

Week 8

Jacobian and Velocities

MRes in Medical Robotics and Instrumentation

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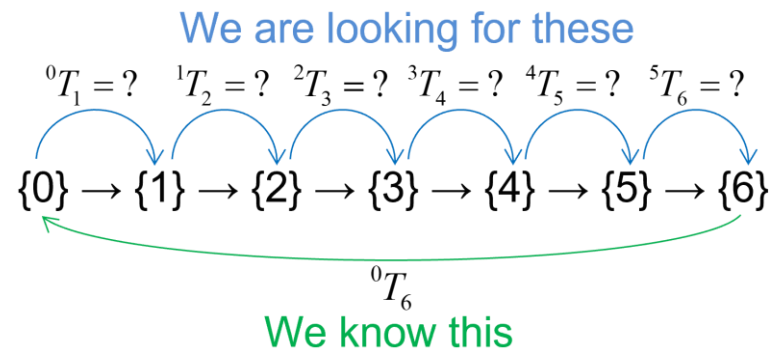
Contents of the lecture

- Revision of main concepts from last week
- Velocity Representation
- Linear and Angular Velocities
- Interlink Velocity Propagation
- Use of Jacobian for velocities mapping

- John J Craig, Introduction to robotics – mechanics and control, 3rd Edition, Pearson Education International, London, ISBN 0-13-123629-6.
- Reza N Jazar, Theory of applied robotics, 2nd Edition, Springer, ISBN 978-1-4419-1749-2.

Revision: Inverse Kinematics (IK) problem

- Inverse kinematics determine the joint configurations of a robot model to achieve a desired end-effect position



The problem of inverse kinematics is **nonlinear**;
the change of the output is not proportional to the change of the input

Revision: Algebraic solution of Inverse Kinematics

Inverse kinematics problem to solve

- End-effector link 3 is located at (x_a, y_a) with an orientation of β relative to frame $\{0\}$
- Problem: Find $\theta_i, (i=1, 2, 3)$?

$${}^0T_3 = \begin{bmatrix} c\beta & -s\beta & 0 & x_a \\ s\beta & c\beta & 0 & y_a \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \quad {}^0T_3 = \begin{bmatrix} c\theta_{123} & -s\theta_{123} & 0 & l_1c\theta_1 + l_2c\theta_{12} \\ s\theta_{123} & c\theta_{123} & 0 & l_1s\theta_1 + l_2s\theta_{12} \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

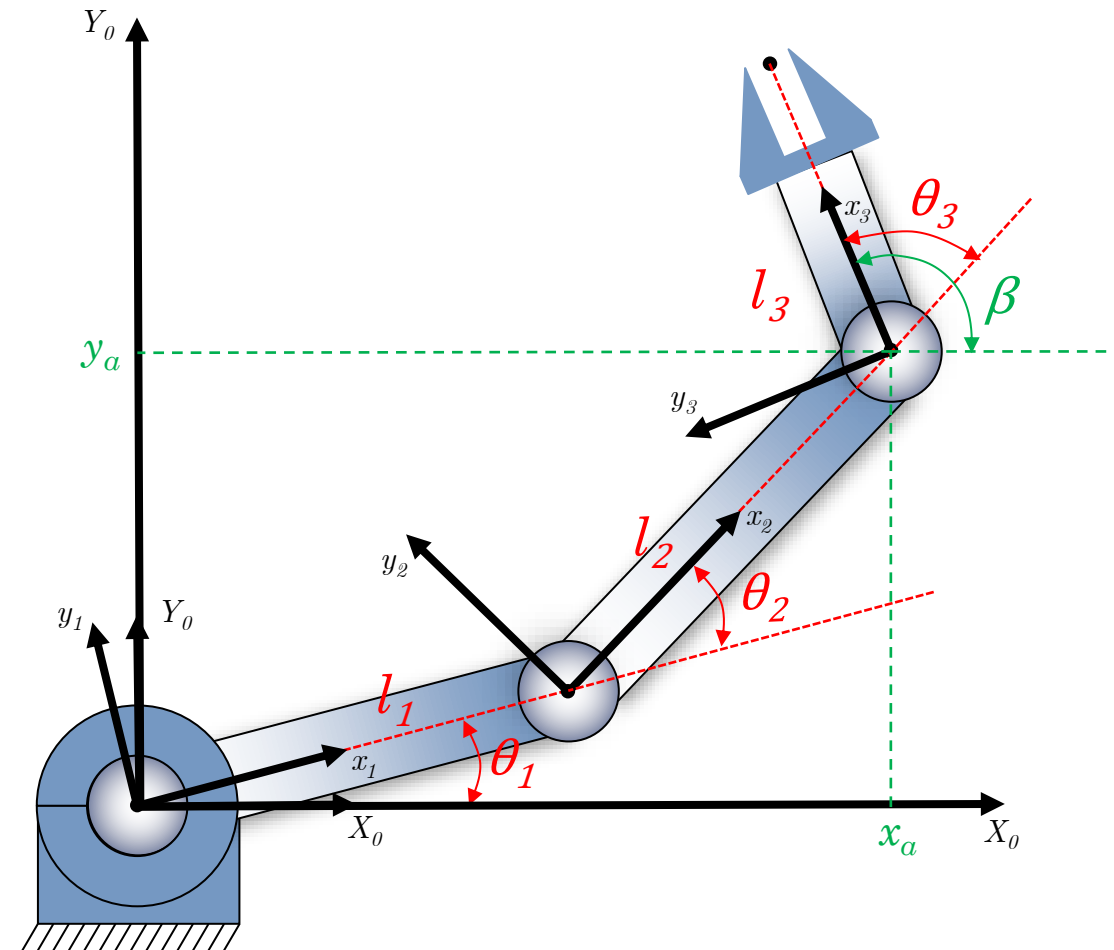
Expressed in
task/cartesian space

Expressed in
joint space

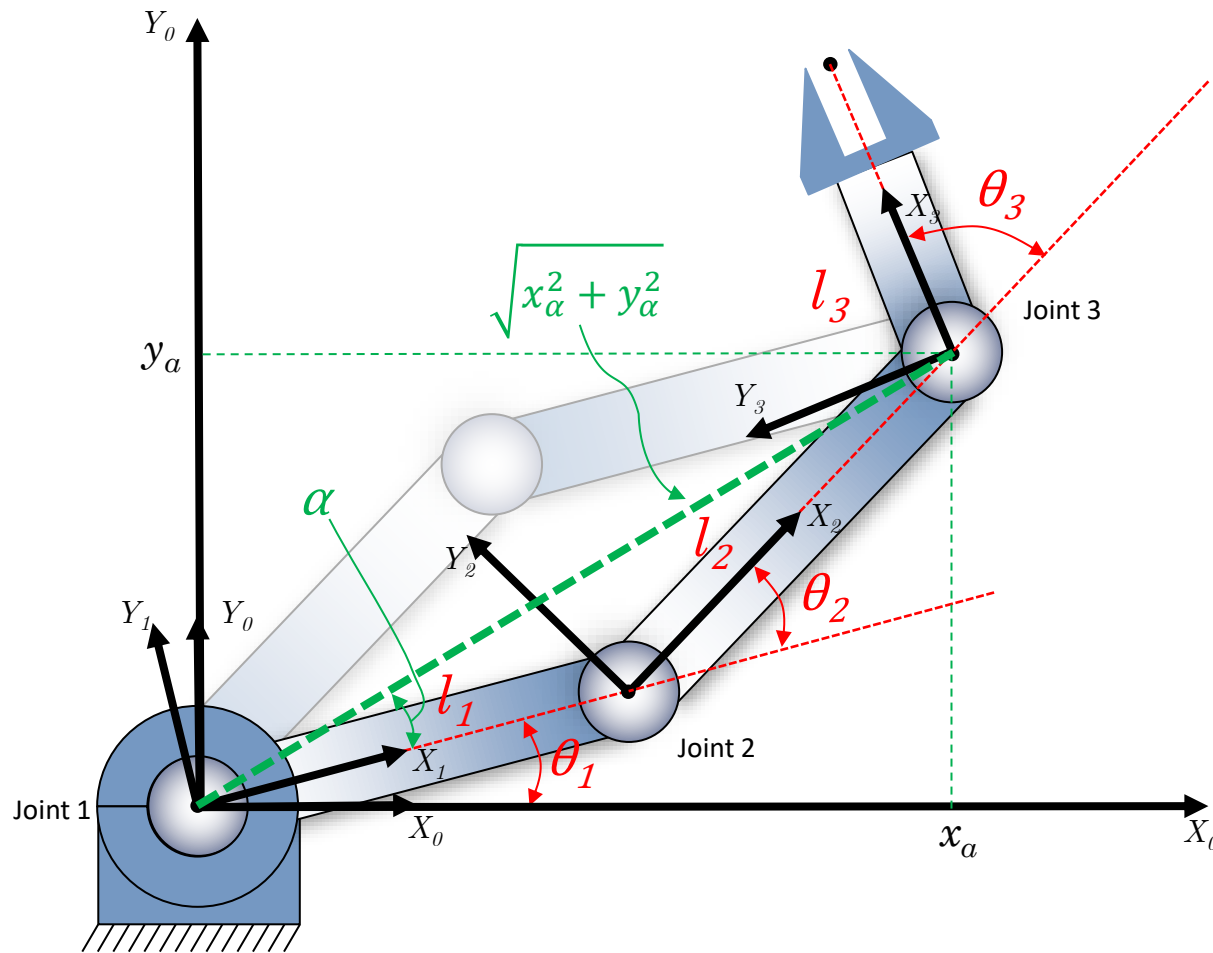
Equate above:

$$\begin{cases} x_a = l_1c\theta_1 + l_2c\theta_{12} \\ y_a = l_1s\theta_1 + l_2s\theta_{12} \\ c\beta = c\theta_{123} \\ s\beta = s\theta_{123} \end{cases}$$

*Kinematic
equations
(non-linear)*



Revision: Geometric solution of Inverse Kinematics



$$\begin{aligned} x_a^2 + y_a^2 &= l_1^2 + l_2^2 - 2l_1l_2c(180 - \theta_2) \\ &= l_1^2 + l_2^2 + 2l_1l_2c\theta_2 \end{aligned}$$

Cosine rule

$$\theta_2 = \pm \cos^{-1} \left[(x_a^2 + y_a^2) - (l_1^2 + l_2^2) / 2l_1l_2 \right]$$

$$\text{Condition: } \left| \left[(x_a^2 + y_a^2) - (l_1^2 + l_2^2) \right] / 2l_1l_2 \right| \leq 1$$

Which basically means that $\sqrt{x_a^2 + y_a^2} \leq l_1 + l_2$

$$l_2^2 = x_a^2 + y_a^2 + l_1^2 - 2l_1\sqrt{x_a^2 + y_a^2}c\alpha$$

Cosine rule in α

$$c\alpha = \frac{x_a^2 + y_a^2 + l_1^2 - l_2^2}{2l_1\sqrt{x_a^2 + y_a^2}}$$

$$\theta_1 = \text{atan2}(y_a, x_a) \pm \alpha$$

$$\theta_1 + \theta_2 + \theta_3 = \beta, \text{ solve for } \theta_3$$

Revision: Newton-Raphson Method

- Solve (find the roots of) $p(\theta) = a$ or $p(\theta) - a = 0$

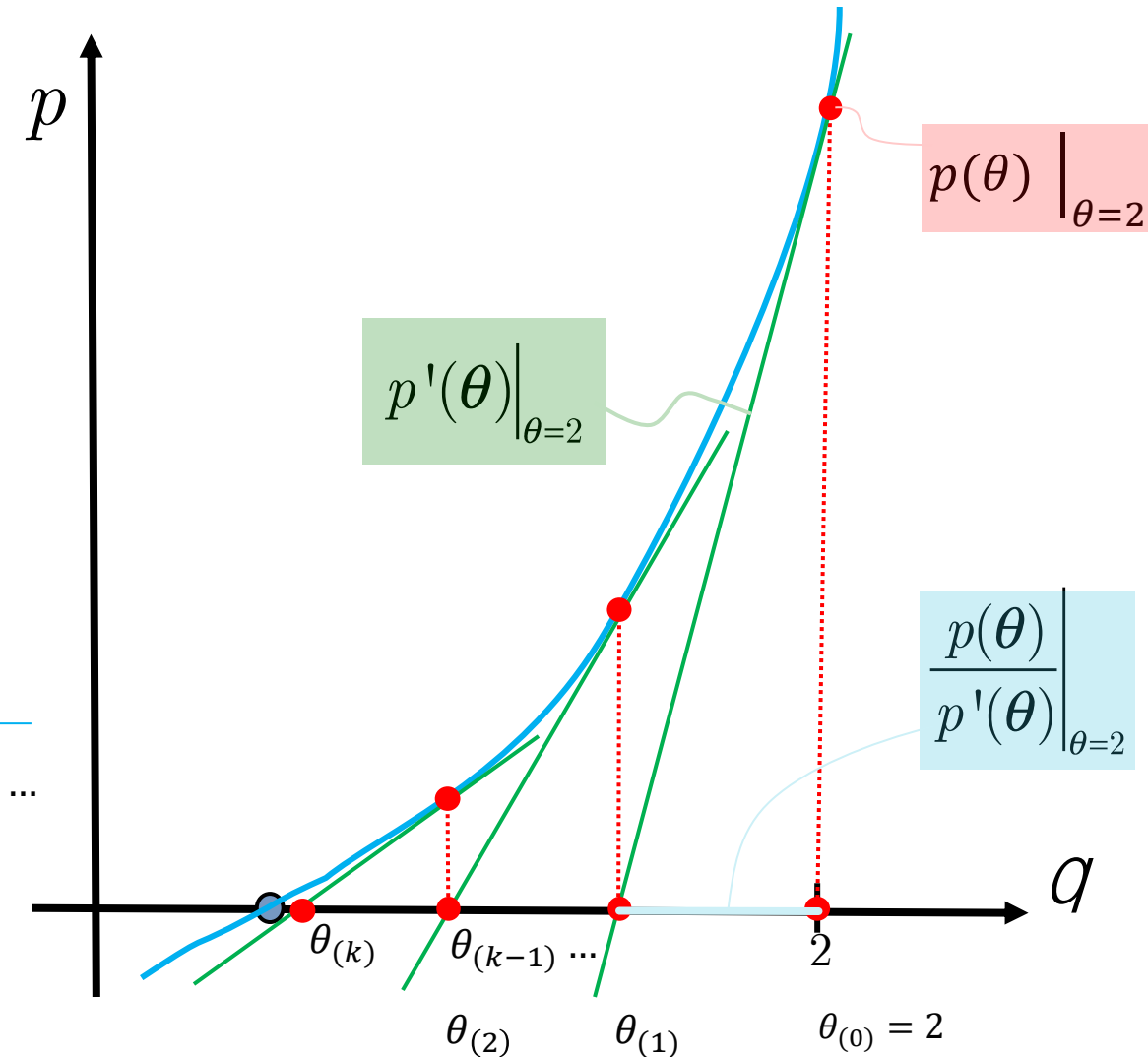
Iterative solution steps

1. Start with an initial value of $\theta_{(0)}$ (**a –hopefully good – guess**)
2. Evaluate p for this guess $p(\theta_{(0)})$
3. Iteratively refine the initial guess of θ values, to approach the true solution for p

$\theta_{(k)}$ represents the result at iteration step k .

$\overbrace{\theta_{(0)} \rightarrow p(\theta_{(0)}) \rightarrow p'(\theta_{(0)})}^{k=0, \theta_0 \text{ is a guess}} \rightarrow \overbrace{\theta_{(1)} \rightarrow p(\theta_{(1)}) \rightarrow p'(\theta_{(1)})}^{k=1} \rightarrow \overbrace{\theta_{(2)} \rightarrow p(\theta_{(2)})}^{k=2}, \dots$

$$\theta_{(k)} = \theta_{(k-1)} + p'(\theta_{(k-1)})^{-1} (a - p(\theta_{(k-1)}))$$



Revision: Iterative numerical solution of Inverse Kinematics

- For a given pose \mathbf{P} of the end effector, we can use the general vector function below to represent the joint kinematics:

$$\mathbf{P} = \mathbf{Q}(\theta_1, \theta_2, \dots, \theta_N) = \mathbf{Q}(\vec{\theta})$$

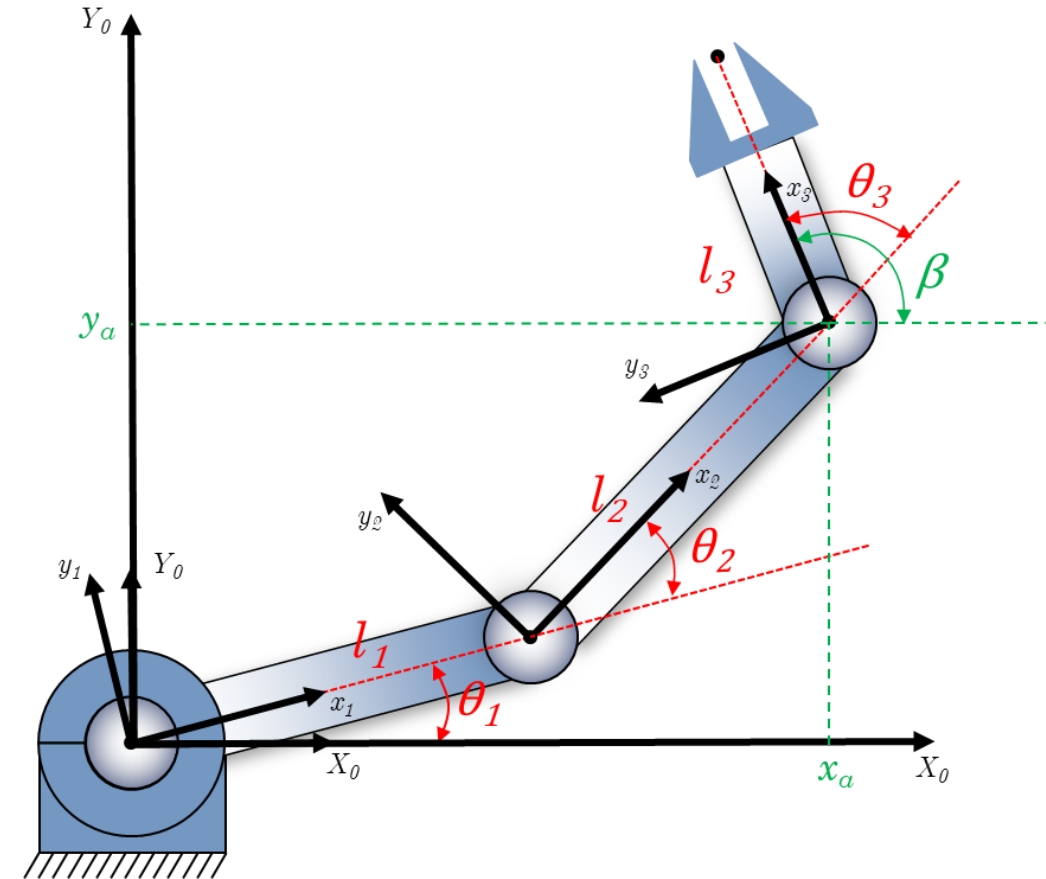
- The inverse kinematics problem then would be:

$$\text{Given } \mathbf{P} = \mathbf{Q}(\vec{\theta}), \text{ find } \vec{q}$$

e.g., here, given $\mathbf{P} = \begin{bmatrix} x_a \\ y_a \\ \beta \end{bmatrix} = \mathbf{Q}(\vec{\theta}_1, \vec{\theta}_2, \vec{\theta}_3)$, find $\vec{\theta}_1, \vec{\theta}_2, \vec{\theta}_3$

$$\theta_{(k)} = \theta_{(k-1)} + p'(\theta_{(k-1)})^{-1} (a - p(\theta_{(k-1)}))$$

$$\vec{\theta}_{(k)} = \vec{\theta}_{(k-1)} + \mathbf{J}^{-1}(\vec{\theta}_{(k-1)}) (\mathbf{P} - \mathbf{Q}(\vec{\theta}_{(k-1)})) \quad \text{where} \quad \mathbf{J}(\vec{\theta}_{(k-1)}) = \frac{\partial \mathbf{Q}(\vec{\theta}_{(k-1)})}{\partial \vec{\theta}_{(k-1)}}$$



Revision: General expression of manipulator Jacobian matrix

$$\mathbf{J}(\vec{\theta}) = \frac{\partial \mathbf{Q}(\vec{\theta})}{\partial \vec{\theta}} = \begin{bmatrix} \frac{\partial q_1(\vec{\theta})}{\partial \theta_1} & \frac{\partial q_1(\vec{\theta})}{\partial \theta_2} & \cdots & \frac{\partial q_1(\vec{\theta})}{\partial \theta_n} \\ \frac{\partial q_2(\vec{\theta})}{\partial \theta_1} & \frac{\partial q_2(\vec{\theta})}{\partial \theta_2} & \cdots & \frac{\partial q_2(\vec{\theta})}{\partial \theta_n} \\ \cdots & \cdots & \cdots & \cdots \\ \frac{\partial q_N(\vec{\theta})}{\partial \theta_1} & \frac{\partial q_N(\vec{\theta})}{\partial \theta_2} & \cdots & \frac{\partial q_N(\vec{\theta})}{\partial \theta_n} \end{bmatrix}$$

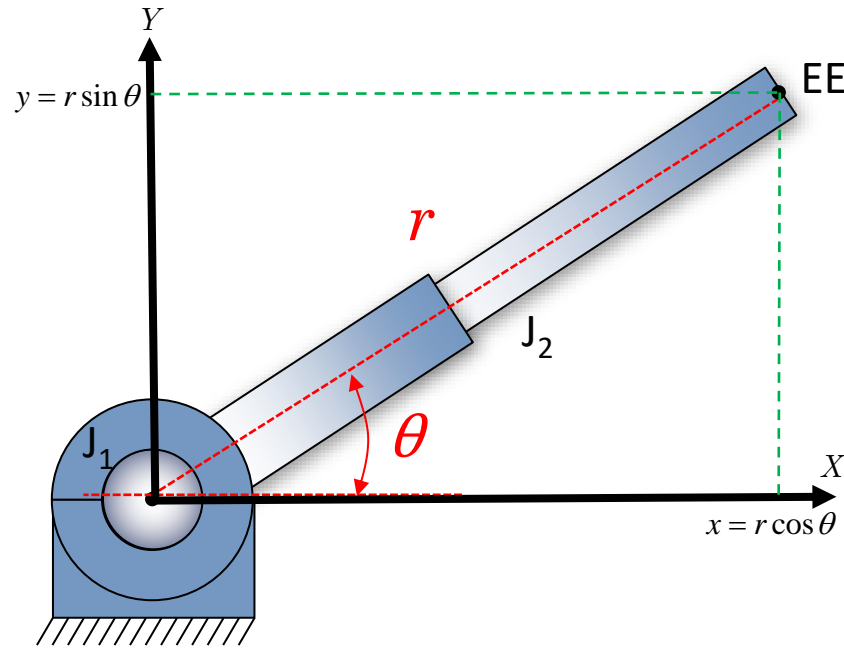
} N rows, task space dimension

} n columns, number of 1 DoF joints

$$\vec{\theta}_{(k)} = \vec{\theta}_{(k-1)} + \mathbf{J}^{-1}(\vec{\theta}_{(k-1)}) (\mathbf{P} - \mathbf{Q}(\vec{\theta}_{(k-1)})) \Rightarrow \underbrace{\vec{\theta}_{(k)} - \vec{\theta}_{(k-1)}}_{\text{Joint space}} = \mathbf{J}^{-1}(\vec{\theta}_{(k-1)}) \underbrace{(\mathbf{P} - \mathbf{Q}(\vec{\theta}_{(k-1)}))}_{\text{Task space}}$$

The **Jacobian** defines a mapping between small (differential) changes in joint space, and how they create small (differential) changes in cartesian space, and vice versa

Revision: Jacobian derivation for a simple manipulator



- Polar manipulator: J_1 revolute, J_2 prismatic
- Joint space size: $n = 2$, as two 1-DoF joints
- Task space dimension: $N=2$, two translations, ignore rotation
- Kinematic equations: $x = r \cos \theta$
 $y = r \sin \theta$

In terms of earlier slides on iterative solution of IK

- Vector of joint variables: $\vec{\theta} = [\theta, r]$
- Pose vector: $\mathbf{P} = [x, y]^T$
- Vector of kinematic equations: $\mathbf{Q} = [r \cos \theta, r \sin \theta]^T$

- End-effector coordinates are $(x, y) = (r \cos \theta, r \sin \theta)$
- The rate of change (velocities) of x and y , using the chain rule to differentiate with respect to time t , is:

$$\frac{dx}{dt} = \frac{\partial(r \cos \theta)}{\partial r} \frac{dr}{dt} + \frac{\partial(r \cos \theta)}{\partial \theta} \frac{d\theta}{dt} \Rightarrow \dot{x} = \cos \theta \dot{r} - r \sin \theta \dot{\theta}$$

$$\frac{dy}{dt} = \frac{\partial(r \sin \theta)}{\partial r} \frac{dr}{dt} + \frac{\partial(r \sin \theta)}{\partial \theta} \frac{d\theta}{dt} \Rightarrow \dot{y} = \sin \theta \dot{r} + r \cos \theta \dot{\theta}$$

$$\begin{bmatrix} \frac{dx}{dt} \\ \frac{dy}{dt} \end{bmatrix} = \begin{bmatrix} \frac{\partial x}{\partial \theta} & \frac{\partial x}{\partial r} \\ \frac{\partial y}{\partial \theta} & \frac{\partial y}{\partial r} \end{bmatrix} \begin{bmatrix} \frac{d\theta}{dt} \\ \frac{dr}{dt} \end{bmatrix}$$

$$\begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} = \begin{bmatrix} -r \sin \theta & \cos \theta \\ r \cos \theta & \sin \theta \end{bmatrix} \begin{bmatrix} \dot{\theta} \\ \dot{r} \end{bmatrix}$$

$$\begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} = \mathbf{J} \begin{bmatrix} \dot{\theta} \\ \dot{r} \end{bmatrix}$$

- We can also use the inverse Jacobian to find the joint velocities for a given cartesian velocity of the EE

$$\mathbf{J}^{-1} \begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} = \begin{bmatrix} \dot{\theta} \\ \dot{r} \end{bmatrix} \quad \text{where} \quad \mathbf{J}^{-1} = \begin{bmatrix} \frac{-\sin \theta}{r} & \frac{\cos \theta}{r} \\ \cos \theta & \sin \theta \end{bmatrix}$$

Chain rule for functions of two independent variables

$$z = F(x, y)$$

$$dz = \frac{\partial F}{\partial x} dx + \frac{\partial F}{\partial y} dy$$

$$\frac{dz}{dt} = \frac{\partial F}{\partial x} \frac{dx}{dt} + \frac{\partial F}{\partial y} \frac{dy}{dt}$$

The **Jacobian** also defines a mapping between joint velocities and end effector velocity

Today's material – Velocities and Jacobian

- So far, we have been dealing with static robotic manipulators –
 - Basically, robot configurations frozen in time
- We need to consider linear and angular velocity concepts –
 - Robot configuration changing in time
 - By considering infinitesimally small steps of time
- We will see how to calculate joint/link velocities
 - Analytical approach
 - Differential approach, i.e., using the Jacobian

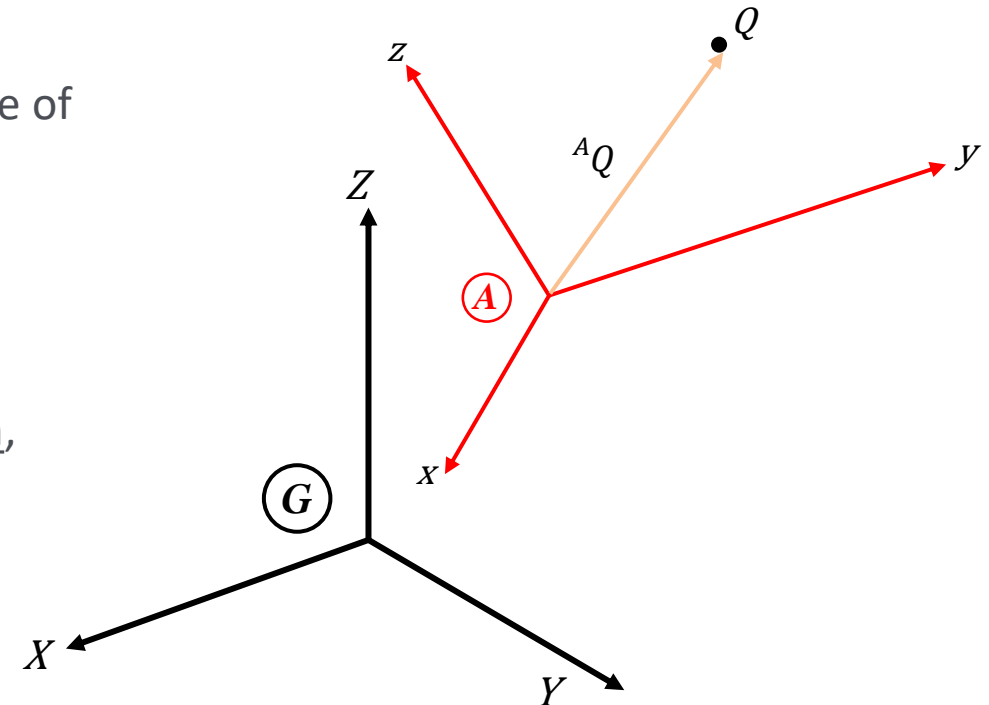
Velocity representation

- The velocity of a position vector (e.g., ${}^A Q$) is the linear velocity of the point in space represented by this vector (e.g., Q).
- The velocity of Q relative to frame $\{A\}$ can be expressed by the derivative of Q relative to $\{A\}$:

$${}^A V(Q) = \frac{d}{dt} {}^A Q = \lim_{\Delta t \rightarrow 0} \frac{{}^A Q(t + \Delta t) - {}^A Q(t)}{\Delta t}$$

- The calculated velocity is written in terms of the frame of differentiation, which is $\{A\}$ in this case.
- However, the velocity vector can be expressed/described in terms of any other frame, e.g., $\{G\}$:
- If the calculated velocity is described in terms of the frame of differentiation, $\{A\}$ in this case, for simplicity:

$${}^A ({}^A V(Q)) \equiv {}^A V(Q)$$



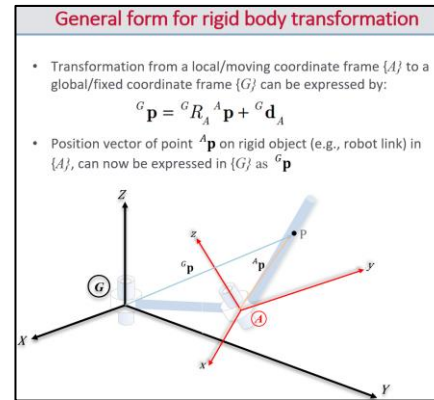
Velocity representation: change in reference frame

- In the case that $\{A\}$ is **stationary** relative to $\{G\}$:

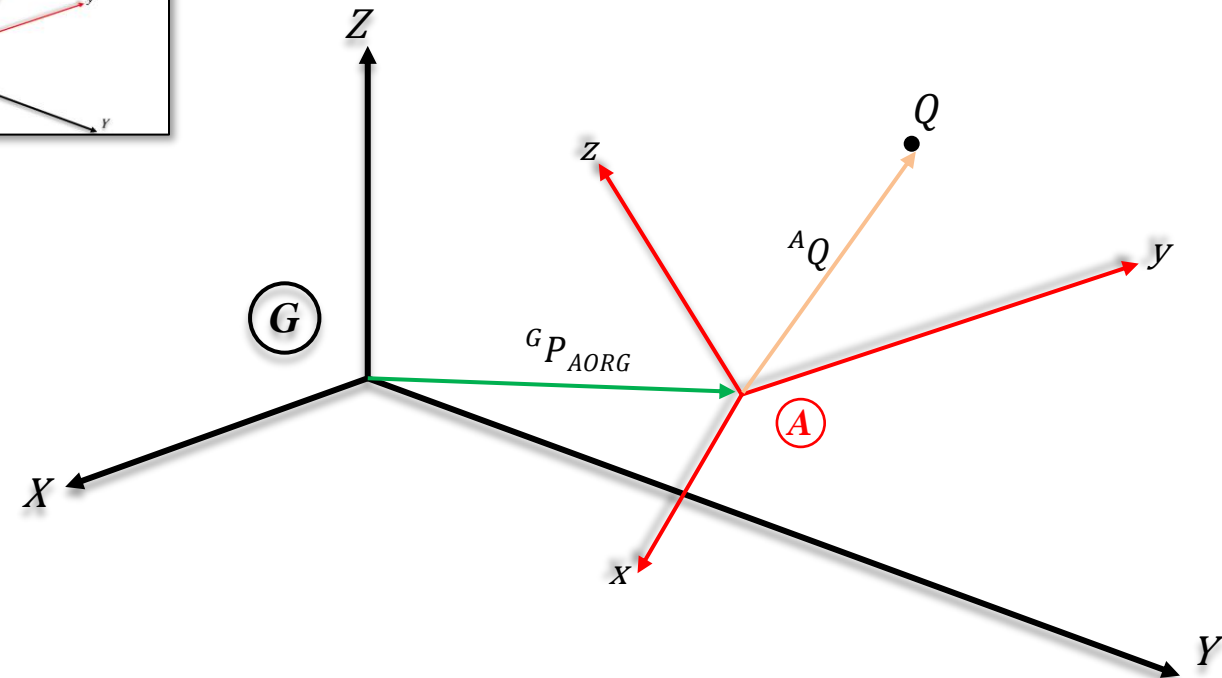
$${}^G Q = {}^G R_A {}^A Q + {}^G P_{AORG}$$

$${}^G \left({}^A V(Q) \right) = \frac{d}{dt} {}^G R_A {}^A Q + \frac{d}{dt} {}^G P_{AORG} \quad \text{with a red arrow pointing to } \frac{d}{dt} {}^G P_{AORG} \text{ and } \equiv 0$$

$${}^G \left({}^A V(Q) \right) = {}^G R_A {}^A V(Q)$$



From Lecture on Kinematics and Transformations



Linear Velocity when Local Frame is Moving (not rotating)

- In the case that $\{A\}$ is **also moving** with a linear velocity V and fixed orientation relative to $\{G\}$:

$${}^A V(Q) = \frac{d}{dt} {}^A Q = \lim_{\Delta t \rightarrow 0} \frac{{}^A Q(t + \Delta t) - {}^A Q(t)}{\Delta t}$$

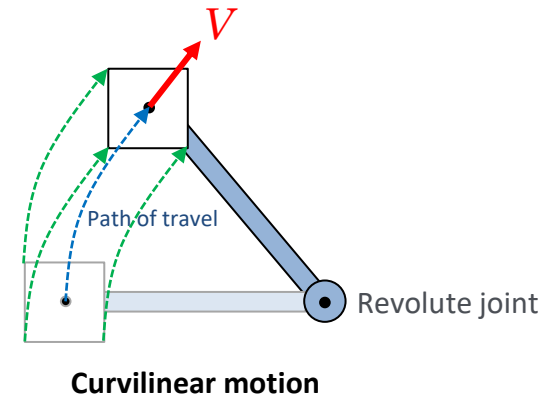
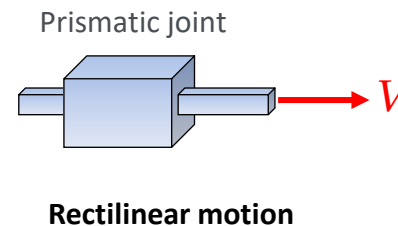
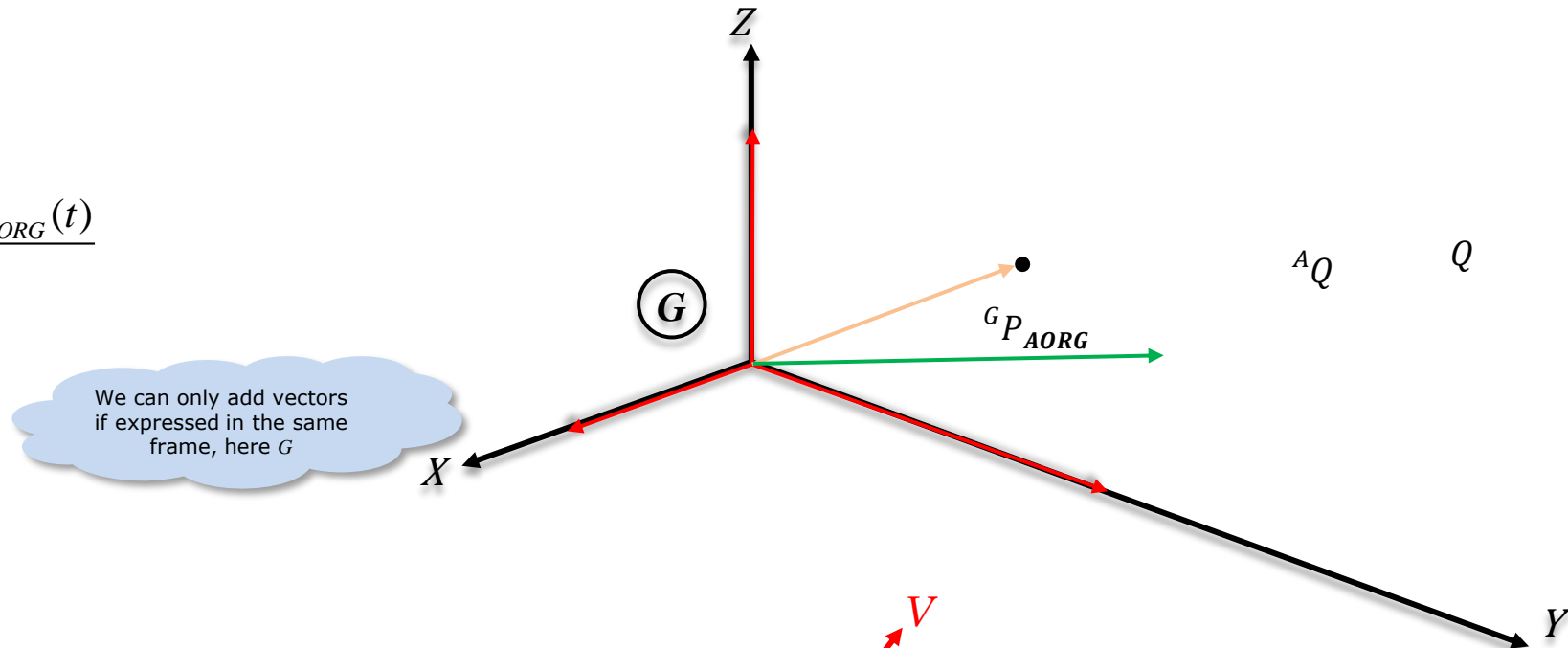
$${}^G V_{AORG} = \frac{d}{dt} {}^G P_{AORG} = \lim_{\Delta t \rightarrow 0} \frac{{}^G P_{AORG}(t + \Delta t) - {}^G P_{AORG}(t)}{\Delta t}$$

- Adding the two velocity vectors in $\{G\}$

$${}^G V(Q) = {}^G ({}^A V(Q)) + {}^G V_{AORG}$$

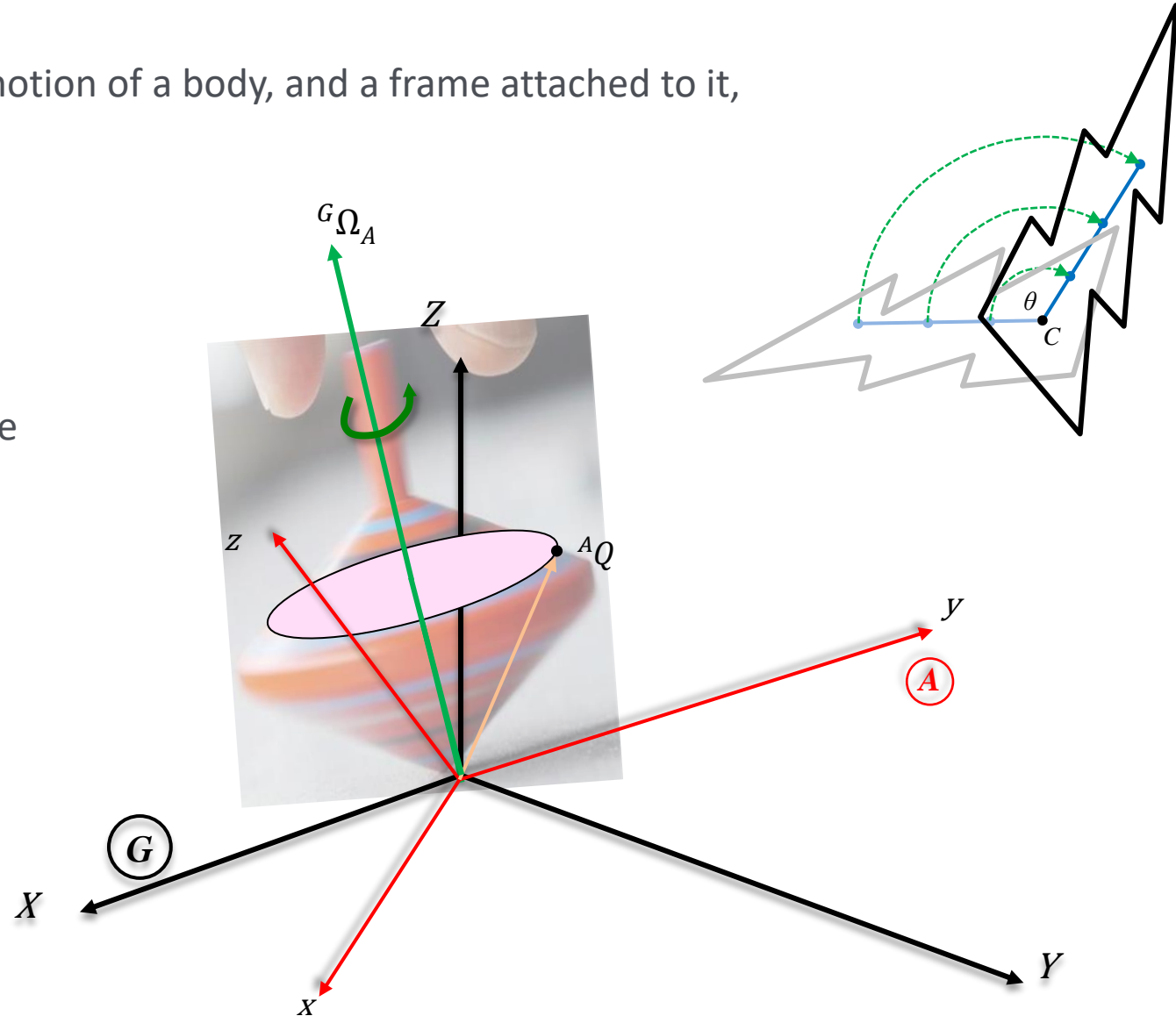
$${}^G V(Q) = {}^G R_A {}^A V(Q) + {}^G V_{AORG}$$

NOTE: this is valid only when the relative orientation of $\{A\}$ and $\{G\}$ remains constant during motion i.e., during **rectilinear or curvilinear** motion

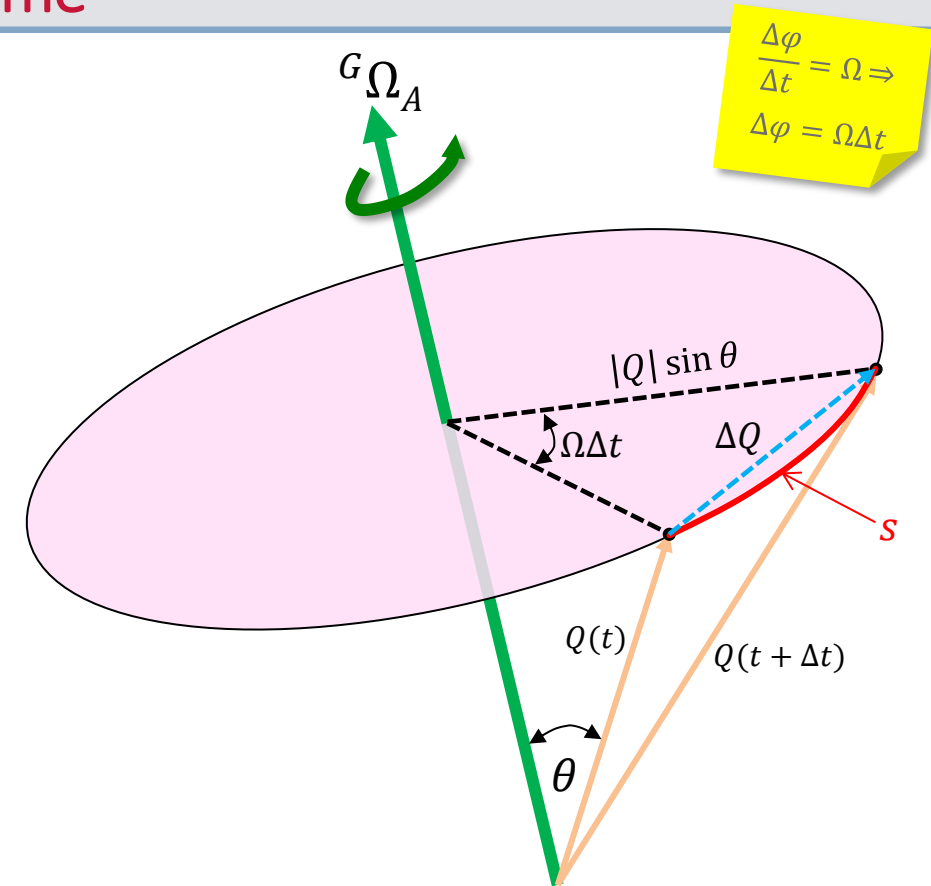
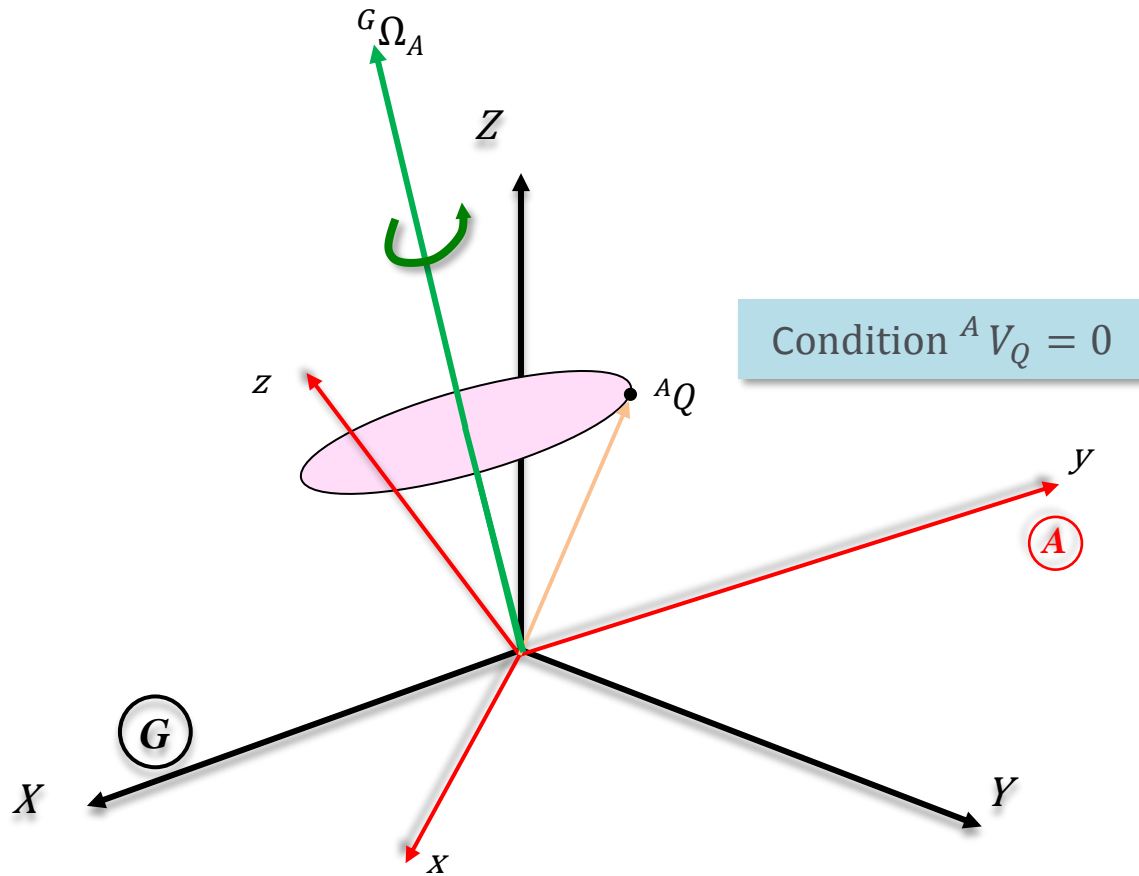


Angular (rotational) velocity: Rotating Local Frame

- Angular (rotational) velocity describes the rotational motion of a body, and a frame attached to it, around a centre of rotation (e.g., centre of mass)
- Frames $\{A\}$ and $\{G\}$ are at same origin, and have zero linear relative velocity
- The rotation of frame $\{A\}$ relative to $\{G\}$ is given by the **angular velocity vector** ${}^G\Omega_A$
- Physically, the **direction** of vector ${}^G\Omega_A$ represents the instantaneous axis of rotation of $\{A\}$ relative to $\{G\}$
- The **magnitude** of vector ${}^G\Omega_A$ indicates the angular speed of rotation
- Assuming linear velocity ${}^A V_Q = 0$, when $\{A\}$ rotates about ${}^G\Omega_A$, point Q is tracing a circular path in $\{G\}$



Solving for the velocity of Q for Rotating Local Frame



- Arc length $s = (|{}^G Q| \sin \theta) (|{}^G \Omega_A| \Delta t)$ When Δt is infinitesimally small: $s = |\Delta Q|$
- Therefore $|\Delta Q| = (|{}^G Q| \sin \theta) (|{}^G \Omega_A| \Delta t) = |{}^G Q| |{}^G \Omega_A| \sin \theta \Delta t \Rightarrow |\Delta Q / \Delta t| = |{}^G Q| |{}^G \Omega_A| \sin \theta$

$${}^G V(Q) = {}^G \Omega_A \times {}^G Q$$

${}^G V(Q)$ = linear velocity
 ${}^G \Omega_A$ = angular velocity

Velocity of Q for Rotating Local Frame and Moving Point Q

- In the previous case that Q was fixed in $\{A\}$:

$${}^G V(Q) = {}^G \Omega_A \times {}^G Q$$

- What if Q is also moving in $\{A\}$, with linear velocity ${}^A V$?
- Simply add velocity components:

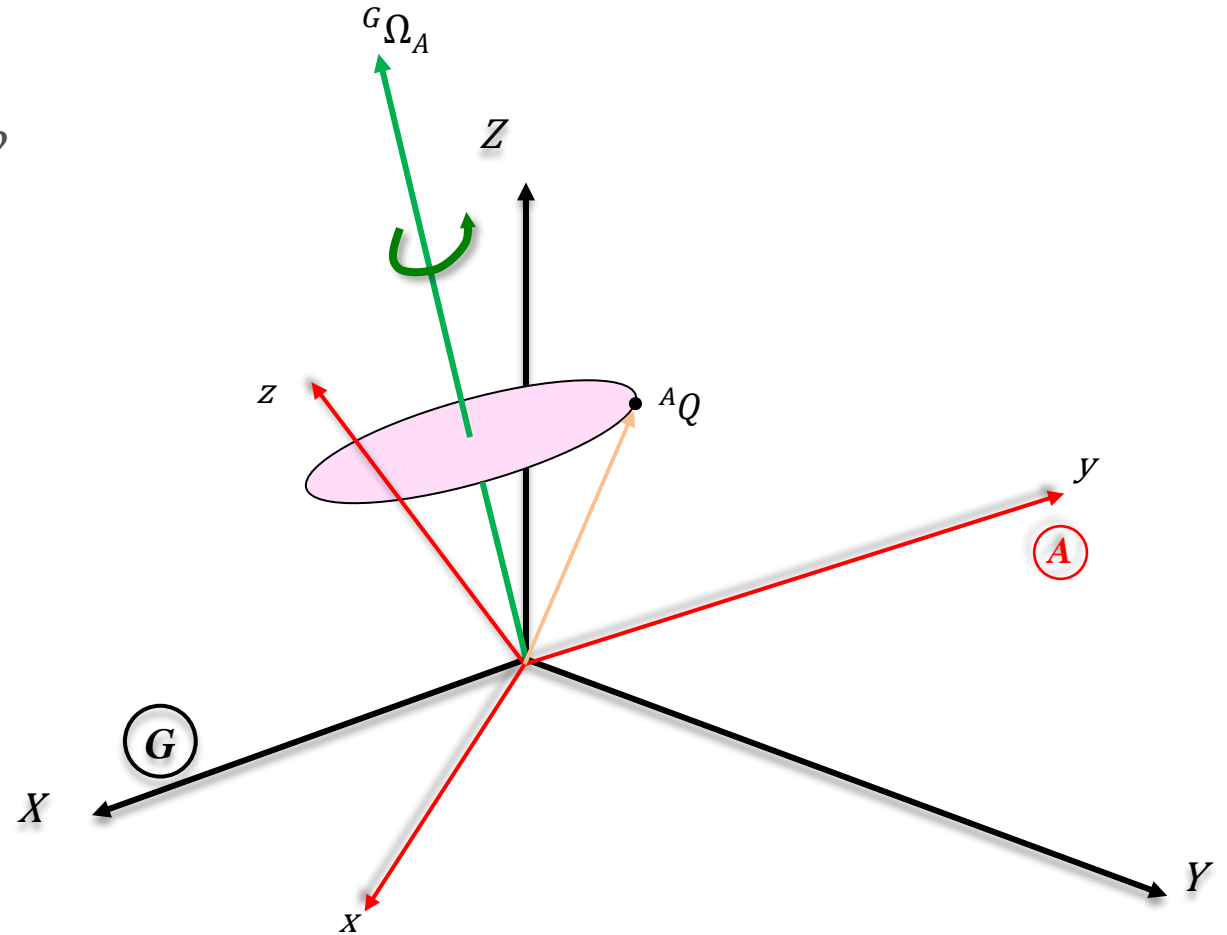
$${}^G V(Q) = \underbrace{{}^G ({}^A V(Q))}_{\text{Linear velocity of point } Q \text{ in } \{G\}, \text{ due to its linear velocity in } \{A\}} + \underbrace{{}^G \Omega_A \times {}^G Q}_{\text{Linear velocity of point } Q \text{ in } \{G\}, \text{ due to angular velocity of } \{A\}}$$

Linear velocity of point Q in $\{G\}$,
due to its linear velocity in $\{A\}$

Linear velocity of point Q in $\{G\}$,
due to angular velocity of $\{A\}$

- Or to remove dual superscript notation and directly express ${}^A V$ in $\{G\}$, use rotation matrix:

$${}^G V(Q) = {}^G R_A {}^A V(Q) + {}^G \Omega_A \times {}^G R_A {}^A Q$$

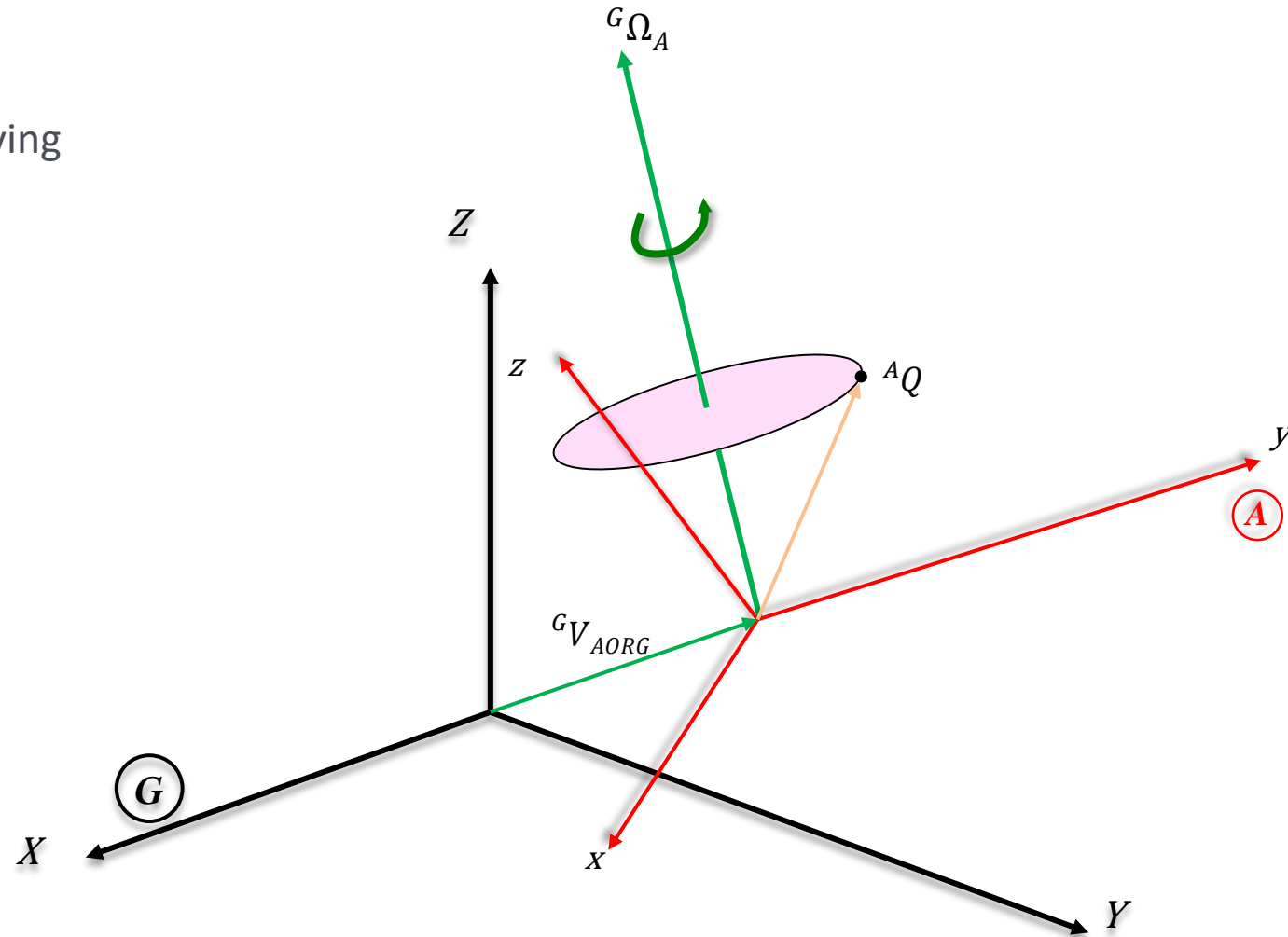


Velocity of Q in simultaneous Rotating and Moving Local Frame and Moving Point

- Now let's make it even more complicated
- Point Q is moving in $\{A\}$, while $\{A\}$ is rotating and moving with a velocity of ${}^G V_{AORG}$
- Again, simply add velocity components

$${}^G V(Q) = \underbrace{{}^G V_{AORG}}_{\text{Linear velocity of } \{A\} \text{ relative to } \{G\}} + \underbrace{{}^G R_A {}^A V(Q)}_{\text{Linear velocity of point } {}^A Q \text{ in } \{G\}, \text{ due to its linear velocity in } \{A\}} + \underbrace{{}^G \Omega_A \times {}^G R_A {}^A Q}_{\text{Linear velocity of point } {}^A Q \text{ in } \{G\}, \text{ due to angular velocity of } \{A\}}$$

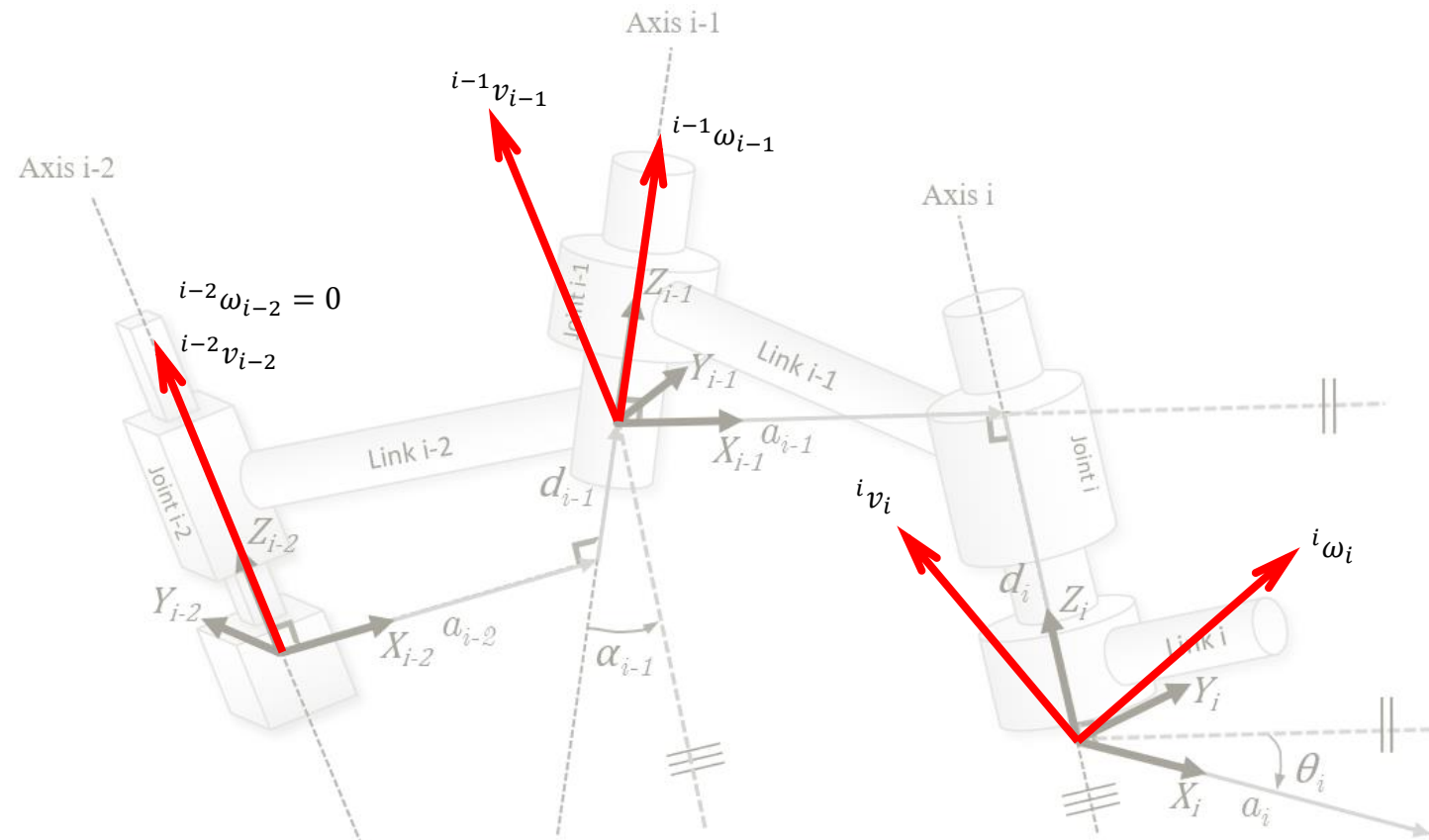
The above equation is the general expression for the derivative of a vector in a moving frame as seen from a stationary frame



Interlink Velocity Propagation

- Problem: Calculate the **linear and angular velocities** of the links of a robotic manipulator
- A manipulator is a chain of bodies; hence velocities **propagate** from the previous to the next links in the chain
- The velocity of **link i** is the velocity of **link $i-1$** plus any additional velocities contributed by **joint i**

- “Velocity of a link” means the **linear velocity of the origin of the link frame** and the **angular velocity of the link**.
- We express link velocities **with respect to the link frame itself**, rather than the base frame
- **Remember:** Velocities can only be added when vectors are expressed in the same frame



Interlink Angular Velocity Propagation

- The angular velocity of link i is that of $i-1$, plus a new component caused by angular velocity of joint i . This can be expressed in frame $\{i-1\}$ as:

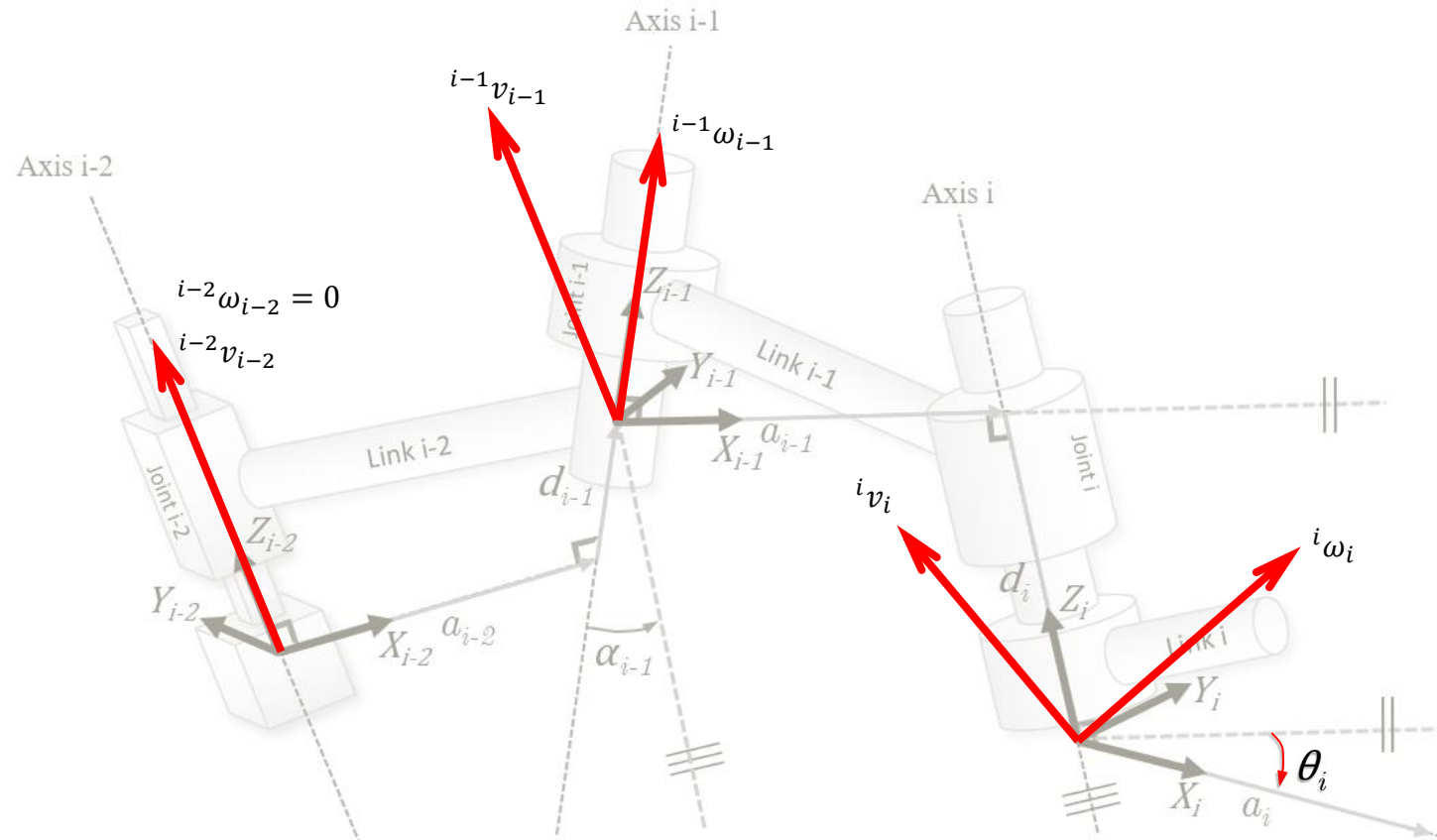
$${}^{i-1}\omega_i = {}^{i-1}\omega_{i-1} + {}^{i-1}R_i \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_i \end{bmatrix}$$

Derived from the link transformation matrix using the DH table of parameters (Lecture on FK, General form of DH transformation)

$${}^{i-1}T_i = \begin{bmatrix} c\theta_i & -s\theta_i & 0 & a_{i-1} \\ s\theta_i c\alpha_{i-1} & c\theta_i c\alpha_{i-1} & -s\alpha_{i-1} & -s\alpha_{i-1}d_i \\ s\theta_i s\alpha_{i-1} & c\theta_i s\alpha_{i-1} & c\alpha_{i-1} & c\alpha_{i-1}d_i \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

- Pre-multiplying both sides by ${}^iR_{i-1}$ allows us to express the angular velocity of link i in its own frame:

$${}^i\omega_i = {}^iR_{i-1} {}^{i-1}\omega_{i-1} + \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_i \end{bmatrix}$$



Interlink Linear Velocity Propagation

- Similarly, the linear velocity of the origin of frame $\{i\}$ is the same as that of the origin of $\{i-1\}$, **plus** a new component caused by angular velocity of link $i-1$.
- Based on the general expression derived earlier

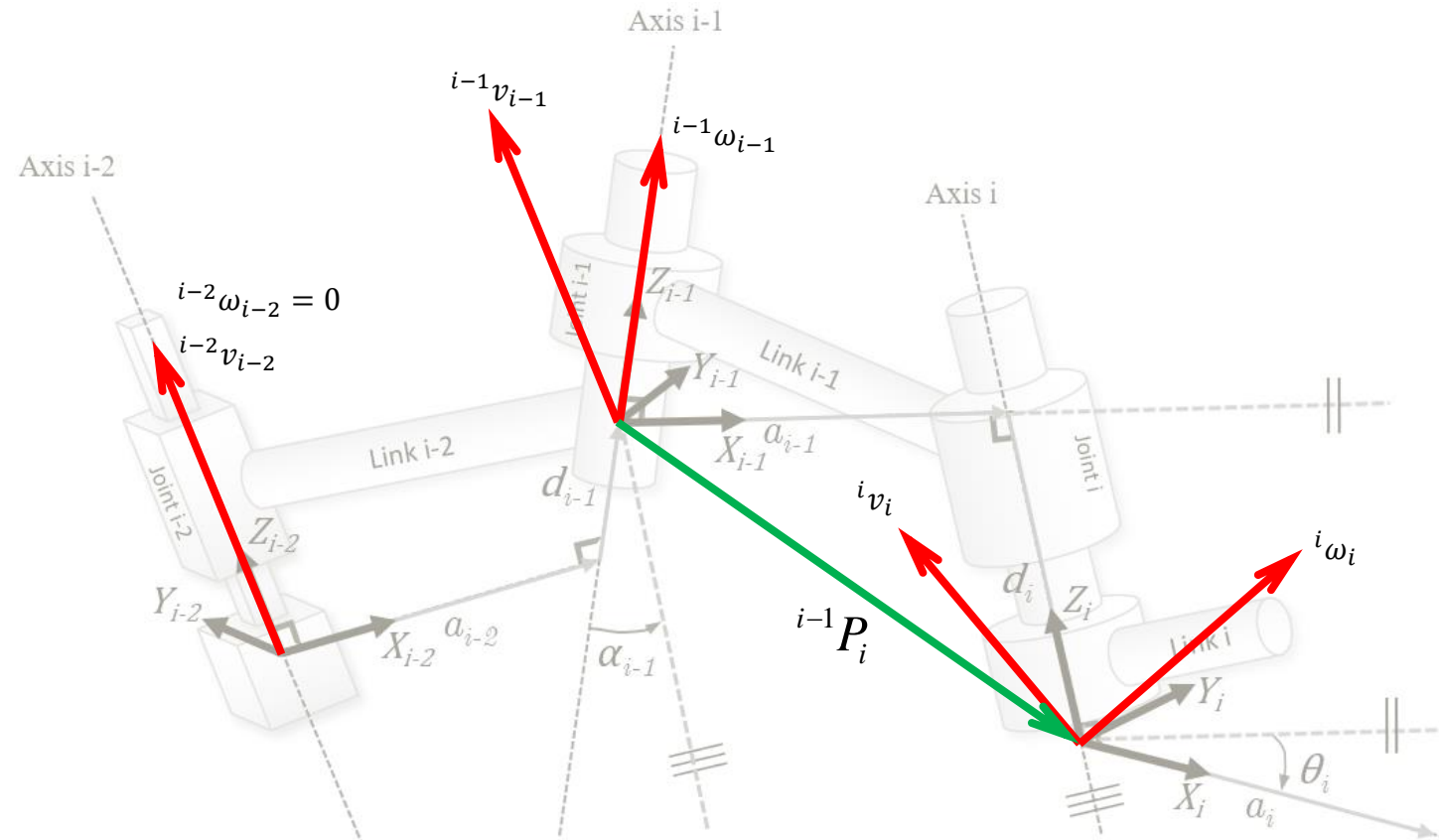
$${}^G V(Q) = {}^G V_{AORG} + {}^G R_A {}^A V(Q) + {}^G \Omega_A \times {}^G R_A {}^A Q$$

- and considering that in this case ${}^{i-1}P_i$ is constant in $\{i-1\}$

$${}^{i-1}v_i = {}^{i-1}v_{i-1} + {}^{i-1}\omega_{i-1} \times {}^{i-1}P_i$$

- Pre-multiplying both sides by ${}^i R_{i-1}$

$${}^i v_i = {}^i R_{i-1} \left({}^{i-1}v_{i-1} + {}^{i-1}\omega_{i-1} \times {}^{i-1}P_i \right)$$



Link to Link Velocity Propagation Summary

- The derived equations can be applied successively to calculate ${}^N\omega_N$ and Nv_N , the angular and linear velocities of the last link, or any other link
- Can be used directly and applied iteratively by computer code
- If velocities are required relative to the base frame, simply multiply by 0R_N

$${}^i\omega_i = {}^iR_{i-1} {}^{i-1}\omega_{i-1} + \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_i \end{bmatrix}$$

$${}^iv_i = {}^iR_{i-1} \left({}^{i-1}v_{i-1} + {}^{i-1}\omega_{i-1} \times {}^{i-1}P_i \right)$$

- If joint i is **prismatic**, the above equations become:

$${}^i\omega_i = {}^iR_{i-1} {}^{i-1}\omega_{i-1}$$

$${}^iv_i = {}^iR_{i-1} \left({}^{i-1}v_{i-1} + {}^{i-1}\omega_{i-1} \times {}^{i-1}P_i \right) + \begin{bmatrix} 0 \\ 0 \\ \dot{d}_i \end{bmatrix}$$

Example of deriving the analytical expression for last link velocities

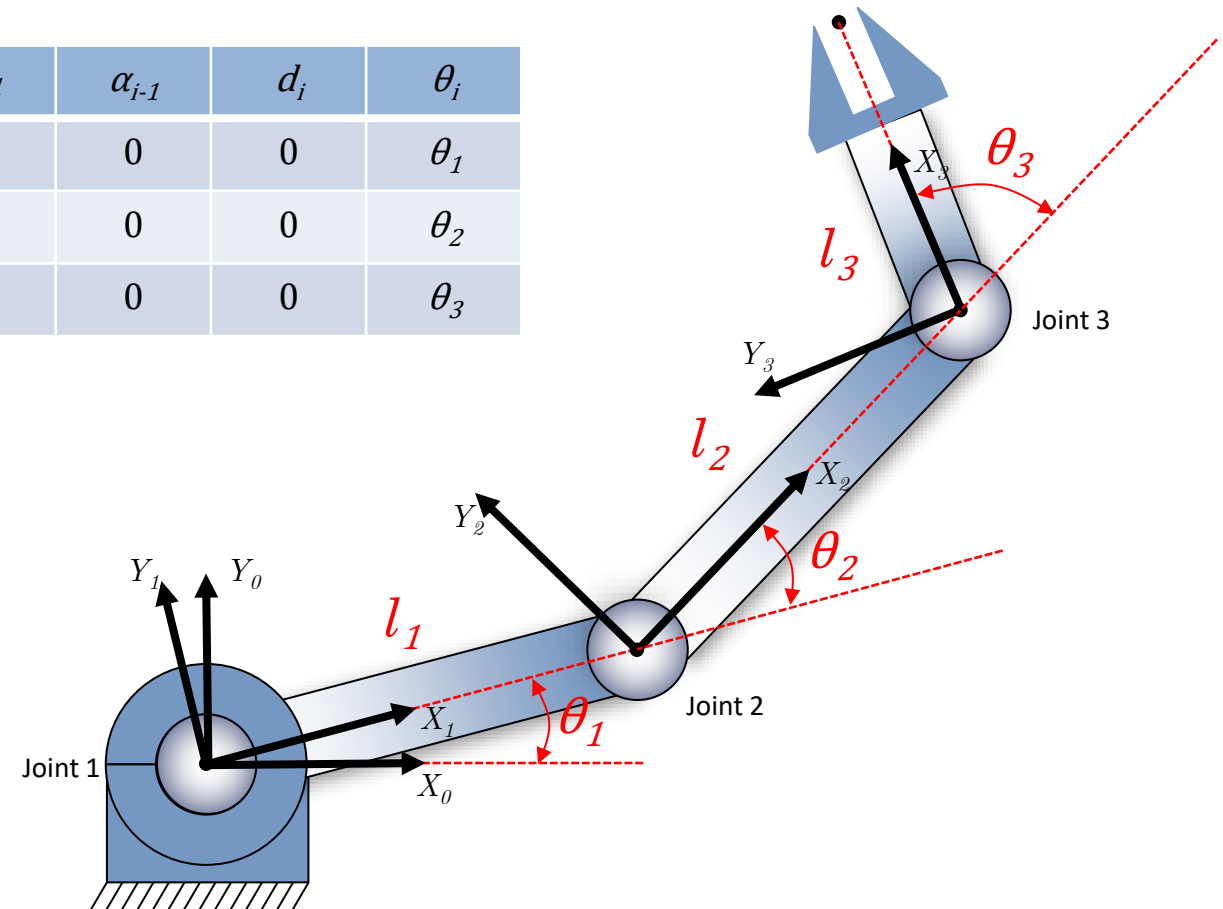
$${}^0T_1 = \begin{bmatrix} c\theta_1 & -s\theta_1 & 0 & 0 \\ s\theta_1 & c\theta_1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

$${}^1T_2 = \begin{bmatrix} c\theta_2 & -s\theta_2 & 0 & l_1 \\ s\theta_2 & c\theta_2 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

$${}^2T_3 = \begin{bmatrix} c\theta_3 & -s\theta_3 & 0 & l_2 \\ s\theta_3 & c\theta_3 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

$${}^0T_3 = {}^0T_1 {}^1T_2 {}^2T_3$$

| Frame | a_{i-1} | α_{i-1} | d_i | θ_i |
|-------|-----------|----------------|-------|------------|
| 1 | 0 | 0 | 0 | θ_1 |
| 2 | l_1 | 0 | 0 | θ_2 |
| 3 | l_2 | 0 | 0 | θ_3 |

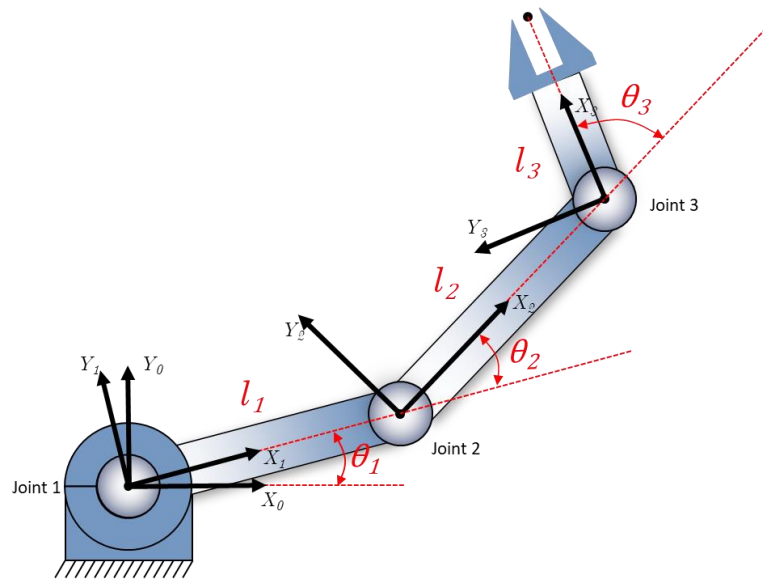


Example of deriving the analytical expression for last link velocities

- Calculate link velocities for the manipulator

$${}^i\omega_i = {}^iR_{i-1} {}^{i-1}\omega_{i-1} + \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_i \end{bmatrix}$$

$${}^i v_i = {}^iR_{i-1} \left({}^{i-1}v_{i-1} + {}^{i-1}\omega_{i-1} \times {}^{i-1}P_i \right)$$



$${}^1\omega_1 = \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_1 \end{bmatrix} \quad {}^1v_1 = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$${}^2\omega_2 = {}^2R_1 {}^1\omega_1 + \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_2 \end{bmatrix} = \begin{bmatrix} c\theta_2 & s\theta_2 & 0 \\ -s\theta_2 & c\theta_2 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_1 \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_1 + \dot{\theta}_2 \end{bmatrix}$$

$$\mathbf{a} \times \mathbf{b} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ a_1 & a_2 & a_3 \\ b_1 & b_2 & b_3 \end{vmatrix}$$

$${}^2v_2 = {}^2R_1 \left({}^1v_1 + {}^1\omega_1 \times \begin{bmatrix} l_1 \\ 0 \\ 0 \end{bmatrix} \right) = \begin{bmatrix} c\theta_2 & s\theta_2 & 0 \\ -s\theta_2 & c\theta_2 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ l_1\dot{\theta}_1 \\ 0 \end{bmatrix} = \begin{bmatrix} l_1\dot{\theta}_1 s\theta_2 \\ l_1\dot{\theta}_1 c\theta_2 \\ 0 \end{bmatrix}$$

$${}^3\omega_3 = {}^3R_2 {}^2\omega_2 + \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_1 + \dot{\theta}_2 + \dot{\theta}_3 \end{bmatrix}$$

$${}^3v_3 = {}^3R_2 \left({}^2v_2 + {}^2\omega_2 \times \begin{bmatrix} l_2 \\ 0 \\ 0 \end{bmatrix} \right) = \begin{bmatrix} l_1\dot{\theta}_1 s\theta_2 \\ l_1\dot{\theta}_1 c\theta_2 + l_2(\dot{\theta}_1 + \dot{\theta}_2) \\ 0 \end{bmatrix}$$

Example using the Jacobian-derived velocities

Kinematics Equations
$$\begin{cases} x_a = l_1 c\theta_1 + l_2 c\theta_{12} \\ y_a = l_1 s\theta_1 + l_2 s\theta_{12} \end{cases}$$

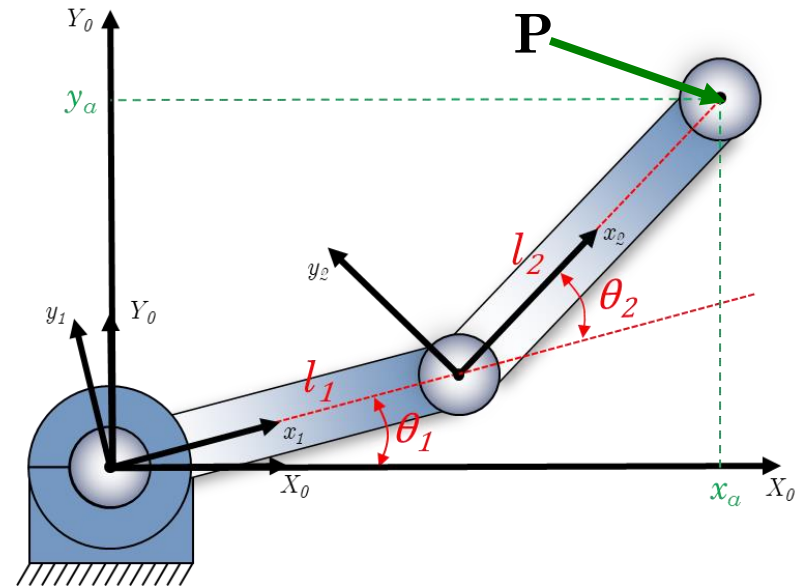
$$\mathbf{P} = \begin{bmatrix} x_a \\ y_a \end{bmatrix} \quad \mathbf{Q} = \begin{bmatrix} l_1 c\theta_1 + l_2 c\theta_{12} \\ l_1 s\theta_1 + l_2 s\theta_{12} \end{bmatrix}$$

$$\mathbf{J}(\vec{\theta}) = \left[\frac{\partial \mathbf{Q}(\vec{\theta})}{\partial \vec{\theta}} \right] = \begin{bmatrix} -l_1 s\theta_1 - l_2 s\theta_{12} & -l_2 s\theta_{12} \\ l_1 c\theta_1 + l_2 c\theta_{12} & l_2 c\theta_{12} \end{bmatrix}$$

$$\dot{\mathbf{P}} = \mathbf{J}(\vec{\theta}) \dot{\vec{\theta}} \Rightarrow \begin{bmatrix} \dot{x}_a \\ \dot{y}_a \end{bmatrix} = \begin{bmatrix} -l_1 s\theta_1 - l_2 s\theta_{12} & -l_2 s\theta_{12} \\ l_1 c\theta_1 + l_2 c\theta_{12} & l_2 c\theta_{12} \end{bmatrix} \begin{bmatrix} \dot{\theta}_1 \\ \dot{\theta}_2 \end{bmatrix}$$

- If for example we want to limit the end-effector \mathbf{P} to move at **1m/s** horizontally, the required joint speed can be found as:

$$\begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} -l_1 s\theta_1 - l_2 s\theta_{12} & -l_2 s\theta_{12} \\ l_1 c\theta_1 + l_2 c\theta_{12} & l_2 c\theta_{12} \end{bmatrix} \begin{bmatrix} \dot{\theta}_1 \\ \dot{\theta}_2 \end{bmatrix} \rightarrow \begin{bmatrix} \dot{\theta}_1 \\ \dot{\theta}_2 \end{bmatrix} = \begin{bmatrix} -l_1 s\theta_1 - l_2 s\theta_{12} & -l_2 s\theta_{12} \\ l_1 c\theta_1 + l_2 c\theta_{12} & l_2 c\theta_{12} \end{bmatrix}^{-1} \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$



Changing a Jacobian's Frame of Reference

- Given a Jacobian in frame $\{A\}$:
$$\begin{bmatrix} {}^A v \\ {}^A \omega \end{bmatrix} = {}^A \mathbf{J}(\vec{\theta}) \dot{\vec{\theta}}$$
- We might be interested in deriving an expressions of the Jacobian in another frame, e.g., $\{G\}$
- As discussed previously, a 6x1 Cartesian velocity vector in $\{A\}$ can be described relative to $\{G\}$ by the transformation:

$$\underbrace{\begin{bmatrix} {}^G v \\ {}^G \omega \end{bmatrix}}_{6 \times 1} = \underbrace{\begin{bmatrix} {}^G R_A & 0 \\ 0 & {}^G R_A \end{bmatrix}}_{6 \times 6} \underbrace{\begin{bmatrix} {}^A v \\ {}^A \omega \end{bmatrix}}_{6 \times 1}$$

Hence, we can write:

$$\begin{bmatrix} {}^G v \\ {}^G \omega \end{bmatrix} = \begin{bmatrix} {}^G R_A & 0 \\ 0 & {}^G R_A \end{bmatrix} {}^A \mathbf{J}(\vec{\theta}) \dot{\vec{\theta}}$$

- Therefore, changing the frame of reference of a Jacobian can be accomplished by the following relationship:

$${}^G \mathbf{J}(\vec{\theta}) = \begin{bmatrix} {}^G R_A & 0 \\ 0 & {}^G R_A \end{bmatrix} {}^A \mathbf{J}(\vec{\theta})$$

Conclusions

- Velocity Representation
- Linear and Angular Velocities
- Interlink Velocity Propagation
- Analytical and Jacobian Velocities
- Changing of Jacobian Frame