

Small post processor

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2. Clear complete workspace

For new Matlab projects best practise is to clear the complete workspace and the command window. It is also a good habit to close all remaining figures since Matlab does not automatically open a new window each time a plot call is invoked.

```
1 path('./functions',path) % add functions directory the Matlab path
2 [datadir,flag]=ClearWs() ;
```

The above mentioned clears are performed in the function **ClearWs**. In addition some basic parameters like the name of the data directory or the dimensionality of the problem are also defined.

```
1 function [datadir,flag]=ClearWs()
2     close all
3     clear all
4     clc
5
```

```

6     datadir='data'; % specify the directory containing the data
7     flag='3D'; % specify the subdirectory of data
8 end

```

3. Read data files

During the evaluation of the function **ReadData** all data files necessary for the calculation of the spectrum and the correlation coefficients are read, namely the velocity components. In addition the import operations are enclosed in a **tic;...;toc** block measuring the time needed for reading the ASCII data. What you should get from the tic/toc block is that most of the time is spent during data I/O (Input/Output operations), nearly 220s. The actual computation needs only about 8s. What you can easily calculate from this is that the computation of the spectrum is nearly 27 times faster than the data import. Why the computation of Fourier transforms is that fast we will come to that later. Although the ASCII data format is not the preferred choice in terms of speed and size, we will use it since other methodologies require additional knowledge of data processing. Just for your information a very famous and highly portable data format is [hdf5](#). It is a software library that runs on a range of computational platforms, from laptops to massively parallel systems and implements a high-level API (Application programming interface) with C, C++, Fortran 90, and Java interfaces. Besides its hierarchical structure it is highly optimized for parallel I/O operations and can be read by nearly all data processing tools.

```

1 [uvel,vvel,wvel,time_read] = ReadData(datadir,flag);

1 function [uvel,vvel,wvel,time] = ReadData(datadir,flag)
2     tic; % enable timer
3     uvel=importdata([datadir,'/',flag,'/uvel']);
4     vvel=importdata([datadir,'/',flag,'/vvel']);
5     wvel=importdata([datadir,'/',flag,'/wvel']);
6     time = toc; % end timer
7 end

```

4. Set necessary parameters

For further computations it is important to define some parameters of the DNS simulation such as

- Number of grid points in on direction n_p ,
- Physical length of one direction L_x ,
- Physical grid spacing Δx ,
- Kinematic viscosity ν .

```
1 [u,v,w,dim,Lx,dx,nu]=Params(uvel,vvel,wvel);

1 function [u,v,w,dim,Lx,dx,nu]=Params(uvel,vvel,wvel)
2     dim=256; % number of points in one dimension
3     Lx=5e-3; % domain size
4     Ly=Lx;
5     Lz=Lx;
6     dx=Lx/dim; % grid spacing
7     dy=dx;
8     dz=dx;
9     nu=1.7e-5; % viscosity
10    u=reshape(uvel,dim,dim,dim); % reshape arrays to have them in 3D
11    v=reshape(vvel,dim,dim,dim);
12    w=reshape(wvel,dim,dim,dim);
13    clear uvel vvel wvel
14 end
```

5. Compute 3D spectrum

The core of the provided code is contained in the function **PowerSpec**. It computes the three dimensional energy spectrum from the given velocity fields, obtained from a direct numerical simulation. Although the theoretical analysis is relatively demanding compared to one dimensional spectra its worth investing the effort. The theory of one dimensional spectra relies on the assumption that the propagation of spectral waves (κ_1) is in the direction of

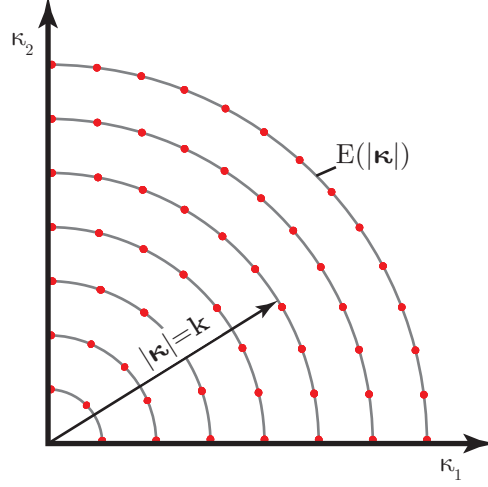


Fig. 1: Illustration of the two dimensional shell integration

the observed velocity fields or to say it differently one dimensional spectra and correlation functions are Fourier transform pairs. The theory of correlation functions will be discussed in a later section. A key drawback of this theory is that the calculated spectrum has contributions from all wavenumbers $\boldsymbol{\kappa}$, so that the magnitude of $\boldsymbol{\kappa}$ can be appreciably larger than κ_1 . This phenomenon is called aliasing.

In order to avoid aliasing effects usually connected with a one dimensional spectrum it is also possible to produce correlations that involve all possible directions. The three dimensional Fourier transformation of such a correlation produces a spectrum that not only depends on a single wavenumber but on the wavenumber vector κ_i . Though the directional information contained in κ_i eliminates the aliasing problem the complexity makes a physical reasoning impossible. For homogeneous isotropic turbulence the situation can be simplified by integrating the three dimensional spectrum over spherical shells. The idea of this integration is illustrated in Fig. 1

$$E(\kappa) = \oint\!\!\!\oint E(\boldsymbol{\kappa}) dS(\boldsymbol{\kappa}) = \oint\!\!\!\oint \frac{1}{2} \Phi_{ii}(\boldsymbol{\kappa}) dS(\boldsymbol{\kappa}) \quad (1)$$

Since the surface of a sphere is completely determined by its radius the surface

integral can be solved analytically.

$$\oint\!\!\!\oint () dS(\kappa) = 4\pi\kappa^2 \cdot () \quad (2)$$

This leads to

$$E(|\kappa|) = \frac{1}{2} \Phi_{ii}(|\kappa|) \quad (3)$$

```
1 [spectrum,k,mu,mv,mw,time_spec] = PowerSpec(u,v,w,Lx);
```

The content of **PowerSpec** reads

```
1 function [spectrum,k,mu,mv,mw,time] = power_spec(u,v,w,L)
2 % nx = size(u,1);
3 % ny = size(u,2);
4 % nz = size(u,3);
5 tic;
6 NFFT = 2.^nextpow2(size(u)); % next power of 2 fitting the length of u
7 u_fft=fftn(u,NFFT);
8 v_fft=fftn(v,NFFT);
9 w_fft=fftn(w,NFFT);
10
11 % Calculate the number of unique points
12 NumUniquePts = ceil((NFFT(1)+1)/2);
13
14 % FFT is symmetric, throw away second half
15 u_fft = u_fft(1:NumUniquePts,1:NumUniquePts,1:NumUniquePts);
16 v_fft = v_fft(1:NumUniquePts,1:NumUniquePts,1:NumUniquePts);
17 w_fft = w_fft(1:NumUniquePts,1:NumUniquePts,1:NumUniquePts);
18
19 mu = abs(u_fft)/length(u)^3;
20 mv = abs(v_fft)/length(v)^3;
21 mw = abs(w_fft)/length(w)^3;
22
23 % Take the square of the magnitude of fft of x.
24 mu = mu.^2;
25 mv = mv.^2;
26 mw = mw.^2;
27
28 % Since we dropped half the FFT, we multiply mx by 2 to keep the same energy.
29 % The Nyquist component, if it exists, is unique and should not be multiplied by 2.
30
```

```

31 if rem(NFFT, 2) % odd nfft excludes Nyquist point
32     mu(2:end,2:end,2:end) = mu(2:end,2:end,2:end)*2;
33     mv(2:end,2:end,2:end) = mv(2:end,2:end,2:end)*2;
34     mw(2:end,2:end,2:end) = mw(2:end,2:end,2:end)*2;
35 else
36     mu(2:end -1,2:end -1,2:end -1) = mu(2:end -1,2:end -1,2:end -1)*2;
37     mv(2:end -1,2:end -1,2:end -1) = mv(2:end -1,2:end -1,2:end -1)*2;
38     mw(2:end -1,2:end -1,2:end -1) = mw(2:end -1,2:end -1,2:end -1)*2;
39 end
40 % Compute the radius vector along which the energies are sumed
41 mx=NumUniquePts;
42 my=NumUniquePts;
43 mz=NumUniquePts;
44
45 dx=pi/L;
46 dy=pi/L;
47 dz=pi/L;
48 for I=1:mx
49     X0(I) = (I) *dx;
50 end
51
52 for J=1:my
53     Y0(J) = (J) *dy;
54 end
55
56 for K=1:mz
57     Z0(K) = (K) *dz;
58 end
59
60 for I=1:mx
61     for J=1:my
62         for K=1:mz
63             R(I, J, K) = sqrt (X0(I) *X0(I) +Y0(J) *Y0(J) +Z0(K) *Z0(K) ) ;
64         end
65     end
66 end
67
68 % P=mod(nx, 2) ;
69 % if (P < 1)
70 %     Nmax=mx-0.5;
71 % else
72 %     Nmax=mx;
73 % end
74 spectrum=zeros(Nmax,1);
75 for N=1:Nmax

```

```

76     Radius1=sqrt(3)*(N-1)*dx; %lower radius bound
77     Radius2=sqrt(3)*N*dx; %upper radius bound
78     % build logical index for selecting values lying on the shell
79     logical = (Radius1 <= R(:, :, :)) & (R(:, :, :) < Radius2);
80     % build summation over shell components
81     T.EVP1=sum(mu(logical))+sum(mv(logical))+sum(mw(logical));
82     % put them at position N in the spectrum
83     spectrum(N)=T.EVP1./2;
84 end
85 k=[1:Nmax].*dx;
86 spectrum = 1./(2*pi)^3.*spectrum;
87 time=toc;
88 end

```

6. Compute dissipation and turbulent kinetic energy

```

1 [Dissipation, kin_E_Sp, kin_E_Ph, up] = SpecProp(spectrum, k, nu, u, v, w);
2 kin_E_Ph
3 kin_E_Sp

1 kin_E_Ph =
2
3     6.8206
4
5
6 kin_E_Sp =
7
8     5.3630

```

The content of SpecProp reads

```

1 function [Dissipation, kin_E_Sp, kin_E_Ph, up] = SpecProp(E, k, nu, u, v, w)
2     kin_E_Sp = trapz(k, E);
3     Dissipation = trapz(k, 2*nu.*k.^2.*E');
4     up = sqrt(1/3*(u.^2+v.^2+w.^2));
5     kin_E_Ph = sum(sum(sum(3/2*up.^2)))/size(up,1)^3;
6 end

```

7. Kolmogorov properties

```

1 [eta,u_eta,tau]=KolmoScale(nu,Dissipation);
2 eta
3 u_eta
4 tau

```

```

1 eta =
2
3     4.5999e-05
4
5
6 u_eta =
7
8     0.3696
9
10
11 tau =
12
13     1.2446e-04

```

The content of `KolmoScale` reads

```

1 function [eta,u_eta,tau]=KolmoScale(nu,Dissipation)
2     eta = (nu^3/Dissipation)^(1/4);
3     u_eta = (nu*Dissipation)^(1/4);
4     tau = (nu/Dissipation)^(1/2);
5 end

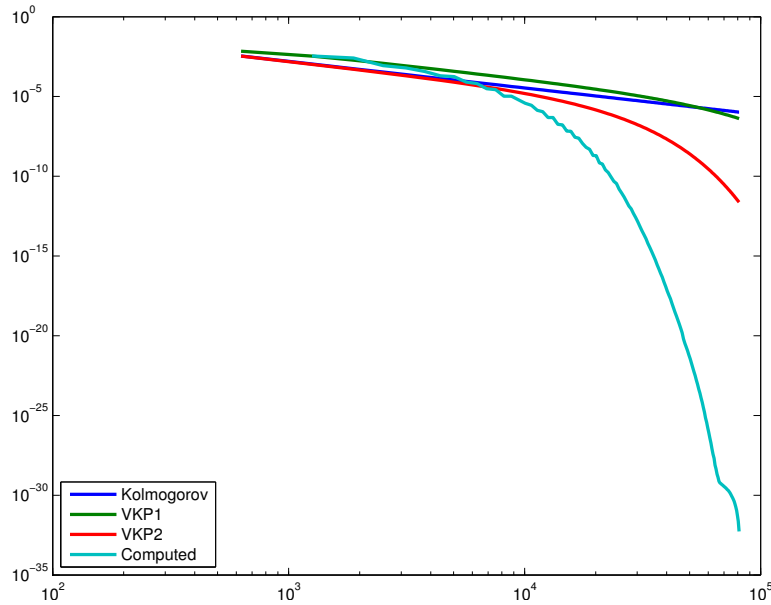
```

8. Compute model spectra

```

1 PlotModelSpec(k,spectrum,Dissipation,up,Lx,eta,nu);

```

The content of PlotModelSpec reads

```

1 function PlotModelSpec(k,spectrum,Dissipation,up,Lx,eta,nu)
2     % Von Karman-Pao Spektren
3     close all
4     kd = k(end);
5     ke = pi/Lx/2;
6     A = 1.5;
7     up = mean2(up);
8
9     VKP1 = A*up^5/Dissipation.*(k./ke).^4./(1+(k./ke).^2).^(17/6).* ...
10         exp(-3/2*A.*(k./kd).^(4/3));
11
12     kd = 1./eta;
13     VKP2 = 1.5*(k./kd).^(-5/3)./(Dissipation*nu^5)^(-1/4).* ...
14         exp(-1.5*1.5.*(k./kd).^(4/3));
15
16     % Kolmogorov Spektrum
17     Kolmo=1.5*Dissipation^(2/3)*(k.^(-5/3));
18
19     % Plot spectra
20     h=loglog(k,Kolmo,k,VKP1,k,VKP2,k,spectrum);
21     set(h,'LineWidth',2);
22
23     h=legend('Kolmogorov','VKP1','VKP2','Computed');
```

```

24         set (h, 'Location', 'SouthWest')
25     end

```

9. Compute correlations

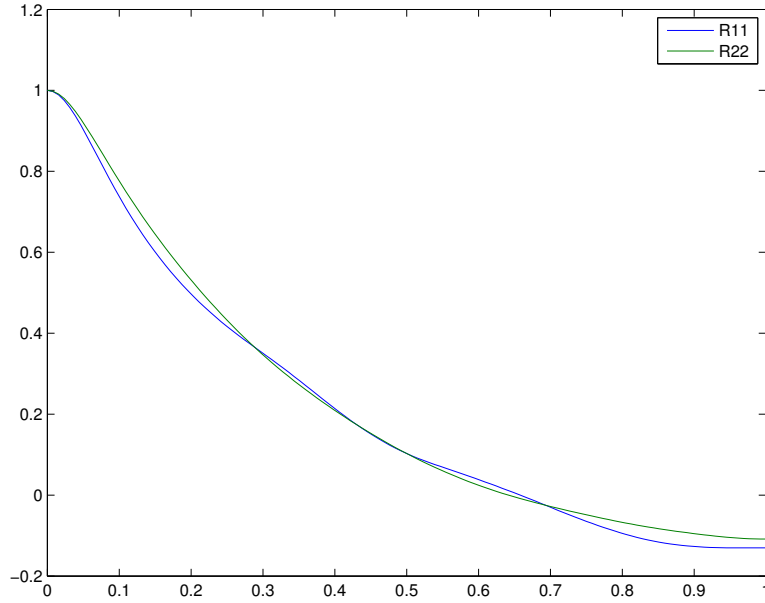
Computing a correlation can be a tedious work (requiring tremendous effort) especially if you have large data sets. From theory it is well known that the multiplication of the transform of a data set and its complex conjugate are an accurate representation of the correlation function. Using the FFT approach this gives an enormous speed advantage. Since we already computed the velocity correlation tensor we may use this result in order to compute the correlation tensor.

$$R_{ij} = \frac{cov(U_i, U_j)}{\sqrt{\sigma_i^2 \sigma_j^2}} = \frac{left < u'_i u'_j right >}{\sqrt{\sigma_i^2 \sigma_j^2}} \quad (4)$$

```

1  [R11,R22,r,R1,R2,R3]=Correlation(u,v,w,Lx);
2  close all
3  figure
4  plot(r,R11,r,R22)
5  legend('R11','R22')

```



The content of Correlation reads

```

1 function [R11,R22,r,R1,R2,R3]=Correlation(u,v,w,Lx)
2     scaling = 1;
3     NFFT = 2.nextpow2(size(u)); % next power of 2 fitting the length of u
4     u_fft=fftn(u,NFFT)./scaling; %2 pi → definition of FFT
5     %
6     NFFT = 2.nextpow2(size(v));
7     v_fft=fftn(v,NFFT)./scaling;
8     %
9     NFFT = 2.nextpow2(size(w));
10    w_fft=fftn(w,NFFT)./scaling;
11
12    Rij_x=(u_fft.*conj(u_fft)); % compute velo. correlation tensor
13    Rij_y=(v_fft.*conj(v_fft));
14    Rij_z=(w_fft.*conj(w_fft));
15
16    % x-component
17    NFFT = 2.nextpow2(size(u_fft));
18    R1=ifftn(Rij_x,NFFT)/std2(u)^2./prod(NFFT);
19    % y-component
20    NFFT = 2.nextpow2(size(v_fft));
21    R2=ifftn(Rij_y,NFFT)/std2(v)^2./prod(NFFT);
22    % z-component
23    NFFT = 2.nextpow2(size(w_fft));

```

```

24     R3=ifftn(Rij_z,NFFT)/std2(w)^2./prod(NFFT);
25
26     R11 = (reshape(R3(1,1,:),NFFT(1),1)+R2(1,:,1)'+R1(:,1,1))/3;
27     R11 = R11(1:size(u_ffft)/2+1);
28     %
29     R1_22 = (R1(1,:,1)+R3(1,:,1))/2;
30     R2_22 = (R2(:,1,1)+R3(:,1,1))/2;
31     R3_22 = (reshape(R1(1,1,:),size(u_ffft,1),1)+...
32             reshape(R2(1,1,:),size(u_ffft,1),1))/2;
33
34     R22 = (R1_22'+R2_22+R3_22)/3;
35     R22 = R22(1:size(u_ffft)/2+1);
36
37     r = linspace(0,Lx/2,size(u_ffft,1)/2+1)/(Lx/2);
38 end

```