows out the other.
- (b) Consider a particle described by the wavefunction $\Psi(x, t) = Ae^{-amx^2/\hbar}e^{-iat}$. Find A for this wavefunction.
- (c) Compute the average velocity $\langle v \rangle = d\langle x \rangle/dt$ for this wavefunction.
- (d) Compute the probability current for this wavefunction.

Proof.

- (a) Since Probability itself has no unit, then J as a consequence of time derivative of probability, endows with the unit of $1/s$ (where s represents seconds for simplicity).
- (b) Given $\Psi(x, t) = Ae^{-amx^2/\hbar}e^{-iat}$, then it satisfies the following:

$$1 = \int_{-\infty}^{\infty} |\Psi(x, t)|^2 dx = \int_{-\infty}^{\infty} |A|^2 e^{-2amx^2/\hbar} dx = \int_{-\infty}^{\infty} |A|^2 e^{-(\sqrt{2am/\hbar}x)^2} dx \quad (33)$$

Takes substitution $u = \sqrt{2am/\hbar}x$, $du = \sqrt{2am/\hbar}dx$, we get:

$$1 = \frac{|A|^2}{\sqrt{2am/\hbar}} \int_{-\infty}^{\infty} e^{-u^2} du = |A|^2 \sqrt{\frac{\pi\hbar}{2am}} \quad (34)$$

Hence, we require $|A| = \left(\frac{\pi\hbar}{2am}\right)^{1/4}$.

- (c) First, we'll compute the expectation value $\langle x \rangle$ (with the same substitution as above):

$$\langle x \rangle = \int_{-\infty}^{\infty} x |\Psi(x, t)|^2 dx = |A|^2 \int_{-\infty}^{\infty} x e^{-(\sqrt{2am/\hbar}x)^2} dx = 0 \quad (35)$$

So, the average velocity $\langle v \rangle = \frac{d\langle x \rangle}{dt} = 0$.

- (d) Given that $\Psi^*(x, t) = A^* e^{-amx^2/\hbar} \cdot e^{iat}$, $\frac{d\Psi^*}{dx} = -A^* \frac{2amx}{\hbar} e^{-amx^2/\hbar} \cdot e^{iat}$, and $\frac{d\Psi}{dx} = -A \frac{2amx}{\hbar} e^{-amx^2/\hbar} \cdot e^{-iat}$. So, the Probability Current is given by:

$$J(x, t) = \frac{i\hbar}{2m} \left(-Ae^{-amx^2/\hbar}e^{-iat} \cdot A^* \frac{2amx}{\hbar} e^{-amx^2/\hbar}e^{iat} + A^* e^{-amx^2/\hbar}e^{iat} \cdot A \frac{2amx}{\hbar} e^{-amx^2/\hbar}e^{-iat} \right) \quad (36)$$

$$= 0 \quad (37)$$

□

Question 5 (More intuition about probability currents). *The wave function of a beam of free particles traveling in one dimension with velocity v is*

$$\Psi(x) = A \exp[ikx]$$

where $k = mv/\hbar$. (You might notice that this wave function is not strictly normalizable, but you can set aside this concern for now.) Argue that the probability current corresponding to this beam of particles has the form

$$J(x) = v \times \text{density of particles.}$$

Proof. Given the definition of probability current $J(x, t) := -\frac{i\hbar}{2m} \left(\Psi^* \frac{\partial \Psi}{\partial x} - \frac{\partial \Psi^*}{\partial x} \Psi \right)$, with $\Psi(x, t) = Ae^{ikx}$, then we have $\frac{\partial \Psi}{\partial x} = Aike^{ikx}$ and $\frac{\partial \Psi^*}{\partial x} = -Aike^{-ikx}$. Hence, J expands to the following form:

$$J(x, t) = -\frac{i\hbar}{2m} (Ae^{-ikx} \cdot Aike^{ikx} + Aike^{-ikx} \cdot Ae^{ikx}) = \frac{\hbar}{m} \cdot A^2 k = A^2 \frac{\hbar}{m} \cdot \frac{mv}{\hbar} = vA^2 \quad (38)$$

Now, given that $\int_a^b |\Psi|^2 dx$ is the probability of finding particles in the range of $[a, b]$ (assume $a < b$), we yield the following (recall that $|e^{ikx}| = 1$):

$$\int_a^b |\Psi|^2 dx = \int_a^b A^2 dx = A^2 \cdot (b - a) \quad (39)$$

Which, classically given a beam of a particle, the probability of finding a particle in the range of $[a, b]$ can be thought of as the density $\lambda \cdot (b - a)$ (which denotes the amount of particles within a certain range). Hence, here if $A^2 \cdot (b - a)$ serves the purpose for arbitrary interval $[a, b]$ (with $a < b$), then A^2 must serve as the density of the beam of free particles.

Hence, $J = vA^2$ is in terms of $v \times$ density of the particles.

□

Question 6 (Classical Particle in a Well (Griffiths Problem 1.11)). *Imagine a particle of mass m and energy E in a potential well $V(x)$, sliding frictionlessly back and forth between the classical turning points (call them a and b). Classically, the probability of finding the particle in the range dx (if, for example, you took a snapshot at a random time t) is equal to the fraction of the time T it takes to get from a to b that it spends in the interval dx :*

$$\rho(x)dx = \frac{dt}{T} = \frac{(dt/dx)dx}{T} = \frac{1}{v(x)T}dx$$

where $v(x)$ is the speed, and

$$T = \int_0^T dt = \int_a^b \frac{1}{v(x)} dx.$$

Thus

$$\rho(x) = \frac{1}{v(x)T}.$$

(a) Use conservation of energy to express $v(x)$ in terms of E and $V(x)$.

(b) As an example, find $\rho(x)$ for the simple harmonic oscillator, $V(x) = kx^2/2$. Plot $\rho(x)$ and check that it is correctly normalized.

(c) For the classical harmonic oscillator in part (b), find $\langle x \rangle$, $\langle x^2 \rangle$, and σ_x .

Proof.

- (a) Assume the conservation of energy, the particle has mass m , total energy E , and the potential is recorded by $V(x)$. Then, the kinetic energy of the particle is $K = \frac{1}{2}mv(x)^2 = E - V(x)$, hence $v(x) = \sqrt{\frac{2(E-V(x))}{m}}$.

- (b) Given $V(x) = \frac{kx^2}{2}$, then based on formula in part (a), $v(x) = \sqrt{\frac{2E-kx^2}{m}}$, hence the time T is:

$$T = \int_a^b \frac{1}{v(x)} dx = \int_a^b \frac{\sqrt{m}}{\sqrt{2E} \cdot \sqrt{1 - \left(\sqrt{\frac{k}{2E}}x\right)^2}} dx = \sqrt{\frac{m}{2E}} \cdot \sqrt{\frac{2E}{k}} \int_{\sqrt{\frac{k}{2E}}a}^{\sqrt{\frac{k}{2E}}b} \frac{1}{\sqrt{1-u^2}} du \quad (40)$$

$$= \sqrt{\frac{m}{k}} \arcsin(u) \Big|_{\sqrt{\frac{k}{2E}}a}^{\sqrt{\frac{k}{2E}}b} \quad (41)$$

Given that $V(x) = kx^2/2$, then the turning point (if assume $a < b$) satisfies $a = -b$ (since the turning points must have the same potential, so $ka^2/2 = kb^2/2$, or $a = \pm b$; with $a < b$, we have $a = -b$, if assuming $b > 0$). Hence, the above equation provides $T = 2\sqrt{\frac{m}{k}} \arcsin\left(\sqrt{\frac{k}{2E}}b\right)$.

Using this information, we have:

$$\rho(x) = \frac{1}{v(x)T} = \frac{1}{2\sqrt{m/k} \arcsin\left(\sqrt{\frac{k}{2E}}b\right)} \cdot \sqrt{\frac{m}{2E-kx^2}} = \frac{1}{2 \arcsin\left(\sqrt{\frac{k}{2E}}b\right)} \cdot \sqrt{\frac{k}{2E-kx^2}} \quad (42)$$

To check that $\rho(x)$ is normalized, consider its integral on $(-b, b)$ (the open interval between the two turning points), based on assumption one has $V(x) = kx^2/2 < E$ for all $x \in (-b, b)$, hence $2E - kx^2 > 0$ on this interval. Perform the integral, one gets:

$$\int_{-b}^b \rho(x) dx = \frac{1}{2 \arcsin\left(\sqrt{\frac{k}{2E}}b\right)} \int_{-b}^b \frac{\sqrt{k}}{\sqrt{2E} \cdot \sqrt{1 - \left(\sqrt{\frac{k}{2E}}x\right)^2}} dx \quad (43)$$

$$= \frac{1}{2 \arcsin\left(\sqrt{\frac{k}{2E}}b\right)} \cdot \sqrt{\frac{k}{2E}} \cdot \sqrt{\frac{2E}{k}} \int_{-\sqrt{\frac{k}{2E}}b}^{\sqrt{\frac{k}{2E}}b} \frac{1}{\sqrt{1-u^2}} du = \frac{1}{2 \arcsin\left(\sqrt{\frac{k}{2E}}b\right)} \cdot \arcsin(u) \Big|_{-\sqrt{\frac{k}{2E}}b}^{\sqrt{\frac{k}{2E}}b} \quad (44)$$

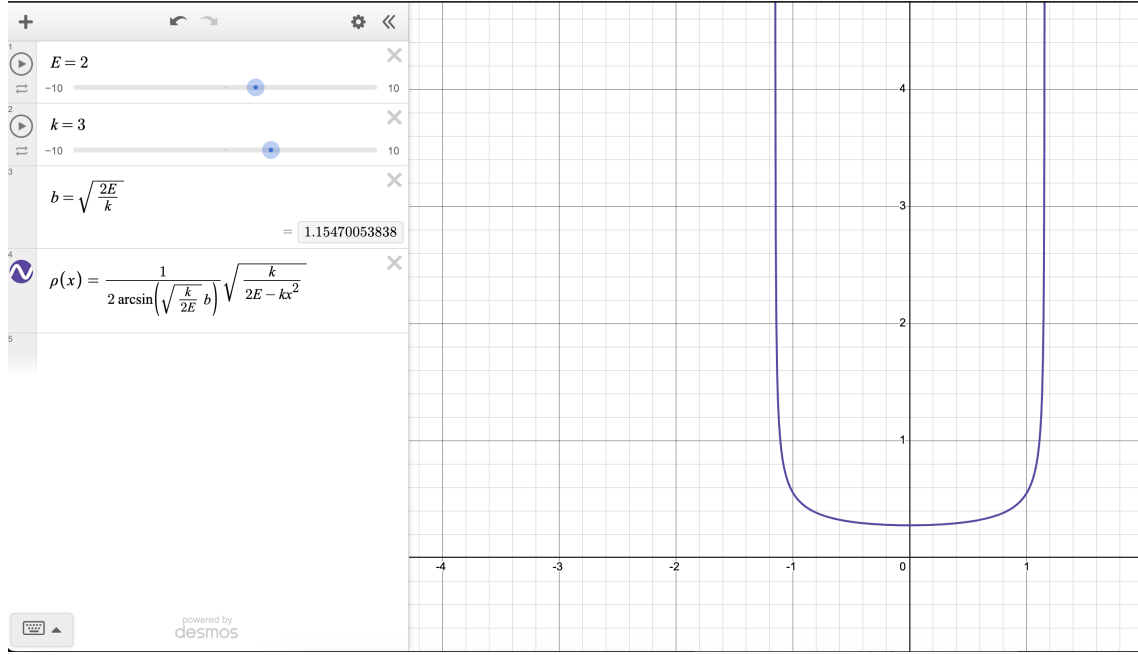
$$= \frac{1}{2 \arcsin\left(\sqrt{\frac{k}{2E}}b\right)} \cdot 2 \arcsin\left(\sqrt{\frac{k}{2E}}b\right) = 1 \quad (45)$$

Hence, $\rho(x)$ is perfectly normalized.

Finally, here is a graph of $\rho(x)$, with the constant given on the side:

- (c) To find $\langle x \rangle, \langle x^2 \rangle$, consider the following integrals (also note that at b , $V(b) = kb^2/2 = E$, so $b = \sqrt{\frac{2E}{k}}$):

$$\langle x \rangle = \int_{-b}^b x \rho(x) dx = \int_{-b}^b x \sqrt{\frac{k}{2E-kx^2}} dx = 0 \quad (46)$$



(The above integral is 0 since eventually it's an integral of odd function with integration bounds of opposite signs).

$$\langle x^2 \rangle = \int_{-b}^b x^2 \rho(x) dx = -\frac{1}{2k} \int_{-b}^b x \cdot (-2kx) \sqrt{\frac{k}{2E - kx^2}} dx \quad (47)$$

$$= -\frac{1}{2k} x \cdot 2\sqrt{k} \cdot \sqrt{2E - kx^2} \Big|_{-b}^b + \frac{1}{2k} \int_{-b}^b 2\sqrt{k} \cdot \sqrt{2E - kx^2} dx \quad (48)$$

$$= -\frac{2}{\sqrt{k}} b \cdot \sqrt{2E - kb^2} + \sqrt{\frac{2E}{k}} \int_{-b}^b \sqrt{1 - \left(\sqrt{\frac{k}{2E}} x\right)^2} dx \quad (49)$$

$$= -\frac{2b}{\sqrt{k}} \sqrt{2E - k \cdot \left(\sqrt{\frac{2E}{k}}\right)^2} + \frac{2E}{k} \int_{-\sqrt{\frac{k}{2E}} b}^{\sqrt{\frac{k}{2E}} b} \sqrt{1 - u^2} du \quad (50)$$

$$= 0 + \frac{2E}{k} \int_{-1}^1 \sqrt{1 - u^2} du = \frac{2E}{k} \cdot \frac{\pi}{2} = \frac{E\pi}{k} \quad (51)$$

With this value, then the final one $\sigma_x = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \sqrt{\frac{E\pi}{k}}$.

□

Question 7 (Unstable Particles (Griffiths Problem 1.17)). Suppose you wanted to describe an unstable particle, that spontaneously disintegrates with a “lifetime” τ . In that case, the total probability of finding the particle somewhere should not be constant, but should decrease at (say) an exponential rate:

$$P(t) = \int_{-\infty}^{\infty} |\Psi(x, t)|^2 dx = e^{-t/\tau}.$$

A crude way of achieving this result is as follows. In Griffiths Eq. 1.24 we tacitly assumed that V (the potential energy) is real. That is certainly reasonable, but it leads to the “conservation of probability” enshrined in Griffiths Eq. 1.27. What if we assign to V an imaginary part,

$$V = V_0 - i\Gamma$$

where V_0 is the true (real) potential energy and Γ is a positive real constant?

(a) Show that (in place of Griffiths Eq. 1.27, i.e. $\frac{d}{dt} \int_{-\infty}^{\infty} |\Psi(x, t)|^2 dx = 0$) we now get

$$\frac{dP}{dt} = -\frac{2\Gamma}{\hbar} P.$$

(b) Solve for $P(t)$, and find the lifetime of the particle in terms of Γ .

Proof.

- (a) Given that $|\Psi(x, t)|^2 = \Psi(x, t) \cdot \Psi^*(x, t)$, and $\frac{d}{dt} \int_{-\infty}^{\infty} |\Psi(x, t)|^2 dx = \int_{-\infty}^{\infty} \frac{\partial}{\partial t} |\Psi(x, t)|^2 dx$ (i.e. Leibniz Rule applies), together with the Schrodinger’s Equation $i\hbar \frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + V\Psi$, we get the following:

$$\frac{\partial}{\partial t} (|\Psi(x, t)|^2) = \frac{\partial}{\partial t} (\Psi\Psi^*) = \Psi^* \frac{\partial \Psi}{\partial t} + \Psi \frac{\partial \Psi^*}{\partial t} \quad (52)$$

$$= \Psi^* \left(\frac{i\hbar}{2m} \frac{\partial^2 \Psi}{\partial x^2} - \frac{i}{\hbar} V\Psi \right) + \Psi \left(-\frac{i\hbar}{2m} \frac{\partial^2 \Psi^*}{\partial x^2} + \frac{i}{\hbar} V^*\Psi^* \right) \quad (53)$$

$$= \frac{i\hbar}{2m} \left(\Psi^* \frac{\partial^2 \Psi}{\partial x^2} - \Psi \frac{\partial^2 \Psi^*}{\partial x^2} \right) - \frac{i}{\hbar} ((V_0 - i\Gamma) - (V_0 + i\Gamma)) |\Psi|^2 \quad (54)$$

$$= \frac{i\hbar}{2m} \frac{\partial}{\partial x} \left(\Psi^* \frac{\partial \Psi}{\partial x} - \Psi \frac{\partial \Psi^*}{\partial x} \right) - \frac{2\Gamma}{\hbar} |\Psi|^2 \quad (55)$$

Perform integration, we get:

$$\frac{dP}{dt} = \frac{d}{dt} \int_{-\infty}^{\infty} |\Psi|^2 dx = \int_{-\infty}^{\infty} \frac{\partial}{\partial t} |\Psi|^2 dx = \int_{-\infty}^{\infty} \frac{i\hbar}{2m} \left(\Psi^* \frac{\partial \Psi}{\partial x} - \Psi \frac{\partial \Psi^*}{\partial x} \right) dx - \frac{2\Gamma}{\hbar} \int_{-\infty}^{\infty} |\Psi|^2 dx \quad (56)$$

$$= -\frac{2\Gamma}{\hbar} P \quad (57)$$

The first integral is 0 because of the boundary condition of probability current, while the second integral (up to some constant) yields P .

- (b) To solve for $P(t)$, since $\frac{dP}{dt} = -\frac{2\Gamma}{\hbar} P$, we get $P(t) = Ke^{-2\Gamma t/\hbar}$ for some undetermined coefficient K . Which, if the lifetime τ of the particle is given as $P(t) = Ke^{-t/\tau}$, then here $\tau = \frac{\hbar}{2\Gamma}$.

□