

Numerical Evaluation of the Arrhenius Integral

C.D. Clark III

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The Arrhenius Integral model for thermal damage requires the evaluation of the integral[2]

$$\Omega(\tau) = \int_0^{\tau} A e^{\frac{-E_a}{RT(t)}} dt. \quad (1)$$

If the temperature, $T(t)$, is predicted using a numerical simulation, i.e. a finite-difference or finite-element model, then the temperature is calculated at a discrete set of times, (t_0, t_1, \dots) . This limits the methods that can be used to evaluate Eq. 1 numerically to those that work on predefined nodes (Gaussian quadrature, for example, cannot be used).

The usual methods for numerically evaluating an integral with pre-defined nodes would be Reimann sum, Trapezoid Rule, and Simpson's Rule. These methods are simple to implement, but approximate the integrand as piece-wise, linear, and parabolic, respectively.

1 Recasting the integral

If we assume that the temperature between two times, t_0 and t_1 can be written as a linear function of time,

$$T(t) = mt + T_0 \quad (2)$$

and inserting into this into Eq. 1 gives,

$$\Omega = \int_{t_0}^{t_1} A e^{\frac{-E_a}{R(mt+T_0)}} dt. \quad (3)$$

Changing the integration variable to T ,

$$dT = m dt \quad (4)$$

$$\Omega = \int_{T_0}^{T_1} \frac{A}{m} e^{\frac{-E_a}{RT}} dT. \quad (5)$$

Where T_0 and T_1 are the temperature at t_0 and t_1 . Next, we substitute $u = \frac{E_a}{RT}$ and change the integration variable again to u .

$$u = \frac{E_a}{R} \frac{1}{T} \quad (6)$$

$$du = -\frac{E_a}{R} \frac{1}{T^2} dT \quad (7)$$

$$dT = -\frac{E_a}{R} \frac{1}{u^2} du \quad (8)$$

$$\Omega = -\frac{A}{m} \frac{E_a}{R} \int_{E_a/RT_0}^{E_a/RT_1} \frac{e^{-u}}{u^2} du \quad (9)$$

$$= \frac{A}{m} \frac{E_a}{R} \int_{E_a/RT_1}^{E_a/RT_0} \frac{e^{-u}}{u^2} du. \quad (10)$$

The integrand can be evaluated “analytically” using the *exponential integral*, E_n , which is defined as[1]

$$E_n(x) = \int_1^\infty \frac{e^{-xt}}{t^n} dt \quad (11)$$

The variable substitution $y = xt$ gives

$$dy = xdt \quad (12)$$

$$dt = \frac{1}{x} dy \quad (13)$$

$$E_n(x) = x \int_x^\infty \frac{e^{-y}}{y^n} dy \quad (14)$$

Comparing this to Eq. 10, it is simple to show that

$$\Omega = \frac{A}{m} \frac{E_a}{R} \left[\frac{E_2\{E_a/RT_1\}}{E_a/RT_1} - \frac{E_2\{E_a/RT_0\}}{E_a/RT_0} \right] \quad (15)$$

$$= \frac{A}{m} [T_1 E_2\{E_a/RT_1\} - T_0 E_2\{E_a/RT_0\}] \quad (16)$$

Finally, we note that the slope, m , can be determined from T_0 and T_1 if t_1 and t_0 are known (rise over run),

$$\Omega = A \frac{t_1 - t_0}{T_1 - T_0} [T_1 E_2\{E_a/RT_1\} - T_0 E_2\{E_a/RT_0\}] \quad (17)$$

This gives the accumulated damage over a time $t_1 - t_0$, when the temperature rise is linear. If the actual temperature rise is not linear, then we can discretize the profile into small segments, such that the temperature is approximatly linear over the segment.

It is reasonable to think that this approximation (treating the temperature as linear over some time interval) is better than assuming that the damage rate is linear over the same time. Figure 1 shows the Arrhenius rate as a function of time for an example thermal profile. The figure shows that the Arrhenius rate is concave on both sides its peak, which means that the error accumulated by the trapezoid rule will have the same sign on each side. The trapezoid rule will *over*-predict the accumulated damage. On the otherhand, if the thermal profile is approximated as linear segments, the accumulated damage will be *under*-predicted while the temperature increases, and then be *over*-predicted when the temperature begins to fall.

2 Numerical Implementation

Equation 17 provides an analytic expression to evaluate the Arrhenius integral over some time period where the temperature rise is linear. To numerically integrate the Arrhenius integral for a full thermal profile, we simply break the integral up into small time intervals, use Equation 17 to evaluate the integral over each time interval, and then them all up. However, we must be careful because there is a $T_1 - T_0$ in the denominator. If the temperature is constant, we will divide by zero.

If $T_0 = T_1 = T$, then Equation 17 gives

$$\Omega = A(t_1 - t_0) \left(\frac{TE_2\{E_a/RT\} - TE_2\{E_a/RT\}}{T - T} \right) = A(t_1 - t_0) \frac{0}{0}. \quad (18)$$

To evaluate, we need to apply L'Hôpital's rule for $T_1 - T_0 \rightarrow 0$. Apparently, this limit will give $A(t_1 - t_0)e^{-E_a/RT}$, since this is the integral for a constant temperature.

A numerical implementation will need to check for the case that $T_0 = T_1$ and use the constant temperature quadrature. If T_0 and T_1 are stored as floating point numbers, then we will probably need to check that they are “close”, rather than “equal”. Note that, for a linear temperature rise from T_0 and T_1 , evaluating Equation 17 will give a numerical value between $Ae^{-E_a/RT_0}(t_1 - t_0)$ and $Ae^{-E_a/RT_1}(t_1 - t_0)$. We can therefore write a limit on the maximum error in incurred by any numerical approximation to integral over duration $t_1 - t_0$.

$$\epsilon \leq \frac{|Ae^{-E_a/RT_1}(t_1 - t_0) - Ae^{-E_a/RT_0}(t_1 - t_0)|}{Ae^{-E_a/RT_0}(t_1 - t_0)} = \left| e^{-\frac{E_a}{R}\left(\frac{1}{T_1} - \frac{1}{T_0}\right)} - 1 \right| \quad (19)$$

This gives a metric for deciding if T_1 and T_0 are “close enough”. Given a tolerance, or maximum allowed error, we

$$\frac{R}{E_a} \ln(\epsilon + 1) \leq \left| \frac{1}{T_1} - \frac{1}{T_0} \right| \quad (20)$$

For small ϵ , $\ln(\epsilon + 1) \approx \epsilon$,

$$\frac{1}{T_1} - \frac{1}{T_0} \geq \frac{R}{E_a} \epsilon \quad (21)$$

References

- [1] Milton Abramowitz and Irene A. Stegun, editors. *Handbook of Mathematical Functions with Formulas, Graphs, and Mathematical Tables*. U.S. Department of Commerce, 1970.
- [2] A.J. Welch and M.J.C van Germert, editors. *Optical-Thermal Response of Laser-Irradiated Tissue*. Plenum Press, 2nd edition, 2011.

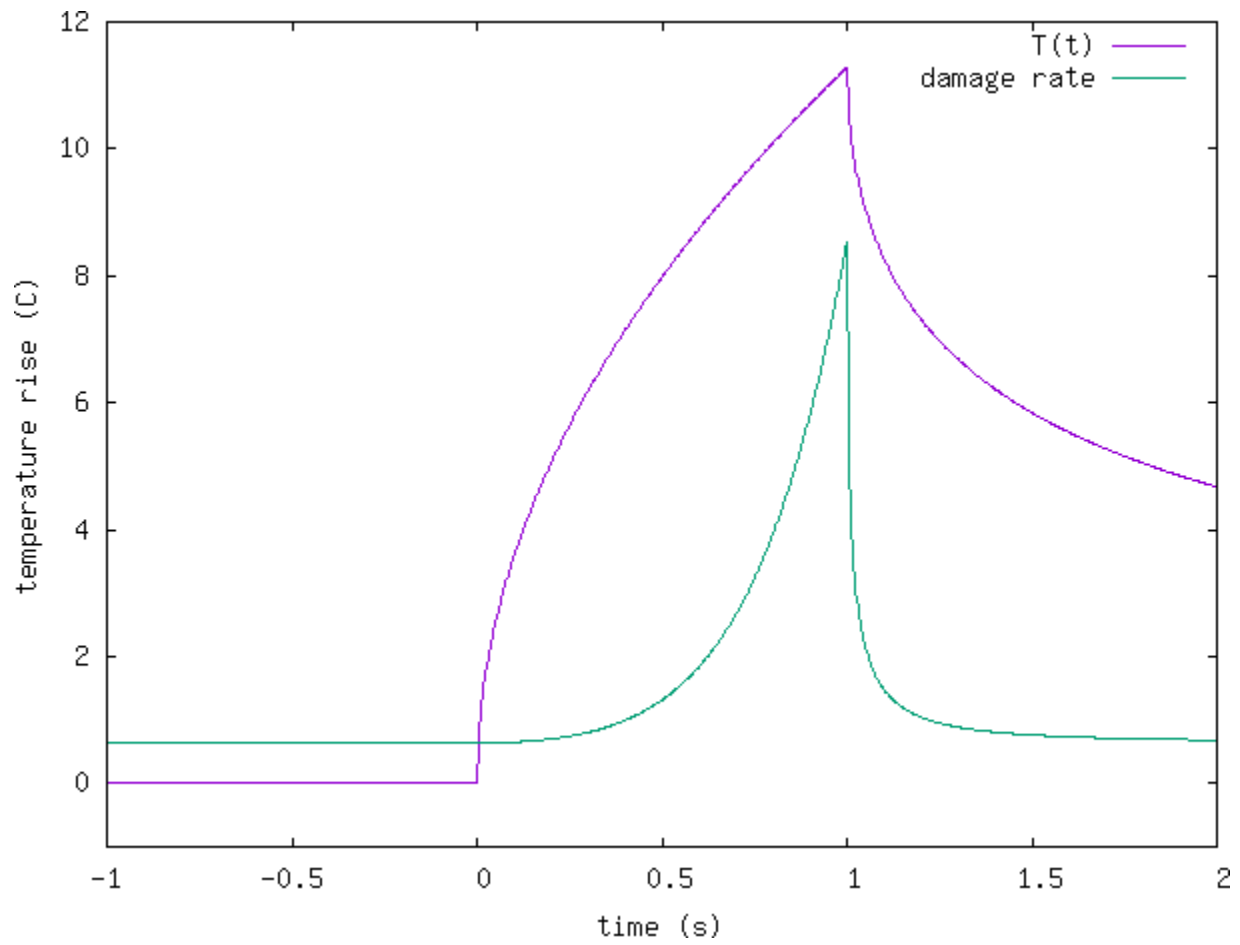


Figure 1: The Arrhenius rate, $Ae^{-E_a/RT(t)}$, plotted for an example thermal profile.