

4 Diagram notation

Notation 4.1. Diagram notation. In diagram notation, particle-hole operators are written as oriented lines extending from a vertex. Particle annihilation operators enter the vertex from below, particle creation operators leave the vertex at the top, and single-excitation operators have both creation and annihilation lines. Contractions are represented by joining particle-hole lines with compatible position and orientation.

$$\begin{array}{cccc} \circ \equiv a_p & \uparrow \equiv a_p^\dagger & \uparrow \circ \equiv a_p^\dagger a_q = a_q^p & \circ \equiv \overline{a_p} a_q^\dagger = a_p^q \bullet \end{array} \quad (4.1)$$

Quasiparticle operators are distinguished by the use of closed-circle vertices, with particle lines pointing upward and with hole lines pointing downward. Single-excitation operators split into four cases representing the virtual and occupied blocks of a_q^p . Internal contractions of single excitations (*bubble contractions*) are implicitly taken to be hole contractions.

$$\begin{array}{cccccc} \bullet \equiv b_a & \uparrow \bullet \equiv b_a^\dagger & \downarrow \bullet \equiv b_i & \downarrow \bullet \equiv b_i^\dagger & \uparrow \bullet \equiv \overline{b_a} b_b^\dagger = a_a^{b\bullet} & \downarrow \bullet \equiv \overline{b_i} b_j^\dagger = a_j^{i\circ} \end{array} \quad (4.2)$$

$$\begin{array}{cccccc} \uparrow \bullet \equiv b_a^\dagger b_b = a_b^a & \swarrow \searrow \bullet \equiv b_a^\dagger b_i^\dagger = a_i^a & \swarrow \bullet \equiv b_i b_a = a_a^i & \downarrow \bullet \equiv b_i b_j^\dagger = a_j^i & \circ \equiv \overline{b_i} b_i^\dagger = a_i^{i\circ} \end{array} \quad (4.3)$$

Normal-ordered products of single excitation operators are joined with a dotted line, forming higher excitation operators.

$$\begin{array}{c} \begin{array}{c} \uparrow \bullet \cdots \uparrow \bullet \\ \downarrow \bullet \cdots \downarrow \bullet \end{array} \equiv :a_{q_1}^{p_1} a_{q_2}^{p_2} \cdots a_{q_m}^{p_m}: = a_{q_1 q_2 \cdots q_m}^{p_1 p_2 \cdots p_m} \\ \text{m times} \end{array} \quad \begin{array}{c} \overbrace{\begin{array}{c} \uparrow \bullet \cdots \uparrow \bullet \\ \downarrow \bullet \cdots \downarrow \bullet \end{array}}^{m \text{ times}} \quad \underbrace{\begin{array}{c} \uparrow \bullet \cdots \uparrow \bullet \\ \downarrow \bullet \cdots \downarrow \bullet \end{array}}_{n \text{ times}} \equiv :a_{q_1 q_2 \cdots q_m}^{p_1 p_2 \cdots p_m} a_{s_1 s_2 \cdots s_n}^{r_1 r_2 \cdots r_n}: \end{array}$$

Contracted operators are implicitly normal ordered together. Φ -normal-ordering is indicated by the use of double-circle vertices, \odot and \bullet instead of \circ and \bullet . These building blocks can be used to construct more complicated products. Products of diagrams are written from top to bottom, representing left-to-right ordering in the corresponding algebraic expression.

Definition 4.1. m -electron operators in diagram notation. The building blocks of a graph are m -electron operators, which can be represented in two equivalent ways. The *Goldstone representation* depicts an operator as a label attached to the corresponding excitation operator, whereas the *Hugenholtz representation* depicts the operator as a single vertex with m outgoing and incoming lines. Note that $(\frac{1}{m!})^2 \sum_{\text{Einstein}}$ is baked into the definition (see def 4.2 and axiom 4.1).

$$\begin{array}{c} \boxed{v} \begin{array}{c} \uparrow \bullet \cdots \uparrow \bullet \\ \downarrow \bullet \cdots \downarrow \bullet \end{array} \equiv \left(\frac{1}{m!}\right)^2 \sum_{\text{Einstein}} \overline{v}_{p_1 \cdots p_m}^{q_1 \cdots q_m} a_{q_1 \cdots q_m}^{p_1 \cdots p_m} \equiv \begin{array}{c} \cdots \\ \swarrow \bullet \cdots \bullet \searrow \\ \bullet \end{array} \end{array} \quad \begin{array}{c} \boxed{v} \begin{array}{c} p_1 \quad p_m \\ \uparrow \bullet \cdots \uparrow \bullet \\ q_1 \quad q_m \end{array} = \overline{v}_{p_1 \cdots p_m}^{q_1 \cdots q_m} a_{q_1 \cdots q_m}^{p_1 \cdots p_m} \end{array} \quad (4.4)$$

$$\begin{array}{c} \boxed{v} \begin{array}{c} \uparrow \bullet \cdots \uparrow \bullet \\ \downarrow \bullet \cdots \downarrow \bullet \end{array} \equiv \left(\frac{1}{m!}\right)^2 \sum_{\text{Einstein}} \overline{v}_{p_1 \cdots p_m}^{q_1 \cdots q_m} \tilde{a}_{q_1 \cdots q_m}^{p_1 \cdots p_m} \equiv \begin{array}{c} \cdots \\ \swarrow \bullet \cdots \bullet \searrow \\ \bullet \end{array} \end{array} \quad \begin{array}{c} \begin{array}{c} p_1 \quad p_m \\ \swarrow \bullet \cdots \bullet \searrow \\ q_1 \quad q_m \end{array} \equiv \overline{v}_{p_{\pi(1)} \cdots p_{\pi(m)}}^{q_{\sigma(1)} \cdots q_{\sigma(m)}} a_{q_{\sigma(1)} \cdots q_{\sigma(m)}}^{p_{\pi(1)} \cdots p_{\pi(m)}} \end{array} \quad (4.5)$$

The labeled diagrams on the right represent just the summand of the operator, which highlights the difference between representations. Both summands correspond to an excitation operator weighted by its antisymmetrized interaction tensor,¹ but whereas the Goldstone summand specifies an ordering for the indices of its corresponding algebraic term, the Hugenholtz summand does not. Since the phases of $\overline{v}_{p_1 \cdots p_m}^{q_1 \cdots q_m}$ and $a_{q_1 \cdots q_m}^{p_1 \cdots p_m}$ cancel under index permutation, the two labeled diagrams are actually equal – a Hugenholtz summand can be expanded into a Goldstone summand by simply choosing an arbitrary ordering for the indices. In practice, the symmetry of the Hugenholtz operator simplifies the enumeration of Wick expansions whereas the Goldstone operator makes it easier to evaluate a graph's overall phase.

¹In the original paper [J. Goldstone, *P. Roy. Soc. A* **239**, (1957)], Goldstone's diagrams were actually defined in terms of non-antisymmetrized integrals. The *antisymmetrized Goldstone diagrams* used here are sometimes called *Brandow diagrams*.

Definition 4.2. Graph. A graph² $G = (O, L, h, t)$ consists of a set of m -electron operators O , a set of lines L , and two mappings $h, t : L \rightarrow O$ that return the *head* $h(l) \in O$ and *tail end* $t(l) \in O$ of each line in L . Two lines are considered *equivalent* if they share the same head and tail. Here, we allow for *external lines* in which one of the ends is e , the *free end*, which is formally considered a member of O . Lines with no free end are termed *internal*. A graph is termed *closed* if it contains no external lines, representing a scalar-valued algebraic term. Otherwise, the graph is *open* and represents an operator-valued algebraic term. The *rules of interpretation* for translating G into an algebraic expression are given below.

Definition 4.3. Connected and linked graphs. Two lines are *adjacent* if they share a non-free end. In a Goldstone diagram, two lines are termed *Goldstone adjacent* if they end on the same single-excitation vertex of an operator. In this context, the broader sense of the adjacency is called *Hugenholtz adjacency*. A *path* is a sequence of lines (l_1, \dots, l_n) such that l_i is adjacent to l_{i+1} and no line is repeated. In a Goldstone diagram, one can distinguish between *Goldstone* and *Hugenholtz paths* depending on whether or not the lines are all Goldstone adjacent. A *Goldstone cycle* is a Goldstone path whose ends are either Goldstone adjacent or free. A Goldstone cycle with free ends is an *open cycle*. Otherwise, it forms a *loop*. A graph is considered *connected* if there is a Hugenholtz path connecting any two of its operators. A closed, disconnected subgraph is considered *unlinked* from the rest of the graph. Note this last distinction: A *disconnected* graph is still classified as *linked*, as long as none of its disconnected subgraphs are closed.

Definition 4.4. Equivalent subgraphs. Repeated copies of the same operator are formally distinguished as elements of a graph's operator set and are termed *identical operators*. Operators that can be interchanged without altering the graph are termed *interchangeable operators*. This occurs when they are connected to the same set of operators in the same way. When identical operators are also interchangeable, we call them *equivalent operators*. More generally, if $G[O']$ denotes the subgraph associated with³ some subset $O' \subset O$ of the operators in G , then two disjoint subgraphs⁴ can also be classified as *identical*, *interchangeable*, or *equivalent*. For our purposes, these terms will be restricted to *connected subgraphs*.

Definition 4.5. Summand graph. A *summand graph*⁵ $\Sigma(G) = (G, S, s)$ associates each line l in G with a symbol $s(l) \in S$ through the *label mapping* $s : L \rightarrow S$. Pictorially, this corresponds to labeling each line in G with spin-orbital index, p, q, r, s , etc. $\Sigma(G)$ translates directly into an algebraic summand according to Def 4.1 and Notation 4.1.

Definition 4.6. Degeneracy. The *line degeneracy* or simply *degeneracy* of G is the number of permutational symmetries in $\Sigma(G)$, a positive integer here denoted $\text{dg}(G)$. Formally, this can be defined as follows. If $S = \{s_1, \dots, s_n\}$ is the label set of $\Sigma(G)$ and $s(l_i) = s_i$ is its label map, we can define a new summand graph $\Sigma_\pi(G) = (G, S, s_\pi)$ with a permuted label map s_π given by $s_\pi(l_i) = s_{\pi(i)}$, where π is a permutation in S_n . Then $\text{dg}(G)$ is the number of $\Sigma_\pi(G)$ that are equal to $\Sigma(G)$. Assuming no identical operators, the degeneracy of G is simply given by $\text{dg}(G) = |L_1|! \cdots |L_h|!$ where $L_1 \cup \cdots \cup L_h$ partitions its line set into subsets of equivalent lines and $|L_i|$ denotes the number of elements in L_i . When the graph does contain identical operators, there is an additional factor of $k!$ for each set $\{G[O_1], \dots, G[O_k]\}$ of equivalent subgraphs.

Axiom 4.1. Rules of interpretation. The algebraic interpretation of a graph G is obtained from $\Sigma(G)$ as follows.

1. Multiply $\Sigma(G)$ by $\text{dg}(G)^{-1}$, the *degeneracy factor*.
2. Sum each index in $\Sigma(G)$ over its range.

Symbolically, these rules can be stated as follows: $G = \frac{1}{\text{dg}(G)} \sum_{\text{Einstein}} \Sigma(G)$.

Definition 4.7. Contraction. Formally, a *graph contraction* is a map $G \mapsto c(G)$ joining one or more compatible external lines in G . For example, c might replace l_1 and l_2 , which have ends $(t(l_1), h(l_1)) = (o_1, e)$ and $(t(l_2), h(l_2)) = (e, o_2)$, with l_{12} , which has ends $(t(l_{12}), h(l_{12})) = (o_1, o_2)$. Two contractions c and c' of G are *equivalent* if they are graphically indistinguishable, i.e. $c(G) = c'(G)$. The number of equivalent ways of achieving a given contraction c is called its *pattern degeneracy*, which is denoted $\text{pat}(c)$. The complete set of unique contraction patterns for a graph is here denoted $\text{Ctr}(G)$. Combined with Wick's theorem, these concepts lead to the following statement

$$\sum_{\text{Einstein}} \Sigma(G) = \sum_{\text{Einstein}} \text{:}\Sigma(G)\text{:} + \sum_{c \in \text{Ctr}(G)} \text{pat}(c) \sum_{\text{Einstein}} \text{:}\Sigma(c(G))\text{:} \quad (4.6)$$

which will be used to prove a more elegant graphical formulation of Wick's theorem below.

²In graph theory jargon this is essentially a *directed multigraph*, except that the vertical ordering of operators matters.

³See https://en.wikipedia.org/wiki/Induced_subgraph

⁴That is, $G[O_1]$ and $G[O_2]$ with $O_1 \cap O_2 = \emptyset$.

⁵In graph theory jargon this is an *edge-labeled directed multigraph*.

Lemma 4.1. *The pattern degeneracy of a contraction c is $\text{pat}(c) = \frac{\text{dg}(G)}{\text{dg}(c(G))}$.*

Proof[†]: Assume G has no identical operators, so that equivalent subgraphs do not contribute to the line degeneracies. If $L = L_1 \cup \dots \cup L_m$ partitions G 's line set into equivalent lines, let $L_{i,j} = L_{j,i}$ denote lines in $c(G)$ produced by contracting lines from L_i and L_j , and let $L_{i,0}$ denote lines from L_i unaffected by c . Then the number of equivalent ways to partition L_i for contraction is $\frac{|L_i|!}{|L_{i,0}|!|L_{i,1}|!\dots|L_{i,m}|!}$, and the number of ways of forming each $L_{i,j}$ from a given partition of L_i and L_j is $|L_{i,j}|!$. Therefore, the total pattern degeneracy of c is given by the following.

$$\text{pat}(c) = \prod_{i=1}^m \frac{|L_i|!}{|L_{i,0}|!|L_{i,1}|! \dots |L_{i,m}|!} \cdot \prod_{i>j}^m |L_{i,j}|! = \frac{|L_1|! \dots |L_m|!}{\prod_{i=1}^m |L_{i,0}|!|L_{i,1}|! \dots |L_{i,i-1}|!} = \frac{\text{dg}(G)}{\text{dg}(c(G))}$$

The final step follows from the fact that, according to our definitions, $\bigcup_{i=1}^m (L_{i,0} \cup L_{i,1} \cup \cdots L_{i,i-1})$ partitions the line set of $c(G)$ into equivalent lines and therefore the denominator equals $\text{dg}(c(G))$.

Remark 4.1. The assumption that G contains no identical operators limits the scope of lemma 4.1 and of the next theorem, which derives from it. We can generalize this result to allow identical operators if we assume that the contraction c does not create any new pairs of equivalent subgraphs.⁶ The proof is as follows: Each set $E = \{O_1, \dots, O_k\}$ of equivalent subgraph generators⁷ for G partitions into $E_1 \cup \dots \cup E_n$, where $E_i \subset E$ generates a set of equivalent subgraphs in $c(G)$. By our assumption, tallying up these partitions for each E exhausts the sets of equivalent subgraph generators for $c(G)$. Then for each E there will be a factor of $|E|!$ in the line degeneracy of G , a factor of $|E_1|! \dots |E_n|!$ in the line degeneracy of $c(G)$, and a factor of $\frac{|E|!}{|E_1|! \dots |E_n|!}$ in the pattern degeneracy of the contraction, so that lemma 4.1 still holds.

Theorem 4.1. Wick’s theorem for Graphs. $G = \mathfrak{z}G\mathfrak{z} + \sum_{c \in \text{Ctr}(G)} \mathfrak{z}c(G)\mathfrak{z}$

Proof: This follows by expanding G according to axiom 4.1 and then using equation 4.6 and lemma 4.1. \square

$$G = \frac{1}{\deg(G)} \sum_{\text{Einstein}} \Sigma(G) = \frac{1}{\deg(G)} \sum_{\text{Einstein}} :\Sigma(G): + \sum_{c \in \text{Ctr}(G)} \frac{1}{\deg(c(G))} \sum_{\text{Einstein}} :\Sigma(c(G)): = :G: + \sum_{c \in \text{Ctr}(G)} :c(G):$$

Derivation 4.1. Phase rule for graphs with contractions. The summand graph of a pair of m -electron operators with a single particle or hole contraction has the following form in the Goldstone representation

$$\begin{array}{c}
\begin{array}{c} p_1 \\ \uparrow \\ \bigcirc \\ \downarrow \\ q_1 \end{array} \cdots \begin{array}{c} p_i \\ \uparrow \\ \bullet \\ \downarrow \\ d \end{array} \cdots \begin{array}{c} p_m \\ \uparrow \\ \bigcirc \\ \downarrow \\ q_m \end{array} \begin{array}{c} \boxed{v} \end{array} \\
\begin{array}{c} \boxed{w} \end{array} \begin{array}{c} \uparrow \\ \bigcirc \\ \downarrow \\ s_1 \end{array} \cdots \begin{array}{c} \uparrow \\ \bullet \\ \downarrow \\ s_j \end{array} \cdots \begin{array}{c} \uparrow \\ \bigcirc \\ \downarrow \\ s_n \end{array}
\end{array}
= \overline{v}_{p_1 \cdots p_i \cdots p_m}^{q_1 \cdots d \cdots q_m} \overline{w}_{r_1 \cdots r_j \cdots r_n}^{s_1 \cdots s_j \cdots s_n} \cdot \\
\vdots a_{q_1 \cdots d \cdots q_m}^{p_1 \cdots p_i \cdots p_m} a_{s_1 \cdots s_j \cdots s_n}^{d \cdots r_n} \vdots
\end{array}
\quad
\begin{array}{c}
\begin{array}{c} p_1 \\ \uparrow \\ \bigcirc \\ \downarrow \\ q_1 \end{array} \cdots \begin{array}{c} q_i \\ \downarrow \\ \bullet \\ \downarrow \\ k \end{array} \cdots \begin{array}{c} p_m \\ \uparrow \\ \bigcirc \\ \downarrow \\ q_m \end{array} \begin{array}{c} \boxed{v} \end{array} \\
\begin{array}{c} \boxed{w} \end{array} \begin{array}{c} \uparrow \\ \bigcirc \\ \downarrow \\ s_1 \end{array} \cdots \begin{array}{c} \uparrow \\ \bullet \\ \downarrow \\ r_j \end{array} \cdots \begin{array}{c} \uparrow \\ \bigcirc \\ \downarrow \\ s_n \end{array}
\end{array}
= \overline{v}_{p_1 \cdots k \cdots p_m}^{q_1 \cdots q_i \cdots q_m} \overline{w}_{r_1 \cdots r_j \cdots r_n}^{s_1 \cdots k \cdots s_n} \cdot \\
\vdots a_{q_1 \cdots q_i \cdots q_m}^{p_1 \cdots k \cdots p_m} a_{s_1 \cdots k \cdots s_n}^{r_1 \cdots r_j \cdots r_n} \vdots
\end{array}
\quad (4.7)$$

where d is a virtual index, k is an occupied index, and we are not using implicit summation. Using the permutational degrees of freedom of the normal-ordered product, we can bring together the contracted pair of single-excitation operators without changing the sign of the expression. The contracted pair of operators can then be eliminated as follows.

$$:a_{d^{\bullet}s_j}^{p_i} a_{s_j}^{d^{\bullet}}:: = -:a_{d^{\bullet}s_j}^{d^{\bullet}p_i}:: = -(\eta_d^d) a_{s_j}^{p_i} = a_{s_j}^{p_i} \quad :a_{q_i}^{k^{\circ}} a_{k^{\circ}q_i}^{r_j}:: = -:a_{k^{\circ}q_i}^{k^{\circ}r_j}:: = -(\gamma_k^k) a_{q_i}^{r_j} = -a_{q_i}^{r_j} \quad (4.8)$$

Applying this rule to each internal line in an open cycle leaves $(-)^h a_q^p$ where h is the number of hole contractions in the cycle and p and q are indices of the free ends. Applying the same procedure to a loop gives $(-)^h a_{d\bullet}^d = (-)^{h(-\eta_d^d)} = (-)^{h+1}$. Since the lines of any Goldstone graph uniquely partition into Goldstone cycles, these rules provide a quick way to resolve the operator component of any diagram into a normal-ordered product of uncontracted operators times a phase factor.

Corollary 4.1. Phase rule for closed graphs. *After eliminating all contracted operator pairs, the sign of a closed graph is $(-)^{h+l}$, where h is the total number of holes in the graph and l is the number of loops.*

Proof: By definition, every Goldstone cycle in a closed graph is a loop. Let h_i be the number of holes in the i^{th} loop. Then, by derivation 4.1, the fully contracted operator product evaluates to $\prod_{i=1}^l (-)^{h_i+1} = (-)^{\sum_i h_i+l} = (-)^{h+l}$.

[‡]Note that this proof is limited in scope. See [rmk 4.1](#).

⁶This encompasses most cases of interest, but it is not completely general, so be cautious about when you apply thm 4.1.

⁷By this we mean that O_1, \dots, O_k are connected subsets of O such that $G[O_1], \dots, G[O_k]$ are equivalent subgraphs.

Definition 4.8. Coefficient graph. A *coefficient graph* is a closed graph in which one of the interaction tensors is an antisymmetrized Kronecker delta, $\bar{\delta}_{p'_1 \dots p'_m}^{q'_1 \dots q'_m} = \hat{P}_{(q_1/\dots/q_m)}^{(p_1/\dots/p_m)} \delta_{p'_1}^{p_1} \dots \delta_{p'_m}^{p_m} \delta_{q'_1}^{q_1} \dots \delta_{q'_m}^{q_m}$.⁸ Using the identity

$$\tilde{a}_{q_1 \dots q_m}^{p_1 \dots p_m} = \left(\frac{1}{m!}\right)^2 \delta_{p'_1 \dots p'_m}^{q'_1 \dots q'_m} \tilde{a}_{q'_1 \dots q'_m}^{p'_1 \dots p'_m} \quad (4.9)$$

this extends the results derived above to allow for bare excitation operators in addition to m -electron operators. The contraction with $\tilde{a}_{q_1 \dots q_m}^{p_1 \dots p_m}$ is depicted by capping the lines connected to it, \rightarrow , and labelling these creation and annihilation lines with p_1, \dots, p_m and q_1, \dots, q_m , respectively. To simplify the algebraic interpretation of a coefficient graph (see rmk 4.2), it becomes convenient to treat *coefficient lines* as distinct from other internal lines. A closed graph with no coefficient lines is sometimes called an *energy graph*.

Remark 4.2. The algebraic interpretation of a coefficient graph with the operator in eq 4.9 has the form

$$\delta_{p'_1 \dots p'_m}^{q'_1 \dots q'_m} T_{q'_1 \dots q'_m}^{p'_1 \dots p'_m} = \hat{P}_{(q_1/\dots/q_m)}^{(p_1/\dots/p_m)} T_{q_1 \dots q_m}^{p_1 \dots p_m} \quad (4.10)$$

where $T_{q'_1 \dots q'_m}^{p'_1 \dots p'_m}$ is the product of interaction tensors contracted with $\bar{\delta}_{p'_1 \dots p'_m}^{q'_1 \dots q'_m}$, scaled by appropriate sign and degeneracy factors according to axiom 4.1. Suppose $\{p_1, \dots, p_m\}$ and $\{q_1, \dots, q_m\}$ can be partitioned as $P_1 \cup \dots \cup P_h$ and $Q_1 \cup \dots \cup Q_k$ where $P_i = \{p_{i,1}, \dots, p_{i,|P_i|}\}$ and $Q_i = \{q_{i,1}, \dots, q_{i,|Q_i|}\}$ label equivalent coefficient lines. Then eq 4.10 simplifies as follows.

$$\hat{P}_{(q_1/\dots/q_m)}^{(p_1/\dots/p_m)} T_{q_1\dots q_m}^{p_1\dots p_m} = |P_1|! \dots |P_h|! |Q_1|! \dots |Q_k|! \hat{P}_{(Q_1/\dots/Q_k)}^{(P_1/\dots/P_h)} T_{q_1\dots q_m}^{p_1\dots p_m} \quad (4.11)$$

Notice that the factorials on the right exactly cancel the contribution of equivalent coefficient lines to the degeneracy factor. This equation defines a simplified rule for interpreting coefficient graphs.

Remark 4.3. Useful results for interpreting graphs. A few results from the preceding discussion are particularly useful for algebraically interpreting a graph according to axiom 4.1. For easy reference, let's summarize them here.

1. Each set of k equivalent internal lines or equivalent subgraphs contributes a factor of $k!$ to the degeneracy. (If rule 4 is used, this excludes equivalent coefficient lines. Otherwise, replace rule 4 with the full antisymmetrizer.)
2. Each open cycle contributes $(-)^{h_i} a_q^p$ to the normal-ordered product, where p and q label the free ends and h_i is the number of hole contractions in the cycle. A closed cycle (loop) contributes $(-)^{h_i+1}$.
3. The overall sign of a closed graph is $(-)^{h+l}$, where h and l denote the total number of hole lines and loops.
4. In a coefficient graph, antisymmetrize over inequivalent coefficient lines using the reduced⁹ antisymmetrizer $\hat{P}_{(Q_1/\dots/Q_k)}^{(P_1/\dots/P_h)}$ where the P_i 's and Q_i 's label subsets of equivalent coefficient lines.

Example 4.1. The one- and two-electron components of H_e expand into occupied/virtual blocks as follows.

$$h_p^q a_q^p = h_a^b a_b^a + h_a^i a_i^a + h_i^a a_a^i + h_i^j a_j^i \quad (4.12)$$

$$\frac{1}{4}\bar{g}_{pq}^{rs}a_{rs}^{pq} = \frac{1}{4}\bar{g}_{ab}^{cd}a_{cd}^{ab} + \frac{1}{2}\bar{g}_{ab}^{ci}a_{ci}^{ab} + \frac{1}{2}\bar{g}_{ai}^{bc}a_{bc}^{ai} + \frac{1}{4}\bar{g}_{ab}^{ij}a_{ij}^{ab} + \bar{g}_{ia}^{bj}a_{bj}^{ia} + \frac{1}{4}\bar{g}_{ij}^{ab}a_{ab}^{ij} + \frac{1}{2}\bar{g}_{ia}^{jk}a_{jk}^{ia} + \frac{1}{2}\bar{g}_{ij}^{ka}a_{ka}^{ij} + \frac{1}{4}\bar{g}_{ij}^{kl}a_{kl}^{ij} \quad (4.13)$$

which, defining $\text{box} \equiv h_p^q a_q^p$ and $\text{wavy} \equiv \frac{1}{4} \bar{g}_{pq}^{rs} a_{rs}^{pq}$, can be expressed in terms of Goldstone diagrams as

$$\begin{array}{c} \text{---} \square \text{---} \bigcirc \text{---} \\ \updownarrow \quad \updownarrow \end{array} = \begin{array}{c} \text{---} \square \text{---} \bullet \text{---} \\ \updownarrow \quad \updownarrow \end{array} + \begin{array}{c} \text{---} \square \text{---} \bullet \text{---} \\ \swarrow \quad \searrow \end{array} + \begin{array}{c} \text{---} \square \text{---} \bullet \text{---} \\ \swarrow \quad \nwarrow \end{array} + \begin{array}{c} \text{---} \square \text{---} \bullet \text{---} \\ \updownarrow \end{array} \quad (4.14)$$

$$\begin{array}{c} \text{---}\circ\text{---} \\ | \quad | \end{array} = \begin{array}{c} \text{---}\bullet\text{---} \\ | \quad | \end{array} + \begin{array}{c} \diagdown \quad \diagup \\ \text{---}\bullet\text{---} \\ | \quad | \end{array} + \begin{array}{c} \diagup \quad \diagdown \\ \text{---}\bullet\text{---} \\ | \quad | \end{array} + \begin{array}{c} \diagdown \quad \diagdown \\ \text{---}\bullet\text{---} \\ | \quad | \end{array} + \begin{array}{c} \diagup \quad \diagup \\ \text{---}\bullet\text{---} \\ | \quad | \end{array} + \begin{array}{c} \diagdown \quad \diagup \\ \text{---}\bullet\text{---} \\ | \quad | \end{array} + \begin{array}{c} \diagup \quad \diagdown \\ \text{---}\bullet\text{---} \\ | \quad | \end{array} + \begin{array}{c} \diagdown \quad \diagdown \\ \text{---}\bullet\text{---} \\ | \quad | \end{array} + \begin{array}{c} \diagup \quad \diagup \\ \text{---}\bullet\text{---} \\ | \quad | \end{array}. \quad (4.15)$$

The degeneracy factors fall into three cases: $\{l_1, l_2\} \cup \{l_3, l_4\} \implies \text{dg}(G)^{-1} = \frac{1}{2.2}$; $\{l_1, l_2\} \cup \{l_3\} \cup \{l_4\} \implies \text{dg}(G)^{-1} = \frac{1}{2.1.1}$; and $\{l_1\} \cup \{l_2\} \cup \{l_3\} \cup \{l_4\} \implies \text{dg}(G)^{-1} = \frac{1}{1.1.1.1}$. In terms of Hugenholtz diagrams, these are written as follows.

[illegible]

Note that in each case the correct scalar factor is built into the definition of the diagram.

⁸The unprimed indices are fixed. This equals the derivative of a generic interaction tensor $\bar{v}_{p'_1 \dots p'_m}^{q'_1 \dots q'_m}$ with respect to $\bar{v}_{p_1 \dots p_m}^{q_1 \dots q_m}$.

⁹The full antisymmetrizer would be $\hat{P}_{(q_1 \dots q_m)}^{(p_1 \dots p_m)}$, where $\{p_1, \dots, p_m\} = P_1 \cup \dots \cup P_h$ and $\{q_1, \dots, q_m\} = Q_1 \cup \dots \cup Q_k$ are the upper and lower indices of the bare excitation operator.

Example 4.2. The Φ -normal Wick expansion of the one- and two-electron components of H_e are as follows.

$$h_p^q a_q^p = h_p^q \tilde{a}_q^p + h_i^i \quad \frac{1}{4} \bar{g}_{pq}^{rs} a_{rs}^{pq} = \frac{1}{4} \bar{g}_{pq}^{rs} \tilde{a}_{rs}^{pq} + \bar{g}_{ip}^{iq} \tilde{a}_q^p + \frac{1}{2} \bar{g}_{ij}^{ij} \quad (4.17)$$

In terms of Goldstone diagrams and Hugenholtz diagrams, these equations are written as follows.

Diagrammatic equations (4.18) showing the expansion of one- and two-electron operators into Goldstone and Hugenholtz diagrams. The first row shows the expansion of a one-electron operator (a box with a vertical line) into a sum of a box with a vertical line and a box with a loop. The second row shows the expansion of a two-electron operator (a box with two vertical lines) into a sum of a box with two vertical lines, a box with two vertical lines and a loop, and a box with two vertical lines and two loops. The third row shows the expansion of a one-electron operator (a box with a vertical line) into a sum of a box with a vertical line and a box with a loop. The fourth row shows the expansion of a two-electron operator (a box with two vertical lines) into a sum of a box with two vertical lines, a box with two vertical lines and a loop, and a box with two vertical lines and two loops.

Again each degeneracy factor is exactly equal to the scalar factor in front of the corresponding algebraic term, in accord with Wick's theorem for graphs. Using these results, the Φ -normal Wick expansion of H_e in Goldstone representation is

Diagrammatic equation (4.19) showing the expansion of H_e into E_0 and H_c . The left side shows H_e as a box with two vertical lines. The right side shows E_0 as a box with two vertical lines and a loop, and H_c as a box with two vertical lines and a loop. The equation is $H_e = E_0 + H_c$.

where \otimes is the Fock operator, $f_p^q \tilde{a}_q^p = h_p^q \tilde{a}_q^p + \bar{g}_{ip}^{iq} \tilde{a}_q^p$.

Notation 4.2. The bra and ket of the reference state Φ are graphically depicted by thick double lines \equiv above and below the graph, respectively. Open graphs with a ket at the bottom are sometimes classified as *wavefunction graphs*.

Remark 4.4. The singles, doubles, triples, etc. configuration interaction (CI) operators will here be denoted as follows.

$$C_1 \equiv c_a^i \tilde{a}_i^a \equiv \text{diagram} \quad C_2 \equiv \left(\frac{1}{2!}\right)^2 c_{ab}^{ij} \tilde{a}_{ij}^{ab} \equiv \text{diagram} \quad C_3 \equiv \left(\frac{1}{3!}\right)^2 c_{abc}^{ijk} \tilde{a}_{ijk}^{abc} \equiv \text{diagram} \quad \dots \quad (4.20)$$

CI theory uses these operators to define a linear parametrization of the wavefunction, $\Psi_{\text{CI}} = (1 + C_1 + \dots + C_n)\Phi$.¹⁰ Coupled-cluster (CC) theory employs a similar series of operators $T_k \equiv \left(\frac{1}{k!}\right)^2 t_{a_1 \dots a_k}^{i_1 \dots i_k} \tilde{a}_{i_1 \dots i_k}^{a_1 \dots a_k}$ to define an exponential parametrization $\Psi_{\text{CC}} = \exp(T_1 + \dots + T_n)\Phi$.

Example 4.3. Derivation of CIS in diagram notation. The CI singles transition energy to an excited state is $\omega \equiv E - E_0 = \langle \Phi | C_1^\dagger H_c C_1 | \Phi \rangle$. Using Wick's theorem for graphs, we can evaluate this expectation value as

Diagrammatic equation for the CIS transition energy ω . The left side shows ω as a diagram with a box and two vertical lines, and a ket at the bottom. The right side shows the expansion of ω into a sum of diagrams, which is then equated to the algebraic expression $c_i^a (f_a^b \delta_j^i - f_j^i \delta_a^b + \bar{g}_{ja}^{bi}) c_b^j$.

where the expression in parentheses on the right defines the CIS matrix, $\langle \Phi_i^a | H_c | \Phi_j^b \rangle$, which we could have also evaluated separately as a coefficient graph. The main consideration in evaluating such expectation values is to identify which terms in the quasiparticle expansion of H_c can form complete contractions with the CI operators in the bra and the ket.

Example 4.4. Derivation of the ground-state CI energy. Projecting the CI Schrödinger equation¹¹ by Φ , the exact ground-state correlation energy can be expressed in CI as $E_c = \langle \Phi | H_c (1 + C_1 + \dots + C_n) | \Phi \rangle$. Using WTG, we find

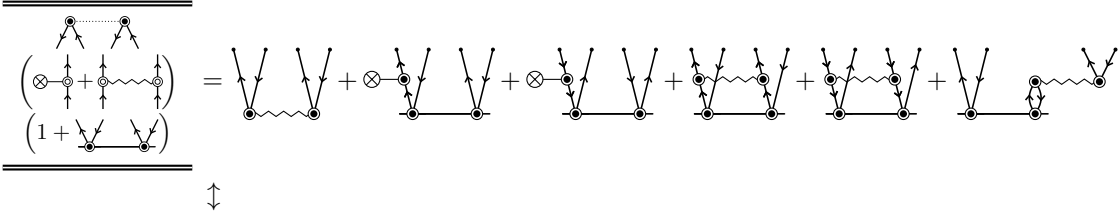
Diagrammatic equation for the ground-state CI energy E_c . The left side shows E_c as a diagram with a box and two vertical lines, and a ket at the bottom. The right side shows the expansion of E_c into a sum of diagrams, which is then equated to the algebraic expression $f_i^a c_a^i + \left(\frac{1}{2}\right)^2 \bar{g}_{ij}^{ab} c_{ab}^{ij}$.

where the first term evaluates to zero whenever Brillouin's theorem applies. This shows that the CI energy depends on triples, quadruples, etc. only indirectly, through their influence on C_1 and C_2 .

¹⁰This expansion assumes "intermediate normalization", in which we normalize the wavefunction to satisfy $\langle \Phi | \Psi \rangle = 1$. In standard normalization, the CI wavefunction is given by $\Psi_{\text{CI}} = (c_0 + C_1 + \dots + C_n)\Phi$, where c_0 is the scalar weight of the reference function.

¹¹By this we mean $H_c \Psi_{\text{CI}} = E_c \Psi_{\text{CI}}$, where $E_c = E - E_0$ is the correlation energy.

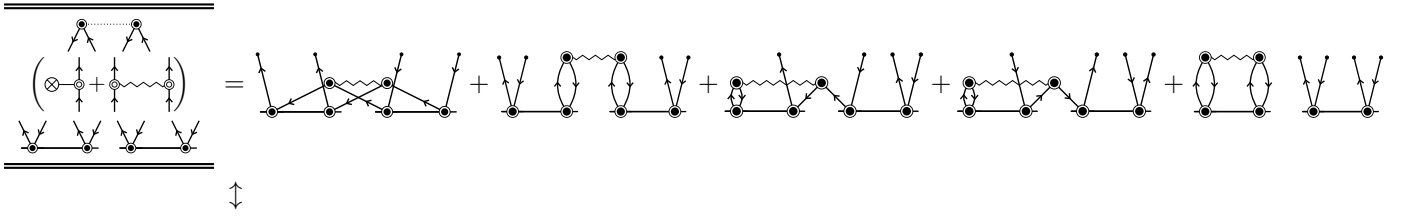
Example 4.5. Derivation of the CID coefficient equation. Projecting the CID Schrödinger equation by Φ_{ij}^{ab} leads to the CID coefficient equation, $\langle \Phi_{ij}^{ab} | H_c (1 + C_2) | \Phi \rangle = E_c c_{ab}^{ij}$. The left-hand side can be evaluated as follows.¹²



$$\langle \Phi_{ij}^{ab} | H_c (1 + C_2) | \Phi \rangle = \bar{g}_{ab}^{ij} + \hat{P}_{(a/b)} f_a^c t_{cb}^{ij} - \hat{P}^{(i/j)} f_k^i c_{ab}^{kj} + \frac{1}{2} \bar{g}_{ab}^{cd} c_{cd}^{ij} + \frac{1}{2} \bar{g}_{kl}^{ij} c_{ab}^{kl} + \hat{P}_{(a/b)}^{(i/j)} \bar{g}_{kb}^{cj} c_{ac}^{ik} \quad (4.21)$$

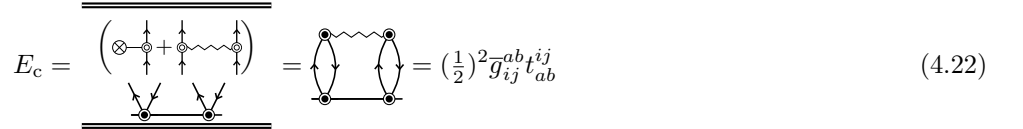
Note that the equivalent coefficient lines in these diagrams don't contribute to the degeneracy factors here, since these have been canceled using item 4 in rmk 4.3. The minus sign in the third term comes from the fact that its diagram has three holes and two loops, $(-1)^{3+2} = -1$.

Example 4.6. Derivation of the CCD amplitude equation. Projecting the CCD Schrödinger equation by Φ_{ij}^{ab} leads to the CCD amplitude equation, $\langle \Phi_{ij}^{ab} | H_c \exp(T_2) | \Phi \rangle = E_c t_{ab}^{ij}$. Note that terms cubic and higher in T_2 cannot form complete contractions with $\tilde{a}_{ab}^{ij} H_c$, reducing the left-hand side to $\langle \Phi_{ij}^{ab} | H_c (1 + T_2 + \frac{1}{2} T_2^2) | \Phi \rangle$. If we exchange CI coefficients for CC amplitudes the linear terms are given by equation 4.21, so only the quadratic term remains to be evaluated.



$$\frac{1}{2} \langle \Phi_{ij}^{ab} | H_c T_2^2 | \Phi \rangle = \left(\frac{1}{2} \right)^2 \bar{g}_{kl}^{cd} t_{ab}^{kl} t_{cd}^{ij} + \frac{1}{2} \hat{P}_{(a/b)}^{(i/j)} \bar{g}_{kl}^{cd} t_{ac}^{ik} t_{db}^{lj} - \frac{1}{2} \hat{P}_{(a/b)} \bar{g}_{kl}^{cd} t_{ca}^{kl} t_{db}^{ij} - \frac{1}{2} \hat{P}^{(i/j)} \bar{g}_{kl}^{cd} t_{cd}^{ki} t_{ab}^{lj} + E_c t_{ab}^{ij}$$

In the last term, we have substituted in the CCD energy expression, which is the same graph as the CID energy expression.



$$E_c = \left(\frac{1}{2} \right)^2 \bar{g}_{ab}^{ij} t_{ab}^{ij} \quad (4.22)$$

Note that the term $E_c t_{ab}^{ij}$ actually cancels with the right-hand side of the projected Schrödinger equation. Since this is the only unlinked graph in the expansion, we can write the amplitude equation as

$$\langle \Phi_{ij}^{ab} | H_c \exp(T_2) | \Phi \rangle_L = 0 \quad (4.23)$$

where the subscript indicates that only linked graphs are included. This cancellation of unlinked terms is a general feature of coupled-cluster theory. Finally, note that when Brillouin's theorem holds the second and third terms in eq 4.21 become

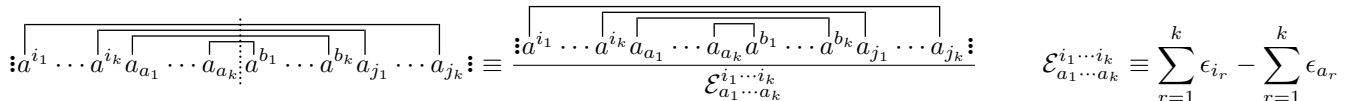
$$\hat{P}_{(a/b)} f_a^c t_{cb}^{ij} - \hat{P}^{(i/j)} f_k^i c_{ab}^{kj} = -(\epsilon_i + \epsilon_j - \epsilon_a - \epsilon_b) t_{ab}^{ij}$$

where $\{\epsilon_p = f_p^p\}$ are canonical Hartree-Fock orbital energies. Subtracting this contribution from both sides leads to

$$t_{ab}^{ij} = (\mathcal{E}_{ab}^{ij})^{-1} \langle \Phi_{ij}^{ab} | (H_c - f_p^p \tilde{a}_p^p) \exp(T_2) | \Phi \rangle_L \quad \mathcal{E}_{ab}^{ij} \equiv \epsilon_i + \epsilon_j - \epsilon_a - \epsilon_b \quad (4.24)$$

which is the standard form for the T_2 amplitude equation.

Definition 4.9. Resolvent line. Complete contractions through a *resolvent line* \vdots are defined as follows.



$$\vdots a_{i_1} \dots a_{i_k} a_{a_1} \dots a_{a_k} \vdots a_{b_1} \dots a_{b_k} a_{j_1} \dots a_{j_k} \vdots \equiv \frac{a_{i_1} \dots a_{i_k} a_{a_1} \dots a_{a_k} a_{b_1} \dots a_{b_k} a_{j_1} \dots a_{j_k}}{\mathcal{E}_{a_1 \dots a_k}^{i_1 \dots i_k}} \quad \mathcal{E}_{a_1 \dots a_k}^{i_1 \dots i_k} \equiv \sum_{r=1}^k \epsilon_{i_r} - \sum_{r=1}^k \epsilon_{a_r}$$

This can be denoted in a graph by drawing a horizontal dotted line through several contraction lines.

¹²Note that we are using Einstein summation here, so c, d, k , and l are summation indices.

Example 4.7. Using the notion of a resolvent line, equation 4.24 can be written in diagram notation as follows.

$$(4.25)$$

This equation is typically solved in an iterative procedure that updates the amplitudes by substituting the previous iteration's amplitudes into the right-hand side, starting with $^{[0]}t_{ab}^{ij} = 0$ as an initial guess.

Remark 4.5. The CEPA₀ and MP2 amplitude equations. If the T_2 amplitudes are small, the CCD exponential, $\exp(T_2) = 1 + T_2 + \frac{1}{2}T_2^2 + \dots$, can be approximated by the linear terms or even just the constant term, leading to the *coupled electron-pair approximation zero* (CEPA₀) and *second-order Møller-Plesset perturbation theory* (MP2) methods. Assuming Brillouin's theorem and writing V for the electron repulsion operator, $\frac{1}{4}\bar{g}_{pq}^{rs}\bar{a}_{rs}^{pq}$, these amplitude equations are

MP2 : $t_{ab}^{ij} = (\mathcal{E}_{ab}^{ij})^{-1} \langle \Phi_{ij}^{ab} | V | \Phi \rangle_L =$ (4.26)

CEPA₀ : $t_{ab}^{ij} = (\mathcal{E}_{ab}^{ij})^{-1} \langle \Phi_{ij}^{ab} | V(1 + T_2) | \Phi \rangle_L =$ (4.27)

where the MP2 amplitude is usually denoted $^{(1)}t_{ab}^{ij}$ because it is correct to first order in perturbation theory. Note that eq 4.26 specifies the MP2 amplitudes directly, $^{(1)}t_{ab}^{ij} = (\mathcal{E}_{ab}^{ij})^{-1} \bar{g}_{ab}^{ij}$, and requires no iterative solution. Substituting these amplitudes into equation 4.22 gives the MP2 energy, which has contributions up to second order in perturbation theory.

$E_c^{(2)} =$ $(1) =$ $= \frac{1}{4} \sum_{ijab} \frac{\bar{g}_{ab}^{ij} \bar{g}_{ij}^{ab}}{\mathcal{E}_{ab}^{ij}}$ (4.28)