

运动学

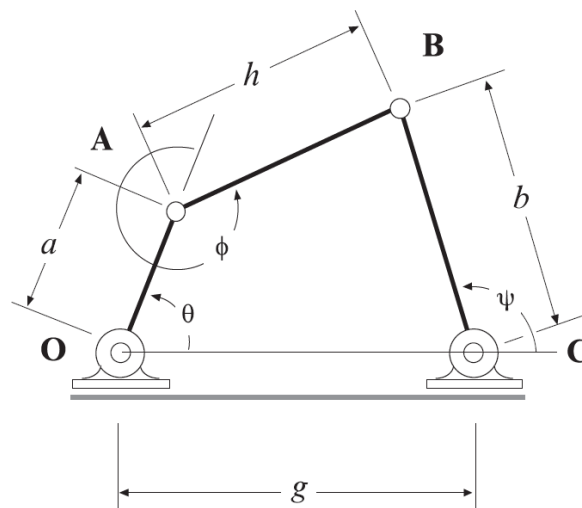
力学：如何实现运动

材料力学：考虑真实变形而非理想刚体情况

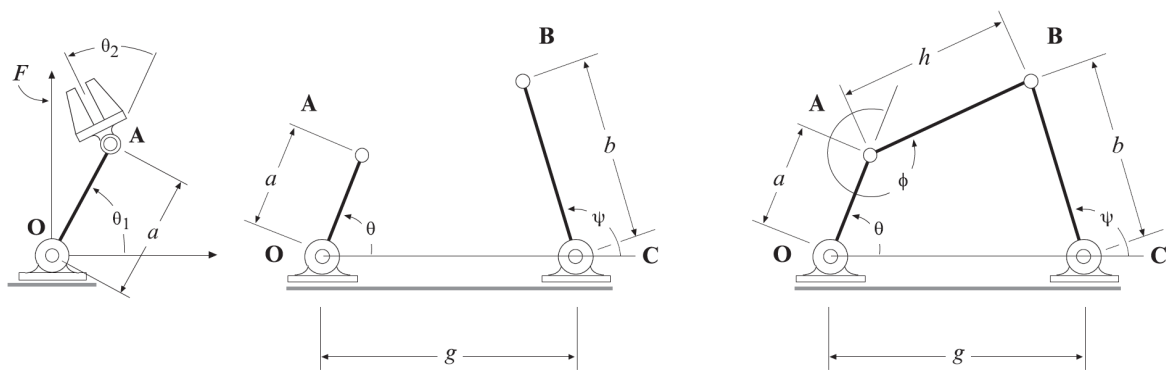
振动

力学

0. 回顾四连杆几何关系推导过程



• Forward kinematics $\theta \mapsto \psi$ (output link angle)



$$\mathbf{A} = \begin{pmatrix} a \cos \theta \\ a \sin \theta \end{pmatrix},$$

$$\mathbf{B} = \begin{pmatrix} g + b \cos \psi \\ b \sin \psi \end{pmatrix}$$

Plug in 

$$(\mathbf{B} - \mathbf{A}) \cdot (\mathbf{B} - \mathbf{A}) - h^2 = 0$$

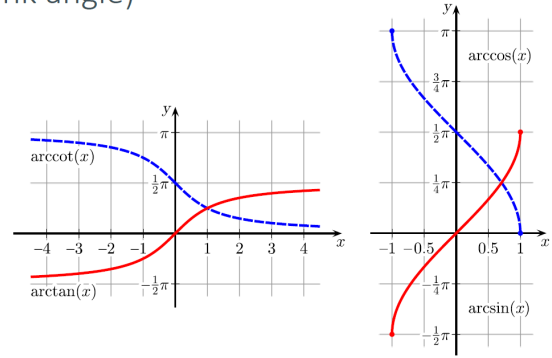
- Forward kinematics $\theta \mapsto \psi$ (output link angle)

$$\psi(\theta) = \arctan\left(\frac{B}{A}\right) \pm \arccos\left(\frac{C}{\sqrt{A^2 + B^2}}\right)$$

$$A(\theta) = 2ab \cos \theta - 2gb$$

$$B(\theta) = 2ab \sin \theta$$

$$C(\theta) = g^2 + b^2 + a^2 - h^2 - 2ag \cos \theta$$

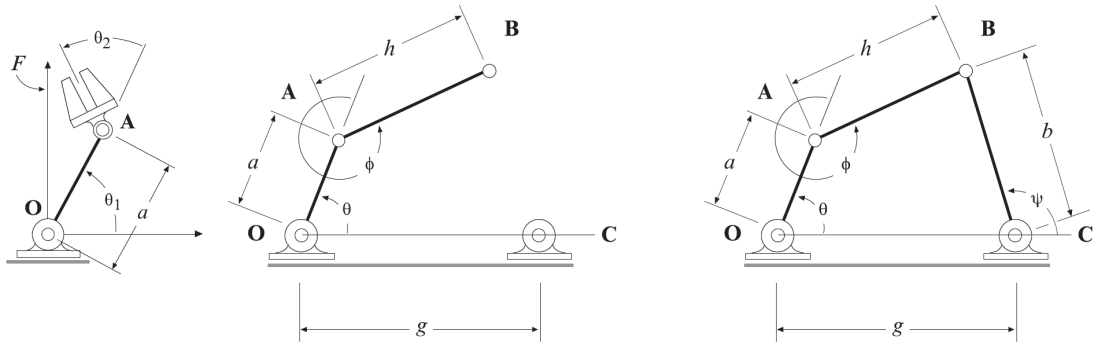


Question: solvable?



$$A^2 + B^2 - C^2 \geq 0$$

- Forward kinematics $\theta \mapsto \phi$ (coupler link angle)



$$\mathbf{A} = \begin{pmatrix} a \cos \theta \\ a \sin \theta \end{pmatrix}, \quad \mathbf{B} = \begin{pmatrix} a \cos \theta + h \cos(\theta + \phi) \\ a \sin \theta + h \sin(\theta + \phi) \end{pmatrix}$$

$$\begin{aligned} a \cos \theta + h \cos(\theta + \phi) &= g + b \cos \psi \\ a \sin \theta + h \sin(\theta + \phi) &= b \sin \psi \end{aligned}$$

$$\phi = \arctan\left(\frac{b \sin \psi - a \sin \theta}{g + b \cos \psi - a \cos \theta}\right) - \theta$$

- Input-output velocity ratio

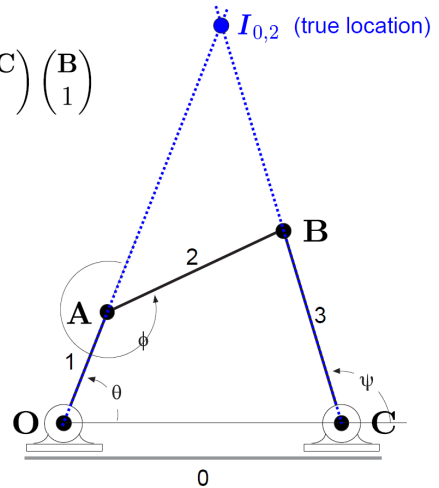
$$\left(\dot{\theta} \begin{pmatrix} \mathbf{J} & \mathbf{0} \\ \mathbf{0}^T & 0 \end{pmatrix} + \dot{\phi} \begin{pmatrix} \mathbf{J} & -\mathbf{J}\mathbf{A} \\ \mathbf{0}^T & 0 \end{pmatrix} \right) \begin{pmatrix} \mathbf{B} \\ 1 \end{pmatrix} = \dot{\psi} \begin{pmatrix} \mathbf{J} & -\mathbf{J}\mathbf{C} \\ \mathbf{0}^T & 0 \end{pmatrix} \begin{pmatrix} \mathbf{B} \\ 1 \end{pmatrix}$$

$$\dot{\theta} \mathbf{J} \mathbf{B} + \dot{\phi} \mathbf{J} (\mathbf{B} - \mathbf{A}) = \dot{\psi} \mathbf{J} (\mathbf{B} - \mathbf{C})$$

$$\dot{\theta} (\mathbf{B} - \mathbf{A})^T \mathbf{J} \mathbf{B} = \dot{\psi} (\mathbf{B} - \mathbf{A})^T \mathbf{J} (\mathbf{B} - \mathbf{C})$$

$$\dot{\theta} (\mathbf{B} - \mathbf{C})^T \mathbf{J} \mathbf{B} = -\dot{\phi} (\mathbf{B} - \mathbf{C})^T \mathbf{J} (\mathbf{B} - \mathbf{A})$$

$$\frac{\dot{\psi}}{\dot{\theta}} = \frac{(\mathbf{B} - \mathbf{A})^T \mathbf{J} \mathbf{B}}{(\mathbf{B} - \mathbf{A})^T \mathbf{J} (\mathbf{B} - \mathbf{C})}, \quad \frac{\dot{\phi}}{\dot{\theta}} = \frac{-(\mathbf{B} - \mathbf{C})^T \mathbf{J} \mathbf{B}}{(\mathbf{B} - \mathbf{C})^T \mathbf{J} (\mathbf{B} - \mathbf{A})}$$



忘记这个是怎么推导的了 ↑

1. Basic Concept

某质点 q $q = (x, y)$

其运动轨迹 $x^2 + y^2 = 1$

其瞬时速度 $\dot{q} = (\dot{x}, \dot{y})$

速度方程 $\frac{d}{dt}(x^2 + y^2) = 2x\dot{x} + 2y\dot{y} = 0$ x 径向 \dot{x} 切向 两者垂直, 点乘结果等于 0

当 x 确定时 (取消一个自由度) $y = \pm\sqrt{1-x^2}$

$$\dot{y} = \mp \frac{x\dot{x}}{\sqrt{1-x^2}}$$

本处做法问题: 1 正负号选哪个 2 $x=1$ 时方程无意义

当 y 确定时情况类似

解决方式: 不必要使用笛卡尔坐标系, 可以使用极坐标

极坐标与笛卡尔坐标的变换方式:

$$x = \cos\theta \quad y = \sin\theta$$

$$\dot{x} = -\sin\theta\dot{\theta} \quad \dot{y} = \cos\theta\dot{\theta}$$

质点可以存在的空间: configuration space 位形空间 $C = \{\theta | \theta \in [0, 2\pi)\}$ 开区间原因: 圆环上 0 与 2π 为同一个点

与质点本身位置 (x, y) 无关的坐标形式叫做广义坐标 generalized coordinates

2. Kinetic energy

牛顿第二定律 $f = \ddot{q}$

定义动能 $T = \frac{1}{2}m\dot{q}^T\dot{q}$

动能函数反推牛顿方程

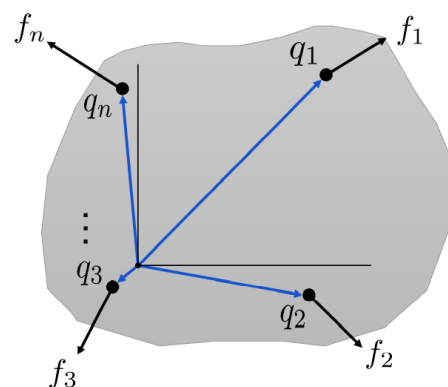
Kinetic energy becomes very useful when you have a whole system of particles (for rigid body, replace summation with integration)

$$T = \frac{1}{2} \sum_{i=1}^n m_i \dot{q}_i^T \dot{q}_i$$

Question: do we have the following equation?

$$\frac{d}{dt} \left(\frac{\partial T}{\partial \dot{q}_i} \right) = \frac{d}{dt} (m_i \dot{q}_i) = f_i, \quad i = 1, \dots, n$$

(Hint: think about the force that keeps a particle on the circle)

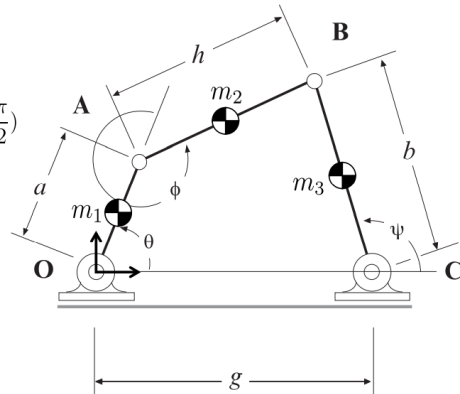


$$\frac{d}{dt} \left(\frac{\partial T}{\partial \dot{q}} \right) = \frac{d}{dt} (m\dot{q}) = f \quad \left(\text{Given } \dot{q} = \begin{bmatrix} \dot{x} \\ \dot{y} \\ \dot{z} \end{bmatrix}, \frac{\partial T}{\partial \dot{q}} = \begin{bmatrix} \partial T / \partial \dot{x} \\ \partial T / \partial \dot{y} \\ \partial T / \partial \dot{z} \end{bmatrix} \right)$$

问题：有内力

Assume center of mass of the links are located at the mid-point of the links, compute the kinetic energy of the 4-bar linkage as a function of θ and $\dot{\theta}$ (assume $\phi(\theta), \psi(\theta), \dot{\phi}(\theta, \dot{\theta}), \dot{\psi}(\theta, \dot{\theta})$ are given)

$$\frac{\dot{\psi}}{\dot{\theta}} = \frac{(\mathbf{B} - \mathbf{A})^T J \mathbf{B}}{(\mathbf{B} - \mathbf{A})^T J (\mathbf{B} - \mathbf{C})}, \quad \frac{\dot{\phi}}{\dot{\theta}} = \frac{-(\mathbf{B} - \mathbf{C})^T J \mathbf{B}}{(\mathbf{B} - \mathbf{C})^T J (\mathbf{B} - \mathbf{A})} \quad J = R\left(\frac{\pi}{2}\right)$$



求三联杆系统的动能

1 直接求导 2 先算出角速度之间的关系再计算

Consider the infinitesimal work done by f :

$$dW = f^T dq$$

If there exists a **work function** $U(q)$ such that:

$$f = \frac{\partial U(q)}{\partial q}$$

we have:

$$\int_{t_0}^{t_f} dW = \int_{t_0}^{t_f} \frac{\partial U}{\partial q} dq = U(q(t_f)) - U(q(t_0))$$

Putting the kinetic energy and work function together:

$$\int_{t_0}^{t_f} m\ddot{q}^T dq = \int_{t_0}^{t_f} \frac{d}{dt} \left(\frac{1}{2} m \dot{q}^T \dot{q} \right) dt = \int_{t_0}^{t_f} f^T dq = \int_{t_0}^{t_f} \left(\frac{\partial U}{\partial q} \right)^T dq$$



$$T(t_f) - T(t_0) = U(t_f) - U(t_0)$$

or

$$T(t_f) - U(t_f) = T(t_0) - U(t_0)$$

We usually define **potential energy** $V(q) = -U(q)$



$$E(t) \triangleq T(t) + V(t) = \text{constant}$$

The total energy E remains constant. $f = -\partial V / \partial q$ is therefore called a **conservative force**.

之前使用动能表示牛顿定律有问题，没表达出内力。应使用动能+势能表示牛顿定律。其中势能对应的为保守力。

3. 虚功原理

约束力

在切向方向上进行一小段微小位移（任取），如果系统处于平衡状态（合力为 0），合力在该微小位移上做的功为 0。如果该功为 0，也能证明力平衡。

虚位移：1 沿切线 2 不真实，用于测试系统是否平衡 3 发生位移时，时间停止 4 性质上类似速度

The **virtual displacements** $\delta q_1, \delta q_2$ act like dq_1, dq_2 , but occurs when the system freezes in time. Recall:

$$\dot{q}_1 = \frac{dq_1}{dt}, \dot{q}_2 = \frac{dq_2}{dt}$$

$$(q_1 - q_2)^T (q_1 - q_2) = h^2 \Rightarrow (q_1 - q_2)^T (\dot{q}_1 - \dot{q}_2) = 0$$

Similarly, we have:

$$(q_1 - q_2)^T (\delta q_1 - \delta q_2) = 0$$

$$\frac{d}{dt}(u^T v) = \dot{u}^T v + u^T \dot{v}$$

虚位移可以是关于时间的函数，不同时间下虚位移不同。

两个重要运算性质:

$$\delta q(t + dt) = \delta q(t) + \delta \dot{q}(t) dt \Rightarrow$$

$$\frac{d}{dt}(\delta q(t)) = \delta \dot{q}(t) \text{ or } d\delta q(t) = \delta dq(t)$$

Similarly, δ commutes with $\int dt$.

$$\delta \int_{t_0}^{t_f} q(t) dt = \int_{t_0}^{t_f} \delta q(t) dt$$

$$\begin{aligned} \delta W &\triangleq (f_1 + f_t + f_{1,g})^T \delta q_1 + (f_2 - f_t + f_{2,g})^T \delta q_2 \\ &= (f_1 + f_{1,g})^T \delta q_1 + (f_2 + f_{2,g})^T \delta q_2 + f_t^T (\delta q_1 - \delta q_2) \\ &= 0 \end{aligned}$$

Recall we also have:

$$(q_1 - q_2)^T (\delta q_1 - \delta q_2) = 0$$

Since f_t is parallel to the vector $q_1 - q_2$, we have:

$$f_t^T (\delta q_1 - \delta q_2) = 0$$



$$\delta W = (f_1 + f_{1,g})^T \delta q_1 + (f_2 + f_{2,g})^T \delta q_2 = 0$$

约束力不做功，对虚位移也不做功，所以使用虚功原理可以不考虑约束力。

两个虚位移线性相关

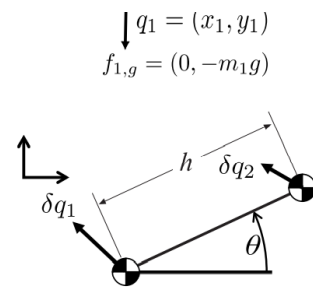
However, if we notice:

$$\delta q_2 = \delta q_1 + h v \delta \theta, \quad v = \begin{bmatrix} -\sin \theta \\ \cos \theta \end{bmatrix}$$

$$\delta W = (f_1 + f_{1,g} + f_2 + f_{2,g})^T \delta q_1 + h(f_2 + f_{2,g})^T v \delta \theta = 0$$

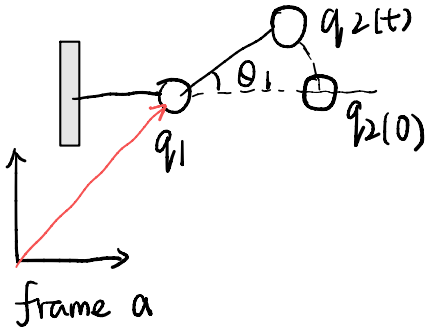
Since $\delta q_1, \delta \theta$ are linear independent, we have:

$$\begin{aligned} f_1 + f_{1,g} + f_2 + f_{2,g} &= 0 \\ (f_2 + f_{2,g})^T v &= 0 \end{aligned}$$



虚功

1. move a point $q_2(0) \rightarrow q_2(t)$

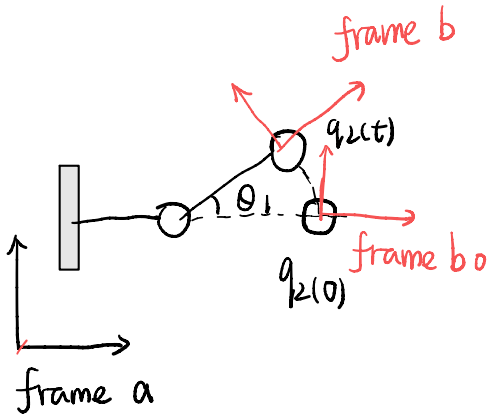


$$q_2(t) = R(q_1, \theta_1) q_2(0)$$

fixed axis rotation
(about axis q_1 for θ_1)

$$R(q_1, \theta_1) = e^{\hat{U}t} = \begin{bmatrix} e^{\hat{\omega}t} & (I - e^{\hat{\omega}t})p \\ 0 & 1 \end{bmatrix} \xrightarrow[p=q_{1a}]{\hat{\omega} = \hat{z} \cdot ||\omega||} \begin{bmatrix} e^{\hat{z}\theta_1} (I - e^{\hat{z}\theta_1})q_1 & \\ 0 & 1 \end{bmatrix}$$

2. move a frame



frame a \rightarrow frame b0 : g_{abo}

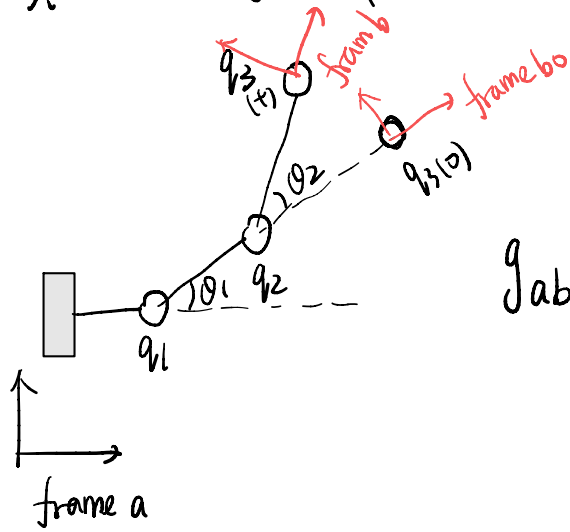
$$\begin{bmatrix} q_{2a}(0) \\ 1 \end{bmatrix} = g_{abo} \begin{bmatrix} q_{2b0} \\ 1 \end{bmatrix}$$

frame a \rightarrow frame b : g_{ab}

$$\begin{bmatrix} q_{2a}(t) \\ 1 \end{bmatrix} = g_{ab} \begin{bmatrix} q_{2b} \\ 1 \end{bmatrix}$$

$$g_{ab} = R(q_1, \theta_1) g_{abo}$$

3. move a frame (2 link robot end)



$$g_{ab} = R(q_{10}, \theta_1) R(q_{20}, \theta_2) \cdot g_{abo}$$

4. g_{ab} (定轴转动情况, $\dot{p}=0$)

$$g_{ab} = \prod_{i=1}^n e^{\hat{E}_i \theta_i} \cdot g_{abo}$$

$$\hat{E}_i = \begin{bmatrix} \hat{w}_{i0} & q_{i0} \times w_{i0} \\ 0 & 0 \end{bmatrix}$$

w_{i0} 为单轴转动所用速度

一般情况:

$$\hat{V} = \begin{bmatrix} \hat{\omega} = \dot{R}R^T & \dot{p}(t) - \hat{\omega} \cdot p(t) \\ 0 & 0 \end{bmatrix}$$

