1. Interpret the following coefficient graph algebraically, denoting the bare excitation operator by \tilde{a}_{abc}^{ijk} .



2. Interpret the following coefficient graph algebraically, denoting the top bare excitation operator by \tilde{a}_{ab}^{ij} and the bottom one by \tilde{a}_{kl}^{cd} and their interaction tensors by ${}^{1}\overline{\delta}_{mn}^{ef}$ and ${}^{2}\overline{\delta}_{ef}^{mn}$.



Solutions:

1. Axiom 4.1 gives

since there are two pairs of equivalent lines on the repulsion operator, and the CI triples operator has a set of three equivalent lines. The interaction tensor for the bare excitation operator is $\bar{\delta}_{lmn}^{def} = \hat{P}_{(a/b/c)}^{(i/j/k)} \delta_a^d \delta_b^e \delta_c^f \delta_l^i \delta_m^j \delta_n^k$, so this can be simplified as

Using item 3 in Remark 4.3, the contracted operator string evaluates as follows

since there are three hole lines and three loops in the graph. At this point, we have simplified our interpretation to the following

where I have relabeled the summation indices $g \mapsto d$, $h \mapsto e$. Finally, using item 4 under Remark 4.3, we can cancel the degeneracy factors 2!3! coming from equivalent coefficient lines by replacing $\hat{P}_{(a/b/c)}^{(i/j/k)}$ with $\hat{P}_{(ab/c)}$. This works because the operand $\bar{g}_{ab}^{de}c_{dec}^{ijk}$ is already antisymmetric with respect to $\{a,b\}$ and $\{i,j,k\}$.

2. Axiom 4.1 gives

$$\otimes - \stackrel{\downarrow}{\bullet} \quad \Big| \quad \Big| \quad \Big| = \frac{1}{2!} \sum_{\substack{ef \\ mno}} \otimes \stackrel{m}{\downarrow} e \Big| n \Big| f \Big| = \frac{1}{2!} \sum_{\substack{ef \\ mno}} {}^1 \overline{\delta}_{mn}^{ef} f_o^m {}^2 \overline{\delta}_{ef}^{on} \\ \vdots \\ a_{e^{\bullet 1}f^{\bullet 2}}^{m^{\circ 1}n^{\circ 2}} a_{m^{\circ 1}}^{o^{\circ 3}} a_{o^{\circ 3}n^{\circ 2}}^{e^{\bullet 1}f^{\bullet 2}} \\ \vdots \\ \Big| \qquad \qquad \\ \vdots \\ \Big| \qquad \qquad \qquad \\ \Big| \qquad \qquad$$

since the two particle lines are equivalent. Using item 3 in Remark 4.3, the operator string evaluates to

$$\mathbf{i} a_{e^{\bullet 1} f^{\bullet 2}}^{m^{\circ 1} n^{\circ 2}} a_{m^{\circ 1}}^{o^{\circ 3}} a_{o^{\circ 3} n^{\circ 2}}^{e^{\bullet 1} f^{\bullet 2}} \mathbf{i} = (-1)^{3+2} = -1$$

since there are three hole lines and two loops in the diagram. Substituting in the definitions of the interaction tensors, we can simplify the result as follows.

$$\bigotimes \frac{1}{l} \quad \frac{1}{l} = -\frac{1}{2!} \sum_{\substack{ef \\ mno}} {}^{1} \overline{\delta}_{mn}^{ef} f_o^{m} {}^{2} \overline{\delta}_{ef}^{on} = -\frac{1}{2!} \sum_{\substack{ef \\ mno}} \left(\hat{P}_{(a/b)}^{(i/j)} \delta_a^e \delta_b^f \delta_m^i \delta_n^j \right) f_o^m \left(\hat{P}_{(k/l)}^{(c/d)} \delta_k^o \delta_l^n \delta_e^c \delta_f^d \right)$$

$$= -\frac{1}{2!} \hat{P}_{(a/b)}^{(i/j)} \hat{P}_{(k/l)}^{(c/d)} f_k^i \delta_l^j \delta_a^c \delta_b^d$$

Finally, using item 4 under Remark 4.3, we can cancel the degeneracy factor 2! by replacing $\hat{P}_{(a/b)}^{(i/j)}$ with $\hat{P}^{(i/j)}$.