

Small post processor

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1. Clear complete workspace

For new Matlab projects best practise is to clear the complete workspace and the command window. It is also a good habit to close all remaining figures since Matlab does not automatically open a new window each time a plot call is invoked.

```
1 display('Clear workspace ...')
2 path('./functions',path) % add functions directory the
   Matlab path
3 close all % close all figures
4 clear all % clear workspace
5 clc % clear command window
6
7 datadir='data/stuttgart/'; % specify directory containg
   input data
8 flag='3D'; % specify dimensionality
9 %set(0,'DefaultFigureWindowStyle','docked')
```

2. Read data files

During the evaluation of the function `ReadData` all velocity files necessary for the calculation of the spectrum and the correlation coefficients are read. In addition the import operations are enclosed in a `tic;...;toc` block which measures the time needed for reading the ASCII data. What you should get from the tic/toc block is that most of the time is spent during data I/O (Input/Output operations). The actual computation needs much less time. Although the ASCII data format is not the preferred choice in terms of speed and size, we will use it since other methodologies require additional knowledge of data processing. Just for your information a very famous and highly portable data format is [hdf5](#). It is a software library that runs on a range of computational platforms, from laptops to massively parallel systems and implements a high-level API (Application programming interface) with C, C++, Fortran 90, and Java interfaces. Besides its hierarchical structure it is highly optimized for parallel I/O operations and can be read by nearly all data processing tools.

```
1 display('Read data ...')
2 [uvel,vvel,wvel,time_read,dim] = ReadData(datadir,flag,'
    uvel','vvel','wvel');
3 % test=importdata('data/3D/CFX_velocity_field.dat');
4 % uvel=reshape(test(:,1),33,33,33);
5 % vvel=reshape(test(:,2),33,33,33);
6 % wvel=reshape(test(:,3),33,33,33);

1 function [uvel,vvel,wvel,time,dim] = ReadData(datadir,flag
    ,...
2                                     u_name,...
3                                     v_name,...
4                                     w_name)
5     tic; % enable timer
6     uvel=importdata([datadir,'/',u_name]);
7     vvel=importdata([datadir,'/',v_name]);
8     wvel=importdata([datadir,'/',w_name]);
9     time = toc; % end timer
10    if strcmp(flag,'3D')
11        dim=round((size(uvel,1))^(1/3));
12    end
13 end
```

3. Set necessary parameters

For further computations it is important to define some parameters of the DNS simulation such as

- Number of grid points in one direction n_p ,
- Physical length of one edge of the domain L_x ,
- Physical grid spacing Δx ,
- Kinematic viscosity ν .

```
1 display('Set parameters ...')
2 [u,v,w,Lx,dx,nu]=Params(uvel,vvel,wvel,dim);

1 function [u,v,w,Lx,dx,nu]=Params(uvel,vvel,wvel,dim)
2 %     Lx=0.1;
3     Lx=2*pi;
4 %     dim=33;
5 %     Lx=3.2e-2; % domain size
6     Ly=Lx;
7     Lz=Lx;
8     dx=Lx/(dim-1); % grid spacing
9     dy=dx;
10    dz=dx;
11    nu=1.7e-5; % viscosity
12    u=reshape(uvel,dim,dim,dim); % reshape arrays to have
    them in 3D
13    v=reshape(vvel,dim,dim,dim);
14    w=reshape(wvel,dim,dim,dim);
15    clear uvel vvel wvel
16 end
```

4. Compute 3D spectrum

The core of the code is contained in the function **PowerSpec**. It computes the three dimensional energy spectrum from the given velocity fields, obtained from a direct numerical simulation. Although the theoretical analysis is

relatively demanding compared to one dimensional spectra its worth investing the effort. The theory of one dimensional spectra relies on the assumption that the propagation of spectral waves (κ_1) is in the direction of the observed velocity fields or to say it differently one dimensional spectra and correlation functions are Fourier transform pairs. The theory of correlation functions will be discussed in section 7. A key drawback of this theory is that the calculated spectrum has contributions from all wavenumbers $\boldsymbol{\kappa}$, so that the magnitude of $\boldsymbol{\kappa}$ can be appreciably larger than κ_1 . This phenomenon is called aliasing. In order to avoid these aliasing effects is also possible to produce correlations that involve all possible directions. The three dimensional Fourier transformation of such a correlation produces a spectrum that not only depends on a single wavenumber but on the wavenumber vector κ_i . Though the directional information contained in κ_i eliminates the aliasing problem the complexity makes a physical reasoning impossible. For homogeneous isotropic turbulence the situation can be considerably simplified. From the knowledge that the velocity field is isotropic it can be shown that the velocity spectrum tensor is fully determined by

$$\Phi_{ij}(\boldsymbol{\kappa}) = A(\kappa)\delta_{ij} + B(\kappa)\kappa_i\kappa_j, \quad (1)$$

where $A(\kappa)$ and $B(\kappa)$ are arbitrary scalar functions. Since we assume incompressible fluids (mathematically expressed by $\nabla \cdot u = 0$ or $\kappa_i u_i = 0$ the following condition holds

$$\kappa_i \Phi_{ij}(\boldsymbol{\kappa}) = 0. \quad (2)$$

It can be shown that this yields a relation between A and B by means of

$$B(\kappa) = -\frac{A(\kappa)}{\kappa^2} \quad (3)$$

In the end this gives a relation between the three dimensional energy spectrum function $E(|\boldsymbol{\kappa}|)$ and the velocity spectrum tensor Φ_{ij} .

$$\Phi_{ij} = \frac{E(|\boldsymbol{\kappa}|)}{4\pi(|\boldsymbol{\kappa}|)^2} \left(\delta_{ij} - \frac{\kappa_i\kappa_j}{(|\boldsymbol{\kappa}|)^2} \right) \quad (4)$$

The question is now how the remaining variable (A or B) can be determined. Regarding the turbulent kinetic energy we know that

$$k = \int_{-\infty}^{\infty} E(|\boldsymbol{\kappa}|) dk = \sum_{\boldsymbol{\kappa}} E(\boldsymbol{\kappa}) = \sum_{\boldsymbol{\kappa}} \frac{1}{2} \langle u^*(\boldsymbol{\kappa}) u(\boldsymbol{\kappa}) \rangle = \iiint_{-\infty}^{\infty} \frac{1}{2} \Phi_{ii}(\boldsymbol{\kappa}) d\boldsymbol{\kappa}. \quad (5)$$

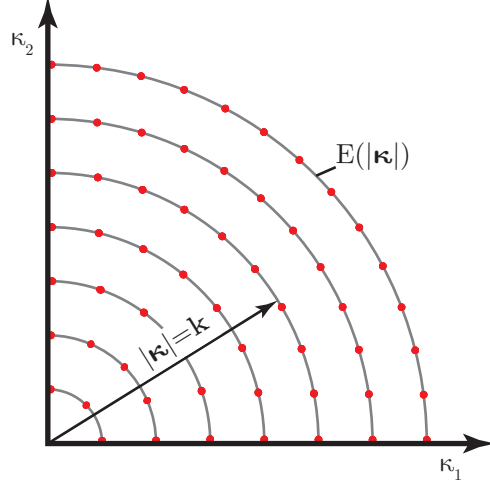


Fig. 1: Illustration of the two dimensional shell integration

Comparing the second and last expression we get

$$E(|\kappa|) = \oint \frac{1}{2} \Phi_{ii}(\kappa) dS(\kappa). \quad (6)$$

This integral can be solved analytically by utilizing again the assumption of isotropy. For these kind of flows the energy spectrum function can be regarded as the sum of kinetic energy (in wave number space) on different energy levels. Each of these energy levels is denoted by a spherical shell in wave number space. Since the surface of a sphere is completely determined by its radius the surface integral can be solved analytically. The idea of this integration is illustrated in Fig. ???. As a result of this one gets

$$E(|\kappa|) = \oint \frac{1}{2} \Phi_{ii}(\kappa) dS(\kappa) = 4\pi(|\kappa|)^2 \Phi_{ii}(|\kappa|). \quad (7)$$

Introducing this relation to equations (1) and (3) one arrives at an expression for the variable B .

$$B = -\frac{E(|\kappa|)}{4\pi(|\kappa|)^2} \quad (8)$$

Together with the approximation of the integral of Φ (equation (5))

$$\iiint_{-\infty}^{\infty} \frac{1}{2} \Phi_{ii}(\kappa) d\kappa \approx \frac{1}{2} \sum_{\kappa} \Phi_{ii}(\kappa) (\Delta\kappa)^3, \quad (9)$$

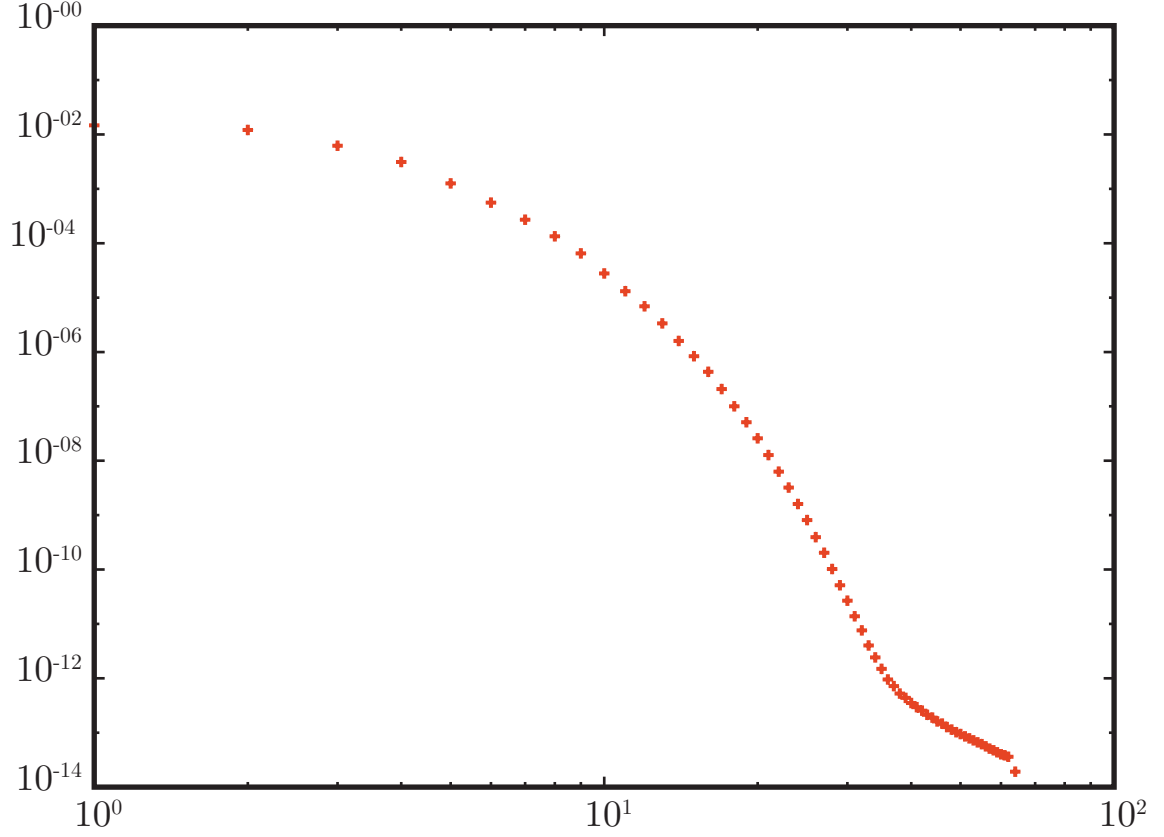


Fig. 2: Computed spectrum

where $\Delta\kappa$ refers to the step size in wave number space, the final expression of the three dimensional discrete energy spectrum can be derived.

$$E(|\boldsymbol{\kappa}|) = 2\pi(|\boldsymbol{\kappa}|)^2 \frac{\langle u^*(\boldsymbol{\kappa}) u(\boldsymbol{\kappa}) \rangle}{(\Delta\kappa)^3} \quad (10)$$

The calling sequence for the computation of the energy spectrum reads

```

1 display('Compute spectrum...')
2 [spectrum,k,bin_counter,time_spec] = PowerSpec(u,v,w,Lx,
    dim);
```

```

1  function [spectrum,k,bin_counter,time] = PowerSpec(u,v,w,
    ...
2                                     L,dim)
3      tic;
4      uu_fft=fftn(u);
5      vv_fft=fftn(v);
6      ww_fft=fftn(w);
7
8      muu = abs(uu_fft)/length(u)^3;
9      mvv = abs(vv_fft)/length(v)^3;
10     mww = abs(ww_fft)/length(w)^3;
11
12     muu = muu.^2;
13     mvv = mvv.^2;
14     mww = mww.^2;
15
16     % %      rx=[0:1:dim-1] - (dim-1)/2;
17     % %      ry=[0:1:dim-1] - (dim-1)/2;
18     % %      rz=[0:1:dim-1] - (dim-1)/2;
19     % %
20     % %      test_x=circshift(rx',[(dim+1)/2 1]);
21     % %      test_y=circshift(ry',[(dim+1)/2 1]);
22     % %      test_z=circshift(rz',[(dim+1)/2 1]);
23     rx=[0:1:dim-1] - (dim)/2+1;
24     ry=[0:1:dim-1] - (dim)/2+1;
25     rz=[0:1:dim-1] - (dim)/2+1;
26
27     test_x=circshift(rx',[(dim)/2+1 1]);
28     test_y=circshift(ry',[(dim)/2+1 1]);
29     test_z=circshift(rz',[(dim)/2+1 1]);
30
31     [X,Y,Z]= meshgrid(test_x,test_y,test_z);
32     r=(sqrt(X.^2+Y.^2+Z.^2));
33
34     dx=2*pi/L;
35     k=[1:(dim)/2].*dx;
36
37     spectrum=zeros(size(k,2),1);
38     bin_counter=zeros(size(k,2),1);
39     for N=2:(dim-1)/2-1
40         picker = (r(:, :, :)*dx <= (k(N+1) + k(N))/2) & ...
41                 (r(:, :, :)*dx > (k(N) + k(N-1))/2);
42         spectrum(N) = sum(muu(picker))+...
43                     sum(mvv(picker))+...
44                     sum(mww(picker));

```

```

45         bin_counter(N) = size(find(picker==1),1);
46     end
47     % compute first value of spectrum
48     picker = (r(:, :, :)*dx <= (k(2) + k(1))/2);
49     spectrum(1) = sum(muu(picker))+...
50                 sum(mvv(picker))+...
51                 sum(mww(picker));
52     bin_counter(1) = size(find(picker==1),1);
53     % compute last value of spectrum
54     picker = (r(:, :, :)*dx > (k(end) + k(end-1))/2 & ...
55             r(:, :, :)*dx <= k(end));
56     spectrum(end) = sum(muu(picker))+...
57                   sum(mvv(picker))+...
58                   sum(mww(picker));
59     bin_counter(end) = size(find(picker==1),1);
60     %compute final spectrum
61     spectrum = spectrum*2*pi.*k'.^2./(bin_counter.*dx.^3);
62     time=toc;
63
64     y = [k; spectrum'];
65     fid = fopen('spectrum.dat', 'w');
66     fprintf(fid, '%10.2E %10.2E\n', y);
67     fclose(fid);
68 end

```

5. Compute dissipation and turbulent kinetic energy

The function **SpecProp** calculates the kinetic energy both from the velocities and the previously computed spectrum. The latter one is calculated by

$$k = \int_{-\infty}^{\infty} E(|\boldsymbol{\kappa}|) d|\boldsymbol{\kappa}| \quad (11)$$

A second integral, also evaluated in this routine, gives the value of the Dissipation

$$\epsilon = 2 \int_{-\infty}^{\infty} \nu(|\boldsymbol{\kappa}|)^2 E(|\boldsymbol{\kappa}|) d|\boldsymbol{\kappa}|, \quad (12)$$

where ν refers to the kinematic viscosity. The calling sequence reads


```

1 display('Compute kinetic energy...')
2 [Dissipation,kin_E_Sp,kin_E_Ph,up] = SpecProp(spectrum,k,
    ...
3                                     nu,u,v,w,dim);

1 function [Dissipation,kin_E_Sp,kin_E_Ph,up] = ...
2                                     SpecProp(E,k,nu,u,v,w,dim)
3     kin_E_Sp = trapz(k,E);
4     Dissipation = trapz(k,2*nu.*k.^2.*E');
5     up = sqrt(1/3/dim^3*sum(sum(sum(u.^2+v.^2+w.^2))));
6     kin_E_Ph = 3/2*up^2;
7 end

```

6. Kolmogorov properties

According to the Kolmogorov hypotheses the length scale η , characteristic velocity u_η and the characteristic time scale τ of the smallest swirls in the flow are computed within the function **KolmoScale**. From a dimensionality analysis Kolmogorov derived

$$\eta = \left(\frac{\nu^3}{\epsilon} \right)^{1/4}, \quad (13)$$

$$u_\eta = (\epsilon \nu)^{1/4}, \quad (14)$$

$$\tau_\eta = \left(\frac{\nu}{\epsilon} \right)^{1/2}. \quad (15)$$

For further reading concerning his theory it is referred to Pope [1], Hinze [2] and Tennekes and Lumley [3].

```

1 display('Compute Kolmogorov scales...')
2 [eta,u_eta,tau]=KolmoScale(nu,Dissipation);

1 function [eta,u_eta,tau]=KolmoScale(nu,Dissipation)
2     eta = (nu^3/Dissipation)^(1/4);
3     u_eta = (nu*Dissipation)^(1/4);
4     tau = (nu/Dissipation)^(1/2);
5 end

```

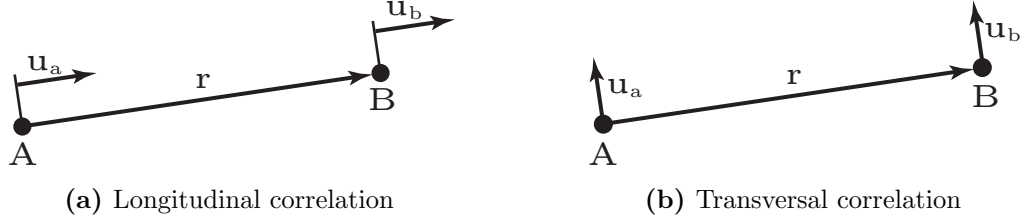


Fig. 3: Illustration of different correlation functions

7. Compute correlations

From a general perspective correlation functions are a measure of how much two physical quantities are connected. So how is this helpful for the analysis of turbulent flows? For seemingly chaotic and random processes it would be beneficial if we had a measure of how the velocity at point A is influenced by the velocity at point B . A maybe more intuitive quantity that can be calculated from the correlation functions is the integral lengthscale which gives a measure of the largest eddies in the flow. In fluid dynamics one generally differentiates between two forms of correlation functions, the longitudinal and the transversal or lateral correlation function. The difference between both forms is illustrated in figure 3. In general the correlation between two components of an isotropic homogeneous velocity field is expressed by

$$R_{ij} = \langle u_i(\mathbf{x} + \mathbf{r}) u_j(\mathbf{x}) \rangle. \quad (16)$$

The correlation coefficients are computed by normalizing R_{ij} with the square root of the product of the two variances σ_i^2 and σ_j^2 .

$$r_{ij} = \frac{\langle u_i(\mathbf{x} + \mathbf{r}) u_j(\mathbf{x}) \rangle}{\sqrt{\sigma_i^2 \sigma_j^2}} \quad (17)$$

As illustrated in figure 3 the longitudinal and lateral correlation depend on the direction of \mathbf{r} , i.e.

$$f(r) = \frac{\langle u_1(\mathbf{x} + r\mathbf{e}_1) u_1(\mathbf{x}) \rangle}{\sigma_1} \quad (18)$$

$$g(r) = \frac{\langle u_2(\mathbf{x} + r\mathbf{e}_1) u_2(\mathbf{x}) \rangle}{\sigma_2} \quad (19)$$

From several theoretical analysis it is well known that correlations can be efficiently computed by means of multiplying the Fourier transform of a quantity with its complex conjugate. Using the FFT approach this gives an enormous speed advantage.

$$R_{ij} = \mathfrak{F}^{-1} \{ \mathfrak{F} \{ u_i \}^* \cdot \mathfrak{F} \{ u_j \} \} \quad (20)$$

```
1 display('Compute Correlations...')
2 [R11,R22,r,R1,R2,R3]=Correlation(u,v,w,Lx,dim);
```

The content of **Correlation** reads

```
1 function [R11,R22,r,R1,R2,R3]=Correlation(u,v,w,Lx,dim)
2     scaling = 1;
3 %     NFFT = 2.^nextpow2(size(u)); %power of 2 fitting
4 %     length of u
5 %     u_fft=fftn(u,NFFT)./scaling;
6 %     NFFT = 2.^nextpow2(size(v));
7 %     v_fft=fftn(v,NFFT)./scaling;
8 %     NFFT = 2.^nextpow2(size(w));
9 %     w_fft=fftn(w,NFFT)./scaling;
10
11 u_fft=fftn(u)./scaling;
12 v_fft=fftn(v)./scaling;
13 w_fft=fftn(w)./scaling;
14
15 Rij_x=(u_fft.*conj(u_fft)); % compute velo.
16      correlation tensor
17 Rij_y=(v_fft.*conj(v_fft));
18 Rij_z=(w_fft.*conj(w_fft));
19
20 % % % x-component
21 % % NFFT = 2.^nextpow2(size(u_fft));
22 % % R1=ifftn(Rij_x,NFFT)/std2(u)^2/dim^3;
23
24 % % % y-component
25 % % NFFT = 2.^nextpow2(size(v_fft));
26 % % R2=ifftn(Rij_y,NFFT)/std2(v)^2/dim^3;
27
28 % % % z-component
29 % % NFFT = 2.^nextpow2(size(w_fft));
30 % % R3=ifftn(Rij_z,NFFT)/std2(w)^2/dim^3;
```

```

29     R1=ifftn(Ri_j_x)/std2(u)^2/dim^3;
30     R2=ifftn(Ri_j_y)/std2(v)^2./dim^3;
31     R3=ifftn(Ri_j_z)/std2(w)^2./dim^3;
32
33     NFFT=size(u_ffft,1);
34     R11 = (reshape(R3(1,1,:),NFFT(1),1)+R2(1,:,1)'+R1
35           (:,1,1))/3;
36     R11 = R11(1:size(u_ffft)/2+1);
37     %
38     R1_22 = (R1(1,:,1)+R3(1,:,1))/2;
39     R2_22 = (R2(:,1,1)+R3(:,1,1))/2;
40     R3_22 = (reshape(R1(1,1,:),size(u_ffft,1),1)+...
41             reshape(R2(1,1,:),size(u_ffft,1),1))/2;
42
43     R22 = (R1_22'+R2_22+R3_22)/3;
44     R22 = R22(1:size(u_ffft)/2+1);
45
46     r = linspace(0,Lx/2,size(u_ffft,1)/2+1)/(Lx/2);
47 end

```

References

- [1] S. B. Pope, Turbulent Flows, Cambridge University Press, 1 edition, 2000.
- [2] J. O. Hinze, Turbulence, Mcgraw-Hill College, 2 edition, 1975.
- [3] H. Tennekes, J. L. Lumley, A first course in turbulence, The MIT Press, 1972.