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# 1 Scientific, Technical, and Management Section

## 1.1 Observables and Baselines

We are proposing to study a probe-scale mission to extract the wealth of cosmological information contained in the spectrum and polarization of the cosmic microwave background (CMB). This information is unique and can not be provided through other astrophysical probes.

The starting points for our study are two space missions, EPIC-IM and Super-PIXIE [? ? ]. EPIC-IM is an imaging instrument that was presented to the 2010 decadal panel as a candidate CMB polarization space mission. It is based on a 2 m aperture telescope and 11,094 bolometric transition edge sensors. Table ?? gives the mission's frequency bands and sensitivities. PIXIE is a proposed Explorer scale mission focused on a measurement of the spectrum of the CMB. Super-PIXIE is a scaled up, more capable version of PIXIE. Improvement in technology by the next decade will enable designing a mission that is more capable compared to both EPIC-IM and Super-PIXIE. Therefore, all quantitative predictions presented here for EPIC-IM and Super-PIXIE represent *minimum* capabilities.

The best measurements of the CMB spectrum – made by COBE/FIRAS approximately 25 years ago – show that the average CMB spectrum is consistent with that of a blackbody to an accuracy of 4 parts in  $10^4$  [1, 2]. Distortions in this spectrum encode a wealth of new information. The distortion shapes are commonly denoted as  $\mu$ - and  $y$ -types. The  $\mu$ -distortion arises from energy release in the early Universe and can only be produced in the hot and dense environment present at redshifts  $z \gtrsim 5 \times 10^4$ . This makes  $\mu$ -distortions a unique messenger from a red-shift range that is not accessible to other probes. The  $y$  distortion is caused by energy exchange between the CMB photons and free electrons through inverse Compton scattering [3, 4]. These are caused at lower redshifts and are sensitive to the evolution of the large scale structure.

Thomson scattering at the surface of last scattering is the source of the polarization of the CMB. It is useful to decompose the polarization field to two modes that are independent over the full sky,  $E$  and  $B$ -modes. Together with the pattern of temperature anisotropy  $T$ , the CMB thus gives three auto- and three cross-spectra. The *Planck* satellite and larger aperture ground-based instruments measured the  $T$  spectrum to cosmic variance limit for  $\ell \leq ??$ . Much information remains encoded in the  $E$  and  $B$  spectra, whose full exploration has just begun [5, 6? ? ].

## 1.2 Science Objectives

### 1.2.1 The Primordial Universe and Cosmic Inflation

The observed temperature and E-mode polarization of the Cosmic Microwave Background require primordial inhomogeneities in the gravitational potential, providing a remarkable observational link to the dynamics of the universe before the hot big bang era. Inflation [7, 8, 9, 10, 11], a primordial era of accelerated expansion, provides a compelling dynamical origin for the observed statistical homogeneity of the primordial perturbations on even the largest scales. But, inflation also predicts an as yet unobserved spectrum of primordial gravitational waves sourced directly by quantum fluctuations of the tensor component of the metric. Any detection of B-mode polarization, whether generated by the primordial gravitational waves of inflation [12, 13] or by any other source of early time vector or tensor perturbations, would reveal completely new information about the primordial era. The results would provide significant constraints and consistency checks for current models or could perhaps even overturn them. A detection would have implications far beyond the difficult question of cosmological dynamics by providing evidence for a new scale of particle physics near the GUT scale. In the context of an inflation, the relationship is particularly clear:

the potential energy  $V$  of the inflaton is related to the tensor-to-scalar ratio  $r$  at the peak of the spectrum by  $V^{1/4} = 3.7 \times 10^{16} r^{1/4}$  GeV.

Figure 2 shows current data, as well as forecasts for the EPIC-IM satellite, for CMB temperature, E-mode polarization, and B-mode polarization from lensing. The large-scale polarization information accessible by a satellite mission is complementary to the smaller scale modes that can be well measured from the ground. The figure also shows the predictions for B-modes from vacuum fluctuations of the metric during an inflationary era for two values of the tensor-to-scalar ratio  $r$ . For testing inflation, the largest scales are particularly important: only a satellite can demonstrate the presence of B-mode correlations on scales that were super-horizon at the time of recombination [14]. In addition, new tests of isotropy on the largest scales can only come from the reionization bump in the E-mode polarization []. In its recent report New Worlds New Horizons (NWNH), the decadal survey committee strongly endorsed sub-orbital searches for the B-mode signal from inflation saying that “The convincing detection of B-mode polarization in the CMB produced in the epoch of reionization would represent a watershed discovery.” [15].

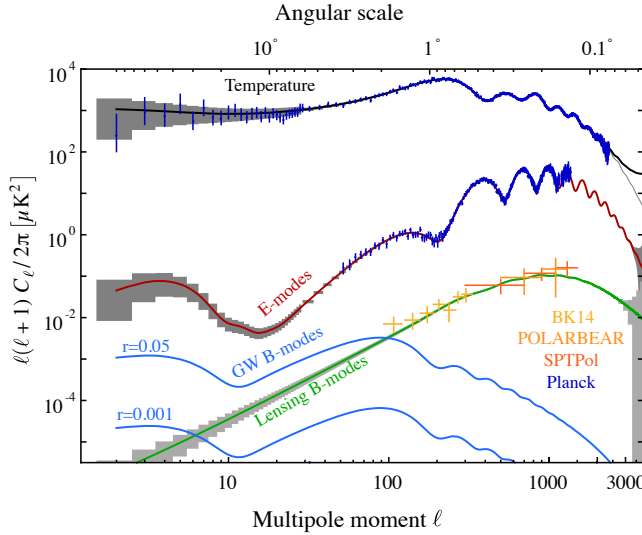


Figure 1: Predicted determination of the CMB power spectra for EPIC-IM (grey boxes) overlaid on theoretical predictions (solid lines) and including Planck measurements of the temperature and E modes (blue) and of several ground-based measurements of the lensing B-modes. The tensor B-mode predictions (blue) are shown for two representative values of the tensor-to-scalar ratio:  $r = 0.001$  and  $r = 0.05$ .

Since all inflation models predict primordial B-modes, the space of inflation models can be used to motivate a target sensitivity for future observations. For example, when the slow-roll criteria is met ( $\epsilon \equiv -\dot{H}/H^2 \ll 1$ ) there are just two observationally viable classes of models that naturally explain the value of the spectral index  $n_s$  (by requiring  $n_s(\mathcal{N}) - 1 \propto -\frac{1}{\mathcal{N}}$ , where  $\mathcal{N}$  is the number of e-folds between the scale  $n_s$  where is observed and end of inflation) [16, 17, 18]. One class is the set of monomial potentials ( $V(\phi) \propto \phi^p$ ), which contains many of the canonical inflation models (eg, a quadratic potential) and is already under significant observational pressure. If the error bars on the spectral index tighten by a factor of about 2, and the 95% C.L. upper limit on  $r$  is pushed to even  $\sim 0.01$ , all such models would be ruled out. The remaining class of natural models includes Starobinsky and Higgs inflation (which both have  $r \sim 0.003$ ). A future mission capable of reaching  $\sigma_r \sim \mathcal{O}(10^{-4})$  would provide significant constraints on nearly every currently favored inflation model. The EPIC-IM configuration is forecasted to achieve  $\sigma(r) \sim 4.8 \times 10^{-4}$  assuming  $r = 0.01$  and no foregrounds. **Check this statement.**

A detection of B-modes consistent with a primordial spectrum of vacuum fluctuations would be the first observation of a phenomena directly related to quantum gravity. In addition, any detection with a next generation satellite would be evidence for “large-field” inflation [19], where a smooth potential that supports inflation extends over a distance in field space  $\Delta\phi$  of order  $M_p$ . Any



Figure 2: Forecasted constraints in the  $n_s$ - $r$  plane for a fiducial model with  $r = 0.01$  for EPIC-IM together with the current measured  $1$  and  $2\sigma$  constraints (blue) [? ]. Also shown are predictions for the models of the inflaton potential discussed in the text: Chaotic inflation for a range of  $N_*$  values (blue lines); Higgs and  $R^2$  (large and small dots, respectively); quartic hilltop (green band); and a sub-class of  $\alpha$ -attractor models [? ]

inflationary model is necessarily an effective theory valid only below the scale of quantum gravity, but it is a generic expectation that additional heavy degrees (mass  $M$ ) required by quantum gravity [ ] can easily generate structure in the low-energy potential that limits inflation to  $\Delta\phi \lesssim M \sim M_p$ . Studies of inflation in string theory have both born out the generic expectation [ ] and suggested mechanisms to realize large-field inflation [ ]. A detection of  $r$  would therefore provide very strong motivation to better understand how large-field inflation can be naturally incorporated in quantum gravity.

Inflation generically predicts a B-mode spectrum with the shape shown in Figure 2, but inflation need not be correct [ ] and does not preclude additional sources of  $B$ -mode polarization either during or after inflation. To be confident of the implications of a detection, the shape and Gaussianity of the B-mode spectrum must be characterized. The vast majority of inflation scenarios predict an extremely Gaussian and very nearly scale-invariant, but slightly red, spectrum for gravitational waves. A future satellite mission, combined with Stage 3 ground based observations, could target  $\sigma(n_t) \lesssim 1$  at  $r = 0.01$ , for example, to rule out non-vacuum inflationary sources [20, 21] and physics completely inconsistent with inflation.

Although the  $B$ -mode polarization is the richest source of new information, deeper mapping of  $E$ -mode polarization will also contribute to testing inflationary models. On the largest scales (accessible only from space),  $E$ -modes will provide new tests of isotropy, while on sufficiently small scales they will allow tighter constraints on the shape of the scalar power spectrum, the amplitude of scalar non-Gaussianities, and isocurvature modes. A satellite mission, combined with next generation ground-based instruments will improve on all measures of the scalar sector by factor of order two.

There is also much to be learned from very small scales. The dissipation of small-scale perturbations through Silk-damping [22, 23, 24, 25] leads to  $\mu$ -distortions in the CMB spectrum, which provide a unique means to place stringent constraints on the amplitude of the small-scale curvature power spectrum. This information is on small scales (wavelength  $0.1 \text{ kpc} \lesssim \lambda \lesssim 1 \text{ Mpc}$ ) and from early times ( $10^4 \lesssim z \lesssim 10^6$ ), inaccessible through any other observation. In  $\Lambda$ CDM [26], distortions are predicted at a level of  $\mu = (2.0 \pm 0.14) \times 10^{-8}$ . A detection would deliver a complementary test for the inflation paradigm [27, 28, 29, 30, 31], and new probes of the particle spectrum of inflation through new tests of non-Gaussianity [32, 33, 34, 35]. Precise measurements of signals at this level will be extremely challenging and requires unprecedented control of systematics and modeling of foregrounds. It would also bring us to the sensitivity level required to detect the cosmological recombination radiation [36, 37] imprinted by the recombination of hydrogen and helium at redshift  $z \simeq 10^3 - 10^4$ , which can be used to probe the physics of recombination. Optimizing next-generation CMB spectrometers for these purposes requires extensive studies.

### 1.2.2 Neutrinos and Light Relics

After inflation, the universe was reheated to temperatures of at least 10 MeV and perhaps as high as  $10^{10}$  GeV. At these high temperatures, even very weakly interacting or very massive particles, such as those arising in extensions of the Standard model of particle physics, can be produced in large abundances. As the universe expands and cools, the particles fall out of equilibrium and leave observable signatures in the CMB power spectra. Through these effects the CMB is a sensitive probe of neutrino properties and of other candidate light particles.

Much of the information about our thermal history and the particle content of the universe is encoded in the  $T$  and  $E$  power spectra. A high-precision measurement of these spectra over the full sky is expected to significantly improve our understanding of the post-inflationary universe. This is particularly true in  $E$ -mode polarization where, to date, far fewer modes have been measured at the level of cosmic variance than in temperature.

One particular compelling target is the effective number of neutrino species  $N_{\text{eff}}$ , which parameterizes the total amount of energy density in radiation at the time of recombination. Despite its common name,  $N_{\text{eff}}$  encodes the total number of light relic particles. The canonical value with three species of neutrinos is  $N_{\text{eff}} = 3.046$ . Additional light particles contribute a change to  $N_{\text{eff}}$  of at least  $\Delta N_{\text{eff}} \geq g 0.027$  where  $g \geq 1$  is the number of degrees of freedom of the new particle [38, 39]. This defines a compelling target of  $\sigma(N_{\text{eff}}) < 0.027$  for future CMB observations. New light particles are a common feature of many approaches to physics beyond the Standard model and are directly tied to some of the most significant problems in the Standard model. Either a limit or detection of  $\Delta N_{\text{eff}}$  at this level would provide a powerful insight into the fundamental laws of nature.

The presence of free-streaming radiation changes the detailed features of the  $TT$ ,  $TE$  and  $EE$  spectra at all  $\ell$ . In particular, it changes the locations of the acoustic peaks and alters the damping tail at high- $\ell$  [40, 41, 42, 43]. Similar changes to the spectra arise from many other compelling targets including the helium fraction  $Y_p$  and more general dark sector physics. For this reason, constraints on  $N_{\text{eff}}$  a useful proxy for the information available in the high- $\ell$  power spectra.

Forecasts for  $N_{\text{eff}}$  are shown in the right hand panel of Figure 3. The two most important quantities for improving constraints on  $N_{\text{eff}}$  and other high- $\ell$  targets are the fraction of sky observed  $f_{\text{sky}}$  and the noise temperature. Both are specific strengths of a space mission. The EPIC-IM mission reaching an effective noise temperature noise of  $0.9 \mu\text{K-arcmin}$  over the full sky gives competitive constraining power when compared other proposals. The full-sky nature of the proposed mission would allow for cosmic variance limited  $E$ -modes over most of the sky and a large range of  $\ell$ .

The main downside of a space-based mission is that we cannot reach the resolutions available from the ground. However, we see that at  $5'$  resolution and  $1 \mu\text{K-arcminute}$  noise the forecasts are less sensitive to the resolution than one might naively expect. In particular we can reach  $\sigma(N_{\text{eff}}) < 0.035$  for temperature noise from  $1\text{--}2 \mu\text{K-arcmin}$  and  $f_{\text{sky}} = 0.6 - 0.8$ . These forecasts are competitive with CMB Stage IV. Specifically, the larger sky fraction and sensitivity available from space appears to compensate for the reduced resolution. In fact, the full sky measurement would provide complimentary information that could be combined with ground based surveys to further improve over the limits available from either experiment. This is particularly important for  $N_{\text{eff}}$  which is tantalizingly close to the target of  $\sigma(N_{\text{eff}}) = 0.027$  and therefore even an apparently modest improvement could have a major scientific impact.

A measurement of the sum of neutrino masses  $\sum m_\nu$  is another compelling target for the CMB Probe. The lower the sum of the masses, the more suppressed is the growth of structures in

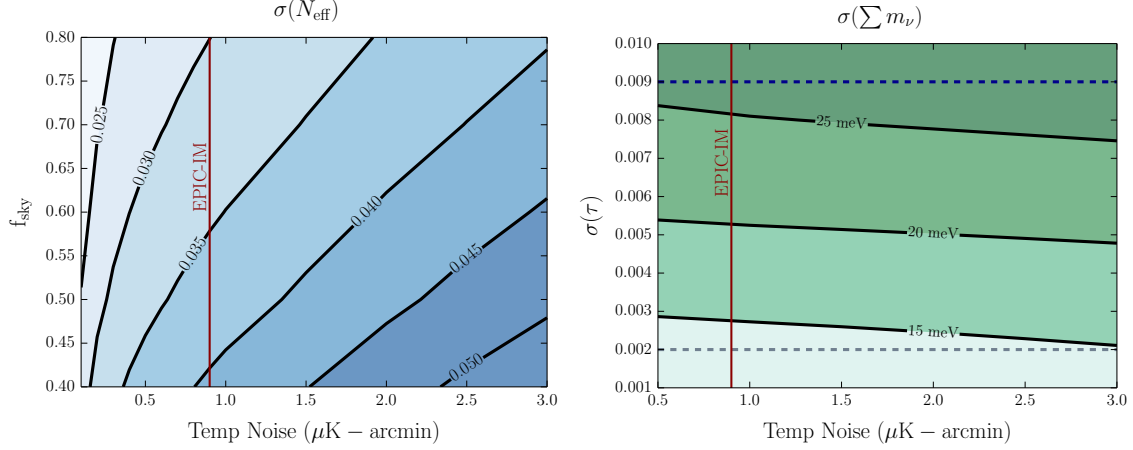


Figure 3:  $N_{\text{eff}}$  as a function of noise and sky fraction assuming  $5'$  resolution (left) and Neutrino mass constraints as a function of uncertainties in measurement of  $\tau$  and noise for a  $5'$  beam and sky fraction of  $f_{\text{sky}} = 0.7$ . Vertical lines denote the expected performance of EPIC-IM. The blue dashed line is the current *Planck* limit; the grey dashed line is the limit from cosmic variance measurement of  $\tau$ . All forecasts assume internal delensing of the  $T$  and  $E$ -maps [45], including residual non-Gaussian covariances.

the Universe on small scales. The level of suppression, and therefore  $\sum m_\nu$  can be measured by comparing the amplitude of the lensing power spectrum to that of the other CMB spectra [44]. But the sensitivity to  $\sum m_\nu$  is limited by our knowledge of the primordial amplitude of fluctuations  $A_s$  which is strongly degenerate with the optical depth  $\tau$ .

The current limit on  $\tau$  from *Planck* of  $\tau = 0.055 \pm 0.009$  [ ] limits  $\sigma(\sum m_\nu) \gtrsim 25$  meV; see Figure 3. Therefore, a detection of the minimum value expected from particle physics  $\sum m_\nu = 58$  meV at  $3\text{-}5\sigma$  depends crucially on an improved measurement of  $\tau$ . To date, the only proven method for such a measurement is from space-based CMB observations. The best constraints on  $\tau$  come from  $E$ -modes with  $\ell < 20$  which requires control over the largest angular scales.

Measurements of the primary CMB and CMB lensing are very sensitive to particles with long-lifetimes that service until the era of recombination. Furthermore, the primary CMB is indirectly sensitive to BBN through the helium fraction,  $Y_p$ , that can be detected energy injections at those times. However, many relics of the early universe are not cosmologically stable but can survive beyond the time of BBN. Energy injection between the two eras is not directly captured by the CMB anisotropies but can be tested by BBN consistency of  $Y_p$  with the measured value of  $N_{\text{eff}}$  [43, 46].

Additional, far more sensitive measurements of this intermediate era are possible through the *distortions* of the CMB black-body spectrum. This will allow us to place stringent bounds on the presence of long-lived decaying particles [29, 47, 48, 49] and other non-standard physics, such as axions, primordial magnetic fields and superconducting strings [e.g., 50, 51, 52, 53, 54].

### 1.2.3 Cosmological structure formation

Understanding the evolution of cosmological structures from the formation of the first stars to the present time is a key goal of cosmology [55].

In particular understanding cosmological reionization, the transformation of neutral hydrogen into an ionized state, and establishing a connection between reionization and early galaxies, will reveal valuable information about the star formation history and the physical processes that formed the galaxies of various luminosities and masses we see today. Though a number of diverse observational probes have rendered a preliminary picture of the reionization epoch, details of the transition

from the neutral to ionized Universe are still the subject of intense activity. In the current picture, early galaxies reionize hydrogen progressively throughout the entire Universe between  $z = 12$  and  $z = 6$ , while quasars take over to reionize helium from  $z \simeq 6$  to  $z \simeq 2$ . However many questions remain. When did the epoch of reionization (EoR) start, and how long did it last? Are early galaxies enough to reionize the entire Universe or is another source required? To answer these questions one can resort to using the traces left by the EoR in the cosmic microwave background (CMB) anisotropies. More specifically, the integrated Thomson scattering optical depth,  $\tau$ , due to the scattering of CMB photons off free electrons after reionization, as inferred from the polarised CMB photons. The optical depth to reionization,  $\tau$ , places an important integral constraint on the extended reionization history. The *Planck* Collaboration [14, 15] reported recently a value of  $\tau = 0.055 \pm 0.009$  significantly lower than previous estimates. This suggests that an early onset of reionization is strongly disfavoured by the *Planck* data. The *Planck* Collaboration [15] showed that this result reduces the tension between CMB-based analyses and constraints from other astrophysical sources. A cosmic variance limited measurement of E-mode polarization on large scales, possible with a probe mission, will render the most accurate determination of  $\tau$  (Figure 3 shows a cosmic variance limit measurement of  $\tau$  along with the current *Planck* limit, break the degeneracy with the neutrino mass, set stringent constraints on models of the reionization epoch, and, finally, help understanding the formation of the cosmological structures we see today.

The anisotropies in the Cosmic Infrared Background (CIB), produced by dusty star-forming galaxies (DSFGs) in a wide redshift range, are an excellent probe of both the history of star formation and the link between galaxies and dark matter across cosmic time. The *Planck* Collaboration [56, 57], derived values of the star formation rate density that, at redshifts  $z \sim 3$ , are about three times higher than constraints from number counts measurements ([58]). The new mission probe, by measuring CIB anisotropies with 100 times more sensitivity than *Planck*, will be able to shed light on this intriguing discrepancy. In particular, it will be possible to constrain the global star formation rate (SFR) with one tenth of *Planck*'s uncertainty at  $z = 3$ . Moreover, it will be possible to identify and constrain a characteristic halo mass  $M_{\text{eff}}$ , which determines the most efficient gas accretion and SFR, and therefore sets the evolution of the galaxies residing within a dark matter halo. Current models and measurements constrain this characteristic halo mass at  $M_{\text{eff}} \sim 10^{12}$  solar masses with about 10% uncertainty, while the new mission probe will constrain this parameter at the percent level.

Moreover, because DSFGs trace the underlying dark matter field in a broad redshift range, the CIB will correlate with multiple dark matter tracers such as catalogs of galaxies and quasars [59, 60], and diffuse maps of the  $\gamma$ -ray and the X-ray background [61]. These cross-correlations will provide an additional probe of the global star formation history, and they will shed light on the interaction between light and matter in a broad wavelength range.

Large-scale structure can also be probed using CMB spectral distortions measurements. In fact, the largest guaranteed distortion is caused by the associated late-time energy release of forming structures and from reionization [62, 63, 64, 65, 66], imprinting a  $y$ -type distortion with  $y \simeq 2 \times 10^{-6}$  [e.g., 66, 67]. This distortion is only one order of magnitude below the current limit from COBE/FIRAS and, even with most pessimistic assumptions about foregrounds, should be clearly detected with the next-generation spectrometers we propose to study. A detection will give information about the total energy output of first stars, AGN and galaxy clusters. In particular, group-size clusters that have masses  $M \simeq 10^{13} M_{\odot}$  contribute significantly to the signal. With temperature  $kT_e \simeq 1$  keV these are still sufficiently hot to create a detectable relativistic temperature correction to the large  $y$ -distortion, which can be used to constrain the currently uncertain feed-



back mechanisms used in hydrodynamical simulations of cosmic structure formation [67]. These two inevitable signals probe the low-redshift Universe and provide clear targets for future spectral distortions measurements and their requirements in the presence of foregrounds.

The CMB spectrum varies spatially across the sky. One source of such anisotropic distortion is related to clusters of galaxies and has already been measured by Planck [68]. A combination of precise CMB imaging and spectroscopic measurements will allow observing the relativistic temperature correction of individual SZ clusters [69, 70, 71], which will calibrate cluster scaling relations and inform our knowledge of the dynamical state of the cluster atmosphere. Finally, resonant scattering signals in the recombination [72, 73, 74] and post-recombination eras [75, 76] can lead to spectral-spatial CMB signals that can be used to constrain the presence of metals in the dark ages and the physics of recombination. For all these applications, instrumental synergies between CMB imaging and spectroscopy need to be studied in detail.

These studies will be key to address three of the seven key questions identified in the Astro2010 report “New Worlds, New Horizons in Astronomy and Astrophysics” (NAS Decadal Survey, p. 47): *What is the fossil record of galaxy assembly from the first stars to present? What are the connections between dark and luminous matter? How do cosmic structures form and evolve?*

#### 1.2.4 Dark Matter

Characterizing the nature of dark matter is one of the most basic problems in cosmology and astroparticle physics. For a conventional WIMP candidate, the CMB places very stringent constraints on dark matter through possible energy injection at the time of recombination. A thermal relic will necessarily annihilate into Standard model particles which distorts the Thomson visibility function and ultimately the shape of the primary  $T$  and  $E$  spectra [77, 78, 79]. A dark matter candidate can also have a slow decay to Standard models particles which is similarly constrained by the CMB [80].

A somewhat more model-independent approach is to constraint dark matter interactions that would effect the evolution of the effective dark matter fluid and its interactions with baryons or photons. The simplest example is to constraint the baryon-dark matter cross section through its effective coupling of the two fluids [81]. These couplings effect the evolution of fluctuations and ultimately the  $T$  and  $E$  spectra. More exotic dark sectors can produce a rich phenomenology in the CMB and in the large-scale structure without necessarily producing an associated signature in direct detection experiments or indirect searches (e.g. [82, 83, 84]).

Interactions of DM with standard model particles can also be constrained through CMB spectral distortions [85]. As the Universe expands, the matter cools. Compton interactions with electrons keep the normal matter at the CMB temperature until well after recombination. This leads to a small distortion of the CMB with a negative chemical potential [86]. In a similar manner, interactions of DM with electron, protons or directly with CMB photons can lead to a distortion. These constraints could be strongly improved with future CMB spectroscopy [85].

### 1.3 The Challenges: Foregrounds, systematics

**3 pages. Discuss the challenge of Foregrounds and Systematics.**

The search for primordial B-modes is one of the main science objectives of any future CMB mission. A satellite mission is the most sensitive probe of this signal from the early universe because it can access the full sky and the choice of frequencies is not limited by the atmosphere.

By the time of launch of a probe class mission substantial progress will already have been made. In the case of a detection at the time of launch, the probe mission should be able to convincingly



detect the expected reionization signal. In the absence of a signal it should significantly improve the upper limits beyond  $\sigma(r) = 0.001$ .

The contribution from reionization to the Stokes  $Q$ -parameter for a tensor-to-scalar ratio  $r = 0.001$  is shown in the left panel of Figure 4. The amplitude of the signal is approximately 10 nK so that any experiment attempting to measure such a signal must control foregrounds and systematics on large angular scales at an unprecedented level of a few nK.

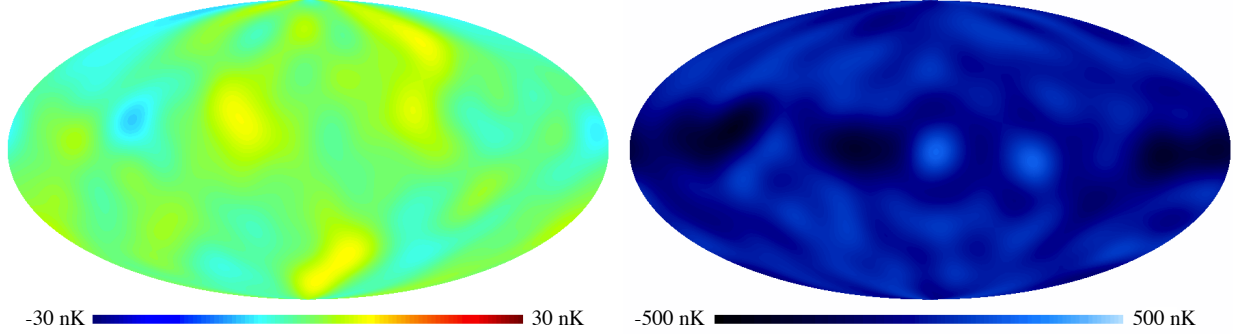


Figure 4: *Left panel:* Contribution to the Stokes  $Q$  parameter from inflationary B-modes for  $\ell < 12$  for  $r = 0.001$ . *Right panel:* Noise in the *Planck* 353 GHz map of the Stokes  $Q$  parameter for  $\ell < 12$  rescaled to 150 GHz assuming the spectral properties of dust.

### 1.3.1 Foregrounds

Data from the *Planck* satellite has significantly improved our understanding of foregrounds in both intensity and polarization. In polarization, the sky is dominated by the expected sources, synchrotron and dust, and *Planck* has provided us with much improved measurements of their amplitude and spectral dependence. The observed spectral dependence is shown in the left panel of Figure 5 for different sky fractions. One of the most important lessons is that foregrounds are larger than putative primordial B-mode signal everywhere.

This is particularly severe on the scales relevant for the search of primordial B-modes as can be seen in the right panel of Figure 5, which shows the power spectrum of foregrounds over 75% of the sky for frequencies between 70 and 200 GHz together with the lensing and inflationary contribution for different values of the tensor-to-scalar ratio.

Perfect knowledge of the foreground components would allow to remove them, but the sensitivity of *Planck* limits our ability. The right panel in Figure 4 shows the noise in the *Planck* 353 GHz map of the Stokes  $Q$ -parameter rescaled to 150 GHz assuming the same spectral dependence as for dust and on the angular scales relevant for the measurement of the inflationary  $B$ -modes. The noise is more than an order of magnitude larger than the inflationary contribution for  $r = 0.001$ , clearly indicating the necessity to measure foregrounds with a potential future space mission. A probe class mission likely is the only way to map foregrounds at a level required to measure primordial gravitational waves for  $r < 0.001$ .

While the search for primordial B-modes leads to the strictest constraints on foreground residuals, exquisit control of foregrounds is also necessary for other science objectives. A satellite mission is likely also the only reliable way to measure the optical depth at a level necessary to

break the degeneracy with the neutrino mass. Furthermore, a cosmic variance limited measurement of E-mode polarization on large scales possible with a probe mission would contain valuable information about the star formation history.

Similarly, a clear objective of the spectral science is to have a robust, foreground-marginalized expectation of detecting the  $\mu$ -distortion generated by the dissipation of small-scale acoustic modes in  $\Lambda$ CDM. While an instrument like PIXIE just falls short of this objective, it seems within reach of a probe mission.

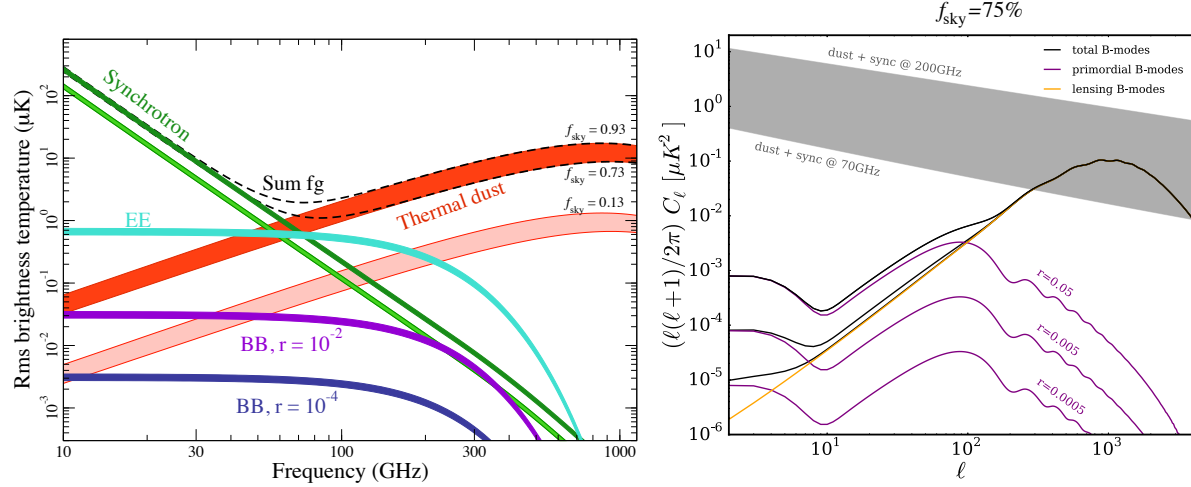


Figure 5: *Left panel:* Brightness temperature as function of frequency for the CMB as well as synchrotron emission (green) and dust emission (red). The darker bands show the brightness temperature for sky fractions between 73% and 93%, the lighter bands show the brightness temperature for the cleanest 13% with the width indicating the uncertainty. *Right panel:* Angular power spectrum for B-mode polarization of the CMB for  $r = 0.0005$ ,  $r = 0.005$ , and  $r = 0.05$  as well as for foreground emission between 70 and 200 GHz.

One of the key ingredients in the design of a CMB experiment is the frequency coverage required to achieve the science goals. Consequently, optimizing frequency coverage in light of the new information from *Planck* and its limitations will be one of key task of the study proposed here.

To achieve these goals we plan to investigate the effect on the measurements of  $r$  and  $\tau$  of the presence of foregrounds residuals in the CMB maps (after foreground separation and/or cleaning) including, for example, the properties of the polarized thermal dust emission, specifically the potential spatial variation of its spectral index, hinted at by the observed decorrelation of the dust between 217 GHz and 353GHz *Planck* channels (??), and the study of the breakdown of the modified black body spectrum model. These aspects are better explored with the help of physically motivated models of the foregrounds (??) and simulations based on these, using either existing simulation tools and/or implementing new simulators when required.

The optimization of frequency channels will be conducted using both traditional (Fisher codes both spectra and map based) and novel techniques currently in development (such as direct Bayesian MCMC inference of cosmological parameters in the presence of foregrounds, (extension of the method presented in ??). Realistic simulations (including time domain based) are essential to fully assess the performance of a given instrument design/concept in view of both foreground residu-

als and the presence of systematics. We plan to generate these simulations, process them as we would do the real data, specifically applying some of the current component separation methods to clean the frequency maps from foreground emission (akin to *Planck* collaboration procedures ??) and propagate to parameter estimation, assessing the impact of all contaminating effects on the accuracy and potential biases of the parameters estimated with particular emphasis on  $r$  and  $\tau$ .

### 1.3.2 Systematic Errors

Advances in detector technology since the formulation of Planck and WMAP will enable huge gains in raw sensitivity for a CMB probe. To fully take advantage of this sensitivity, systematic errors must be controlled to detect polarization signals at nano-Kelvin levels. The proposed study will invest heavily in designing an instrument, test plan, and observation strategy to address systematic errors, gathering decades of experience of ground-, balloon-, and space-based CMB polarimetry. The latest analyses of Planck HFI and LFI data in particular show the effects likely to be important to a future space mission(<https://arxiv.org/abs/1605.02985>). These systematic errors can be considered in three broad categories: 1. Intensity-to-polarization leakage, 2. stability, and 3. straylight. Each of these is considered in light of differential polarimetry; the instantaneous signals measured by polarization-sensitive detectors at different times and orientations are combined to recover the maximum likelihood polarization signal from each point on the sky.

**Leakage.** The CMB temperature signal is many orders of magnitude larger than the polarization B-mode signal (see, e.g. Fig. 2). Therefore, instrumental mismatch effects that can leak even a very small fraction of an intensity fluctuation into a spurious polarization signal must be addressed. The main effects are relative gain and bandpass calibration, differential pointing error, and differential beam shapes.

Relative calibration requirements are likely to exceed those of Planck, whose High Frequency Instrument achieved of order 0.01% (cite <https://arxiv.org/abs/1605.02985>). Bandpass mismatch is the imperfect matching of the frequency dependant gain of the two detectors or acquisition channels used to estimate the polarized signal. Any asymmetry turns into a spurious polarization signal when viewing sources with SED different from the calibrator. For CMB-probe, this includes every foreground contaminant on the sky. The effect can be minimized by matching the bandpasses, and residual mismatch can be estimated by measuring the bandpasses on the ground, however, for both instruments of Planck the ground based measurements of the bandpasses were found to be insufficient to determine the impact on the polarized maps. In both cases estimates of the bandpass mismatch needed to be determined as part of the mapmaking pipeline, and bandpass mismatch remains one of the most problematic systematics in the pipelines. CMB-probe will need a more robust program to mitigate this problem on ground, simulate the impact, and to deal with it in the flight data.

A second order complication of unmatched bandpasses is unmatched far sidelobe contamination, leading to a spurious polarized component from bright unpolarized signals far from the main lobe.

These systematics are likely to drive the instrumental requirements on the optical system as well as the uniformity of the bandpass of each polarimeter. Calibration requirements will also be set by limiting these systematics: particularly on the knowledge of the polarization parameters (such as cross-polar leakage and the angle of polarization sensitivity), as well as measurement of the beam shape (in general a function of the SED of the observed source). These systematic effects can potentially be mitigated by modulation of the sky signal in such a way that allows complete reconstruction of the polarized sky signal using each photometer, for example, using a half-wave

plate.

**Stability.** Given the need to avoid light from the Sun, Earth, and Moon, the full reconstruction of the polarized sky will necessarily involve combination of measurements made at times separated by months, requiring stability of the response of the instrument on corresponding time scales. This systematic error puts requirements on control of thermal drifts of spacecraft temperatures, to mitigate thermal emissivity changes and thermoelastic deformation of telescope structures. The cryogenic operating temperatures of detectors or reference calibration loads must be controlled adequately as well. Careful design of the scan strategy can shorten the time scales needed for stringent stability, for example Planck’s scan strategy traced out great circles which overlapped on 1 minute timescales, giving a shorter effective time scale for stability requirements.

Additionally, the space radiation environment is modulated by the solar activity and can introduce drifts in the cryogenic thermal environment as well as introducing correlated transients in detectors and readout electronics. The design of the instrument must account this environment, which following Planck is much better understood.

**Straylight.** The brightest cm-wave and mm-wave sources in the sky (such as the Sun, Moon, planets, and Galactic center) passing into the far sidelobes of the telescope (defined as the response of a detector from a source more than a few degrees from the optical axis) in a sky-synchronous way can create a spurious polarization signal. The far sidelobe response can be reduced by the optical design and baffling, but will always be present at a non-trivial level. The Planck experience is instructive here as well. The detailed GRASP model of the telescope, convolved with the sky model, predicted sidelobe contamination at a visible (10’s of microK) level in the 30 GHz maps, which was observed in difference maps. As a result the LFI maps have had an estimate of the sidelobe contamination removed from the timelines as part of the mapmaking process. The more stringent requirements for CMB-probe will necessitate at least this level of mitigation. A major limitation to the analysis of far sidelobe contamination in the Planck data has been the lack of bright enough on-orbit sources to validate the GRASP simulation as adjusted by optical and bandpass parameters estimated on-orbit. Finding a way to better validate the FSL model on-orbit for CMB-probe may be critical to successfully removing FSL contamination.

## 1.4 Current and Forthcoming Efforts and the CMB Probe

2.5 pages. S3 experiments, forthcoming S4, Baselines CMB Probe options and their complementarity with S4.

Bill Jones is writing this. Due by Oct. 28

## 1.5 State of Technologies

2 pages. Discuss the technologies, their TRL, and what will be studied

A fourth generation CMB satellite targeting a map sensitivity of  $\sim 1\mu\text{K} - \text{arcmin}$  will require, extremely sensitive detector arrays, tight control over systematics, and ability to reject polarized foregrounds as is described in Section 1.2. Given the frequency dependence of synchrotron and dust foregrounds, this last requirement translates into the need for a large number of spectral bands covering the approximate frequency range from  $\sim 30$  GHz to  $\sim 800$  GHz. Development of the CMB technologies needed to meet these requirements is actively being pursued by many groups who are also demonstrating these technologies on ground, balloon, and satellite platforms. We describe the status and needs in the areas of telescopes, optics, detector coupling, detectors, and readout.

**Telescopes:** Carbon Fiber Reinforced Polymer (CFRP) mirrors are at TRL 9 as they have flown on the Planck satellite. Their 1.9x1.5 m mirror weighted only 28 lbs and met all surface quality requirements. However, small deformations in the mirror caused by its structural supports had a measurable impact to the beam far-sidelobes that was not captured by preflight measurements or the corresponding beam modeling. Future CMB satellites will require improved pre-flight characterization of *polarized beam* at operating temperature augmented by improved simulation tools to meet even more systematics requirements. Current ground and balloon born optical designs achieve large field of views (FOVs) with reflective and refractive designs; related designs and their implementation should be studied in the context of a satellite mission as the sensitivity requirements lead to the need to maximize the size of the FOV while fitting within the tight mass and size constraints imposed by a space mission. Given the heritage of past satellite missions it will be possible to develop a telescope design that meets the requirements for a future mission and uses high TRL components.

**Optical Coupling:** The need for sensitivity drives the push for high efficiency optics; wide bandwidth to compliment multichroic detector; infrared filters to maximize cryogenic performance; and polarization modulators to suppress  $1/f$  noise and mitigate instrument systematics. The CMB field has made tremendous progress recently by drawing on advances in materials, processing techniques, and developments in electrical engineering including meta-material research. Single crystals such as silicon and sapphire are attractive since they offer extremely low dielectric losses and high indices of refraction to better manipulate light. New coating techniques have been developed for silicon and sapphire that span 2:1 bandwidth (TRL 5+ for silicon) and can realize up to 5:1 bandwidth. EBEX deployed broadband cryogenic polarization modulator with a superconducting bearing that covered 150 GHz band to 410 GHz band raising the modulators to TRL 5+ for space. Meta-material metal-mesh optical filters were deployed with the Planck satellite and they are extensively used by ground based and balloon experiments making these TRL6 optical elements. It is necessary to develop a plan for a satellite mission that will cover  $\sim 30$  GHz to  $\sim 800$  GHz. Two configurations could be considered: multiple optical paths with  $< 3 : 1$  bandwidth and a potentially simpler design with only two optical paths with  $\sim 5 : 1$  bandwidth. These studies include evaluating the design tradeoffs inherent to these approaches, developing the new coatings needed, and evaluating the promise of hybrid approaches where filters and lenses are implemented in the same optical elements. In addition, the cryogenic rotation mechanism should be demonstrated at the robustness (eg lifetime) needed for a satellite mission.

**Detector Coupling:** The focal-plane feed determines the shape and polarization properties of the pixel beams and therefore plays a strong role in controlling systematic errors. The feed design also can determine the total bandwidth and number of photometric bands of each pixel which is important for the efficient use of a telescope's focal plane area. CMB experiments developed broadband multichroic *detector* to increase optical throughput of a focal plane. Broadband feed captures signal over wide frequency range. Then on-chip superconducting filter partitions signal into multiple frequency bands prior to detection. Broadband detectors were realized with spline profiled horn and lenslet coupled antenna. Broadband horn detector deployed a pixel that covers 2.3:1 bandwidth with on going development for extending bandwidth to 6:1. Broadband lenslet coupled antenna will deploy 3:1 bandwidth detector this year. Lenslet coupled antenna demonstrated 5:1 bandwidth in laboratory. RF-techniques to partition broadband signals into multiple band are mature. For a

future CMB polarization satellite mission, broadband feed should be demonstrated at high frequency where alignment and line width for micro-fabrication becomes challenging. Detectors for CMB satellite mission were hand picked one by one for optimal performance. Next generation of detector array will be fabricated on a silicon wafer. Micro-fabrication process should demonstrate high yield and uniformity across a wafer that can meet tight requirement of satellite mission. Also detector test need to be able to characterize detector beyond the level of systematic required by next generation CMB satellite experiment.

The Planck HFI deployed Neutron Transmutation Doped Germanium high-resistance bolometer at 100 milli-Kelvin to achieve photon noise limited detector performance. A Transition Edge Sensor (TES) bolometer uses a steep transition of superconducting metal to improve linearity of the detector. TES bolometers have been deployed on 100 milli-Kelvin and 250 milli-Kelvin platform. TES bolometers have been deployed across ground based and balloon CMB experiments spanning 40 GHz-410 GHz with detectors achieving NEPs of  $20\text{-}50\text{ aW}/\sqrt{\text{Hz}}$ , nearly background limited at CMB frequencies. TES bolometers deployed at low optical frequencies ( $\sim 40\text{ GHz}$ ) and balloon-borne payloads should realized even lower NEPs of  $\sim 10\text{ aW}/\sqrt{\text{Hz}}$ . Emerging detector technology for CMB experiment is kinetic-inductance detector (KIDs). The KIDs detector detects signal as change in kinetic inductance. KIDs detectors can be frequency multiplexed easily to  $\sim 1,000$  detectors. Recently, on-sky demonstration at 150 GHz and 230 GHz was done with lumped element KID detector. Noise performance of KID detector at low frequency channels ( $< 40\text{ GHz}$ ) need some improvement to be photon-noise limited. Currently there is no CMB polarization power spectrum data produced with KID detector. Coupling between RF (100 GHz) signal to micro-wave KID (MKID) detector is in a development stage. Planck detectors experienced unexpectedly high rate of cosmic ray events. Data was successfully cleaned with analysis technique. Study of impact of cosmic rays on a detector is crucial for next CMB satellite mission.

Multiplexed readout is being used by CMB experiments to readout thousands of TES bolometers, and readout multiplexing is built into KID detector architecture. Voltage bias and low impedance of a TES bolometer facilitates multiplexing readout by Superconducting Quantum Interference Device (SQUID). Time domain multiplexing uses a SQUID at milli-Kelvin as a switch to rapidly cycle through bolometers. Highest achieved multiplexing factor is 64 channels. Frequency domain multiplexing uses superconducting resonators to assign bolometers to different frequency channels. Highest achieved multiplexing factor is 68 channels. New readout scheme, such as microwave SQUID readout, is emerging to increase multiplexing factor for TES bolometer. MKID detector architecture has multiple resonators coming off from a transmission line. A resonator is both a detector and multiplexer. MKID demonstrated multiplexing factor that exceeds 1,000 channels. For next generation satellite experiment that will readout over thousands detectors require high multiplexing factor. Multiplexing factor is directly related to readout complexity and power consumption. Also the Planck mission experienced ADC non-linearity, thus extensive characterization of end to end readout architecture should be performed pre-flight.

A future CMB satellite mission offers exciting opportunity for millimeter wave polarization science. Experience from Planck mission will be studied to learn lessons for the future mission. Development for CMB instrumentation is an active field with many institutions developing new technologies for ground based, balloon, and proposed satellite missions. For a satellite instrumentation, there is a difficult trade off between desire to have high performance instrument and desire to keep cost manageable. Many developments that is going on for ground based and balloon experiment have similar goal as satellite mission that collaborative development across all platform will be beneficial.

## **1.6 Mission Study, and Management Plan**

1.5 pages; Describe what we want to do, who is doing what, what we are funding

Shaul is writing this. Due by Tuesday Nov. 1



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## **2 Curriculum Vitae**



**3 Summary of Work Effort**

**4 Current and Pending Support**

## **5 Letters of Support**

## **6 Budget Details - Narrative**

### **6.1 Team, and Work Effort**

#### **6.1.1 Funded Team Members**

#### **6.1.2 Non-Funded Team Members**

### **6.2 Costing Principles**

- **Summer Salaries:**
  - **Workshop:**

### **6.3 University of Minnesota Budget**

#### **6.3.1 Direct Labor**

#### **6.3.2 Supplies**

#### **6.3.3 Travel**

#### **6.3.4 Other Direct Costs**

#### **Publications and Teleconferencing**

#### **Other Subcontracts**

#### **6.3.5 Facilities and Administrative Costs**

## 7 Budget Sheets

**ACS** attitude control system

**ADC** analog-to-digital converters

**ADS** attitude determination software

**AHWP** achromatic half-wave plate

**AMC** Advanced Motion Controls

**ARC** anti-reflection coatings

**ATA** advanced technology attachment

**BRC** bolometer readout crates

**BLAST** Balloon-borne Large-Aperture Submillimeter Telescope

**CANbus** controller area network bus

**CMB** cosmic microwave background

**CMM** coordinate measurement machine

**CSBF** Columbia Scientific Balloon Facility

**CCD** charge coupled device

**DAC** digital-to-analog converters

**DASI** Degree Angular Scale Interferometer

**dGPS** differential global positioning system

**DfMUX** digital frequency domain multiplexer

**DLFOV** diffraction limited field of view

**DSP** digital signal processing

**EBEX** E and B Experiment

**EBEX2013** EBEX2013

**ELIS** EBEX low inductance striplines

**ETC** EBEX test cryostat

**FDM** frequency domain multiplexing

**FPGA** field programmable gate array

**FCP** flight control program

**FOV** field of view

**FWHM** full width half maximum

**GPS** global positioning system

**HDPE** high density polyethylene

**HIM** high index materials

**HWP** half-wave plate

**IA** integrated attitude

**IP** instrumental polarization

**JSON** JavaScript Object Notation

**LDB** long duration balloon

**LED** light emitting diode

**LCS** liquid cooling system

**LC** inductor and capacitor

**LZH** Lazer Zentrum Hannover

**MCP** multi-color pixel

**MSM** millimeter and sub-millimeter

**MLR** multilayer reflective

**MAXIMA** Millimeter Anisotropy eXperiment IMaging Array

**NASA** National Aeronautics and Space Administration

**NDF** neutral density filter

**PCB** printed circuit board

**PE** polyethylene

**PME** polarization modulation efficiency

**PSF** point spread function

**PV** pressure vessel

**PWM** pulse width modulation

**RMS** root mean square

**SLR** single layer reflective

**SMB** superconducting magnetic bearing

**SQUID** superconducting quantum interference device

**SQL** structured query language

**STARS** star tracking attitude reconstruction software

**SWS** sub-wavelength structures

**TES** transition edge sensor

**TDRSS** tracking and data relay satellites

**TM** transformation matrix

**UHMWPE** ultra high molecular weight polyethylene

**UMN** University of Minnesota