

COMPUTATION OF FIELDS IN MULTI-LAYERED STRUCTURES

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Multi-layered structures are the basis of modern semiconductor based integrated circuits, and they have been a subject of considerable interest in theoretical studies to gain an understanding of electromagnetic fields [?]. Most commonly, the fields are found in terms of the dyadic Green functions (DGFs) of the environment, which defines the vector field distribution due to a vector point source [1]. Due to a current \mathbf{J} distributed in a region defined by a surface, r' , the surface fields are expressed as in terms of an inner product,

$$\mathbf{E} = \int_{r'} \underline{\underline{\mathbf{G}}}^{\text{EJ}}(\mathbf{r}|\mathbf{r}') \mathbf{J}(\mathbf{r}') d\mathbf{r}', \quad (1a)$$

$$\mathbf{H} = \int_{r'} \underline{\underline{\mathbf{G}}}^{\text{HJ}}(\mathbf{r}|\mathbf{r}') \mathbf{J}(\mathbf{r}') d\mathbf{r}'. \quad (1b)$$

where, $\underline{\underline{\mathbf{G}}}^{\text{EJ}}$ and $\underline{\underline{\mathbf{G}}}^{\text{HJ}}$ are the spatial domain DGF's due to an electric current source located at r' , that define the electric and magnetic fields at a point r respectively. In the presence of magnetic sources, (1a) and (1b) can be augmented using the superposition principle, by adding an inner product containing magnetic DGFs, $\underline{\underline{\mathbf{G}}}^{\text{EM}}$ and $\underline{\underline{\mathbf{G}}}^{\text{HM}}$.

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1 Theory

We follow the well-established approach of field computation for planar multi-layered media [2, 3], in which, an equivalent transmission line (TL) network is first set up for the structure, and then TL based GFs are computed.

1.1 Transmission line representation of Maxwell's equations

As shown in Fig. 3a, it is assumed that the structure is unbounded in the lateral direction, and excited only by an electric source. The electric and magnetic fields are given by the Maxwell's equations in frequency-domain,

$$\nabla \times \mathbf{E} = -j\omega\mu\mathbf{H}, \quad (2a)$$

$$\nabla \times \mathbf{H} = j\omega\varepsilon\mathbf{E} + \mathbf{J}. \quad (2b)$$

For boundary-value problems displaying symmetry along the z direction, it is desirable to decompose the ∇ operator into two components, one d/dz and the other a transverse (to z) operator, $\nabla_{\mathbf{t}}$ [4, p. 64]. By taking the Fourier transform,

$$\mathcal{F}[f(\rho, z)] \equiv \tilde{F}(\mathbf{k}_\rho, z) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(\rho, z) e^{-j\mathbf{k}_\rho \cdot \boldsymbol{\rho}} dx dy \quad (3)$$

the field computation is considerably simplified by switching to the spectral frequency domain \mathbf{k}_ρ , which reduces the complexity of the vector differential operator, ∇ to $-jk_x\hat{\mathbf{x}} - jk_y\hat{\mathbf{y}} + d/dz\hat{\mathbf{z}}$, which contains a derivative term only in z -direction. In (3), the cylindrical coordinates are expressed as,

$$\rho = x\hat{\mathbf{x}} + y\hat{\mathbf{y}}, \quad \text{and} \quad \mathbf{k}_\rho = k_x\hat{\mathbf{x}} + k_y\hat{\mathbf{y}}, \quad (4)$$

and the notation \sim above the capital-letter terms indicates the Fourier transform with respect to the transverse coordinates and from here on, will be used to denote the spectral domain quantities.

As stated earlier, it is advantageous to separate the fields in transverse and longitudinal coordinates because, as we shall see shortly, the longitudinal part of the field can be completely expressed in terms of the transverse component. Applying the Fourier transform (3) on the Maxwell's equations (2), we obtain:

$$\left(-j\mathbf{k}_\rho + \hat{\mathbf{z}}\frac{d}{dz}\right) \times (\tilde{\mathbf{E}}_t + \tilde{\mathbf{E}}_z) = -j\omega\mu(\tilde{\mathbf{H}}_t + \tilde{\mathbf{H}}_z), \quad (5a)$$

$$\left(-j\mathbf{k}_\rho + \hat{\mathbf{z}}\frac{d}{dz}\right) \times (\tilde{\mathbf{H}}_t + \tilde{\mathbf{H}}_z) = j\omega\varepsilon(\tilde{\mathbf{E}}_t + \tilde{\mathbf{E}}_z) - (\tilde{\mathbf{J}}_t + \tilde{\mathbf{J}}_z). \quad (5b)$$

The transverse and longitudinal components of the magnetic field can be separately expressed in (5a) as,

$$-j\mathbf{k}_\rho \times \tilde{\mathbf{E}}_z + \frac{d}{dz}\hat{\mathbf{z}} \times \tilde{\mathbf{E}}_t = -j\omega\mu\tilde{\mathbf{H}}_t, \quad (6a)$$

$$-j\mathbf{k}_\rho \times \tilde{\mathbf{E}}_t = -j\omega\mu\tilde{\mathbf{H}}_z. \quad (6b)$$

Using the vector product property [5, p. 117],

$$\mathbf{A} \times \mathbf{B} = \mathbf{A} \cdot (\mathbf{B} \times \hat{\mathbf{n}}) \hat{\mathbf{n}}, \quad (7)$$

where the unit vector $\hat{\mathbf{n}}$ is normal to the plane containing vectors \mathbf{A} and \mathbf{B} . A scalar form of the longitudinal component of the electric field is obtained by applying (7) on (6b),

$$-j\mathbf{k}_\rho \cdot (\tilde{\mathbf{E}}_t \times \hat{\mathbf{z}}) \hat{\mathbf{z}} = -j\omega\mu\tilde{\mathbf{H}}_z \quad (8)$$

which can be written in the scalar form,

$$-j\tilde{H}_z = \frac{-j}{\omega\mu}\mathbf{k}_\rho \cdot (\tilde{\mathbf{E}}_t \times \hat{\mathbf{z}}). \quad (9)$$

Now taking the vector product with unit vector $\hat{\mathbf{z}}$ on both sides of (6a), the transverse electric field component is expressed as:

$$\begin{aligned} \frac{d\tilde{\mathbf{E}}_t}{dz} &= -j(\mathbf{k}_\rho \times \tilde{\mathbf{E}}_z) \times \hat{\mathbf{z}} - j\omega\mu\tilde{\mathbf{H}}_t \times \hat{\mathbf{z}} \\ &= -j\mathbf{k}_\rho \tilde{E}_z - j\omega\mu\tilde{\mathbf{H}}_t \times \hat{\mathbf{z}} \end{aligned} \quad (10)$$

where the BAC-CAB vector triple product identity, $(\mathbf{A} \times \mathbf{B}) \times \mathbf{C} \equiv \mathbf{B}(\mathbf{A} \cdot \mathbf{C}) - \mathbf{C}(\mathbf{A} \cdot \mathbf{B})$ has been applied.

Following a similar procedure starting from (5b), we obtain the transverse component of magnetic field, along with the scalar longitudinal component of electric field:

$$\begin{aligned} \frac{d\tilde{\mathbf{H}}_t}{dz} &= -j(\mathbf{k}_\rho \times \tilde{\mathbf{H}}_z) \times \hat{\mathbf{z}} + j\omega\varepsilon\tilde{\mathbf{E}}_t \times \hat{\mathbf{z}} + \tilde{\mathbf{J}}_t \times \hat{\mathbf{z}} \\ &= -j\mathbf{k}_\rho\tilde{H}_z + j\omega\varepsilon\tilde{\mathbf{E}}_t \times \hat{\mathbf{z}} + \tilde{\mathbf{J}}_t \times \hat{\mathbf{z}}, \end{aligned} \quad (11)$$

and,

$$-j\omega\varepsilon\tilde{E}_z = j\mathbf{k}_\rho \cdot (\tilde{\mathbf{H}}_t \times \hat{\mathbf{z}}) + \tilde{J}_z. \quad (12)$$

By substituting (12) in (10), we get the transverse component of electric field,

$$\frac{d\tilde{\mathbf{E}}_t}{dz} = \frac{1}{j\omega\varepsilon} (k^2 - \mathbf{k}_\rho\mathbf{k}_\rho \cdot) (\tilde{\mathbf{H}}_t \times \hat{\mathbf{z}}) + \mathbf{k}_\rho \frac{\tilde{J}_z}{\omega\varepsilon}. \quad (13)$$

Similarly, from (9) and (11), the transverse component of magnetic field,

$$\frac{d\tilde{\mathbf{H}}_t}{dz} = \frac{1}{j\omega\mu} (k^2 - \mathbf{k}_\rho\mathbf{k}_\rho \cdot) (\hat{\mathbf{z}} \times \tilde{\mathbf{E}}_t) + \tilde{\mathbf{J}}_t \times \hat{\mathbf{z}} \quad (14)$$

where $k = \omega\sqrt{\mu\varepsilon}$ in (13) and (14) is the medium wavenumber.

The fields in (13) and (14) for arbitrarily aligned sources lie in the plane of a spectral coordinate system as illustrated in Fig. 1, where the arrowheads in color correspond to spectral-domain quantities. A rotational transformation of the coordinate system such that the axes align with the vectors $\mathbf{k}_\rho, \hat{\mathbf{z}} \times \mathbf{k}_\rho$ [6], simplifies the procedure of finding the transmission line equivalent, which allows the TE and TM mode analysis to be made separately. The coordinate transformation can be expressed as:

$$\begin{bmatrix} \hat{\mathbf{u}} \\ \hat{\mathbf{v}} \end{bmatrix} = \begin{bmatrix} \cos \psi & \sin \psi \\ -\sin \psi & \cos \psi \end{bmatrix} \begin{bmatrix} \hat{\mathbf{x}} \\ \hat{\mathbf{y}} \end{bmatrix} \quad (15)$$

where ψ is the angle between \mathbf{k}_ρ and the positive x-axis. A transmission line analogue

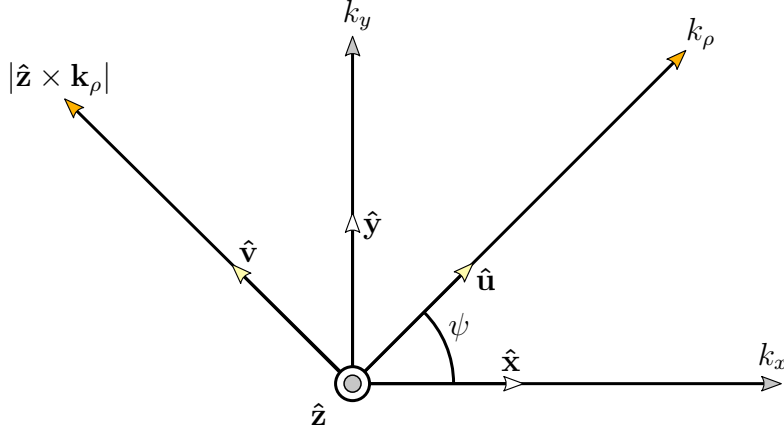


Figure 1: Coordinate System transformation in the spectral domain

for the spectral fields, expressed in terms of modal voltages and currents can therefore, be written as [7, 3],

$$\begin{bmatrix} \tilde{\mathbf{E}}_t \\ \tilde{\mathbf{H}}_t \end{bmatrix} = \begin{bmatrix} V^{\text{TM}} & V^{\text{TE}} \\ -I^{\text{TE}} & I^{\text{TM}} \end{bmatrix} \begin{bmatrix} \hat{\mathbf{u}} \\ \hat{\mathbf{v}} \end{bmatrix}. \quad (16)$$

Using the results of (16) in (13) and noting that $\hat{\mathbf{u}} = \mathbf{k}_\rho / k_\rho$, we get,

$$\frac{d(\hat{\mathbf{u}} V^{\text{TM}} + \hat{\mathbf{v}} V^{\text{TE}})}{dz} = \frac{1}{j\omega\varepsilon} (k^2 - \mathbf{k}_\rho \cdot \mathbf{k}_\rho) (\hat{\mathbf{u}} I^{\text{TM}} + \hat{\mathbf{v}} I^{\text{TE}}) + \hat{\mathbf{u}} \frac{k_\rho}{\omega\varepsilon} \tilde{J}_z \quad (17)$$

By separating the $\hat{\mathbf{u}}$ and $\hat{\mathbf{v}}$ components, we obtain the TM and TE equivalent voltage equations respectively,

$$\frac{dV^{\text{TM}}}{dz} = \frac{1}{j\omega\varepsilon} (k^2 - k_\rho^2) I^{\text{TM}} + \frac{k_\rho}{\omega\varepsilon} \tilde{J}_z, \quad (18a)$$

$$\frac{dV^{\text{TE}}}{dz} = \frac{k^2}{j\omega\varepsilon} I^{\text{TE}}. \quad (18b)$$

Similarly, from (16) and (14), the equivalent current equations can be written as:

$$\frac{dI^{\text{TM}}}{dz} = \frac{k^2}{j\omega\mu} V^{\text{TM}} - \tilde{J}_u, \quad (19a)$$

$$\frac{dI^{\text{TE}}}{dz} = \frac{-1}{j\omega\mu} (k^2 - k_\rho^2) V^{\text{TE}} + \tilde{J}_v. \quad (19b)$$

Equations (18)-(19) can be conveniently written in a compact form as a set of telegrapher's equations [4, p. 190]:

$$\frac{dV^\alpha}{dz} = -jk_z Z^\alpha I^\alpha + v^\alpha \quad (20a)$$

$$\frac{dI^\alpha}{dz} = -jk_z Y^\alpha V^\alpha + i^\alpha \quad (20b)$$

where α is either TE or TM, the propagation constant in the transverse direction is $k_z = \pm\sqrt{k^2 - k_\rho^2}$, for which the sign must be chosen so that the fields decay away from the source. The modal impedances in (20) are,

$$Z^{\text{TM}} = \frac{1}{Y^{\text{TM}}} = \frac{k_z}{\omega\varepsilon}, \quad (21a)$$

$$Z^{\text{TE}} = \frac{1}{Y^{\text{TE}}} = \frac{\omega\mu}{k_z}. \quad (21b)$$

Using the expressions that relate the transverse electric and magnetic to the equivalent transmission line currents and voltages (16), and combining with the longitudinal field expressions (6b) and (12), we obtain the total fields in the spectral domain,

$$\begin{bmatrix} \tilde{\mathbf{E}}(\mathbf{k}_\rho, z) \\ \tilde{\mathbf{H}}(\mathbf{k}_\rho, z) \end{bmatrix} = \begin{bmatrix} V^{\text{TM}} & V^{\text{TE}} & -\frac{k_\rho}{\omega\varepsilon(z)} I^{\text{TM}}(z) \\ -I^{\text{TE}} & I^{\text{TM}} & \frac{k_\rho}{\omega\mu} V^{\text{TE}}(z) \end{bmatrix} \begin{bmatrix} \hat{\mathbf{u}} \\ \hat{\mathbf{v}} \\ \hat{\mathbf{z}} \end{bmatrix} + \hat{\mathbf{z}} \begin{bmatrix} \frac{j}{\omega\varepsilon(z)} \tilde{\mathbf{J}}_z(\mathbf{k}_\rho, z) \\ 0 \end{bmatrix}, \quad (22)$$

where $\varepsilon(z)$ may vary from one layer to another.

Assuming only electric sources existing in space, the corresponding TL sources, v^α and i^α , defined in (20), are illustrated in 2. A horizontally oriented (x-directed) electric dipole is represented by a current source in an equivalent TM transmission line network. Likewise, the equivalent configuration of a vertical (y-directed) electric dipole is a TE network with a current source. A z-directed dipole corresponds to voltage source in a TM transmission line. For an arbitrarily directed source, the

equivalent TL model consists of a superposition of the three representations.

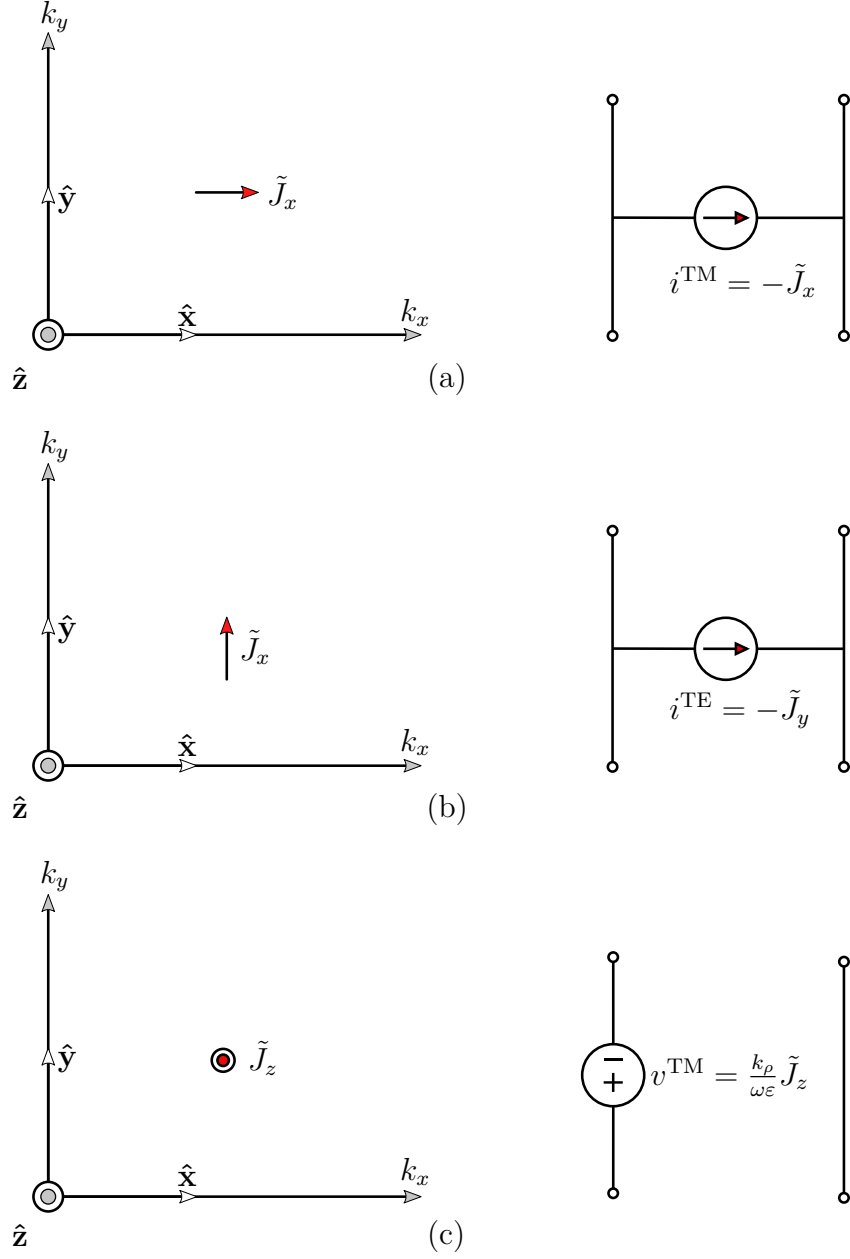


Figure 2: Electric Source representation in a transmission line network

1.2 Green functions for the TL equations

To obtain the transmission line voltages and currents that define the spectral fields in (22), we introduce the one-dimensional transmission line Green functions (TLGFs) that are analogous to the spatial domain DGFs in (1). Following [4, 2], we define $V_i(z, z')$ and $I_i(z, z')$ as the voltage and current respectively, at a point z along the transmission line due to a unit-strength current source located at z' . Similarly, $V_v(z, z')$ and $I_v(z, z')$ are the respective voltage and current due to a unit-strength voltage source. Analogous to (1) which is augmented by magnetic sources, we write:

$$V(z) = \int_{z'} [V_i(z, z')i(z') + V_v(z, z')v(z')] dz', \quad (23a)$$

$$I(z) = \int_{z'} [I_i(z, z')i(z') + I_v(z, z')v(z')] dz'. \quad (23b)$$

The telegrapher's equations (20), are rewritten for a voltage excited line as:

$$\frac{dV_v}{dz} = -jk_z Z_v I_v + \delta(z - z') \quad (24a)$$

$$\frac{dI_v}{dz} = -jk_z Y_v V_v, \quad (24b)$$

and for a current excited line, we obtain:

$$\frac{dV_i}{dz} = -jk_z Z_i I_i \quad (25a)$$

$$\frac{dI_i}{dz} = -jk_z Y_i V_i + \delta(z - z') \quad (25b)$$

1.3 Spectral Domain Dyadic Green Functions

By substituting (24) and (25) into (22) and referring to Fig. 2, we obtain the spectral-domain versions of DGFs [3]:

$$\begin{aligned} \underline{\underline{\tilde{\mathbf{G}}}}^{\text{EJ}}(\mathbf{k}_\rho, z|z') = & -\hat{\mathbf{u}}\hat{\mathbf{u}}V_i^{\text{TM}} - \hat{\mathbf{v}}\hat{\mathbf{v}}V_i^{\text{TE}} + \hat{\mathbf{z}}\hat{\mathbf{u}}\frac{k_\rho}{\omega\varepsilon(z)}I_i^{\text{TM}} \\ & + \hat{\mathbf{u}}\hat{\mathbf{z}}\frac{k_\rho}{\omega\varepsilon(z')}V_v^{\text{TM}} + \hat{\mathbf{z}}\hat{\mathbf{z}}\frac{1}{j\omega\varepsilon(z')}\left[\frac{k_\rho^2}{j\omega\varepsilon(z)}I_v^{\text{TM}} - \delta(z-z')\right], \end{aligned} \quad (26a)$$

$$\underline{\underline{\tilde{\mathbf{G}}}}^{\text{HJ}}(\mathbf{k}_\rho, z|z') = \hat{\mathbf{u}}\hat{\mathbf{v}}I_i^{\text{TE}} - \hat{\mathbf{v}}\hat{\mathbf{u}}I_i^{\text{TE}} - \hat{\mathbf{z}}\hat{\mathbf{v}}\frac{k_\rho}{\omega\mu}V_i^{\text{TE}} + \hat{\mathbf{v}}\hat{\mathbf{z}}\frac{k_\rho}{\omega\varepsilon(z')}I_v^{\text{TM}}. \quad (26b)$$

1.4 Spatial Domain Dyadic Green Functions

By taking an inverse Fourier transform (27),

$$\mathcal{F}^{-1}[\tilde{F}(\mathbf{k}_\rho, z)] \equiv f(\rho, z) = \frac{1}{(2\pi)^2} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \tilde{F}(\mathbf{k}_\rho, z) e^{j\mathbf{k}_\rho \cdot \boldsymbol{\rho}} d\mathbf{k}_x d\mathbf{k}_y. \quad (27)$$

we find the spatial domain analogues of the spectral domain DGFs, which are defined in (26). In case of rotational symmetry which implies that \tilde{F} only depends on one spectral co-ordinate k_ρ , the double integral in (27) can be simplified to an integral of only one variable. Using (4) and the coordinate transformation shown in Fig. 1 where,

$$k_x = k_\rho \cos \psi \quad \text{and} \quad k_y = k_\rho \sin \psi, \quad (28)$$

we rewrite (27) in cylindrical coordinates,

$$\mathcal{F}^{-1}[\tilde{F}(\mathbf{k}_\rho, z)] \equiv f(\rho, z) = \frac{1}{(2\pi)^2} \int_0^\infty \int_0^{2\pi} \tilde{F}(\mathbf{k}_\rho, z) e^{j\mathbf{k}_\rho \cdot \boldsymbol{\rho}} k_\rho dk_\rho d\psi. \quad (29)$$

Applying the Fourier-Bessel transform (FBT) to (29) which states that,

$$\frac{1}{2\pi} \int_0^{2\pi} e^{j\mathbf{k}_\rho \cdot \boldsymbol{\rho}} d\psi = J_0(k_\rho \rho), \quad (30)$$

where $J_0(\cdot)$ is the Bessel function of zero order we obtain the Sommerfeld integral (SI), $\mathcal{S}_0\{\cdot\}$:

$$f(\rho) = \mathcal{S}_0\{\tilde{F}\} \equiv \frac{1}{2\pi} \int_0^\infty J_0(k_\rho \rho) \tilde{F}(k_\rho) k_\rho dk_\rho. \quad (31)$$

In some cases, the Bessel function may be up to an order of 2, therefore, generalized expression for SI is,

$$\mathcal{S}_n\{\tilde{F}\} \equiv \frac{1}{2\pi} \int_0^\infty J_n(k_\rho \rho) \tilde{F}(k_\rho) k_\rho dk_\rho. \quad (32)$$

Following the above procedure, we write the spatial domain DGFs as:

$$\underline{\underline{\mathbf{G}}}^\kappa(\rho, z|z') = \mathcal{F}^{-1} \left[\underline{\underline{\tilde{\mathbf{G}}}}^\kappa(\mathbf{k}_\rho, z|z') \right] \equiv \mathcal{S}_n \left\{ \underline{\underline{\tilde{\mathbf{G}}}}^\kappa(\mathbf{k}_\rho, z|z') \right\} \quad (33)$$

where the superscript κ denotes either ‘EJ’ or ‘HJ’.

Starting with (26), we now discuss each step in detail to obtain the spatial domain versions of the DGFs. First, the spectral domain dyads are converted to their spatial domain counterparts using the relations listed in Table 1 which are obtained from Fig. 1 and (15). In general, a spatial domain DGF can be expressed in a matrix form:

$$\underline{\underline{\mathbf{G}}}^\kappa = \begin{bmatrix} G_{xx}^\kappa & G_{xy}^\kappa & G_{xz}^\kappa \\ G_{yx}^\kappa & G_{yy}^\kappa & G_{yz}^\kappa \\ G_{zx}^\kappa & G_{zy}^\kappa & G_{zz}^\kappa \end{bmatrix}, \quad (34)$$

As an example, we consider the spatial domain component, G_{xx}^{EJ} of the DGF, which

Spectral domain dyad	Spatial domain dyad
$\hat{\mathbf{u}}\hat{\mathbf{u}}$	$\hat{\mathbf{x}}\hat{\mathbf{x}} \cos^2 \psi + \hat{\mathbf{x}}\hat{\mathbf{y}} \cos \psi \sin \psi + \hat{\mathbf{y}}\hat{\mathbf{x}} \cos \psi \sin \psi + \hat{\mathbf{y}}\hat{\mathbf{y}} \sin^2 \psi$
$\hat{\mathbf{v}}\hat{\mathbf{v}}$	$\hat{\mathbf{x}}\hat{\mathbf{x}} \sin^2 \psi - \hat{\mathbf{x}}\hat{\mathbf{y}} \cos \psi \sin \psi - \hat{\mathbf{y}}\hat{\mathbf{x}} \cos \psi \sin \psi + \hat{\mathbf{y}}\hat{\mathbf{y}} \sin^2 \psi$
$\hat{\mathbf{u}}\hat{\mathbf{v}}$	$-\hat{\mathbf{x}}\hat{\mathbf{x}} \cos \psi \sin \psi + \hat{\mathbf{x}}\hat{\mathbf{y}} \cos^2 \psi - \hat{\mathbf{y}}\hat{\mathbf{x}} \sin^2 \psi + \hat{\mathbf{y}}\hat{\mathbf{y}} \cos \psi \sin \psi$
$\hat{\mathbf{v}}\hat{\mathbf{u}}$	$-\hat{\mathbf{x}}\hat{\mathbf{x}} \cos \psi \sin \psi - \hat{\mathbf{x}}\hat{\mathbf{y}} \sin^2 \psi + \hat{\mathbf{y}}\hat{\mathbf{x}} \cos^2 \psi + \hat{\mathbf{y}}\hat{\mathbf{y}} \cos \psi \sin \psi$
$\hat{\mathbf{z}}\hat{\mathbf{u}}$	$\hat{\mathbf{z}}\hat{\mathbf{x}} \cos \psi + \hat{\mathbf{z}}\hat{\mathbf{y}} \sin \psi$
$\hat{\mathbf{u}}\hat{\mathbf{z}}$	$\hat{\mathbf{x}}\hat{\mathbf{z}} \cos \psi + \hat{\mathbf{y}}\hat{\mathbf{z}} \sin \psi$
$\hat{\mathbf{z}}\hat{\mathbf{v}}$	$-\hat{\mathbf{z}}\hat{\mathbf{x}} \sin \psi + \hat{\mathbf{z}}\hat{\mathbf{y}} \cos \psi$
$\hat{\mathbf{v}}\hat{\mathbf{z}}$	$-\hat{\mathbf{x}}\hat{\mathbf{z}} \sin \psi + \hat{\mathbf{y}}\hat{\mathbf{z}} \cos \psi$

Table 1: Conversion of spectral domain dyad to spatial domain

can be expressed as:

$$\begin{aligned}
G_{xx}^{\text{EJ}} &= \mathcal{F}^{-1} \left\{ -\cos^2 \psi V_i^{\text{TM}} - \sin^2 \psi V_i^{\text{TE}} \right\} \\
&= -\mathcal{F}^{-1} \left\{ \frac{1 + \cos 2\psi}{2} V_i^{\text{TM}} + \frac{1 - \cos 2\psi}{2} V_i^{\text{TE}} \right\} \\
&= -\frac{1}{2} \mathcal{S}_0 \left\{ V_i^{\text{TM}} + V_i^{\text{TE}} \right\} - \mathcal{F}^{-1} \left\{ \frac{\cos 2\psi}{2} V_i^{\text{TM}} - V_i^{\text{TE}} \right\}
\end{aligned} \tag{35}$$

Using the formula [2],

$$\mathcal{F}^{-1} \left\{ \cos 2\psi \tilde{F} \right\} = -\cos 2\phi \mathcal{S}_2 \left\{ \tilde{F} \right\} \tag{36}$$

and the recurrence relation for Bessel-type functions [8],

$$J_{n+1}(z) = \frac{2n}{z} J_n(z) - J_{n-1}(z) \tag{37}$$

(35) is simplified to:

$$\begin{aligned}
G_{xx}^{\text{EJ}} &= -\frac{1}{2}\mathcal{S}_0 \left\{ V_i^{\text{TM}} \right\} - \frac{1}{2}\mathcal{S}_0 \left\{ V_i^{\text{TE}} \right\} + \frac{\cos 2\phi}{2}\mathcal{S}_2 \left\{ V_i^{\text{TM}} - V_i^{\text{TE}} \right\} \\
&= -\frac{1}{2}\mathcal{S}_0 \left\{ V_i^{\text{TM}} \right\} - \frac{1}{2}\mathcal{S}_0 \left\{ V_i^{\text{TE}} \right\} + \frac{\cos 2\phi}{2} \left[\frac{2}{\rho}\mathcal{S}_1 \left\{ V_i^{\text{TM}} - V_i^{\text{TE}} \right\} - \mathcal{S}_0 \left\{ V_i^{\text{TM}} - V_i^{\text{TE}} \right\} \right].
\end{aligned} \tag{38}$$

After some algebraic manipulation and knowing the fact that Sommerfeld integrals obey the linearity principle, we obtain the final expression:

$$G_{xx}^{\text{EJ}} = -\cos^2 \phi \mathcal{S}_0 \left\{ V_i^{\text{TM}} \right\} - \sin^2 \phi \mathcal{S}_0 \left\{ V_i^{\text{TE}} \right\} + \frac{\cos 2\phi}{2}\mathcal{S}_2 \left\{ V_i^{\text{TM}} - V_i^{\text{TE}} \right\}. \tag{39}$$

Following similar procedure, the remaining components of both electric and magnetic fields DGFs in the spatial domain can be obtained, which are listed in Table 2 [2].

1.5 Introducing Potential functions

As evident in Tables 2 and 3, the spatial domain DGFs that directly describe fields are tedious. An alternative formulation using vector potential functions considerably reduces the computation complexity. The approach is commonly known as mixed potentials integral equation (MPIE) formulation. From the Maxwell's equations, we know that,

$$\nabla \cdot \mu \mathbf{H} = 0, \tag{40}$$

and using the vector identity, $\nabla \cdot (\nabla \times \mathbf{A}) \equiv 0$, we define the magnetic vector potential, \mathbf{A} such that:

$$\mathbf{H} = \frac{1}{\mu} \nabla \times \mathbf{A}. \tag{41}$$

Through (2a) and (41), we write:

$$\nabla \times (\mathbf{E} + j\omega \mathbf{A}) = 0, \tag{42}$$

GF	$\kappa = \text{EJ}$
G_{xx}	$-\cos^2 \psi \mathcal{S}_0 \{V_i^{\text{TM}}\} - \sin^2 \psi \mathcal{S}_0 \{V_i^{\text{TE}}\} + \frac{\cos(2\psi)}{\rho} \mathcal{S}_1 \left\{ \frac{V_i^{\text{TM}} - V_i^{\text{TE}}}{k_\rho} \right\}$
G_{xy}	$-\frac{\sin(2\psi)}{2} \mathcal{S}_0 \{V_i^{\text{TM}} - V_i^{\text{TE}}\} + \frac{\sin(2\psi)}{\rho} \mathcal{S}_1 \left\{ \frac{V_i^{\text{TM}} - V_i^{\text{TE}}}{k_\rho} \right\}$
G_{xz}	$\frac{\eta_0}{j k_0 \varepsilon(z')} \cos \psi \mathcal{S}_1 \{k_\rho V_v^{\text{TM}}\}$
G_{yx}	$-\frac{\sin(2\psi)}{2} \mathcal{S}_0 \{V_i^{\text{TM}} - V_i^{\text{TE}}\} + \frac{\sin(2\psi)}{\rho} \mathcal{S}_1 \left\{ \frac{V_i^{\text{TM}} - V_i^{\text{TE}}}{k_\rho} \right\}$
G_{yy}	$-\sin^2 \psi \mathcal{S}_0 \{V_i^{\text{TM}}\} - \cos^2 \psi \mathcal{S}_0 \{V_i^{\text{TE}}\} - \frac{\cos(2\psi)}{\rho} \mathcal{S}_1 \left\{ \frac{V_i^{\text{TM}} - V_i^{\text{TE}}}{k_\rho} \right\}$
G_{yz}	$\frac{\eta_0}{j k_0 \varepsilon(z')} \sin \psi \mathcal{S}_1 \{k_\rho V_v^{\text{TM}}\}$
G_{zx}	$\frac{\eta_0}{j k_0 \varepsilon(z)} \cos \psi \mathcal{S}_1 \{k_\rho I_i^{\text{TM}}\}$
G_{zy}	$\frac{\eta_0}{j k_0 \varepsilon(z)} \sin \psi \mathcal{S}_1 \{k_\rho I_i^{\text{TM}}\}$
G_{zz}	$-\frac{\eta_0^2}{j k_0^2 \varepsilon(z) \varepsilon(z')} \mathcal{S}_0 \{k_\rho^2 I_v^{\text{TM}}\} - \frac{\eta_0}{j k_0 \varepsilon(z)} \delta(\rho) \delta(z - z')$

Table 2: Scalar Green functions for computation of electric field due to an electric current source [2]

and knowing the vector identity, $\nabla \times (-\nabla \phi) \equiv 0$, the electric field is expressed as:

$$\mathbf{E} = j\omega \mathbf{A} - \nabla \phi. \quad (43)$$

The scalar potential function ϕ can be related to A by the Lorenz gauge, which states that:

$$\nabla \cdot \mathbf{A} = -j\omega \mu \varepsilon \phi. \quad (44)$$

The electric field in (43) is then expressed only as a function of \mathbf{A} ,

$$\mathbf{E} = \frac{j\omega}{k^2} (k^2 + \nabla \nabla \cdot) \mathbf{A}. \quad (45)$$

GF	$\kappa = \text{HJ}$
G_{xx}	$-\frac{\sin(2\psi)}{2}\mathcal{S}_0\{I_i^{\text{TE}} - I_i^{\text{TM}}\} + \frac{\sin(2\psi)}{\rho}\mathcal{S}_1\left\{\frac{I_i^{\text{TE}} - I_i^{\text{TM}}}{k_\rho}\right\}$
G_{xy}	$\cos^2\psi\mathcal{S}_0\{I_i^{\text{TE}}\} + \sin^2\psi\mathcal{S}_0\{I_i^{\text{TM}}\} - \frac{\cos(2\psi)}{\rho}\mathcal{S}_1\left\{\frac{I_i^{\text{TE}} - I_i^{\text{TM}}}{k_\rho}\right\}$
G_{xz}	$-\frac{\eta_0}{jk_0\varepsilon(z)}\sin\psi\mathcal{S}_1\{k_\rho I_v^{\text{TM}}\}$
G_{yx}	$-\sin^2\psi\mathcal{S}_0\{I_i^{\text{TE}}\} - \cos^2\psi\mathcal{S}_0\{I_i^{\text{TM}}\} - \frac{\cos(2\psi)}{\rho}\mathcal{S}_1\left\{\frac{I_i^{\text{TE}} - I_i^{\text{TM}}}{k_\rho}\right\}$
G_{yy}	$\frac{\sin(2\psi)}{2}\mathcal{S}_0\{I_i^{\text{TE}} - I_i^{\text{TM}}\} - \frac{\sin(2\psi)}{\rho}\mathcal{S}_1\left\{\frac{I_i^{\text{TE}} - I_i^{\text{TM}}}{k_\rho}\right\}$
G_{yz}	$\frac{\eta_0}{jk_0\varepsilon(z')}\sin\psi\mathcal{S}_1\{k_\rho I_v^{\text{TM}}\}$
G_{zx}	$\frac{1}{jk_0\eta_0\mu}\sin\psi\mathcal{S}_1\{k_\rho V_i^{\text{TE}}\}$
G_{zy}	$-\frac{1}{jk_0\eta_0\mu}\cos\psi\mathcal{S}_1\{k_\rho V_i^{\text{TE}}\}$
G_{zz}	.

Table 3: Scalar Green functions for computation of magnetic field due to an electric current source [2]

For a current distribution \mathbf{J} , the vector potential \mathbf{A} can also be described by a DGF in a similar fashion to (1):

$$\mathbf{A} = \int_{r'} \underline{\underline{\mathbf{G}}}^{\text{A}}(\mathbf{r}|\mathbf{r}')\mathbf{J}(\mathbf{r}') d\mathbf{r}', \quad (46)$$

where $\underline{\underline{\mathbf{G}}}^{\text{A}}$ is the vector potential DGF. Following the formulation of [3], the spectral domain form is,

$$\underline{\underline{\mathbf{G}}}^{\text{A}} = (\hat{\mathbf{u}}\hat{\mathbf{u}} + \hat{\mathbf{v}}\hat{\mathbf{v}})\tilde{G}_{vv}^{\text{A}} + \hat{\mathbf{z}}\hat{\mathbf{u}}\tilde{G}_{zu}^{\text{A}} + \hat{\mathbf{z}}\hat{\mathbf{z}}\tilde{G}_{zz}^{\text{A}}. \quad (47)$$

The vector potential DGF is related to its fields counterpart through the magnetic

field in (41) and written as:

$$\underline{\underline{\tilde{\mathbf{G}}}}^{\text{HJ}} = \frac{1}{\mu} \nabla \times \underline{\underline{\tilde{\mathbf{G}}}}^{\text{A}} \quad (48)$$

where isotropic permeability is assumed. To find the components of $\underline{\underline{\tilde{\mathbf{G}}}}^{\text{A}}$, we first note that the ∇ operator used in the curl operation, which was initially defined in rectangular coordinates, is re-written in the spectral domain form as $\tilde{\nabla} = -jk_\rho \hat{\mathbf{u}} + \hat{\mathbf{z}} d/dz$. Therefore, (48) becomes,

$$\begin{aligned} \nabla \times \underline{\underline{\tilde{\mathbf{G}}}}^{\text{A}} &= (-jk_\rho \hat{\mathbf{u}} + \hat{\mathbf{z}} d/dz) \times \left[(\hat{\mathbf{u}}\hat{\mathbf{u}} + \hat{\mathbf{v}}\hat{\mathbf{v}}) \tilde{G}_{vv}^{\text{A}} + \hat{\mathbf{z}}\hat{\mathbf{u}} \tilde{G}_{zu}^{\text{A}} + \hat{\mathbf{z}}\hat{\mathbf{z}} \tilde{G}_{zz}^{\text{A}} \right] \\ &= -jk_\rho \tilde{G}_{vv}^{\text{A}} (\hat{\mathbf{u}} \times \hat{\mathbf{v}}) \hat{\mathbf{v}} + \frac{d\tilde{G}_{vv}^{\text{A}}}{dz} (\hat{\mathbf{z}} \times \hat{\mathbf{u}}) \hat{\mathbf{u}} - jk_\rho \tilde{G}_{zu}^{\text{A}} (\hat{\mathbf{u}} \times \hat{\mathbf{z}}) \hat{\mathbf{u}} - jk_\rho \tilde{G}_{zz}^{\text{A}} (\hat{\mathbf{u}} \times \hat{\mathbf{z}}) \hat{\mathbf{z}} \\ &= -jk_\rho \tilde{G}_{vv}^{\text{A}} \hat{\mathbf{z}} \hat{\mathbf{v}} + \frac{d\tilde{G}_{vv}^{\text{A}}}{dz} \hat{\mathbf{v}} \hat{\mathbf{u}} + jk_\rho \tilde{G}_{zu}^{\text{A}} \hat{\mathbf{v}} \hat{\mathbf{u}} + jk_\rho \tilde{G}_{zz}^{\text{A}} \hat{\mathbf{v}} \hat{\mathbf{z}}. \end{aligned} \quad (49)$$

Using (48) and comparing the dyads of (26b) and (49), we obtain the scalar components of the vector DGF, $\underline{\underline{\tilde{\mathbf{G}}}}^{\text{A}}$:

$$j\omega \tilde{G}_{vv}^{\text{A}} = V_i^{\text{TE}}, \quad (50a)$$

$$j\omega \mu \varepsilon(z') \tilde{G}_{zz}^{\text{A}} = I_v^{\text{TM}}, \quad (50b)$$

$$\frac{d\tilde{G}_{vv}^{\text{A}}}{dz} + jk_\rho \tilde{G}_{zu}^{\text{A}} = -\mu I_v^{\text{TM}}. \quad (50c)$$

Substituting (50a) into (50c) and invoking the telegrapher's equation, (20a), we obtain:

$$\begin{aligned} \frac{1}{j\omega} \left[-jk_z^{\text{TE}} Z^{\text{TE}} I_i^{\text{TE}} \right] + jk_\rho \tilde{G}_{zu}^{\text{A}} &= -\mu I_v^{\text{TM}} \\ \tilde{G}_{zu}^{\text{A}} &= \frac{\mu}{jk_\rho} \left(I_i^{\text{TE}} - I_i^{\text{TM}} \right), \end{aligned} \quad (51)$$

where $Z^{\text{TE}} = \omega \mu / k_z^{\text{TE}}$.

1.6 Spatial domain DGFs

The spectral domain vector potentials defined in (47) can be conveniently written in a matrix form containing scalar elements,

$$\underline{\underline{\tilde{\mathbf{G}}}}^A = \begin{bmatrix} \tilde{G}_{vv}^A & & \\ & \tilde{G}_{vv}^A & \\ \frac{k_x}{k_\rho} \tilde{G}_{zu}^A & \frac{k_y}{k_\rho} \tilde{G}_{zu}^A & \tilde{G}_{zz}^A \end{bmatrix}. \quad (52)$$

A Fourier inversion following a similar procedure used to convert fields DGFs (35)-(39) yields the spatial domain potential DGFs expressed in terms of TLGFs, which are listed in Table 4.

$G_{xx}^A = G_{yy}^A$	$\frac{1}{j\omega} \mathcal{S}_0 \{V_i^{\text{TE}}\}$
G_{zx}^A	$j\mu \cos \phi \mathcal{S}_1 \left\{ \frac{I_i^{\text{TE}} - I_i^{\text{TM}}}{k_\rho} \right\}$
G_{zy}^A	$j\mu \sin \phi \mathcal{S}_1 \left\{ \frac{I_i^{\text{TE}} - I_i^{\text{TM}}}{k_\rho} \right\}$
G_{zz}^A	$\frac{1}{j\omega\mu\epsilon(z')} \mathcal{S}_0 \{I_v^{\text{TM}}\}$

Table 4: Scalar potential Green functions expressed in terms of TLGFs [9]

The above formulation is applicable to only multi-layered structure in which the layers are stacked along the x-y plane. The formulation allows insertion of arbitrarily oriented electric current source that can be placed anywhere in the structure. Once the magnetic vector potential, \mathbf{A} , is computed using (46), the corresponding electric and magnetic fields can be found through (45) and (41) respectively.

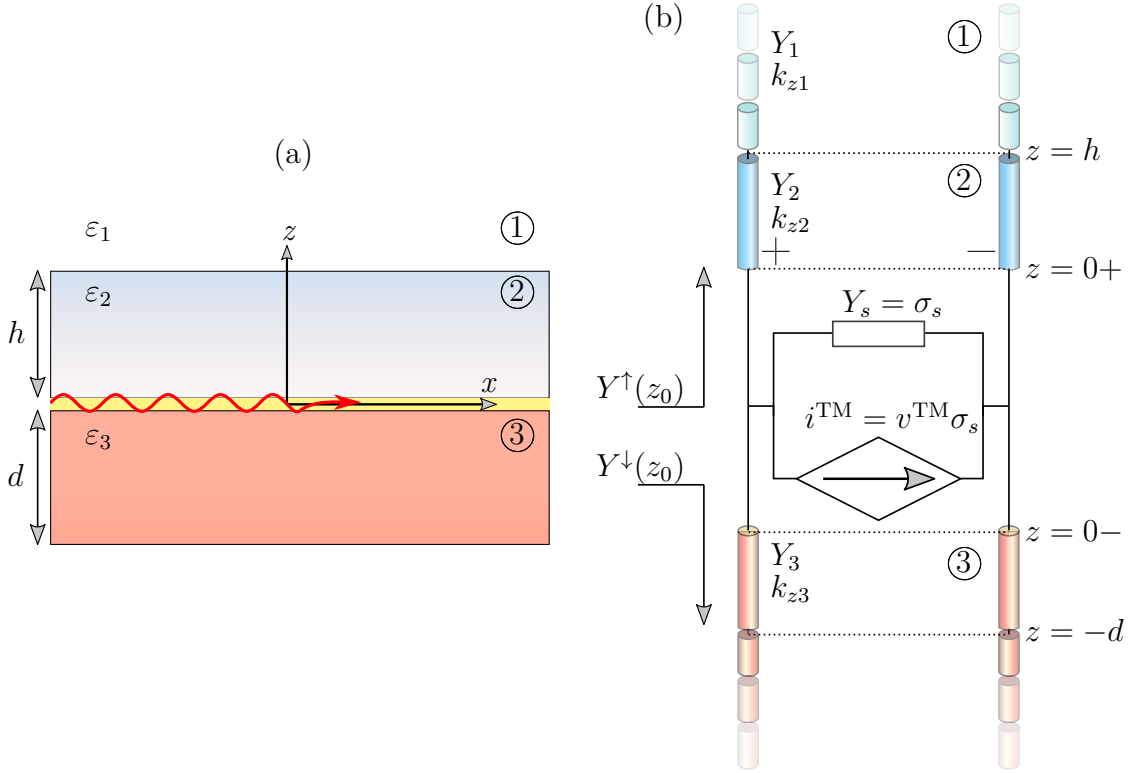


Figure 3: (a) Multilayer structure typically found in a high electron mobility transistor, (b) Equivalent transmission line network

2 Computation of Transmission line Green functions

Following network analysis techniques, the voltage and current at any point in a transmission line can be obtained [4]. Each layer in the multi-layered structure is represented by a section in the transmission line which is specified by its characteristic impedance Z and the propagation constant k_z . As an example, an equivalent TL network for a semiconductor heterostructure that forms the substrate of the modern transistors is shown in Fig. 3. The layers numbered 2 and 3 in Fig. 3a are commonly made up of group III-V materials such as gallium nitride (GaN) and their subsequent alloys such aluminum gallium nitride (AlGaN). Based on the device type, the top layer labeled 1, can represent either freespace or a perfectly conducting material.

The thin layer which is sandwiched between the semiconductor layers is known as a two-dimensional electron gas (2DEG) which can be visualized as a highly conductive sheet of electrons that is known to exhibit extraordinary electromagnetic properties such as highly confined surface waves [10, 11]. In this section, we formulate the TLGFs for various configuration of the 2DEG and compute the respective potential Green functions.

2.1 Lossy Half-space

We start with a case in which the 2DEG occupies the whole lower half-space. Such a case is an example of the classical Sommerfeld half-space problem that has been actively researched in the past century. [?, ?, ?] An electric dipole of arbitrary orientation is placed in air at a height z' from the interface where the origin is also centered. The observation point, at first also lies in air. To find the component G_{xx}^A , we require TE-mode voltage of an equivalent, current excited TL network which is expressed using (25a) and (25b) [4, Sec. 2.4] [3],

$$V_i^{\text{TE}}(z, z') = \frac{Z_1^{\text{TE}}}{2} \left[e^{-jk_{z1}|z-z'|} + \Gamma^{\downarrow, \text{TE}} e^{-jk_{z1}(z+z')} \right], \quad (53)$$

where $\Gamma^{\downarrow, \text{TE}} = (Z_{2\text{DEG}}^{\text{TE}} - Z_1^{\text{TE}})/(Z_{2\text{DEG}}^{\text{TE}} + Z_1^{\text{TE}})$, $Z_{2\text{DEG}}^{\text{TE}}$ and Z_1^{TE} are the TE mode impedances of the 2DEG and freespace regions respectively, given by (21b). The scalar potential GF, G_{xx}^A is then expressed as:

$$G_{xx}^A = \mathcal{S}_0 \left\{ \frac{e^{-jk_{z1}|z-z'|}}{2jk_{z1}} \right\} + \mathcal{S}_0 \left\{ \frac{\Gamma^{\downarrow, \text{TE}} e^{-jk_{z1}(z+z')}}{2jk_{z1}} \right\}. \quad (54)$$

Following similar steps, G_{zz}^A is found using the TM-mode current of a voltage excited network. For the components G_{zx}^A and G_{zy}^A , both the TM and TE-mode currents need to be considered for a current excited network. The final expressions are:

$$G_{zz}^A = \mathcal{S}_0 \left\{ \frac{e^{-jk_{z1}|z-z'|}}{2jk_{z1}} \right\} - \mathcal{S}_0 \left\{ \frac{\Gamma^{\downarrow, \text{TM}} e^{-jk_{z1}(z+z')}}{2jk_{z1}} \right\} \quad (55)$$

and,

$$G_{zx}^A = \frac{j\mu}{2} \cos \phi \mathcal{S}_1 \left\{ \frac{\Gamma^{\downarrow, \text{TM}} - \Gamma^{\downarrow, \text{TE}}}{k_\rho} e^{-jk_{z1}(z+z')} \right\}. \quad (56)$$

The first term in (54) and (55) is the Sommerfeld identity which can be written as the freespace Green function of a point source,

$$4\pi\mathcal{S}_0 \left\{ \frac{e^{-jk_{z1}|z-z'|}}{2jk_{z1}} = \frac{e^{-jk_1 R}}{4\pi R} \right\}. \quad (57)$$

The plasma frequency ω_p of the 2DEG lies in the terahertz frequency at which the dielectric function vanishes to zero. At frequencies lower than ω_p , the real part is negative and the imaginary part, which expresses loss is highly dependent on temperature. Here, we assume a dielectric constant of $-8.114 - j8.159$ for a GaN/AlGaN heterostructure at 5.600 THz, and observed at room temperature (295 K).

For a horizontally oriented electric dipole, only the spatial domain GFs in G_{xx}^A (54), G_{zx}^A (56) exist, and are plotted in Fig. 4 by excluding the freespace GF part. To

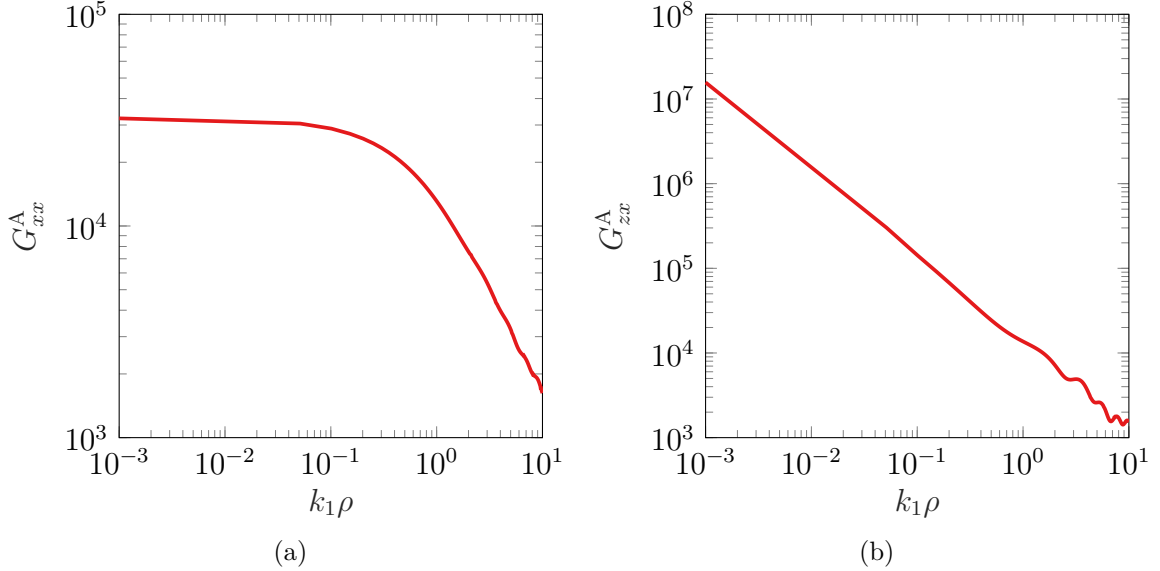


Figure 4: Components of potential DGF for a 2DEG half-space (a) G_{xx}^A (54), (b) G_{zx}^A (56)

compute the SIs, real-axis integration using a combination of *tanh-sinh* quadrature

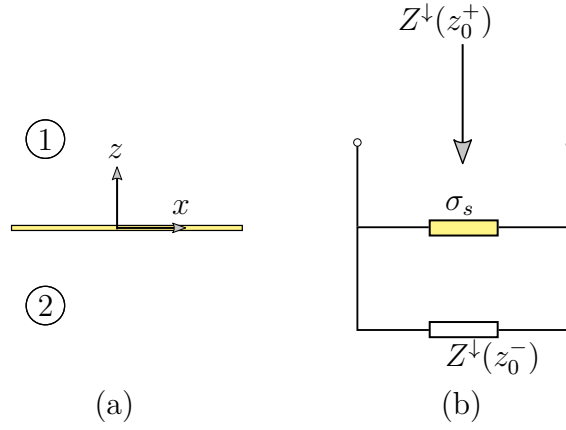


Figure 5: (a) 2DEG sheet in surrounded by freespace, (b) Equivalent transmission line circuit consisting of a shunt impedance

[?, ?] and partition extrapolation [?] for the tail part is used by implementing the algorithms in [?].

2.2 Thin Sheet

Next we consider the case of an infinitesimally thin sheet of 2DEG located on the x-y plane and surrounded by freespace. A horizontal electric dipole is placed at a height h above the sheet. The surface conductivity of the thin sheet is expressed in terms of a Drude-type model:

$$\sigma_s = \frac{N_s e^2 \tau}{m^*} \frac{1}{1 + j\omega\tau}, \quad (58)$$

where N_s is the free electron concentration at the interface, e is the electron charge, m^* is the effective electron mass in the heterostructure, τ is the scattering time of electrons. The dielectric function is subsequently expressed as,

$$\varepsilon(\omega) = 1 - j \frac{\sigma_s}{\omega \Delta \varepsilon_r \varepsilon_0}, \quad (59)$$

in which Δ is the thickness of the 2DEG region, ε_r is the dielectric constant of the GaN layer which acts as a substrate and ε_0 is the freespace permittivity. For a 5 nm

thick GaN/AlGaIn based 2DEG, where $N_s = 7.5 \times 10^{12} \text{ cm}^{-2}$ and $\varepsilon_r = 9.6$ [?], the surface conductivity and dielectric function are plotted in Fig. 6 at room temperature. The equivalent transmission line network used to calculate the reflection coefficients

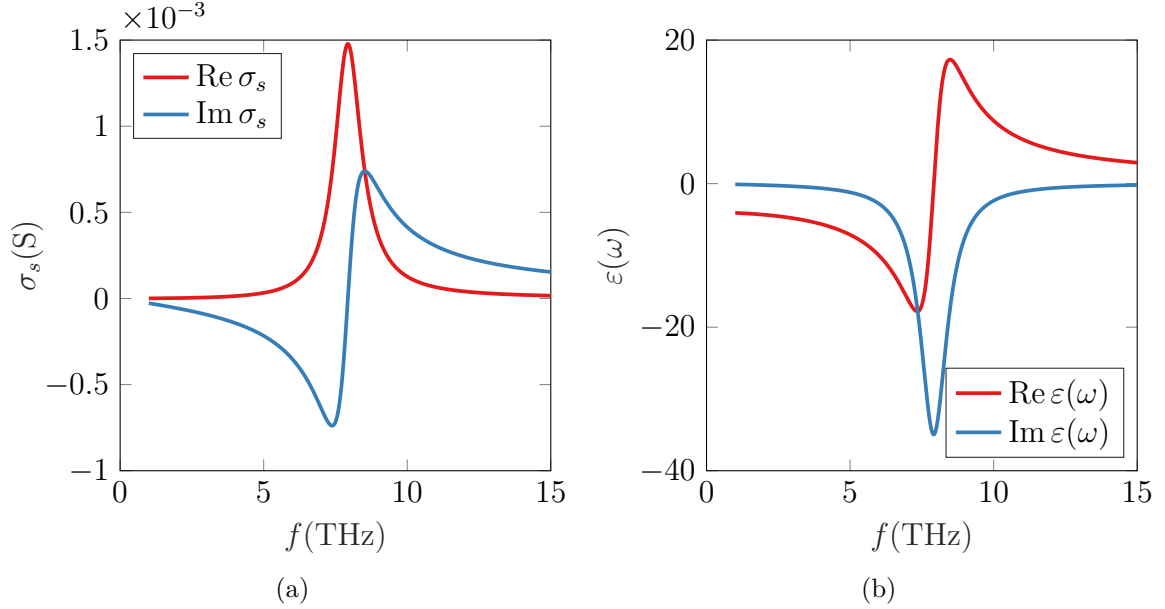


Figure 6: Electrical properties of GaN/AlGaIn based 2DEG at room temperature (a) Surface conductivity (58) (b) Dielectric function (59)

consists of a simple circuit of two parallel impedances [12, 13, 14] which is shown in Fig. 5.

Assuming the source lies on the top of the 2DEG, the down-looking impedance of the shunt configuration circuit is,

$$Z^\downarrow(z_0^+) = \frac{Z^\downarrow(z_0^-)}{1 + \sigma_s Z^\downarrow(z_0^-)} \quad (60)$$

where it is noted that σ_s has the dimensions of admittance, and $Z^\downarrow(z_0^-)$ is the down-looking effective impedance as seen from the bottom of the sheet. In the case considered, $Z^\downarrow(z_0^-) = Z_0$ where Z_0 is the freespace intrinsic impedance. The reflections needed in (54)-(56) are listed in Table 5.

$\Gamma^{\downarrow, \text{TE}}$	$\Gamma^{\downarrow, \text{TM}}$
$\frac{k_{z1} - \omega\mu_1\sigma_s}{k_{z1} + \omega\mu_1\sigma_s}$	$\frac{\omega\varepsilon_1 - \sigma_s k_{z1}}{\omega\varepsilon_1 + \sigma_s k_{z1}}$

Table 5: Reflection coefficients for 2DEG sheet in freespace

2.2.1 Location of poles

The integrands of the SIs in (54)-(56) contain poles that are associated with the zeros of the integrand denominator. For a thin sheet in free-space, the poles can be analytically found by applying using the transverse resonance (TR) condition [4, Sec. 2.4], which requires that the total impedance at any point ($z = 0$ in this case) must be zero. Using the TM parameters given in Table 5 and (60), we write:

$$Z^{\downarrow} + Z^{\uparrow} = 0, \quad (61a)$$

$$Z_1 + Z_1(Z_1\sigma_s + 1) = 0. \quad (61b)$$

The expression is obtained by substituting (21a) into (61b) [?],

$$k_P^{\text{TM}} = k_1 \sqrt{1 - \left(\frac{2}{\eta_0\sigma_s}\right)^2} \quad (62)$$

Similarly the TE-mode pole is expressed as,

$$k_P^{\text{TE}} = k_1 \sqrt{1 - \left(\frac{\eta_0\sigma_s}{2}\right)^2} \quad (63)$$

A mode is termed as proper only when the pole lies in the fourth quadrant of the complex plane. As evident from Figs. 6 and 7a, proper TM modes only exist when the real part of dielectric function of the 2DEG is negative which also corresponds to a negative imaginary part of the conductivity. On the other hand, a proper TE mode exists when the real part of dielectric function is positive which is shown in

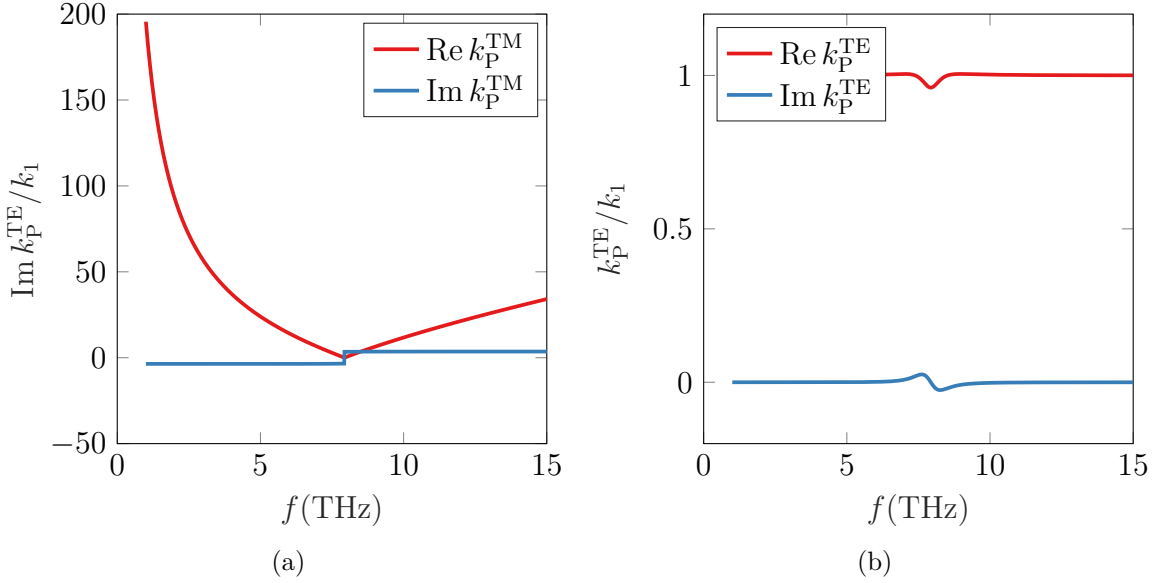


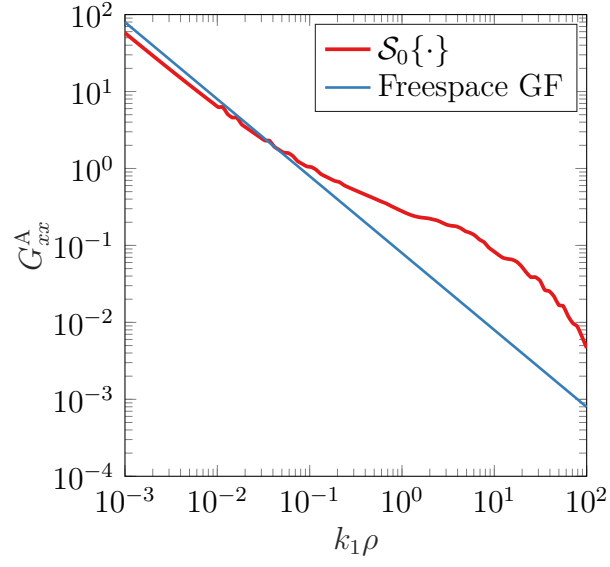
Figure 7: Distribution of analytically calculated poles in the complex plane (a) TM mode (62) (b) TE mode (63)

Fig. 7b.

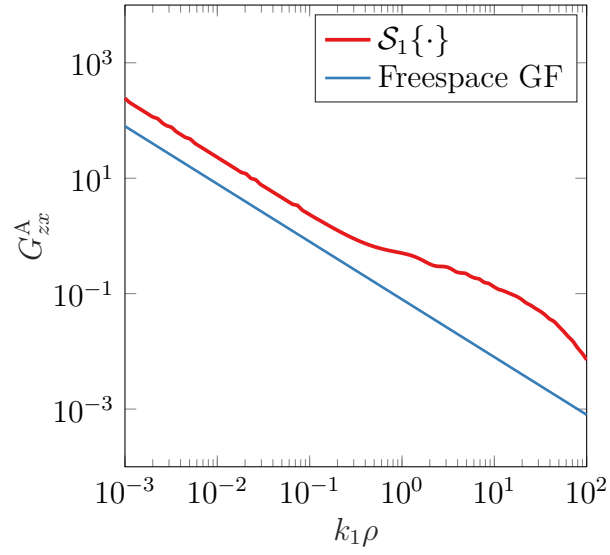
Figure ?? shows the three components of the potential GF, G^A , which are compared with the freespace GF.

The effect of the plasmon mode on the curves is the oscillations that are present at average horizontal distances. For small horizontal distances, the near field effect dominates, while for larger horizontal distances, the losses of the plasmonic mode are not any more present. Moreover, at relatively low frequencies, the plasmon mode has extremely large losses [113], and consequently the effect of the plasmon in the curves of $f = 1$ THz is not so intense and a behavior closer to the one of free space is obtained.

In this chapter, a transmission line network based technique is discussed through which spectral and subsequently spatial domain potential GFs have been derived for a multilayer structure. Of particular interest is an infinitesimally thin sheet of free-charges known as a 2DEG, for which components of the potential GFs are computed. Existence of surface waves supported by the 2DEG is explored. Based

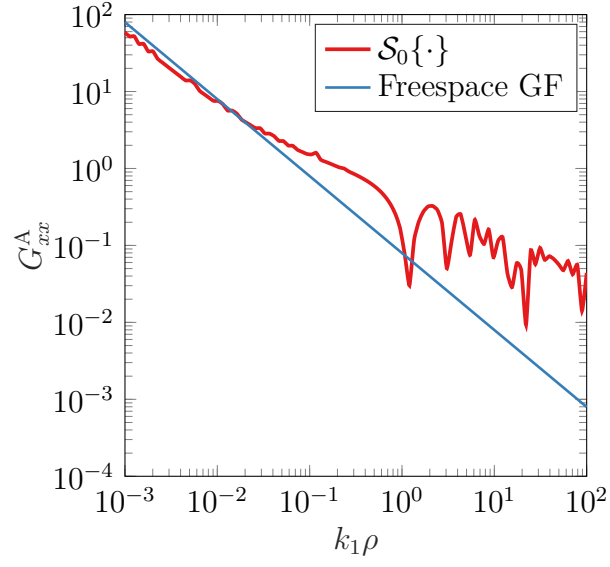


(a)

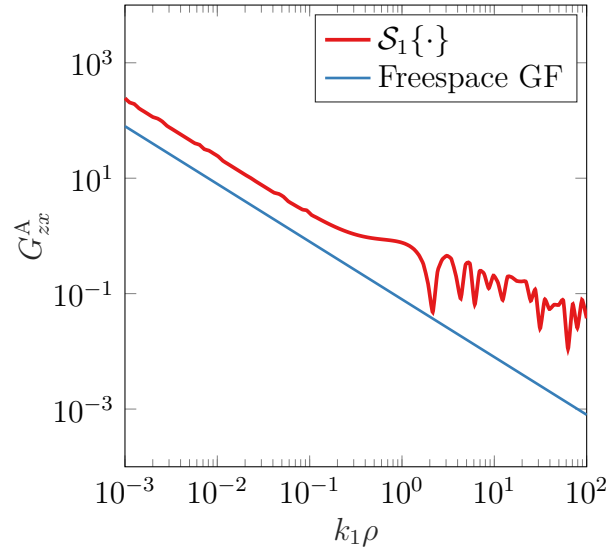


(b)

Figure 8: Components of G^A computed for a GaN/AlGaIn based 2DEG sheet suspended in freespace. The surface conductivity of the sheet is $\sigma_s = 7.6 \times 10^{-5} - j2.98 \times 10^{-3} \text{ S}$ corresponding to 5.6 THz at room temperature (295 K). (a) G_{xx}^A , (b) G_{zx}^A



(a)



(b)

Figure 9: Components of G^A computed for a GaN/AlGaIn based 2DEG sheet suspended in freespace. The cryogenic surface conductivity of the sheet is $\sigma_s = 7.6 \times 10^{-8} - j2.98 \times 10^{-3} \text{ S}$ corresponding to 5.6 THz at a temperature (3 K). (a) G_{xx}^A , (b) G_{zx}^A

on the location of poles that are analytically calculated, it is concluded that such structure do support plasmonic behavior in the terahertz frequency region.

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