1. In solving the canonical Hartree-Fock equations, it is common to use eigenfunctions of the core Hamiltonian as starting guesses for the canonical molecular orbitals. This eigenvalue equation looks as follows

$$\hat{h}\phi_p = \epsilon_p \phi_p$$

where  $\phi_p$  is a spatial function. Project this equation by  $\chi_{\mu}$  and then expand  $\phi_p$  as a linear combination of atomic orbitals,  $\phi_p = \sum_{\nu} \chi_{\nu} C_{\nu p}$ , in order to arrive at the following matrix equation for the core guess.

$$\mathbf{HC} = \mathbf{SC}\boldsymbol{\epsilon} \qquad (\mathbf{H})_{\mu\nu} = \langle \chi_{\mu} | \hat{h} | \chi_{\nu} \rangle \qquad (\mathbf{S})_{\mu\nu} = \langle \chi_{\mu} | \chi_{\nu} \rangle \qquad (\mathbf{C})_{\mu p} = C_{\mu p} \qquad (\boldsymbol{\epsilon})_{pq} = \epsilon_{p} \delta_{pq}$$

For extra credit, you may also briefly answer the following question in words: For an open-shell system, why do we end up with different spatial orbitals for different spins, even though the initial core-guess orbitals are the same for each spin?

2. Derive the expansion of H in terms of  $\Phi$ -normal-ordered excitations.

3. Derive the second Slater rule using Wick's theorem, and explain why this matrix element evaluates to zero for canonical Hartree-Fock orbitals.

$$\langle \Phi | H | \Phi_i^a \rangle = ?$$