

### 3 Kutzelnigg-Mukherjee tensor notation

**Notation 3.1. Kutzelnigg-Mukherjee tensor notation.** In *Kutzelnigg-Mukherjee (KM) tensor notation*, creation operators are denoted using superscripts,  $a^p \equiv a_p^\dagger$ . Consequently, the one-particle and one-hole density matrices are written as  $\gamma_q^p \equiv \langle \Phi | a^p a_q | \Phi \rangle$  and  $\eta_p^q \equiv \langle \Phi | a_p a^q | \Phi \rangle$ , respectively. The one- and two-electron integrals are also written with upper and lower indices: lower indices denote spin-orbitals in the bra and upper ones refer to spin-orbitals in the ket.

$$h_p^q \equiv \langle \psi_p | \hat{h} | \psi_q \rangle \quad g_{pq}^{rs} \equiv \langle pq | rs \rangle \quad \bar{g}_{pq}^{rs} \equiv \langle pq || rs \rangle \quad (3.1)$$

Vacuum-normal-ordered excitations are given the compact notation  $a_{q_1 \dots q_m}^{p_1 \dots p_m} \equiv a^{p_1} \dots a^{p_m} a_{q_m} \dots a_{q_1}$  and  $\Phi$ -normal-ordered excitations are written as  $\tilde{a}_{q_1 \dots q_m}^{p_1 \dots p_m} \equiv :a^{p_1} \dots a^{p_m} a_{q_m} \dots a_{q_1}:$ . Using upper and lower indices enables one to employ the *Einstein summation convention*, in which any index that appears twice in a product, once as a lower index and once as an upper one, is implicitly summed over. As an example, consider the electronic Hamiltonian as expressed in KM notation.

$$H_e = h_p^q a_q^p + \frac{1}{4} \bar{g}_{pq}^{rs} a_{rs}^{pq} = E_0 + H_c \quad H_c = f_p^q \tilde{a}_p^q + \frac{1}{4} \bar{g}_{pq}^{rs} \tilde{a}_{rs}^{pq} \quad E_0 \equiv h_p^q \gamma_q^p + \frac{1}{2} \bar{g}_{pr}^{qs} \gamma_q^p \gamma_s^r \quad f_p^q \equiv h_p^q + \bar{g}_{pr}^{qs} \gamma_s^r \quad (3.2)$$

Here,  $E_0$  is the Hartree-Fock reference energy,  $f_p^q$  denotes a matrix element of the Fock operator, and  $H_c$  denotes the correlation component of the Hamiltonian. More generally, if  $\hat{v}$  is an  $m$ -electron operator, i.e. an operator that acts on  $m$  electronic coordinates, its second quantized form is expressed in KM notation as

$$\hat{v} |_{\mathcal{F}(\mathcal{H})} = \frac{1}{m!} v_{p_1 \dots p_m}^{q_1 \dots q_m} a_{q_1 \dots q_m}^{p_1 \dots p_m} \quad v_{p_1 \dots p_m}^{q_1 \dots q_m} \equiv \int d(1 \dots m) \psi_{p_1}^*(1) \dots \psi_{p_m}^*(m) \hat{v}(1, \dots, m) \psi_{q_1}(1) \dots \psi_{q_m}(m) \quad (3.3)$$

where  $v_{p_1 \dots p_m}^{q_1 \dots q_m}$  is the *interaction tensor* of  $\hat{v}$ . Equivalently,  $\hat{v}$  can also be expressed as

$$\hat{v} |_{\mathcal{F}(\mathcal{H})} = \left(\frac{1}{m!}\right)^2 \bar{v}_{p_1 \dots p_m}^{q_1 \dots q_m} a_{q_1 \dots q_m}^{p_1 \dots p_m} \quad \bar{v}_{p_1 \dots p_m}^{q_1 \dots q_m} \equiv \sum_{\pi \in S_m} \varepsilon_{\pi} v_{p_1 \dots p_m}^{q_{\pi(1)} \dots q_{\pi(m)}} \quad (3.4)$$

where  $\bar{v}_{p_1 \dots p_m}^{q_1 \dots q_m}$  is an *antisymmetrized interaction tensor*. Ordinary interaction tensors are symmetric under simultaneous permutation of upper and lower indices, which is equivalent to changing integration variables in equation 3.3. Antisymmetrized interaction tensors allow for independent permutations of upper and lower indices, with a sign factor corresponding to the parity of the permutation. The same permutational symmetries are shared by the excitation operators

$$\tilde{a}_{q_1 \dots q_m}^{p_1 \dots p_m} = \varepsilon_{\pi} \tilde{a}_{q_{\pi(1)} \dots q_{\pi(m)}}^{p_{\pi(1)} \dots p_{\pi(m)}} = \varepsilon_{\pi} \tilde{a}_{q_1 \dots q_m}^{p_1 \dots p_m} = \tilde{a}_{q_{\pi(1)} \dots q_{\pi(m)}}^{p_{\pi(1)} \dots p_{\pi(m)}} \quad \text{for all } \pi \in S_m \quad (3.5)$$

since creation operators anticommute with each other, as do annihilation operators. Note also the following rearrangements

$$\tilde{a}_{q_1 \dots q_m}^{p_1 \dots p_m} = :a_{q_1}^{p_1} \dots a_{q_m}^{p_m}: \quad :a_{q_1 \dots q_m}^{p_1 \dots p_m} a_{s_1 \dots s_n}^{r_1 \dots r_n}: = :a_{q_1}^{p_1} \dots a_{q_m}^{p_m} a_{s_1}^{r_1} \dots a_{s_n}^{r_n}: = \tilde{a}_{q_1 \dots q_m s_1 \dots s_n}^{p_1 \dots p_m r_1 \dots r_n} \quad (3.6)$$

which follow from the fact that the normal-ordering mapping is antisymmetric with respect to its operator string.

**Notation 3.2. Dot notation for contractions.** To make the notation more flexible, we here augment the traditional KM notation with the following definitions of *particle • contractions* and *hole ◦ contractions*.<sup>1</sup>

$$a_p \bullet a^q \equiv \overline{a_p} a^q \quad a^q \bullet a_p \equiv -\overline{a_p} a^q \quad a^q \circ a_p \equiv \overline{a^q} a_p \quad a_p \circ a^q \equiv -\overline{a^q} a_p \quad (3.7)$$

Note that  $a_q^{\bullet} = -\eta_q^p$  and  $a_q^{\circ} = \gamma_q^p$ . In vac-normal ordering, the hole contractions vanish and the particle contractions become Kronecker deltas. For multiply contracted strings, different contractions will be distinguished with repeated dots,  $a_{p \bullet \bullet} a^{p \bullet \bullet}$ , or dots with numbers,  $a_{p \bullet 2} a^{p \bullet 2}$ . This notation allows for normal ordered strings with contractions to keep all of the permutational symmetries shown in equations 3.5 and 3.6.

**Example 3.1.** Using notations 3.1 and 3.2, the Wick expansions for vacuum-normal single and double excitations in terms of  $\Phi$ -normal ones look as follows. The signs are determined by pairing up contracted indices using eqs (3.5) and (3.6).

$$\begin{aligned} a_q^p &= \tilde{a}_q^p + \tilde{a}_{q^{\circ}}^{p^{\circ}} = \tilde{a}_q^p + \gamma_q^p \\ a_{rs}^{pq} &= \tilde{a}_{rs}^{pq} + \tilde{a}_{r^{\circ}s}^{p^{\circ}q} + \tilde{a}_{r^{\circ}s^{\circ}}^{p^{\circ}q^{\circ}} + \tilde{a}_{r^{\circ}s}^{p^{\circ}q^{\circ}} + \tilde{a}_{r^{\circ}s^{\circ}}^{p^{\circ}q^{\circ}} + \tilde{a}_{r^{\circ}s^{\circ}}^{p^{\circ}q^{\circ}} \\ &= \tilde{a}_{rs}^{pq} + \gamma_r^p \tilde{a}_s^q - \gamma_s^p \tilde{a}_r^q - \gamma_r^q \tilde{a}_s^p + \gamma_s^q \tilde{a}_r^p + \gamma_r^p \gamma_s^q - \gamma_s^p \gamma_r^q \end{aligned}$$

<sup>1</sup>The dot notation is borrowed from physics: [https://en.wikipedia.org/wiki/Wick's\\_theorem#Definition\\_of\\_contraction](https://en.wikipedia.org/wiki/Wick's_theorem#Definition_of_contraction)

**Definition 3.1. Riffle shuffle permutations.** Let  $S_R$  be the symmetric group on  $R \equiv (r_1, \dots, r_k)$ , comprising all  $k!$  permutations of this tuple. For each integer composition<sup>2</sup>  $(k_1, \dots, k_m)$  of  $k$ , the  $(k_1, \dots, k_m)$ -shuffles of  $R$  partition the tuple into blocks  $R_1 \cup \dots \cup R_m$  of the form  $R_i \equiv (r_{h_i+1}, \dots, r_{h_i+k_i})$  with  $h_i \equiv \sum_{j=1}^{i-1} k_j$  and interleave them in all possible ways. The cardinality of this subset  $S_R^{(k_1, \dots, k_m)} \subseteq S_R$  is given by the following multinomial coefficient

$$|S_R^{(k_1, \dots, k_m)}| \equiv \frac{k!}{k_1! \dots k_m!} \quad (3.8)$$

since for each shuffle there are  $k_1! \dots k_m!$  permutations in  $S_R$  derived by permuting the elements within each block. These are termed *riffle shuffle permutations*,<sup>3</sup> by analogy with the process of cutting and interleaving a deck of cards.

**Definition 3.2. Index antisymmetrizers.** Let  $\hat{P}_{(r_1/\dots/r_k)}$  be the full antisymmetrizer for the indices in  $R$ .<sup>4 5</sup> More generally, let  $\hat{P}_{(R_1/\dots/R_m)}$  be a reduced antisymmetrizer, which antisymmetrizes a term in  $R$  that it is already antisymmetric in each block,  $R_i$ . Reduced antisymmetrization is achieved by summing over riffle shuffles.

$$\hat{P}_{(R_1/\dots/R_m)} \equiv \sum_{\pi \in S_R^{(k_1, \dots, k_m)}} \varepsilon_\pi \hat{\pi} \quad \hat{\pi} t_{r_1 \dots r_k} \equiv t_{\pi(r_1) \dots \pi(r_k)} \quad (3.9)$$

To antisymmetrize multiple disjoint sets of indices, we use the compact notation  $\hat{P}_{(X_1|\dots|X_l)}^{(Y_1|\dots|Y_m)} \equiv \hat{P}^{(X_1)} \dots \hat{P}^{(X_l)} \hat{P}^{(Y_1)} \dots \hat{P}^{(Y_m)}$ , where each  $X_i$  or  $Y_i$  stands for a reduced antisymmetrizer argument  $R_1/\dots/R_m$  for a different set of indices.

**Example 3.2.** Permutation operators  $\hat{\pi}$  can be written as products of transpositions  $(pq)t_{pq} \equiv t_{qp}$ . In this notation,

$$\hat{P}_{(p/q)} = 1 - (pq) \quad \hat{P}_{(p/q/r)} = 1 - (pq) - (pr) - (qr) + (pq)(qr) + (pr)(qr) \quad \hat{P}_{(p/qr)} = 1 - (pq) + (pr)(qr)$$

are the unique antisymmetrizers for two and three indices. Note that  $\hat{P}_{(p/qr)}\hat{P}_{(q/r)} = \hat{P}_{(p/q/r)}$ , which is an immediate consequence of the definition of the reduced antisymmetrizer.

**Example 3.3.** Using index antisymmetrizers, the second Wick expansion in ex 3.1 can be expressed even more compactly.

$$a_{rs}^{pq} = \tilde{a}_{rs}^{pq} + \hat{P}_{(r/s)}^{(p/q)} \tilde{a}_{r^\circ s}^{p^\circ q} + \hat{P}_{(r/s)} \tilde{a}_{r^\circ s^\circ}^{p^\circ q^\circ} = \tilde{a}_{rs}^{pq} + \hat{P}_{(r/s)}^{(p/q)} \gamma_r^p \tilde{a}_{rs}^{pq} + \hat{P}_{(r/s)} \gamma_r^p \gamma_s^q$$

Note that permuting  $p$  and  $q$  in the second term, for example, is equivalent to “moving the contraction dot” from one to the other, since the excitation operators are antisymmetric in their upper and lower indices:  $(pq)\tilde{a}_{r^\circ s}^{p^\circ q} = -\tilde{a}_{r^\circ s}^{q^\circ p} = \tilde{a}_{r^\circ s}^{p^\circ q^\circ}$ . In general, this allows us to reduce Wick expansions to a sum over unique contraction “patterns”, with the remaining contractions generated from these by index antisymmetrizers. The next two examples show how this works in practice.

**Example 3.4.** Reduced antisymmetrizers come into play in the Wick expansion of a triple excitation operator.

$$a_{stu}^{pqr} = \tilde{a}_{stu}^{pqr} + \hat{P}_{(s/tu)}^{(p/qr)} \tilde{a}_{s^\circ tu}^{p^\circ qr} + \hat{P}_{(st/u)}^{(p/q/r)} \tilde{a}_{s^\circ t^\circ u}^{p^\circ q^\circ r} + \hat{P}_{(p/q/r)} \tilde{a}_{s^\circ t^\circ u}^{p^\circ q^\circ r^\circ} \quad (3.10)$$

The permutation factors are chosen as follows. In the second term, permutations of  $qr$  and  $tu$  are omitted because they do not produce unique contractions. In the third term, we can include permutations of either  $pq$  or  $st$  but not both. Otherwise, we would double-count terms like  $\tilde{a}_{s^\circ t^\circ u}^{p^\circ q^\circ r}$  by also including  $\tilde{a}_{s^\circ t^\circ u}^{q^\circ p^\circ r}$ , which represents the same contraction. The same reasoning applies to the last term, where we can either antisymmetrize  $pqr$  or  $stu$ .

**Example 3.5. Derivation of CIS equations in KM notation.** A programmable expression for the CI singles Hamiltonian matrix elements,  $\langle \Phi_i^a | H_c | \Phi_j^b \rangle$ , can be derived in KM notation as follows.

$$\begin{aligned} \langle \Phi | \tilde{a}_a^i \tilde{a}_q^p \tilde{a}_j^b | \Phi \rangle &= :a_a^i \bullet a_q^p \bullet a_j^b : + :a_a^i \bullet a_q^{p^\circ} a_j^{b^\circ} : = \gamma_j^i \eta_a^p \eta_q^b - \gamma_q^i \gamma_j^p \eta_a^b \\ \langle \Phi | \tilde{a}_a^i \tilde{a}_{rs}^{pq} \tilde{a}_j^b | \Phi \rangle &= \hat{P}_{(r/s)}^{(p/q)} :a_a^i \bullet a_r^{p^\circ} a_s^{q^\circ} a_j^{b^\circ} : = \hat{P}_{(r/s)}^{(p/q)} \gamma_r^i \eta_a^p \gamma_s^q \eta_j^b \\ \implies \langle \Phi_i^a | H_c | \Phi_j^b \rangle &= f_p^a (\gamma_j^i \eta_a^p \eta_q^b - \gamma_q^i \gamma_j^p \eta_a^b) + \frac{1}{4} \bar{g}_{pq}^{rs} \left( \hat{P}_{(r/s)}^{(p/q)} \gamma_r^i \eta_a^p \gamma_s^q \eta_j^b \right) = f_a^b \gamma_j^i - f_j^i \eta_a^b + \bar{g}_{aj}^{ib} \end{aligned}$$

which simplifies to  $\langle \Phi_i^a | H_c | \Phi_j^b \rangle = f_a^b \delta_j^i - f_j^i \delta_a^b + \bar{g}_{aj}^{ib}$ . Here, we have used the fact that  $\tilde{a}_i^a = a_i^a$  and  $\tilde{a}_a^i = a_a^i$ .

<sup>2</sup>See [https://en.wikipedia.org/wiki/Composition\\_\(combinatorics\)](https://en.wikipedia.org/wiki/Composition_(combinatorics))

<sup>3</sup>See [https://en.wikipedia.org/wiki/Riffle\\_shuffle\\_permutation](https://en.wikipedia.org/wiki/Riffle_shuffle_permutation)

<sup>4</sup>Note that we are borrowing some definitions from def 3.1, such as  $R$  and  $R_i$ .

<sup>5</sup>For example, we can express the position-space Slater determinant  $\Phi_{(p_1 \dots p_n)}(1, \dots, n)$  as  $\frac{1}{\sqrt{n!}} \hat{P}_{(p_1/\dots/p_n)} \psi_{p_1}(1) \dots \psi_{p_n}(n)$ .

**Example 3.6. Derivation of CID equations in KM notation.** Projecting the CI doubles Schrödinger equation,  $H_c \Psi = E_c \Psi$  where  $\Psi = (1 + \frac{1}{4} c_{cd}^{kl} \tilde{a}_{kl}^{cd}) \Phi$ , by  $\Phi$  and  $\Phi_{ij}^{ab}$  gives a system of linear equations

$$\begin{aligned} E_c &= \langle \Phi | H_c (1 + \frac{1}{4} c_{cd}^{kl} \tilde{a}_{kl}^{cd}) | \Phi \rangle & \implies & E_c = \frac{1}{4} \langle \Phi | H_c | \Phi_{kl}^{cd} \rangle c_{cd}^{kl} \\ E_c c_{ab}^{ij} &= \langle \Phi_{ij}^{ab} | H_c (1 + \frac{1}{4} c_{cd}^{kl} \tilde{a}_{kl}^{cd}) | \Phi \rangle & \implies & E_c c_{ab}^{ij} = \langle \Phi_{ij}^{ab} | H_c | \Phi \rangle + \frac{1}{4} \langle \Phi_{ij}^{ab} | H_c | \Phi_{kl}^{cd} \rangle c_{cd}^{kl} \end{aligned}$$

which can be simplified into programmable expressions as follows.

$$\begin{aligned} \langle \Phi | \tilde{a}_{rs}^{pq} \tilde{a}_{kl}^{cd} | \Phi \rangle &= \hat{P}_{(r/s)}^{(p/q)} \gamma_k^p \gamma_l^q \eta_r^c \eta_s^d \\ \langle \Phi | \tilde{a}_{ab}^{ij} \tilde{a}_q^p \tilde{a}_{kl}^{cd} | \Phi \rangle &= \hat{P}_{(a/b|k/l)}^{(c/d)} \mathbf{:} a_{a \bullet 1 b \bullet 3}^{i \circ 1 j \circ 2} a_{q \bullet 2}^{p \bullet 1} a_{k \circ 1 l \circ 2}^{c \bullet 2 d \bullet 3} \mathbf{:} + \hat{P}_{(k/l)}^{(i/j|c/d)} \mathbf{:} a_{a \bullet 1 b \bullet 2}^{i \circ 1 j \circ 3} a_{q \circ 1}^{p \circ 2} a_{k \circ 2 l \circ 3}^{c \bullet 1 d \bullet 2} \mathbf{:} \\ &= \hat{P}_{(a/b|k/l)}^{(c/d)} \eta_a^p \eta_q^c \eta_b^d \gamma_k^i \gamma_l^j - \hat{P}_{(k/l)}^{(i/j|c/d)} \gamma_q^i \gamma_k^p \gamma_l^j \eta_a^c \eta_b^d \\ \langle \Phi | \tilde{a}_{ab}^{ij} \tilde{a}_{rs}^{pq} \tilde{a}_{kl}^{cd} | \Phi \rangle &= \hat{P}_{(a/b|k/l)}^{(c/d)} \mathbf{:} a_{a \bullet 1 b \bullet 2}^{i \circ 1 j \circ 2} a_{r \bullet 3 s \bullet 4}^{p \bullet 1 q \bullet 2} a_{k \circ 1 l \circ 2}^{c \bullet 3 d \bullet 4} \mathbf{:} + \hat{P}_{(k/l)}^{(i/j|c/d)} \mathbf{:} a_{a \bullet 1 b \bullet 2}^{i \circ 1 j \circ 2} a_{r \circ 1 s \circ 2}^{p \circ 3 q \circ 4} a_{k \circ 3 l \circ 4}^{c \bullet 1 d \bullet 2} \mathbf{:} + \hat{P}_{(r/s|k/l|a/b)}^{(p/q|i/j|c/d)} \mathbf{:} a_{a \bullet 1 b \bullet 3}^{i \circ 1 j \circ 3} a_{r \circ 1 s \bullet 2}^{p \circ 2 q \bullet 1} a_{k \circ 2 l \circ 3}^{c \bullet 2 d \bullet 3} \mathbf{:} \\ &= \hat{P}_{(a/b|k/l)}^{(c/d)} \eta_a^p \eta_b^q \eta_r^c \eta_s^d \gamma_k^i \gamma_l^j + \hat{P}_{(k/l)}^{(i/j|c/d)} \gamma_r^i \gamma_s^j \gamma_k^p \gamma_l^q \eta_a^c \eta_b^d - \hat{P}_{(r/s|k/l|a/b)}^{(p/q|i/j|c/d)} \gamma_r^i \gamma_k^p \gamma_l^q \eta_a^c \eta_b^d \\ \implies \quad E_c &= \frac{1}{4} \bar{g}_{kl}^{cd} c_{cd}^{kl} & E_c c_{ab}^{ij} &= \bar{g}_{ab}^{ij} + \hat{P}_{(a/b)} f_a^c c_{cb}^{ij} - \hat{P}^{(i/j)} f_k^i c_{ab}^{kj} + \frac{1}{2} \bar{g}_{ab}^{cd} c_{cd}^{ij} + \frac{1}{2} \bar{g}_{kl}^{ij} c_{ab}^{kl} + \hat{P}_{(a/b)}^{(i/j)} \bar{g}_{ak}^{ic} c_{bc}^{jk} \end{aligned}$$