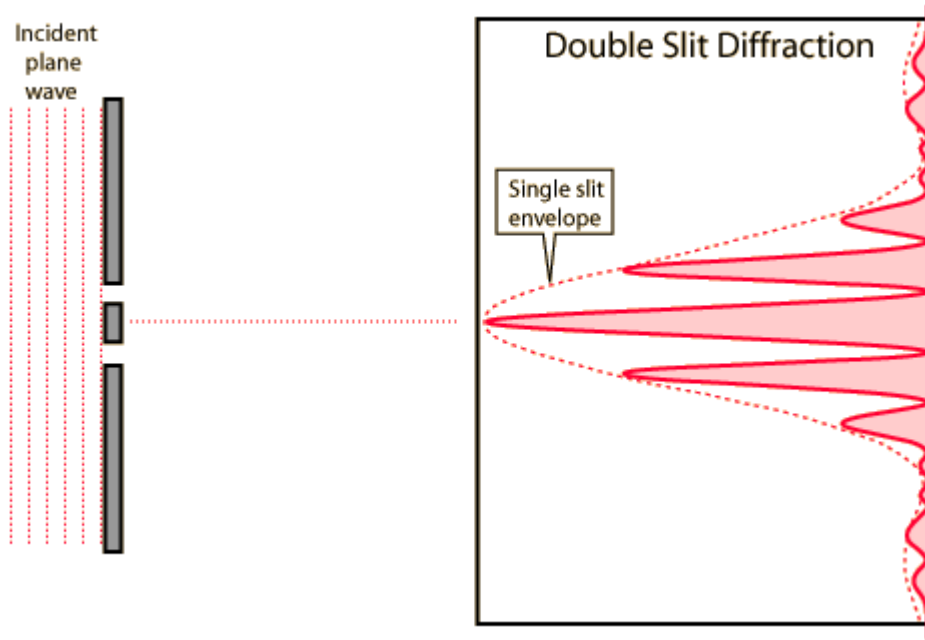


Young's Double Slit Experiment

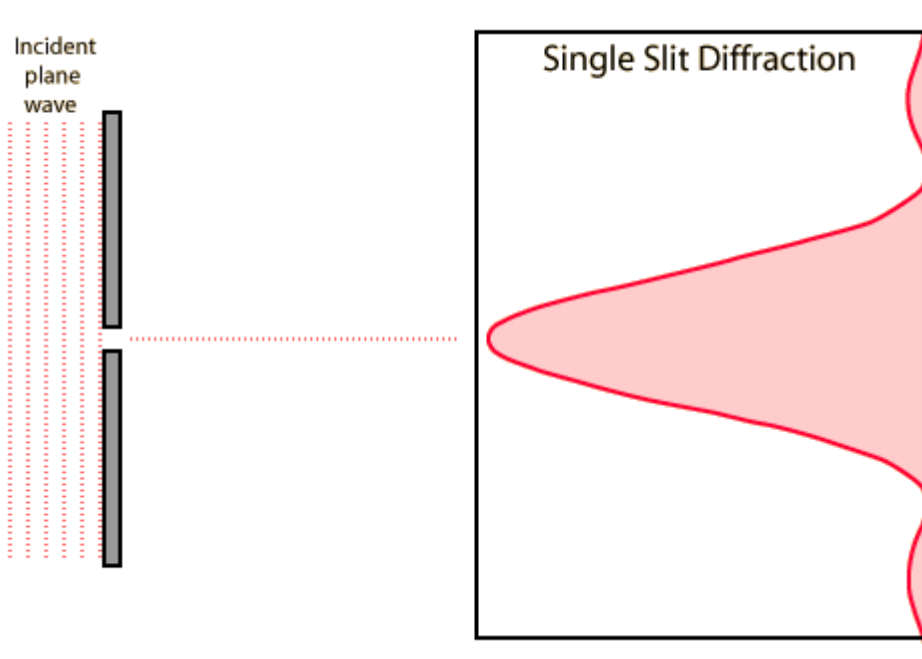
Bryan Turo

Particle Wave Duality of Photons



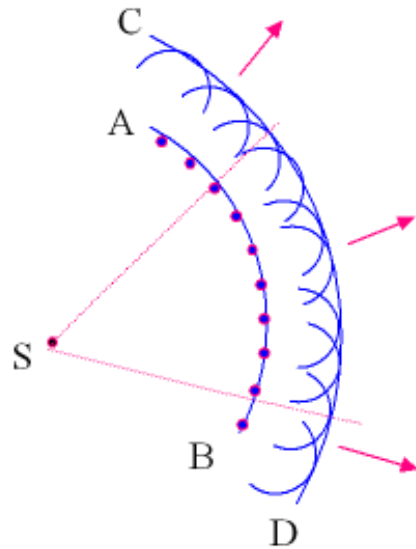
- During this time there were discussions for light being a particle or a wave
- Evidence for the wave like nature of light came from interference in Young's double slit experiment

Wave Properties



- Interference in the double slit is obvious; it's not in the single slit.
- We saw interference in single slits, this led to the development

Huygens' Principle



Huygens' Principle:

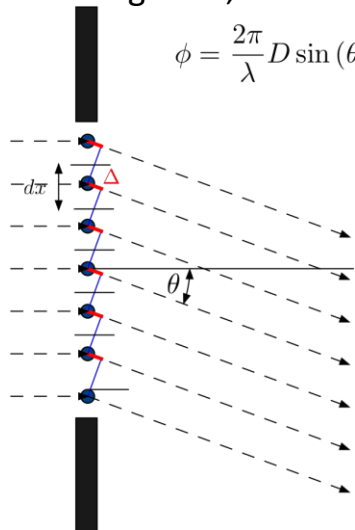
Each wavefront is the envelope of the wavelets. Each point on a wavefront acts as an independent source to generate wavelets for the next wavefront. AB and CD are two wavefronts.

- This works the same way for planar wavefronts
- You could think of this instead by imagining infinitesimal small slits

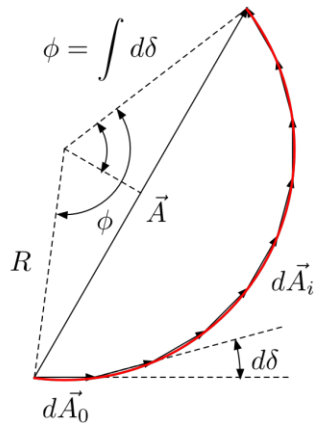
Fraunhofer diffraction

For single slit,

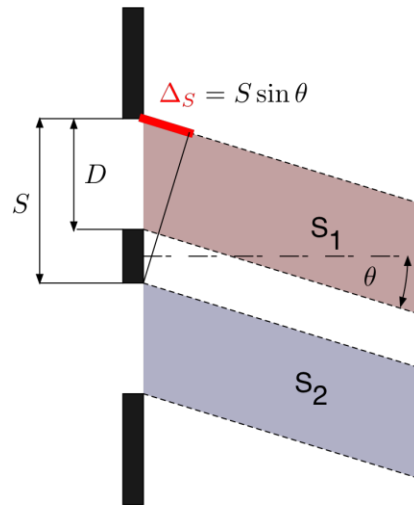
$$\phi = \frac{2\pi}{\lambda} D \sin(\theta)$$



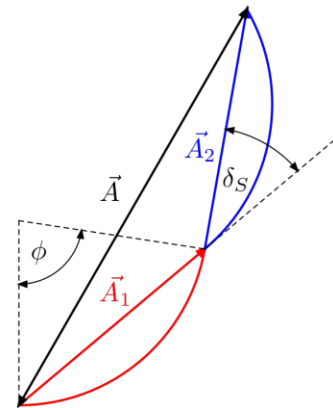
$$\Delta = dx \sin \theta$$



Under the assumption that the source is very far away... (This way the rays are paraxial)



A Theoretical Model for far-field diffraction



$$\vec{A} = \sum_i^N d\vec{A}_i$$

$$|d\vec{A}_i| = \frac{dx}{D} |\vec{A}_0|$$

As $N \Rightarrow \infty$

$$|\vec{A}| = 2R \sin(\phi/2)$$

$$\phi = \frac{2\pi}{\lambda} D \sin(\theta)$$

$$R = |\vec{A}_0|/\phi$$

Finally,

$$|\vec{A}| = |\vec{A}_0| \frac{\sin(\phi/2)}{\phi/2}$$

$$I(\theta) = I_0 \left(\frac{\sin(\phi/2)}{\phi/2} \right)^2$$

For double slit, add a phase shift,

$$|\vec{A}| = |\vec{A}_1 + \vec{A}_2| = 2|\vec{A}_0| \frac{\sin(\phi/2)}{\phi/2} \cos(\psi/2)$$

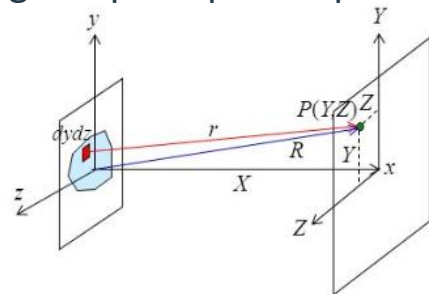
$$I(\theta) = |\vec{A}|^2$$

A quick introduction to Fourier optics

We can represent a plane wave propagating through a slit using Huygens' principle as spherical waves propagating from the source of infinitesimal slits ($dydz$ away)

Where $b-a$ is the total slit width

$$E(x, y, z) = \int_a^b E_0 \frac{e^{jkr}}{r} dz$$



Using some geometry you can find in “Introduction to Fourier optics” by Joseph W. Goodman

$$r = \sqrt{(Z - z)^2 + D^2}$$

Let's assume a paraxial approximation (small angle approximation) (far field) and other approximations to get the Fresnel integral

$$e^{-jkD} \frac{1}{D} \int_a^b e^{-\frac{1}{2}jk \frac{(Z-z)^2}{D}} dz$$

A continuation of Fourier optics

Making the Fraunhofer assumptions, $(b-a) \ll Z$ and $z \ll Z$ (Basically, we are at the far-field) we get,

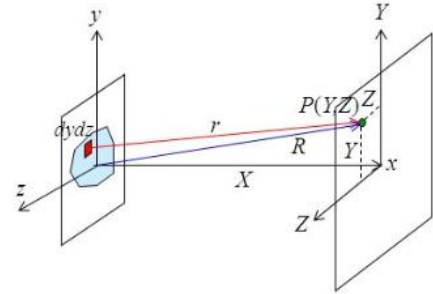
$$\frac{e^{-jkd} e^{-\frac{jkZ^2}{d}}}{d} \int_a^b e^{-jk_z z} dz$$

Now, this is exactly like a Fourier transform,

$$E(x, y, z) = E_0 \int_a^b e^{-jk_z z} dz$$

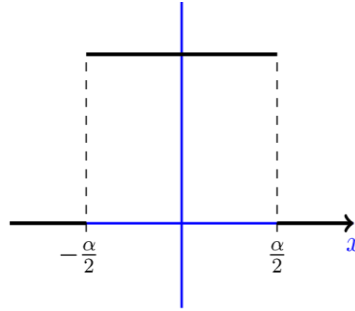
To get the expected intensity function of a slit we could use an Aperture function and replace the function to be,

$$E(x, y, z) = E_0 \int_{-\infty}^{\infty} g(x) e^{-jk_z z} dz = E_0 \mathcal{F}(g(x))$$



Example of Fourier Optics

Let's say we have an aperture function defining a single slit of width α , it should look something like this,



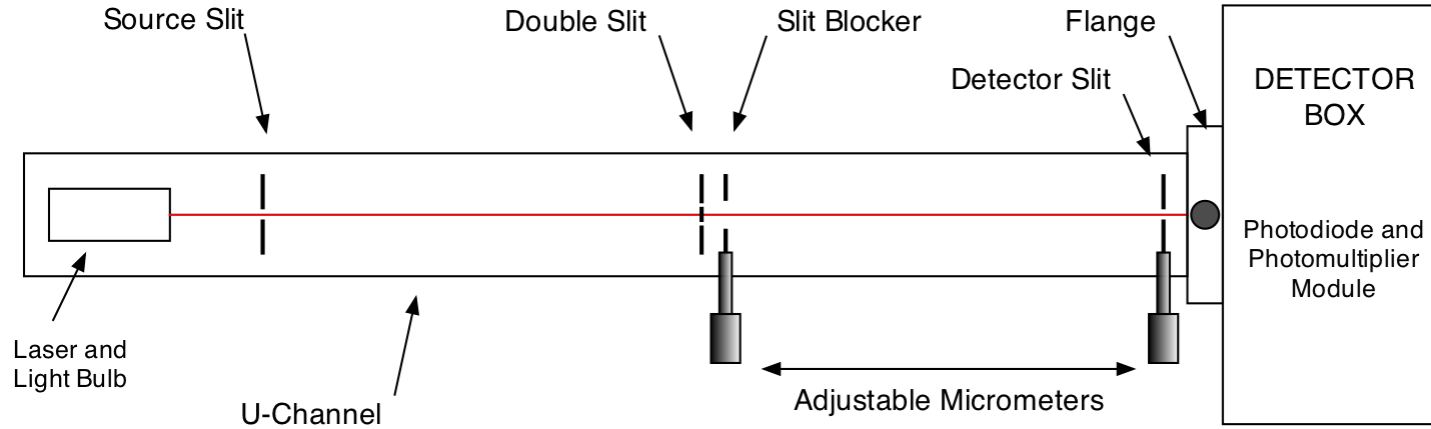
You will see that the Fourier transform of this function is

$$\text{sinc}\left(\frac{k_x \alpha}{2}\right)$$

Which means we get the following relationship when squaring the electric field,

$$I = I_0 \text{sinc}^2\left(\frac{k_x \alpha}{2}\right) = I_0 \left(\frac{\sin\left(\frac{k_x \alpha}{2}\right)}{\frac{k_x \alpha}{2}} \right)^2$$

Experimental Set-up



Specifications

Photomultiplier tube: Hamamatsu R 212

Preamplifier-Discriminator: Amptek A-111

Interference Filter: 546 nm, 10 nm FWHM

All Slit Widths: 0.09 mm

Double Slit Separations: 0.35, 0.40, 0.45 mm

Laser (Class II): 670 ± 20 nm < 1.0 mW

Universal Power Supply

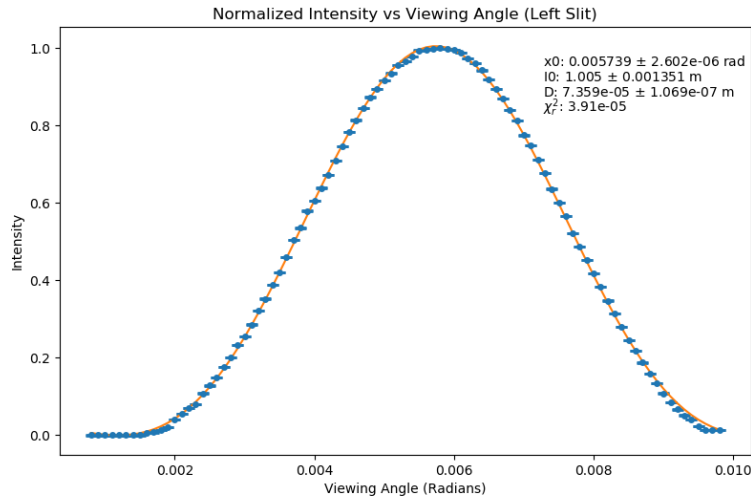
Input: 110 - 230 VAC, 50 - 60 Hz

Output: 15 V DC

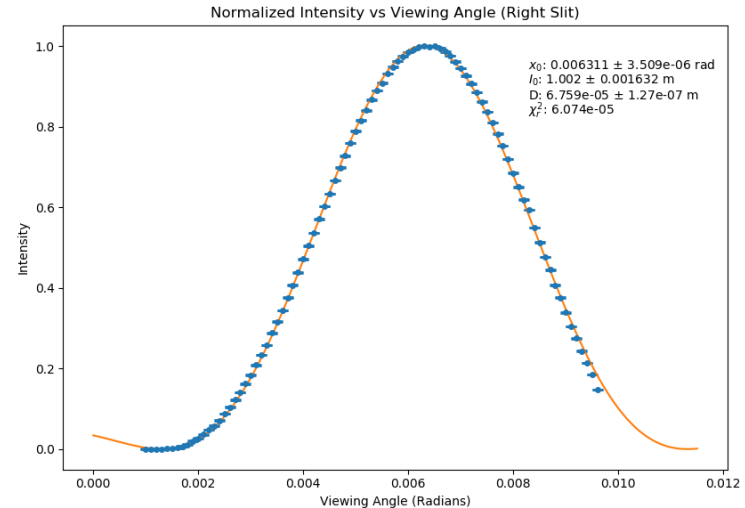
Optical Light Path: 1 meter

Overall Instrument Length: 52 inches

Single Slit Plots



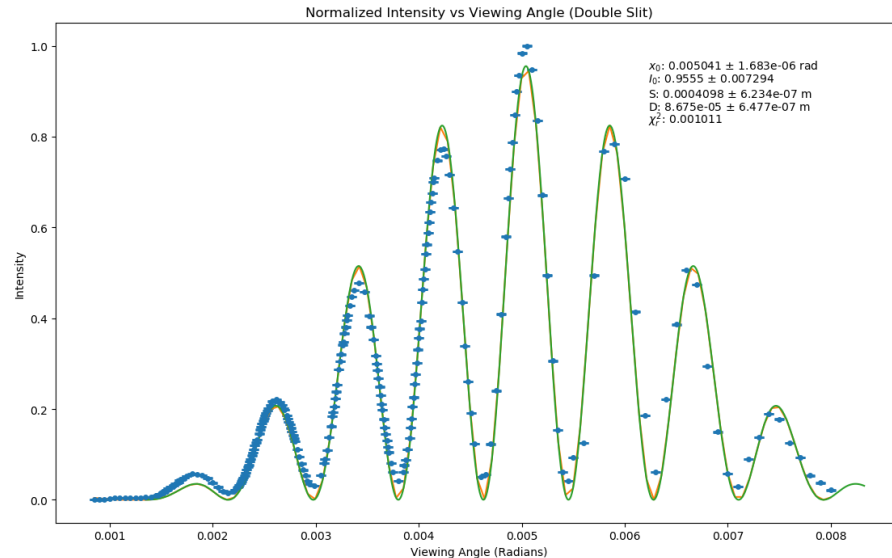
$$V_{max} = 0.802 \pm 0.0005 \text{ V}$$



$$V_{max} = 0.611 \pm 0.0005 \text{ V}$$

Double Slit Plot

- Background noise was subtracted from both single and double slit.



Data Fit

$$\theta \approx \frac{x - x_0}{L}$$

$$\phi = \frac{2\pi}{\lambda} D \sin(\theta)$$

Single Slit,

$$I(\theta) = I_0 \text{sinc} \left(\frac{\phi}{2} \right)^2$$

Double Slit,

$$\psi = \frac{2\pi}{\lambda} S \sin(\theta)$$

$$I(\theta) = I_0 \text{sinc}^2 \left(\frac{\phi}{2} \right) \cos^2 \left(\frac{\psi}{2} \right)$$

Parameters for Single Slit:

D : Slit Width

I_0 : Max intensity

x_0 : Location of Maximum (rad)

Parameters for Double Slit:

D

I_0

x_0

S : Slit Separation

Results

The weighted mean of the slit widths resulted in:

$$7.13 \cdot 10^{-5} \pm 8 \cdot 10^{-8} \text{ m}$$

The slit separation gotten was:

$$0.00041 \pm 6 \cdot 10^{-7} \text{ m}$$

Uncertainty Propagation

For the voltage we recorded an instrumental uncertainty of about $\pm 0.0005 \text{ V}$ for the double slit and about $\pm 0.005 \text{ V}$ for the single slit.

This can be propagated through the normalization function by applying the same ratios to the uncertainty

Uncertainty can be further analyzed in the x-direction

$$\theta = \frac{x - x_0}{L}$$

$$\frac{\partial \theta}{\partial x_0} = \frac{x}{L}$$

$$\frac{\partial \theta}{\partial x} = -\frac{x_0}{L}$$

$$\frac{\partial \theta}{\partial L} = \frac{-(x - x_0)}{L^2}$$

$$\sigma_\theta = \sqrt{\left(\frac{\partial \theta}{\partial x_0} \sigma_{x_0}\right)^2 + \left(\frac{\partial \theta}{\partial x} \sigma_x\right)^2 + \left(\frac{\partial \theta}{\partial L} \sigma_L\right)^2}$$

References

[“Introduction to Fourier Optics” Joseph W. Goodman, 2ed, Pg. 73](#)

http://wanda.fiu.edu/boeglinw/courses/Modern_lab_manual3/interference.html

<https://www.teachspin.com/two-slit>

http://www.pstcc.edu/departments/natural_behavioral_sciences/Web%20Physics/Chapter038.htm

<http://hyperphysics.phy-astr.gsu.edu/hbase/phyopt/slits.html>