



# Templated CUDA Lattice Boltzmann Method: generic CFD solver for single and multi-phase problems

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Michał Dzikowski<sup>a</sup>, Grzegorz Gruszczyński<sup>a,b</sup>

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<sup>a</sup>University of Warsaw, Interdisciplinary Centre for Mathematical and Computational Modelling

<sup>b</sup>Warsaw University of Technology, Faculty of Power and Aeronautical Engineering

## Introduction

- Problem Statement and Governing Equations

- LBM - Theory

- Discrete Boltzmann equation

- LBM - Algorithm

## Multiphysics models

- Advection-Diffusion Equation revisited

## Sample simulations

## Solver Structure

## Memory requirements

# Introduction

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## Goal

Simulate an incompressible flow coupled with advection - diffusion equation of some scalar field.

The continuity and momentum equations are:

$$\begin{cases} \frac{\partial \rho}{\partial t} + \nabla \cdot \rho \mathbf{u} = 0 \\ \rho \left( \frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} \right) = -\nabla p + \nabla \cdot (\mu [\nabla \mathbf{u} + (\nabla \mathbf{u})^T]) + \mathbf{F} \end{cases}$$

And the advection - diffusion equation of some scalar field  $\phi$ :

$$\frac{\partial \phi}{\partial t} + \nabla \cdot (\phi \mathbf{u}) = \nabla \cdot (M \nabla \phi)$$

# Lattice Boltzmann Method

Probability of finding a particle in the phase space:

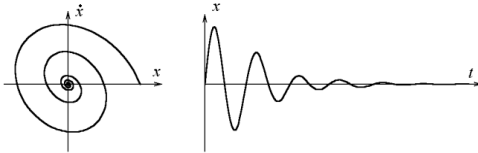


Figure 1:  $\Psi = \Psi(t, \mathbf{x}, \dot{\mathbf{x}})$

In an infinitesimally small volume of the phase space  $d\mathbf{x}d\mathbf{u}$ :

$$\Psi_{no\ collisions}(t + dt, \mathbf{x} + d\mathbf{x}, \mathbf{u} + d\mathbf{u})d\mathbf{x}d\mathbf{u} = \Psi_{no\ collisions}(t, \mathbf{x}, \mathbf{u})d\mathbf{x}d\mathbf{u}$$

Now, include the collision term  $\mathbb{C}(\Psi)$ :

$$\Psi(t + dt, \mathbf{x} + d\mathbf{x}, \mathbf{u} + d\mathbf{u})d\mathbf{x}d\mathbf{u} = \Psi(t, \mathbf{x}, \mathbf{u})d\mathbf{x}d\mathbf{u} + \mathbb{C}(\Psi)d\mathbf{x}d\mathbf{u}dt$$

Taylor series expansion:

$$\Psi(t + dt, \mathbf{x} + d\mathbf{x}, \mathbf{u} + d\mathbf{u}) = \Psi(t, \mathbf{x}, \mathbf{u}) + \frac{\partial \Psi}{\partial t} dt + \nabla_{\mathbf{x}} \Psi d\mathbf{x} + \nabla_{\mathbf{u}} \Psi d\mathbf{u}$$

Plug in:

$$\left[ \Psi(t, \mathbf{x}, \mathbf{u}) + \frac{\partial \Psi}{\partial t} dt + \nabla_{\mathbf{x}} \Psi d\mathbf{x} + \nabla_{\mathbf{u}} \Psi d\mathbf{u} \right] d\mathbf{x} d\mathbf{u} = \left[ \Psi(t, \mathbf{x}, \mathbf{u}) + \mathbb{C}(\Psi) dt \right] d\mathbf{x} d\mathbf{u}$$

Reformulate velocity  $\mathbf{u} = \frac{d\mathbf{x}}{dt}$  and acceleration  $\frac{d\mathbf{u}}{dt} = \frac{\mathbf{F}}{\rho}$ :

$$\frac{\partial \Psi}{\partial t} + (\mathbf{u} \cdot \nabla_{\mathbf{x}}) \Psi + \left( \frac{\mathbf{F}}{\rho} \cdot \nabla_{\mathbf{u}} \right) \Psi = \mathbb{C}(\Psi)$$

Streaming and Collision:

$$\underbrace{\Psi(t + dt, \mathbf{x} + d\mathbf{x}, \mathbf{u} + d\mathbf{u})d\mathbf{x}d\mathbf{u}}_{\text{Streaming}} = \underbrace{\Psi(t, \mathbf{x}, \mathbf{u})d\mathbf{x}d\mathbf{u} + \mathbb{C}(\Psi)d\mathbf{x}d\mathbf{u}dt}_{\text{Collision}}$$

The Boltzmann equation can be viewed as a substantial derivative (of an intensive quantity  $\Psi$ ) which is equal to the collision term  $\mathbb{C}$  applied to the distribution function of  $\Psi$ :

$$\frac{\partial \Psi}{\partial t} + (\mathbf{u} \cdot \nabla_{\mathbf{x}})\Psi + \left(\frac{\mathbf{F}}{\rho} \cdot \nabla_{\mathbf{u}}\right)\Psi = \mathbb{C}(\Psi)$$



$$\underbrace{f_i(\mathbf{x} + \mathbf{e}_i \Delta \mathbf{x}, t + \Delta t)}_{\text{Streaming}} = f_i(\mathbf{x}, t) - \underbrace{\frac{1}{\tau} (f_i - f_i^{eq})}_{\text{Collision}} + F_i(\mathbf{x}, t)$$

- $\tau = \tau(\nu)$  relaxation parameter,  $\nu$  is the kinematic viscosity
- $f_i$  - discrete probability distribution function
- $F_i$  - source term (ex. gravity force)

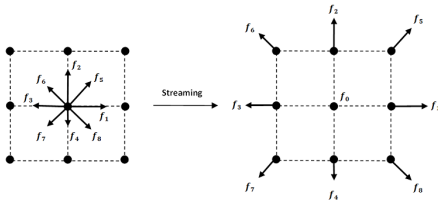


Figure 2: D2Q9: Streaming

1 Initialize  $f_i^n$

2 Compute

$$\rho = \sum_{i=0}^8 f_i^n(x, t) \quad \text{and} \quad u(x, t) = \frac{1}{\rho} \sum_{i=0}^8 f_i^n(x, t) e_i + \frac{F}{2\rho} \delta t$$

3 Compute  $f_i^{eq}(x, t) = w_i \rho \left[ 1 + \frac{e_i u}{c_s^2} + \frac{(e_i u)^2}{2c_s^4} - \frac{u^2}{2c_s^2} \right]$  where  $c_s^2 = \frac{1}{3}$

4 Collision  $f_i^{out}(x, t) = f_i^n(x, t) - \frac{1}{\tau_f} \left[ f_i^n(x, t) - f_i^{eq}(x, t) \right] + F_i(x, t)$

5 Streaming  $f_i^n(x + e_i, t + 1) = f_i^{out}(x, t)$

# Multiphysics models

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How to re-use the same algorithm to track temperature field or interface between phases?

Consider Advection-Diffusion Equation of scalar field  $\phi$ .

$$\frac{\partial \phi}{\partial t} + \nabla \cdot (\phi \mathbf{u}) = \nabla \cdot (M \nabla \phi)$$

The  $\phi$  field can be viewed as a temperature ... or a phase indicator.

Imagine a separation flux,  $\mathbf{j}_S$ , supposed to counteract the diffusion and reach a predefined interface profile in the equilibrium state:

$$\frac{\partial \phi}{\partial t} + \nabla \cdot (\phi \mathbf{u}) = \nabla \cdot (\underbrace{M \nabla \phi}_{\mathbf{j}_D} - \mathbf{j}_S)$$

Let us use a *tanh* to smooth the step interface:

$$\phi^{eq} = \frac{1}{2} \tanh \left( \frac{2(\mathbf{x} - \mathbf{x}_0)}{\gamma} \right)$$

Evaluate diffusive flux in the equilibrium:

$$\begin{aligned}
 j_D^{eq} &= M \nabla \left[ \overbrace{\frac{1}{2} \tanh \left( \frac{2(\mathbf{x} - \mathbf{x}_0)}{\gamma} \right)}^{\phi^{eq}} \right] = \frac{M}{2} \mathbf{n} \frac{\partial}{\partial \mathbf{x}_n} \tanh \left[ \left( \frac{2(\mathbf{x} - \mathbf{x}_0)}{\gamma} \right) \right] \\
 &= \frac{M}{\gamma} \mathbf{n} \left[ \underbrace{1 - \tanh^2 \left( \frac{2(\mathbf{x} - \mathbf{x}_0)}{\gamma} \right)}_{1 - 4(\phi^{eq})^2} \right]
 \end{aligned}$$

Therefore:

$$j_S = M \mathbf{n} \frac{1 - 4\phi^2}{\gamma} \quad \text{where} \quad \mathbf{n} = \frac{\nabla \phi}{|\nabla \phi|}$$

## Conclusion

Modifying the advection-diffusion equation:

$$\frac{\partial \phi}{\partial t} + \nabla \cdot (\phi \mathbf{u}) = \nabla \cdot (M \nabla \phi)$$

The interface tracking equation can be obtained:

$$\frac{\partial \phi}{\partial t} + \nabla \cdot \phi \mathbf{u} = \nabla \cdot M \left( \nabla \phi - \frac{\nabla \phi}{|\nabla \phi|} \frac{[1 - 4(\phi - \phi_0)^2]}{\gamma} \right)$$

## Double distribution function approach



'Advection - Diffusion' of  $\phi$  is solved on a separate D2Q9 lattice

1 Initialize  $h_i^{in}(x, t)$

2 Compute  $\phi = \sum_{i=0}^9 h_i^{in}(x, t)$

3 Compute  $h_i^{eq}(x, t) = \phi w_i \left[ 1 + \frac{e_i u}{c_s^2 e^2} + \frac{(e_i u)^2}{2c_s^2 e^4} - \frac{u^2}{2c_s^2 e^2} \right]$

4 Collision  $h_i^{out}(x, t) = h_i^{in}(x, t) - \frac{1}{\tau_\phi} \left[ h_i^{in}(x, t) - h_i^{eq}(x, t) \right] + F_i^\phi(x, t)$

5 Streaming  $h_i^{in}(x + e_i, t + 1) = h_i^{out}(x, t)$

## Sample simulations

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Figure 3:  $Re = 1000$ ,  $Pr = 0.71$ ,  $D = 128$  [lu]

Case-ID	Lattice Size	Velocity set	Blockage Ratio	D	U	Pr	Re	$\nu$	$k$
Pr10 <sub>small</sub>	1000x150x3	D3Q27Q27	1/5	30	0.01	10	10	3E-02	3E-03
Pr10 <sub>medium</sub>	2000x300x3	D3Q27Q27	1/5	60	0.005	10	10	3E-02	3E-03
Pr10 <sub>large</sub>	4000x600x3	D3Q27Q27	1/5	120	0.0025	10	10	3E-02	3E-03
Pr100 <sub>small</sub>	1000x150x3	D3Q27Q27	1/5	30	0.01	100	10	3E-02	3E-04
Pr100 <sub>medium</sub>	2000x300x3	D3Q27Q27	1/5	60	0.005	100	10	3E-02	3E-04
Pr100 <sub>large</sub>	4000x600x3	D3Q27Q27	1/5	120	0.0025	100	10	3E-02	3E-04
Pr1000 <sub>small</sub>	1000x150x3	D3Q27Q27	1/5	30	0.01	1000	10	3E-02	3E-05
Pr1000 <sub>medium</sub>	2000x300x3	D3Q27Q27	1/5	60	0.005	1000	10	3E-02	3E-05
Pr1000 <sub>large</sub>	4000x600x3	D3Q27Q27	1/5	120	0.0025	1000	10	3E-02	3E-05

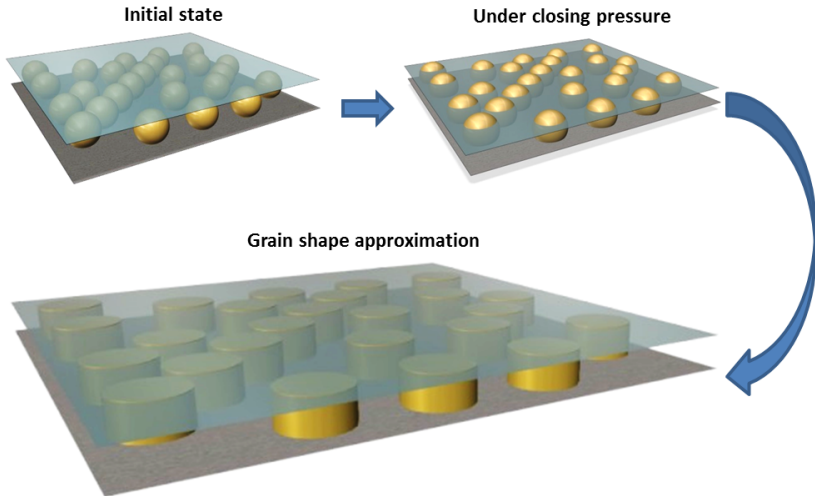
**Table 1:** Case-ID: lookup table.

Case-ID	$Nu_{CM\ HIGHER}^{1st\ order\ bc}$	$Nu_{CM\ HIGHER}^{2nd\ order\ bc}$	$Nu_{Cumulants}^{1st\ order\ bc}$	$Nu_{Cumulants}^{2nd\ order\ bc}$	$Nu_{FEM}$
Pr10 <sub>small</sub>	4.91	4.81	5.04	4.83	4.82
Pr10 <sub>medium</sub>	4.86	4.81	4.92	4.82	4.82
Pr10 <sub>large</sub>	4.83	4.81	4.87	4.81	4.82
Pr100 <sub>small</sub>	10.66	10.27	14.50	11.52	10.1
Pr100 <sub>medium</sub>	10.32	10.13	11.82	10.36	10.1
Pr100 <sub>large</sub>	10.19	10.08	10.83	10.13	10.1
Pr1000 <sub>small</sub>	27.09	24.58	94.31	58.33	21.43
Pr1000 <sub>medium</sub>	22.73	21.84	53.97	34.19	21.43
Pr1000 <sub>large</sub>	21.84	21.37	34.04	24.56	21.43

**Table 2:** Influence of kernel and BC on Nu number.

1st order bc = BB (hydrodynamics) + EQ (thermodynamics)

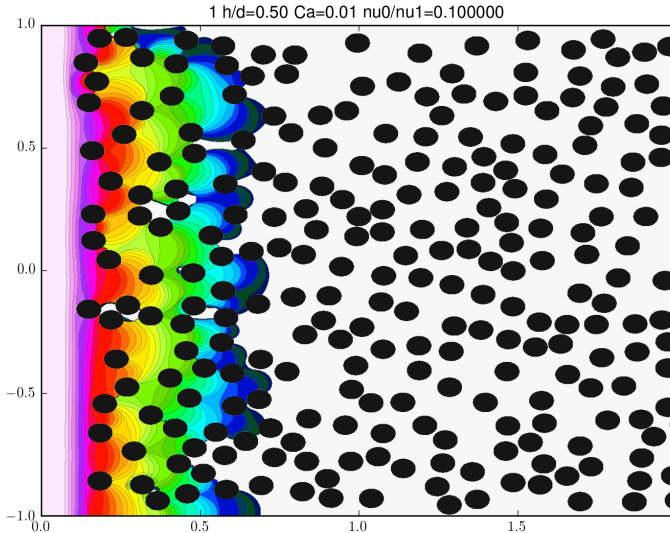
2nd order bc = IBB (hydrodynamics) + IABB (thermodynamics)



Movie:

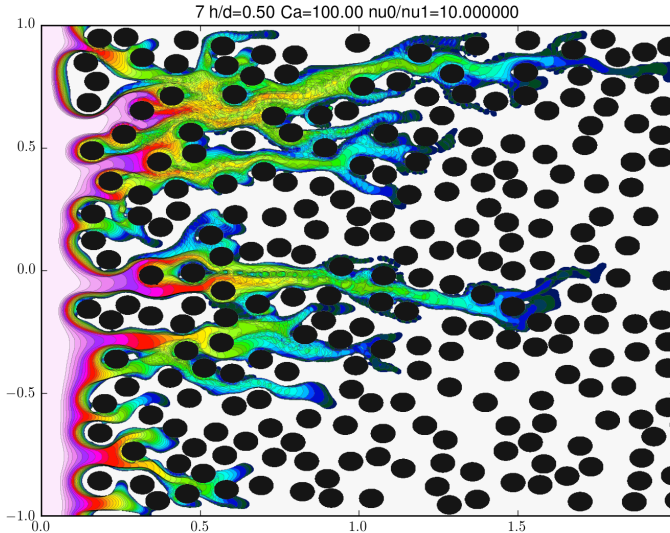
[https://www.youtube.com/watch?v=JoMFy2M\\_RDI](https://www.youtube.com/watch?v=JoMFy2M_RDI)

And multiphase flow (strong Capillary intraction,  $>10\text{mln}$  LBM cells):





And multiphase flow (weak Capillary intraction,  $>10\text{mln}$  LBM cells):



Small system - detailed mesh

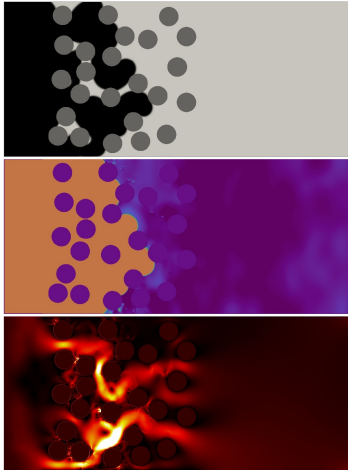


Figure 4: Viscous invasion.

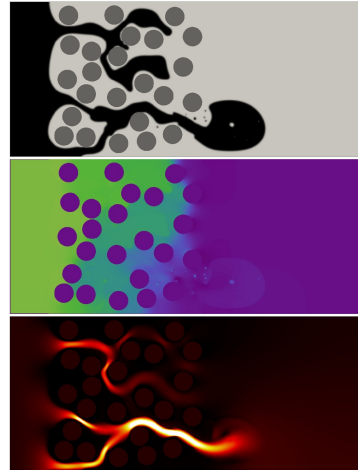
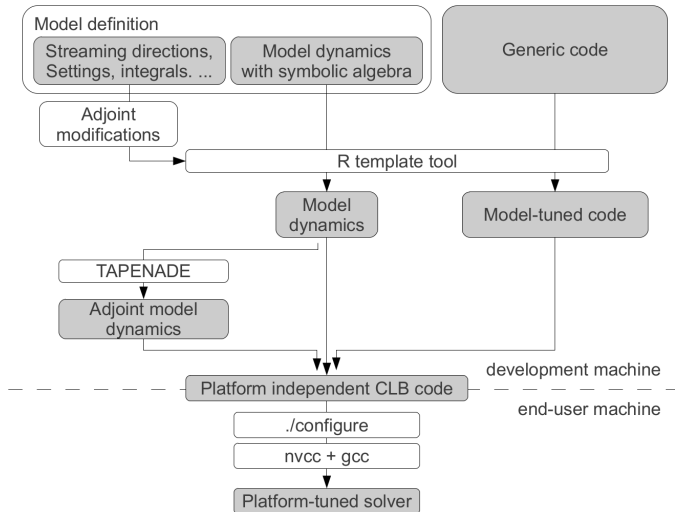


Figure 5: Capillary invasion.

## Solver Structure

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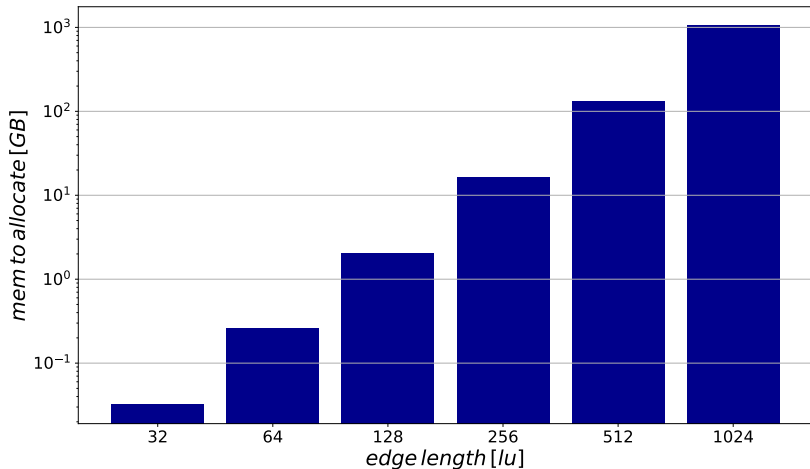


## Memory requirements

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	No [-]	[double]	[Bytes]
DF	2x27	54	432
DF temp	2x27	54	432
q	2x27	13.5	108
flag	1	0.5	4
memory per node		122	976

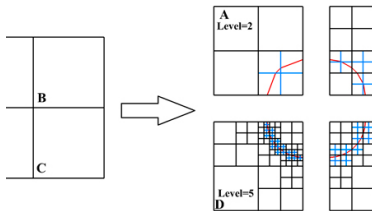
**Table 3:** Theoretical memory requirements for:  
3D, DDF model with interpolated BC.



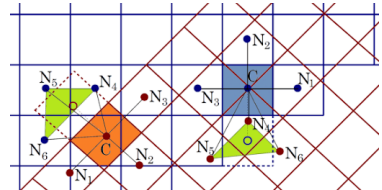
**Figure 6:** Theoretical memory requirements for a 3D cube domain, D3Q27Q27 lattice with interpolated BC.

**Q:** Does the high memory requirement limits applicability of LBM?

Not necessary. Similar problems occur in other CFD methods like VOF, FEM, FD.



**Figure 7:** Mesh refinement.



**Figure 8:** Overlapping mesh.



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<https://github.com/CFD-GO/TCLB>