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# 1 Scientific, Technical, and Management Section

## 1.1 Baselines and Observables

We are proposing to study a probe-scale mission to extract the wealth of cosmological information contained in the spectrum and polarization of the cosmic microwave background (CMB). The starting points for our study of this CMB Probe are two current-decade space missions, EPIC-IM and Super-PIXIE [1, 2]. EPIC-IM was presented to the 2010 decadal panel as a candidate CMB imaging polarization space mission. It was based on a 2 m aperture telescope and 11,094 bolometric transition edge sensors. PIXIE is a proposed Explorer-scale mission focused on a measurement of the spectrum and polarization of the CMB on large angular scales. Super-PIXIE is envisioned to be a scaled up, more capable version of PIXIE. It consists of 4 spectrometers, each operating between 30 and 6015 GHz with 400 15 GHz-wide bands. Improvements in technology by the next decade will enable the design of a mission that is more capable compared to EPIC-IM and Super-PIXIE. Therefore, all quantitative predictions presented in this proposal, which are based on EPIC-IM and Super-PIXIE, represent *minimum* capabilities for the CMB Probe.

The best measurements of the CMB spectrum – made by COBE/FIRAS approximately 25 years ago – show that the average CMB spectrum is consistent with that of a blackbody to an accuracy of 4 parts in  $10^4$  [3, 4]. Distortions in this spectrum encode a wealth of new information. The distortion shapes are commonly denoted as  $\mu$ - and  $y$ -types [5, 6]. The  $\mu$ -distortion arises from energy release in the early Universe and can only be produced in the hot and dense environment present at high redshifts. This makes  $\mu$ -distortions a novel messenger from a redshift range  $z \geq 5 \times 10^4$ . The  $y$  distortions are caused by energy exchange between CMB photons and free electrons through inverse Compton scattering. These originate at lower redshifts and are sensitive to the evolution of the large scale structure of the Universe.

Thomson scattering at the surface of last scattering is the source of the polarization of the CMB. It is useful to decompose the polarization field to two modes that are independent over the full sky,  $E$  and  $B$  modes. Together with the pattern of temperature anisotropy  $T$ , the CMB thus gives three auto- and three cross-spectra. The *Planck* satellite and larger aperture ground-based instruments measured the  $T$  spectrum to cosmic variance limit for  $\ell \leq ??$ . Much information remains encoded in the  $E$  and  $B$  spectra, whose full exploration has just begun [7, 8? ? ].

A future CMB Probe-scale mission will address the physics of the big bang and of quantum gravity; it will measure the sum of the neutrino masses, and constrain the effective number of light particle species and the nature of dark matter; it will probe the existence of new forms of matter at the early Universe; it will give new insights on the star-formation history across cosmic times, and it will provide information about the processes that control structure formation. In addressing these broad array of fundamental questions the Probe firmly fits into NASA’s strategic plan as articulated by its Strategic Goal 1 “Expand the frontiers of knowledge”, and specifically Objective 1.6 “Discover how the Universe works, [and] explore how it began and evolved”.

## 1.2 Science Objectives

### 1.2.1 The Primordial Universe and Cosmic Inflation

The observed temperature and  $E$ -mode polarization of the CMB require primordial inhomogeneities in the gravitational potential, providing a remarkable observational link to the dynamics of the Universe near the big bang. Inflation, a primordial era of accelerated expansion, provides a compelling dynamical origin for the observed nearly scale-invariant spectrum of the primordial perturbations [9, 10, 11, 12, 13]. But, Inflation also predicts an as yet unobserved spectrum of primordial

gravitational waves sourced directly by quantum fluctuations of the tensor component of the metric. These gravitational waves make a distinct B-mode imprint on the polarization of the CMB. Any detection of B-mode polarization, whether generated by the primordial gravitational waves of Inflation [14, 15] or by any other source of early time vector or tensor perturbations, would reveal completely new information about the primordial era. The results would provide significant constraints and consistency checks for current models or could perhaps even overturn them. A detection would have implications for fundamental physics by providing evidence for a new energy scale near the GUT scale. In the context of Inflation, the relationship is particularly clear: the potential energy  $V$  of the inflaton is related to the tensor-to-scalar ratio  $r$  at the peak of the spectrum by  $V^{1/4} = 3.7 \times 10^{16} r^{1/4}$  GeV.

Figure 1 shows current CMB data, B-modes from vacuum fluctuations of the metric during an Inflationary era for two values of  $r$ , as well as forecasts for the determination of the CMB spectra for EPIC-IM. The most recent constraint on the tensor to scalar ratio gives  $r < 0.07$  (95%); see Figure 2 [18]. For testing inflation, the largest scales  $\ell \leq 10$  are particularly important because they reveal the presence of B mode correlations on scales that were super-horizon at the time of recombination [16], and because the signal is strongest relative to the B-mode from lensing. No sub-orbital platform has yet produced B-mode measurements at  $\ell < 80$ , and a satellite is by far the most suitable platform to making the all sky observations necessary to reach the lowest modes,  $\ell < 20$ . In its recent report New Worlds New Horizons (NWNH), the decadal survey committee strongly endorsed searches for the B-mode signal from inflation saying that “The convincing detection of B-mode polarization in the CMB produced in the epoch of reionization would represent a watershed discovery.” [17].

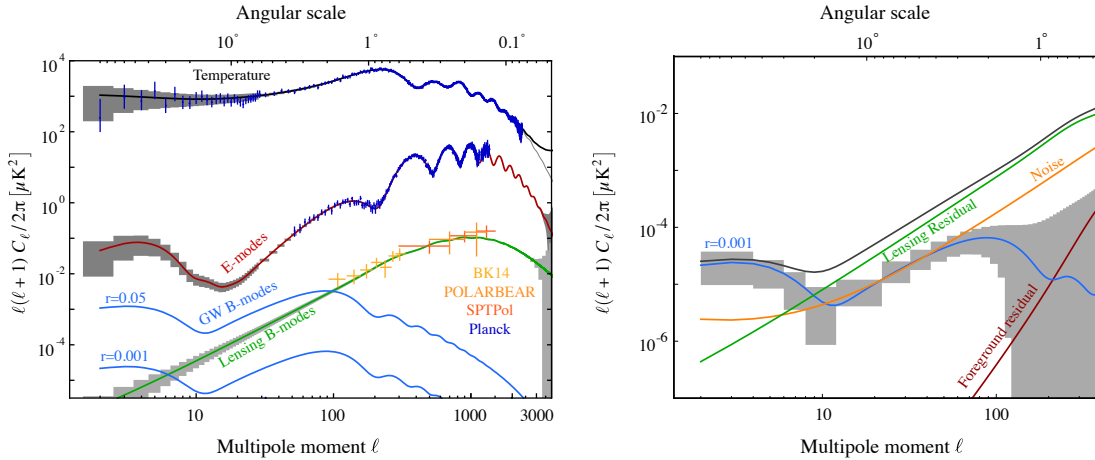


Figure 1: Predicted determination of the CMB power spectra for EPIC-IM (grey boxes) overlaid on theoretical predictions (solid lines) and including Planck measurements of the temperature and E modes (blue) and of several ground-based measurements of the lensing B-modes. The tensor B-mode predictions (blue) are shown for two representative values of the tensor-to-scalar ratio:  $r = 0.001$  and  $r = 0.05$ .

In slow roll Inflation there are just two observationally viable classes of models that naturally explain the measured value of the spectral index  $n_s$ . One is the set of potentials  $V(\phi) \propto \phi^p$ , which contains many of the canonical inflation models. This set is already under significant observational pressure. If the error bars on the spectral index tighten by a factor of about 2, and the 95% C.L. upper limit on  $r$  is pushed to even  $\sim 0.01$ , all such models would be ruled out. The other class of models includes Starobinsky and Higgs inflation, which both have  $r \sim 0.003$ . A future mission capable of reaching  $\sigma_r \sim \mathcal{O}(10^{-4})$  would provide significant constraints on nearly every currently

favored inflation model. The EPIC-IM configuration is forecasted to achieve **Raphael will update the following to match figures**  $\sigma(r) \sim 4.8 \times 10^{-4}$  assuming  $r = 0.01$  and no foregrounds.

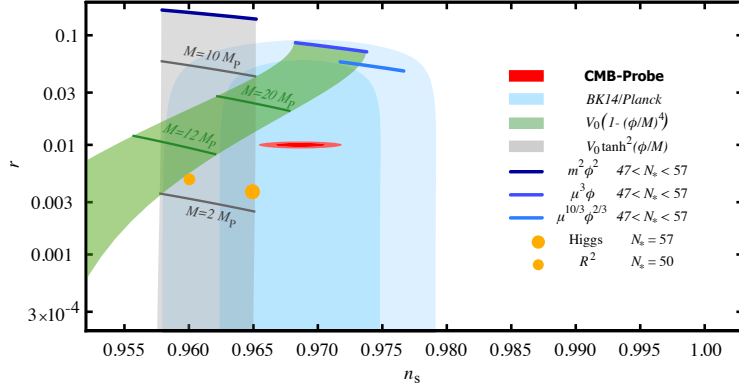


Figure 2: Current 1 and  $2\sigma$  limits on  $r$  and  $n_s$  (blue) and forecasted constraints for a fiducial model with  $r = 0.01$  for EPIC-IM [18]. Also shown are predictions for the models of the inflaton potential discussed in the text: Chaotic inflation for a range of  $N_*$  values (blue lines); Higgs and  $R^2$  (large and small dots, respectively); quartic hilltop (green band); and a sub-class of  $\alpha$ -attractor models [19]

A detection of  $B$ -modes consistent with a primordial spectrum of vacuum fluctuations would be the first observation of a phenomena directly related to quantum gravity. In addition, any detection with a next generation satellite would be evidence for *large-field* Inflation [20], in which a smooth potential that supports Inflation extends over a distance in field space  $\Delta\phi \gtrsim M_p$ . Quantum gravity studies of inflation give a generic expectation  $\Delta\phi \lesssim M_p$  [21, 22, 23, 24], although there are some mechanisms to realize large-field inflation [25, 26, 27, 28]. A detection of  $r$  would therefore provide strong motivation to better understand how large-field inflation can be naturally incorporated into quantum gravity.

All inflation models predict a  $B$ -mode spectrum with the shape shown in Figure 1, but inflation need not be correct [29, 30, 31] and does not preclude additional sources of  $B$ -mode polarization either during or after inflation. To be confident of the implications of a detection, the shape and Gaussianity of the  $B$ -mode spectrum must be characterized. The vast majority of inflation scenarios predict an extremely Gaussian and nearly scale-invariant spectrum for gravitational waves. A target constraint of  $\sigma(n_t) < 1$  at  $r = 0.01$ , driven by the information in the reionization bump, would significantly constrain non-vacuum inflationary sources [32, 33] and rule out physics completely inconsistent with inflation.

Deeper mapping of  $E$ -mode polarization will also contribute to testing inflationary models. Large scale  $E$ -modes will provide new tests of isotropy, a prediction of most models of Inflation; for example, the observations can reject at 99% confidence models in which low multipoles are aligned in the temperature maps [34]. Together with continued improvements at high  $\ell$  from the ground, these modes will also improve constraints on the scalar spectral index and its changes with scale by factors of about two.

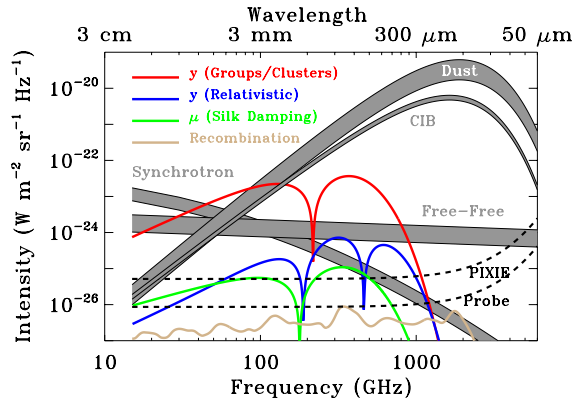


Figure 3: Anticipated  $y$  and  $\mu$  spectral distortions (solid), the signature of resonant recombination lines (solid), and anticipated foreground signal levels relevant for spectral distortion measurements (grey bands). The simplest extension of a proposed Explorer class mission (Probe, dash grey) gives approximately 10 times the Explorer sensitivity (PIXIE). A better optimized Probe may give detections of all anticipated distortions.

Spectral distortion measurements give additional tests of Inflation. The dissipation of small-scale perturbations through Silk-damping leads to  $\mu$ -distortions [35, 36, 37, 38]. In  $\Lambda$ CDM the distortions are predicted at a level of  $\mu = (2.0 \pm 0.14) \times 10^{-8}$ , a level that is readily accessible to a Probe class mission, see Fig. 3 [38, 39].

A better optimized probe may also give the sensitivity to detect the signature of recombination radiation imprinted by cosmological recombination of hydrogen and helium at redshift  $z \simeq 10^3 - 10^4$ ; see Fig. 3 [40, 41]. The detailed physics is sensitive to the values of  $n_s$ , which is a direct probe of Inflation.

### 1.2.2 Light Relics and Dark Matter

After inflation, the universe was reheated to temperatures of at least 10 MeV and perhaps as high as  $10^{10}$  GeV. At these high temperatures, even very weakly interacting or very massive particles, such as those arising in extensions of the Standard model of particle physics, can be produced in large abundances [42, 43]. As the universe expands and cools, the particles fall out of equilibrium and leave observable signatures in the CMB power spectra. Through these effects the CMB is a sensitive probe of neutrino and of other particles' properties.

One particularly compelling target is the effective number of light relic particle species  $N_{\text{eff}}$ , also called the effective number of neutrinos. The canonical value with three neutrino families is  $N_{\text{eff}} = 3.046$ . Additional light particles contribute a change to  $N_{\text{eff}}$  of  $\Delta N_{\text{eff}} \geq 0.027 g$  where  $g \geq 1$  is the number of degrees of freedom of the new particle [44, 45]. This defines a target of  $\sigma(N_{\text{eff}}) < 0.027$  for future CMB observations. Either a limit or detection of  $\Delta N_{\text{eff}}$  at this level would provide a powerful insight into the basic constituents of matter.

Forecasts for  $N_{\text{eff}}$  are shown in Figure 4. The two most important parameters for improving constraints are the fraction of sky observed  $f_{\text{sky}}$  and the noise. Achieving both larger  $f_{\text{sky}}$  and lower noise are strengths of the CMB Probe compared to other platforms. Our baseline mission nearly reaches the target constraint with  $g = 1$ , already exceeding constraints from other astrophysical probes and planned CMB observations. A newly designed mission is likely to reach  $\sigma(N_{\text{eff}}) < 0.027$  with high signal-to-noise ratio.

Many light relics of the early universe are not stable. They decay, leaving faint evidence of their past existence on other tracers. The relics with sufficiently long lifetime to survive few minutes, past the epoch of light element synthesis, leave a signature on the helium fraction  $Y_p$ . If they decay by the time of recombination, their existence through this period is best measured through the ratio of  $N_{\text{eff}}$  to  $Y_p$ . The Probe's cosmic variance limited determination of the  $E$  power spectra will improve current limits for these quantities by a factor of five thus eliminating sub-MeV mass thermal relics. Spectrum distortion measurements give additional constraints on the lifetime and abundance of such relics [47, 48, 49, 50]. A future Probe's  $\mu$ -distortion constraint gives a two orders of magnitude improvement on the abundance and life time of early Universe relics [51, 52] compared to current constraints derived from measurements of light element abundances [53, 54].

Cosmological measurements have already confirmed the existence of one relic that lies beyond the Standard Model: dark matter. For a conventional WIMP candidate, the CMB places very stringent constraints on its properties through the signature of its annihilation on the  $T$  and  $E$  spectra [55, 56, 57]. Planck currently excludes WIMPs with mass  $m_{\text{dm}} < 16$  GeV and a future CMB mission could reach  $m_{\text{dm}} < 45$  GeV for  $f_{\text{sky}} = 0.8$ . The CMB provides the most stringent constraints on the dark matter annihilation cross section for dark matter in this mass range. The CMB is complimentary to direct detection experiments which probe the scattering cross-section of dark matter with Standard model particles.

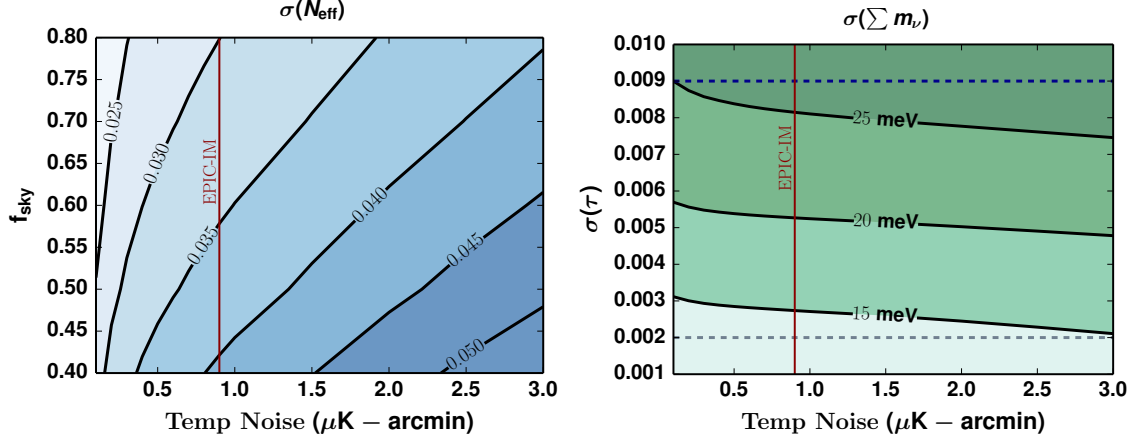


Figure 4:  $N_{\text{eff}}$  as a function of noise and sky fraction (left) and Neutrino mass constraints as a function of uncertainties in measurement of  $\tau$ , noise, and sky fraction of  $f_{\text{sky}} = 0.7$ . The resolution assumed is  $5'$ . Vertical lines denote the expected performance of EPIC-IM. The blue dashed line is the current *Planck* limit; the grey dashed line is the limit from cosmic variance measurement of  $\tau$ . All forecasts assume internal delensing of the  $T$  and  $E$ -maps [46], including residual non-Gaussian covariances. The  $\Sigma m_\nu$  forecasts includes DESI BAO.

A particle-independent approach is to constrain dark matter interactions that would affect the evolution of the effective dark matter fluid and its interactions with baryons or photons. The simplest example is to constrain the baryon-dark matter cross section through its effective coupling of the two fluids [58]. These couplings affect the evolution of fluctuations and ultimately the  $T$  and  $E$  spectra. The current limits of  $\sigma \gtrsim 10^{-31} - 10^{-34} \text{ cm}^2 \times (m_{\text{dm}}/\text{MeV})$  can be competitive with direct detection for sub-GeV masses. More exotic dark sectors that include long-range forces can produce an even richer phenomenology in the CMB and in the large-scale structure without necessarily producing an associated signature in direct detection experiments or indirect searches (e.g. [59, 60, 61]).

Interactions of dark matter with standard model particles can also be constrained through measurements of spectral distortions [62]. Current constraints from FIRAS are most sensitive to small dark matter mass,  $m_\chi \lesssim 0.2 \text{ MeV}$ , but these could be extended to  $m_\chi \lesssim 1 \text{ GeV}$  with a Probe-class mission, testing DM interaction down to cross-sections  $\sigma \simeq 10^{-39} - 10^{-35} \text{ cm}^2$  [62]. This provides new constraints on the low mass end,  $m_\chi \lesssim 10 \text{ MeV}$  and improve existing limits [63, 64] by up to a factor of  $\simeq 50$ . Distortion measurements furthermore open a new avenue for testing dark matter-proton interactions [62].

A host of other physical phenomena including the existence and properties of axions, primordial magnetic fields, and superconducting strings, leave signatures on the spectrum of the CMB and can therefore be constrained by the sensitive measurements of a future Probe [e.g., 65, 66, 67, 68, 69].

### 1.2.3 Neutrino Mass

One of the last unknowns of the Standard model of particle physics is the absolute mass scale of the neutrinos. Cosmology presents a unique opportunity to measure the sum of neutrino masses  $\Sigma m_\nu$  through the suppression of the growth of structures in the Universe on small scales. The sensitivity to  $\Sigma m_\nu$  from suppression of power is limited by our knowledge of the primordial amplitude of fluctuations  $A_s$ , which is strongly degenerate with the optical depth  $\tau$ . The current limit on  $\tau$  from *Planck* of  $\sigma(\tau) = 0.009$  [ ] limits  $\sigma(\Sigma m_\nu) \gtrsim 25 \text{ meV}$ . Forecasts for an internal CMB measurement

of  $\sum m_\nu$  via CMB lensing [70] are shown Figure 4 but the conclusion is the same for any proposed cosmological probe. Therefore, a cosmological detection of the minimum value expected from particle physics  $\sum m_\nu = 58$  meV at more than  $2\sigma$  will require a better measurement of  $\tau$ . The best constraints on  $\tau$  come from  $E$ -modes with  $\ell < 20$  which require measurements over the largest angular scales. To date, the only proven method for such a measurement is from space. The CMB Probe will reach the cosmic variance limit of  $\tau \sim 0.002$  and will therefore reach  $\sigma(\sum m_\nu) < 15$  meV when combined with DESI’s measurements of baryon acoustic oscillations [71]. A detection of  $\sum m_\nu$  at this level is not possible with any other existing survey.

#### 1.2.4 Cosmological structure formation

Understanding the evolution of cosmological structures from small density perturbations through the formation of the first stars to present day galaxies and cluster is a key goal of cosmology [72]. Cosmological reionization, the transition of the Universe from dominated by neutral to ionized hydrogen, is a cornerstone of this evolution because it encodes information about the star formation history and the physical processes that formed the galaxies of various luminosities and masses we see today. But when did the epoch of reionization start? How long did it last? Are early galaxies enough to reionize the entire Universe or is another source required?

Measurements of the CMB  $E$  mode power spectrum over large angular scales are sensitive to the optical depth to reionization  $\tau$ , a key parameter for all reionization models that attempt to answer these questions. The *Planck* team reported recently a value of  $\tau = 0.055 \pm 0.009$  [? ? ]. The level is significantly lower than previous estimates and reduces the tension between CMB-based analyses and constraints from other astrophysical sources. The CMB Probe’s cosmic variance limited measurement of  $E$ -mode polarization will improve the  $1\sigma$  error by a factor of 4.5 to reach a cosmic variance limited measurement of  $\tau$ , thus setting stringent constraints on models of the reionization epoch.

The anisotropy in the cosmic infrared background (CIB) produced by dusty star-forming galaxies in a wide redshift range, are an excellent probe of both the history of star formation and the link between galaxies and dark matter across cosmic time. The *Planck* collaboration derived values of the star formation rate that, at redshifts  $z \sim 3$ , are three times larger than constraints from number counts measurements ([73, 74, 75]). The new mission probe, By measuring CIB anisotropy with 100 times higher signal-to-noise ratio the CMB Probe will shed light on this intriguing discrepancy. Specifically, it will constrain the star formation rate with one tenth of *Planck*’s uncertainty.

A key parameter in simulations of the angular power spectrum of the CIB is  $M_{\text{eff}}$ , the galaxy halo mass that is most efficient in producing star formation activity. Comparing measurements of the power spectrum to simulations constrains this parameter, which informs structure formation models. Current models and measurements find  $M_{\text{eff}} \sim 10^{12}$  solar masses with about 10% uncertainty. The CMB Probe will constrain this parameter at the percent level.

The transition to reionized Universe and the onset of structure formation inject energy into the sea of CMB photons. This injection is detectable through a distinct spectral distortion. This is the largest expected distortion – marked ‘ $y$  Groups/Clusters’ in Figure 3 – and will be clearly detected by the CMB Probe. A detection will give information about the total energy output of the first stars, AGNs, and galaxy clusters, an important parameter in structure formation models.

Group-size clusters that have masses  $M \simeq 10^{13} M_\odot$  contribute significantly to the signal. With temperature  $kT_e \simeq 1$  keV these are sufficiently hot to create a relativistic temperature correction to the large  $y$ -distortion. This relativistic correction, denoted ‘ $y$  relativistic’ in Figure 3, will also be detected with high signal-to-noise ratio by the CMB Probe, and will be used to constrain the

currently uncertain feedback mechanisms used in hydrodynamical simulations of cosmic structure formation [76].

The CMB spectrum varies spatially across the sky. One source of such anisotropic distortion is due to the spatial distribution clusters of galaxies and has already been measured by Planck [77]. A combination of precise CMB imaging and spectroscopic measurements will allow observing the relativistic temperature correction of individual SZ clusters [78, 79, 80], which will calibrate cluster scaling relations and inform our knowledge of the dynamical state of the cluster atmosphere.

Resonant scattering of the CMB photons during and post last scattering leads to spectral-spatial signals that can be used to constrain the abundance of metals in the dark ages and therefore the make-up of the first, and subsequent generations of stars [81, 82, 83, 84, 85].

### 1.3 The Challenges: Foregrounds and Systematics

A tentative target for the CMB Probe is to constrain the tensor to scalar ratio with an uncertainty that is a factor of 100 smaller than the current upper limit  $r < 0.07$ , that is, to reach  $\sigma(r) < 0.0007$ . Data from *Planck* and sub-orbital experiments reveal that foregrounds already dominate the signal. A 100-fold improvement on the final error will require exquisite measurement and accounting for foreground sources of confusion.

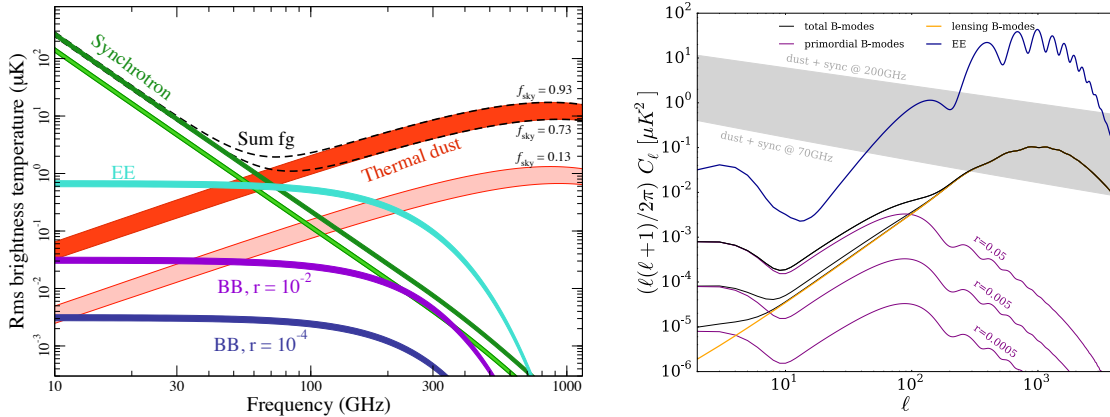


Figure 5: *Left:* Brightness temperature as function of frequency for the polarized CMB (cyan, purple, blue) and galactic foreground signals: dust (red) and synchrotron (green). The darker bands correspond to sky fractions between 73% and 93%; the lighter bands to the cleanest 13%, with the width indicating the uncertainty. *Right:* Angular power spectrum for B-mode polarization several  $r$  values, and for foreground emission between 70 and 200 GHz.

to ascertain that the Probe's uncertainty on the measurement of  $r$  is dominated by statistic rather than systematic error, the mission's design, execution, and data analysis will have to be dominated by the need to controlling systematic uncertainties to an unprecedented, and certainly not yet achieved, levels.

#### 1.3.1 Foregrounds

Whereas the CMB temperature anisotropy signal dominates galactic sources of emission, this is not the case for polarization. Figure 5 compares the expected RMS brightness temperature of polarized emission from galactic sources to the level of  $E$  and  $B$  mode signals as a function of frequency. It also gives the expected signal levels as a function of angular scale  $\ell$ . The conclusions, some of which are also borne out by *Planck* measurements, are that:

- over the largest angular scales (lowest  $\ell$ s), where the inflationary  $B$  mode signal is comparable or larger than that from lensing, foreground sources of confusion will need to be measured and



subtracted to a level better than 1 part in 100 in temperature;

- foregrounds dominate the inflationary  $B$  mode signal on *all* angular scales by an order of magnitude or more.

Known signals can be accounted for and removed, even if their amplitude is large. But the best measurements to date, from *Planck*, fall far shorter than needed for the fidelity envisioned for the Probe. This is visually demonstrated by Figure 6, which compares the level of  $B$  mode signal at low  $\ell$  for  $r = 0.001$  to the *Planck* noise, extrapolated to 150 GHz, a frequency band in which the signal is among the strongest.

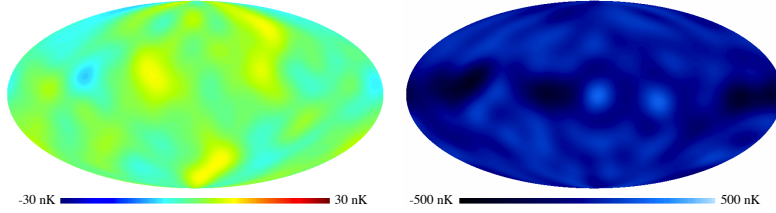


Figure 6: *Left:* Stokes  $Q$  for inflationary  $B$ -modes for  $\ell < 12$  and  $r = 0.001$ . *Right:* Noise in the *Planck* 353 GHz map of Stokes  $Q$  for  $\ell < 12$  extrapolated to 150 GHz assuming the sky average spectral properties of dust.

The extrapolation assumes the properties of dust as measured by *Planck* on average across the sky. We now know that these average properties are, as expected from fundamental physics, only approximate. **say something about correlations, spatial distribution, variations in the spectral index.**

While the search for primordial B-modes leads to the strictest constraints on foreground residuals, exquisite control of foregrounds is also necessary for the other science objectives. The critical measurements of  $\tau$ , available by a cosmic variance limited measurement of the  $E$  power spectrum, are buried below the foregrounds at  $\ell < 10$ . **check what else?** Recovering *all* of the spectral distortions signals from the raw data will require proper accounting for emission by dust grains, synchrotron emission from electrons spiraling in the galactic magnetic field, and coulomb scattering of charged particles (‘free-free’); see Figure 3.

### 1.3.2 Systematic Errors

The latest experience with *Planck* points to the following systematic error categories likely to be important for the CMB Probe, or for that matter, for any instrument striving to map the polarization over large portions of the sky to the levels targeted by the CMB Probe [? ]: 1. Intensity-to-polarization leakage, 2. stability, and 3. straylight. Each of these is considered in light of polarimetry measurements through differencing the signals of two detectors that are sensitive to orthogonal polarization states.

**Leakage** The CMB anisotropy signal is a factor of 1000 larger than the strongest possible inflationary B-mode signal (see for example Fig. 1). Therefore instrumental effects that can leak even a small fraction of an intensity fluctuation into spurious polarization must be understood and controlled. The main effects are differences between gains of detectors, their frequency bandpass mismatch, their differential pointing on the sky, and their differential antenna patterns. These differential effects need to be controlled, through instrument design, characterization, and data analysis to 1 part in  $10^4$  to give a negligible contribution for  $\sigma(r) < 0.001$ . This level of control represents a factor of 100 improvement on *Planck*’s performance [? ]. **what about Bicep/Keck?**

Leakage-related effects will drive: requirements on the optical system, and the uniformity of the bandpass of each polarimeter; calibration requirements on the level of cross-polar leakage and its angle; and measurements of the the beam shape as a function of source spectrum. These systematic effects can potentially be mitigated by modulation of the sky signal in such a way that

allows complete reconstruction of the polarized sky signal using each photometer, for example, using a half-wave plate.

**Stability.** need more number(s) in this paragraph The reconstruction of deep, full sky polarization maps involves a combination of measurements made at times separated by months, requiring stability of the response of the instrument on corresponding time scales. Random deviations from stability are a source of noise; systematic deviations are a source of systematic error. This type of systematic error puts requirements on control of thermal drifts of spacecraft temperatures, to mitigate thermal emissivity changes and thermoelastic deformation of telescope structures. The cryogenic operating temperatures of detectors or reference calibration loads must be controlled adequately as well. Careful design of the scan strategy can shorten the time scales needed for stringent stability, for example Planck’s scan strategy traced out great circles which overlapped on 1 minute timescales, giving a shorter effective time scale for stability requirements.

The spacecraft’s ambient radiation environment is modulated by the solar activity and can introduce temperature drifts in the cryogenic stages as well as introducing correlated transients in detectors and readout electronics. The design of the instrument must account for these effects.

**Straylight.** When the brightest sources in the sky – the Sun, Moon, planets, and Galaxy – are passing through the far sidelobes of the telescope they create a spurious polarization signal. If they are passing in repeated, scan synchronous pattern, the spurious signal becomes a source of systematic error. This far sidelobe response can be reduced through careful optical design and baffling, but will always be present at a non-trivial level. Detailed modeling of the *Planck* telescope, convolved with sky sources, gave a predicted sidelobe contamination at a detectable level of tens of micro-Kelvin in the 30 GHz maps. This contamination has been observed in *Planck* difference maps. As a result an estimate of the sidelobe contamination was removed from some of the *Planck* time ordered data as part of the mapmaking process. The more stringent requirements for CMB-probe will necessitate at least this level of mitigation.

**Study Plan.** Understanding and controlling the effects of systematic errors in a next-generation CMB probe is critical. During the study period we plan to build an end-to-end simulation pipeline that will generate simulated measurements at an instrumental level, and feed them into the notional analysis pipeline, including foreground/CMB component separation and power spectral analysis. This simulation pipeline will allow us to explore 1. mitigation of systematic errors by design, for example exploring modulation schemes and modulator technologies, and 2. mitigation of systematic errors by analysis techniques. This pipeline will be used eventually to define requirements for a notional mission, and could be helpful to prioritize the technology development needs of a future mission.

## **1.4 Current, Forthcoming, and Proposed Efforts and the CMB Probe**

### **1.4.1 Current and Forthcoming Sub-Orbital Efforts**

The remarkable forthcoming scientific yield has motivated significant agency investments in current and future sub-orbital experiments which are designed to realize the full potential of this unique probe of fundamental physics and astrophysics. These experiments are designed to exploit the comparative advantages of the sub-orbital platforms, while providing the design heritage and experience necessary to maximize the probability of success of an orbital mission.

For the ground-based efforts, these include combinations of *i*) provision for large apertures and therefore high angular resolution, *ii*) extremely long observation times, *iii*) flexibility to rapidly deploy new technologies, and *iv*) allowance for detector formats that are relatively unconstrained

by mass and power limitations. In aggregate, these have demonstrated low noise measurements of small and intermediate angular scale  $E$  polarization structures over areas of up to 10% of the sky (check). The largest sky fraction over which  $B$  mode has been measured to date is ??%.

The balloon-borne missions *i)* extend the frequency reach of the ground based telescopes, *ii)* enable high fidelity measurements on larger angular scales than can be probed from the ground, and *iii)* grant inexpensive access to a space-like environment, providing heritage for future space missions as well as experience in dealing with the analysis of data that are representative of a space mission. In this way, the sub-orbital programs complement and multiply the scientific return of the proposed orbital mission, while reinforcing its technical preparedness.

The 2010 Decadal Panel strongly recommended supporting sub-orbital efforts in preparation for a possible space mission to follow sub-orbital detections of inflationary gravity waves. As a result, the US has clear leadership in the field, both in terms of ground- and balloon-based experiments and results.

This leadership will continue into the foreseeable future. In aggregate, funded, now-being-built ‘Stage 3’ CMB experiments will deploy approximately 50,000 detectors on various sub-orbital experiments within the next 3-5 years. Ground-based experiments plan to extend measurements from few percent of the sky to few tens, although in a limited frequency range between 30 and 300 GHz. Balloon-borne payloads operating at even higher frequencies strive to cover even larger fractions.

#### 1.4.2 Proposed Efforts: LiteBIRD, CORE, and CMB-S4

Japan, in collaboration with NASA, is now considering whether to proceed with LiteBIRD, a space mission designed to search for  $B$ -modes from inflation. The US Team has submitted its Phase A report to NASA; Phase A in Japan will conclude in about a year (check). LiteBIRD is a smaller, more focused mission compared to the CMB Probe. It is an imager based on a 0.5 m aperture telescope. Therefore it has a resolution 4 times lower compared to the 2 m aperture EPIC-IM. Its reach in  $\ell$  space is correspondingly 4 times lower making the science available at  $\ell$ ’s above few hundred in both  $E$  and  $B$  modes unreachable. It has no spectroscopic capabilities and thus not sensitive to any of the spectral distortion science goals.

For the Japanese space agency JAXA, LiteBIRD is meant to fit within a cap of \$350M (check). Although there are uncertainties about comparing JAXA’s cost calculations to NASA’s, LiteBIRD’s overall size and more limited science reach is commensurate with it being below, or just at the lower margin of the Probe financial scope.

A collaboration of scientists in Europe has just recently proposed CORE to ESA as part of the M5 round of space mission proposals. The team includes a number of US collaborators; the PI of this proposal is a member of CORE’s executive board (check). CORE is a CMB polarization imager that is based on a 1.2 m aperture telescope and thus intended to reach 2.5 times the resolution of LiteBIRD. ESA has capped the M5 proposals to EU550M, the equivalent of \$610M. Member countries are expected to contribute an additional ~\$100M making the total cost close to \$700M. Selection of missions for Phase A studies is expected in fall 2017, and end of Phase A selection in fall 2019.

The US CMB community has proposed, and the Particle Physics Project Prioritization Panel (P5) has recommended to the DOE, the establishment of a 4th generation CMB experiment called CMB-S4. This is an ambitious program to field 10 times the number of detectors fielded by Stage 3 experiments. If and when funded CMB-S4 will enable unprecedented sensitivity at frequency bands accessible from the ground, and with telescopes that enable high resolution.

### 1.4.3 Why Study a CMB Probe?

Learning from the successes of COBE/FIRAS, COBE/DMR, WMAP, and *Planck*, a CMB Probe is the single most suitable vehicle to deliver complete sky coverage and therefore information on the largest angular scales, comprehensive frequency coverage, and exquisite control of systematic effects. The mission will provide a single self-consistent and self-calibrated data set. It will provide legacy maps at many frequency bands that will become the basis for hundreds of new contributions to science. We should consider launching this mission to reach all the science goals described in Section ???. Some of the science goals are not reachable from ground-based observatories, or make take many years to achieve. If the Inflationary signal is detected by sub-orbital experiments any time soon, a space mission to characterize the signal in full detail is equally compelling. The existence of ambitious sub-orbital programs is a complementary strength. How to make the best use of this complementarity is an explicit goal of our study; see Sectionsec:management.

### 1.4.4 Does the CMB Probe Fit Within the Cost Window?

When EPIC-IM

## 1.5 State of Technologies

2 pages. Discuss the technologies, their TRL, and what will be studied

A fourth generation CMB satellite targeting a map sensitivity of  $\sim 1\mu\text{k} - \text{arcmin}$  will require, extremely sensitive detector arrays, tight control over systematics, and ability to reject polarized foregrounds as is described in Section 1.2. Given the frequency dependance of synchrotron and dust foregrounds, this last requirement translates into the need for a large number of spectral bands covering the approximate frequency range from  $\sim 30$  GHz to  $\sim 800$  GHz. Development of the CMB technologies needed to meet these requirements is actively being pursued by many groups who are also demonstrating these technologies on ground, balloon, and satellite platforms. We describe the status and needs in the areas of telescopes, optics, detector coupling, detectors, and readout.

**Telescopes:** Carbon Fiber Reinforced Polymer (CFRP) mirrors are at TRL 9 as they have flown on the Planck satellite. Their 1.9x1.5 m mirror weighted only 28 lbs and met all surface quality requirements. However, small deformations in the mirror caused by its structural supports had a measurable impact to the beam far-sidelobes that was not captured by preflight measurements or the corresponding beam modeling. Future CMB satellites will require improved pre-flight characterization of *polarized beam* at operating temperature augmented by improved simulation tools to meet even more systematics requirements. Current ground and balloon born optical designs achieve large field of views (FOVs) with reflective and refractive designs; related designs and their implementation should be studied in the context of a satellite mission as the sensitivity requirements lead to the need to maximize the size of the FOV while fitting within the tight mass and size constraints imposed by a space mission. Given the heritage of past satellite missions it will be possible to develop a telescope design that meets the requirements for a future mission and uses high TRL components.

**Optical Coupling:** The need for sensitivity drives the push for high efficiency optics; wide bandwidth to complement multichroic detector; infrared filters to maximize cryogenic performance; and polarization modulators to suppress  $1/f$  noise and mitigate instrument systematics. The CMB field has made tremendous progress recently by drawing on advances in materials, processing

techniques, and developments in electrical engineering including meta-material research. Single crystals such as silicon and sapphire are attractive since they offer extremely low dielectric losses and high indices of refraction to better manipulate light. New coating techniques have been developed for silicon and sapphire that span 2:1 bandwidth (TRL 5+ for silicon) and can realize up to 5:1 bandwidth. EBEX deployed broadband cryogenic polarization modulator with a superconducting bearing that covered 150 GHz band to 410 GHz band raising the modulators to TRL 5+ for space. Meta-material metal-mesh optical filters were deployed with the Planck satellite and they are extensively used by ground based and balloon experiments making these TRL6 optical elements. It is necessary to develop a plan for a satellite mission that will cover  $\sim 30$  GHz to  $\sim 800$  GHz. Two configurations could be considered: multiple optical paths with  $< 3 : 1$  bandwidth and a potentially simpler design with only two optical paths with  $\sim 5 : 1$  bandwidth. These studies include evaluating the design tradeoffs inherent to these approaches, developing the new coatings needed, and evaluating the promise of hybrid approaches where filters and lenses are implemented in the same optical elements. In addition, the cryogenic rotation mechanism should be demonstrated at the robustness (eg lifetime) needed for for a satellite mission.

**Detector Coupling:** The focal-plane feed determines the shape and polarization properties of the pixel beams and therefore plays a strong role in controlling systematic errors. The feed design also can determine the total bandwidth and number of photometric bands of each pixel which is important for the efficient use of a telescope's focal plane area. CMB experiments developed broadband multi-chroic *detector* to increase optical throughput of a focal plane. Broadband feed captures signal over wide frequency range. Then on-chip superconducting filter partitions signal into multiple frequency bands prior to detection. Broadband detectors were realized with spline profiled horn and lenslet coupled antenna. Broadband horn detector deployed a pixel that covers 2.3:1 bandwidth with on going development for extending bandwidth to 6:1. Broadband lenslet coupled antenna will deploy 3:1 bandwidth detector this year. Lenslet coupled antenna demonstrated 5:1 bandwidth in laboratory. RF-techniques to partition broadband signals into multiple band are mature. For a future CMB polarization satellite mission, broadband feed should be demonstrated at high frequency where alignment and line width for micro-fabrication becomes challenging. Detectors for CMB satellite mission were hand picked one by one for optimal performance. Next generation of detector array will be fabricated on a silicon wafer. Micro-fabrication process should demonstrate high yield and uniformity across a wafer that can meet tight requirement of satellite mission. Also detector test need to able to characterize detector beyond the level of systematic required by next generation CMB satellite experiment.

The Planck HFI deployed Neutron Transmutation Doped Germanium high-resistance bolometer at 100 milli-Kelvin to achieve photon noise limited detector performance. A Transition Edge Sensor (TES) bolometer uses a steep transition of superconducting metal to improve linearity of the detector. TES bolometers have been deployed on 100 milli-Kelvin and 250 milli-Kelvin platform. TES bolometers have been deployed across ground based and balloon CMB experiments spanning 40 GHz-410 GHz with detectors achieving NEPs of  $20\text{-}50 \text{ aW}/\sqrt{\text{Hz}}$ , nearly background limited at CMB frequencies. TES bolometers deployed at low optical frequencies ( $\sim 40$  GHz) and balloon-borne payloads should realized even lower NEPs of  $\sim 10 \text{ aW}/\sqrt{\text{Hz}}$ . Emerging detector technology for CMB experiment is kinetic-inductance detector (KIDs). The KIDs detector detects signal as change in kinetic inductance. KIDs detectors can be frequency multiplexed easily to  $\sim 1,000$  detectors. Recently, on-sky demonstration at 150 GHz and 230 GHz was done with lumped ele-

ment KID detector. Noise performance of KID detector at low frequency channels ( $< 40$  GHz) need some improvement to be photon-noise limited. Currently there is no CMB polarization power spectrum data produced with KID detector. Coupling between RF (100 GHz) signal to microwave KID (MKID) detector is in a development stage. Planck detectors experienced unexpectedly high rate of cosmic ray events. Data was successfully cleaned with analysis technique. Study of impact of cosmic rays on a detector is crucial for next CMB satellite mission.

Multiplexed readout is being used by CMB experiments to readout thousands of TES bolometers, and readout multiplexing is built into KID detector architecture. Voltage bias and low impedance of a TES bolometer facilitates multiplexing readout by Superconducting Quantum Interference Device (SQUID). Time domain multiplexing uses a SQUID at milli-Kelvin as a switch to rapidly cycle through bolometers. Highest achieved multiplexing factor is 64 channels. Frequency domain multiplexing uses superconducting resonators to assign bolometers to different frequency channels. Highest achieved multiplexing factor is 68 channels. New readout scheme, such as microwave SQUID readout, is emerging to increase multiplexing factor for TES bolometer. MKID detector architecture has multiple resonators coming off from a transmission line. A resonator is both a detector and multiplexer. MKID demonstrated multiplexing factor that exceeds 1,000 channels. For next generation satellite experiment that will readout over thousands detectors require high multiplexing factor. Multiplexing factor is directly related to readout complexity and power consumption. Also the Planck mission experienced ADC non-linearity, thus extensive characterization of end to end readout architecture should be performed pre-flight.

A future CMB satellite mission offers exciting opportunity for millimeter wave polarization science. Experience from Planck mission will be studied to learn lessons for the future mission. Development for CMB instrumentation is an active field with many institutions developing new technologies for ground based, balloon, and proposed satellite missions. For a satellite instrumentation, there is a difficult trade off between desire to have high performance instrument and desire to keep cost manageable. Many developments that is going on for ground based and balloon experiment have similar goal as satellite mission that collaborative development across all platform will be beneficial.

## **1.6 Mission Study and Management Plan**

Being Written.

## References

- [1] J. Bock, A. Aljabri, A. Amblard, D. Baumann, M. Betoule, T. Chui, L. Colombo, A. Cooray, D. Crumb, P. Day, C. Dickinson, D. Dowell, M. Dragovan, S. Golwala, K. Gorski, S. Hanany, W. Holmes, K. Irwin, B. Johnson, B. Keating, C.-L. Kuo, A. Lee, A. Lange, C. Lawrence, S. Meyer, N. Miller, H. Nguyen, E. Pierpaoli, N. Ponthieu, J.-L. Puget, J. Raab, P. Richards, C. Satter, M. Seiffert, M. Shimon, H. Tran, B. Williams, and J. Zmuidzinas. Study of the Experimental Probe of Inflationary Cosmology (EPIC)-Intermediate Mission for NASA's Einstein Inflation Probe. *ArXiv e-prints*, June 2009.
- [2] A. Kogut, D. J. Fixsen, D. T. Chuss, J. Dotson, E. Dwek, M. Halpern, G. F. Hinshaw, S. M. Meyer, S. H. Moseley, M. D. Seiffert, D. N. Spergel, and E. J. Wollack. The Primordial Inflation Explorer (PIXIE): a nulling polarimeter for cosmic microwave background observations. *JCAP*, 7:25–+, July 2011.
- [3] J. C. Mather, E. S. Cheng, D. A. Cottingham, R. E. Eplee, Jr., D. J. Fixsen, T. Hewagama, R. B. Isaacman, K. A. Jensen, S. S. Meyer, P. D. Noerdlinger, S. M. Read, and L. P. Rosen. Measurement of the cosmic microwave background spectrum by the COBE FIRAS instrument. *Ap. J.*, 420:439–444, January 1994.
- [4] D. J. Fixsen, E. S. Cheng, J. M. Gales, J. C. Mather, R. A. Shafer, and E. L. Wright. The Cosmic Microwave Background Spectrum from the Full COBE FIRAS Data Set. *Ap. J.*, 473:576–+, December 1996.
- [5] Y. B. Zeldovich and R. A. Sunyaev. The Interaction of Matter and Radiation in a Hot-Model Universe. *ApSS*, 4:301–316, July 1969.
- [6] R. A. Sunyaev and Y. B. Zeldovich. The interaction of matter and radiation in the hot model of the Universe, II. *ApSS*, 7:20–30, April 1970.
- [7] R. Keisler, C. L. Reichardt, K. A. Aird, B. A. Benson, L. E. Bleem, J. E. Carlstrom, C. L. Chang, H. M. Cho, T. M. Crawford, A. T. Crites, T. de Haan, M. A. Dobbs, J. Dudley, and E. M. George. A Measurement of the Damping Tail of the Cosmic Microwave Background Power Spectrum with the South Pole Telescope. *Ap. J.*, 743:28, December 2011.
- [8] <http://bolo.berkeley.edu/polarbear/>.
- [9] A. H. Guth. Inflationary universe: A possible solution to the horizon and flatness problems. *Phys. Rev. D.*, 23:347–356, January 1981.
- [10] A. D. Linde. A New Inflationary Universe Scenario: A Possible Solution of the Horizon, Flatness, Homogeneity, Isotropy and Primordial Monopole Problems. *Phys. Lett.*, B108:389–393, 1982.
- [11] A. Albrecht and P. J. Steinhardt. Cosmology for grand unified theories with radiatively induced symmetry breaking. *Phys. Rev. Lett.*, 48:1220–1223, 1982.
- [12] K. Sato. First-order phase transition of a vacuum and the expansion of the Universe. *MNRAS*, 195:467–479, May 1981.



- [13] E. W. Kolb and M. S. Turner. *The Early Universe*. Addison-Wesley, Redwood City, CA, 2nd. edition, 1994.
- [14] M. Kamionkowski, A. Kosowsky, and A. Stebbins. A Probe of Primordial Gravity Waves and Vorticity. *Phys. Rev. Lett.*, 78:2058–2061, March 1997. astro-ph/9609132.
- [15] M. Zaldarriaga and U. Seljak. All-sky analysis of polarization in the microwave background. *Phys. Rev. D.*, 55:1830–1840, 1997.
- [16] Hayden Lee, S. C. Su, and Daniel Baumann. The Superhorizon Test of Future B-mode Experiments. *JCAP*, 1502(02):036, 2015.
- [17] Committee for a Decadal Survey of Astronomy and Astrophysics. *New Worlds, New Horizons in Astronomy and Astrophysics*. National Academy Press, 2010.
- [18] P. A. R. Ade et al. Improved Constraints on Cosmology and Foregrounds from BICEP2 and Keck Array Cosmic Microwave Background Data with Inclusion of 95 GHz Band. *Phys. Rev. Lett.*, 116:031302, 2016.
- [19] Renata Kallosh and Andrei Linde. Universality Class in Conformal Inflation. *JCAP*, 1307:002, 2013.
- [20] David H. Lyth. What would we learn by detecting a gravitational wave signal in the cosmic microwave background anisotropy? *Phys.Rev.Lett.*, 78:1861–1863, 1997.
- [21] Tom Banks, Michael Dine, Patrick J. Fox, and Elie Gorbatov. On the possibility of large axion decay constants. *JCAP*, 0306:001, 2003.
- [22] Daniel Baumann and Liam McAllister. *Inflation and String Theory*. Cambridge University Press, 2015.
- [23] Jon Brown, William Cottrell, Gary Shiu, and Pablo Soler. Fencing in the Swampland: Quantum Gravity Constraints on Large Field Inflation. *JHEP*, 10:023, 2015.
- [24] Tom Rudelius. Constraints on Axion Inflation from the Weak Gravity Conjecture. *JCAP*, 1509(09):020, 2015.
- [25] Eva Silverstein and Alexander Westphal. Monodromy in the CMB: Gravity Waves and String Inflation. *Phys.Rev.*, D78:106003, 2008.
- [26] Nemanja Kaloper and Lorenzo Sorbo. A Natural Framework for Chaotic Inflation. *Phys. Rev. Lett.*, 102:121301, 2009.
- [27] Fernando Marchesano, Gary Shiu, and Angel M. Uranga. F-term Axion Monodromy Inflation. *JHEP*, 09:184, 2014.
- [28] Ralph Blumenhagen, Cesar Damian, Anamaria Font, Daniela Herschmann, and Rui Sun. The Flux-Scaling Scenario: De Sitter Uplift and Axion Inflation. 2015.
- [29] Justin Khoury, Burt A. Ovrut, Nathan Seiberg, Paul J. Steinhardt, and Neil Turok. From big crunch to big bang. *Phys. Rev.*, D65:086007, 2002.



- [30] Robert H. Brandenberger. The Matter Bounce Alternative to Inflationary Cosmology. 2012.
- [31] Anna Ijjas and Paul J. Steinhardt. Implications of Planck2015 for inflationary, ekpyrotic and anamorphic bouncing cosmologies. *Class. Quant. Grav.*, 33(4):044001, 2016.
- [32] Ryo Namba, Marco Peloso, Maresuke Shiraishi, Lorenzo Sorbo, and Caner Unal. Scale-dependent gravitational waves from a rolling axion. *JCAP*, 1601(01):041, 2016.
- [33] Marco Peloso, Lorenzo Sorbo, and Caner Unal. Rolling axions during inflation: perturbativity and signatures. 2016.
- [34] Cora Dvorkin, Hiranya V. Peiris, and Wayne Hu. Testable polarization predictions for models of CMB isotropy anomalies. *Phys. Rev.*, D77:063008, 2008.
- [35] R. A. Sunyaev and Y. B. Zeldovich. Small scale entropy and adiabatic density perturbations - Antimatter in the Universe. *ApSS*, 9:368–382, December 1970.
- [36] R. A. Daly. Spectral distortions of the microwave background radiation resulting from the damping of pressure waves. *Ap. J.*, 371:14–28, April 1991.
- [37] W. Hu, D. Scott, and J. Silk. Power spectrum constraints from spectral distortions in the cosmic microwave background. *Ap. J. Lett.*, 430:L5–L8, July 1994.
- [38] J. Chluba, R. Khatri, and R. A. Sunyaev. CMB at 2 x 2 order: the dissipation of primordial acoustic waves and the observable part of the associated energy release. *MNRAS*, 425:1129–1169, September 2012.
- [39] J. Chluba. Which spectral distortions does  $\Lambda$ CDM actually predict? *MNRAS*, 460:227–239, July 2016.
- [40] R. A. Sunyaev and J. Chluba. Signals from the epoch of cosmological recombination (Karl Schwarzschild Award Lecture 2008). *Astronomische Nachrichten*, 330:657–+, 2009.
- [41] J. Chluba and Y. Ali-Haïmoud. COSMOSPEC: fast and detailed computation of the cosmological recombination radiation from hydrogen and helium. *MNRAS*, 456:3494–3508, March 2016.
- [42] G. Steigman. Cosmology confronts particle physics. *Annual Review of Nuclear and Particle Science*, 29:313–338, 1979.
- [43] M. Bolz, A. Brandenburg, and W. Buchmuller. Thermal production of gravitinos. *Nucl. Phys.*, B606:518–544, 2001. [Erratum: *Nucl. Phys.*B790,336(2008)].
- [44] Christopher Brust, David E. Kaplan, and Matthew T. Walters. New Light Species and the CMB. *JHEP*, 12:058, 2013.
- [45] Daniel Baumann, Daniel Green, and Benjamin Wallisch. A New Target for Cosmic Axion Searches. *Phys. Rev. Lett.*, 117(17):171301, 2016.
- [46] Daniel Green, Joel Meyers, and Alexander van Engelen. CMB Delensing Beyond the B Modes. 2016.

- [47] S. Sarkar and A. M. Cooper. Cosmological and experimental constraints on the tau neutrino. *Physics Letters B*, 148:347–354, November 1984.
- [48] M. Kawasaki and K. Sato. The effect of radiative decay of massive particles on the spectrum of the microwave background radiation. *Physics Letters B*, 169:280–284, March 1986.
- [49] W. Hu and J. Silk. Thermalization constraints and spectral distortions for massive unstable relic particles. *Physical Review Letters*, 70:2661–2664, May 1993.
- [50] J. Chluba and R. A. Sunyaev. The evolution of CMB spectral distortions in the early Universe. *MNRAS*, 419:1294–1314, January 2012.
- [51] J. Chluba. Distinguishing different scenarios of early energy release with spectral distortions of the cosmic microwave background. *MNRAS*, 436:2232–2243, December 2013.
- [52] J. Chluba and D. Jeong. Teasing bits of information out of the CMB energy spectrum. *MNRAS*, 438:2065–2082, March 2014.
- [53] M. Kawasaki, K. Kohri, and T. Moroi. Big-bang nucleosynthesis and hadronic decay of long-lived massive particles. *Phys. Rev. D.*, 71(8):083502, April 2005.
- [54] K. Jedamzik. Big bang nucleosynthesis constraints on hadronically and electromagnetically decaying relic neutral particles. *Phys. Rev. D.*, 74(10):103509, November 2006.
- [55] P. J. E. Peebles, S. Seager, and W. Hu. Delayed Recombination. *Ap. J. Lett.*, 539:L1–L4, August 2000.
- [56] X. Chen and M. Kamionkowski. Particle decays during the cosmic dark ages. *Phys. Rev. D.*, 70(4):043502–+, August 2004.
- [57] N. Padmanabhan and D. P. Finkbeiner. Detecting dark matter annihilation with CMB polarization: Signatures and experimental prospects. *Phys. Rev. D.*, 72(2):023508–+, July 2005.
- [58] Cora Dvorkin, Kfir Blum, and Marc Kamionkowski. Constraining Dark Matter-Baryon Scattering with Linear Cosmology. *Phys. Rev.*, D89(2):023519, 2014.
- [59] Francis-Yan Cyr-Racine, Roland de Putter, Alvise Raccanelli, and Kris Sigurdson. Constraints on Large-Scale Dark Acoustic Oscillations from Cosmology. *Phys. Rev.*, D89(6):063517, 2014.
- [60] Manuel A. Buen-Abad, Gustavo Marques-Tavares, and Martin Schmaltz. Non-Abelian dark matter and dark radiation. *Phys. Rev.*, D92(2):023531, 2015.
- [61] Julien Lesgourgues, Gustavo Marques-Tavares, and Martin Schmaltz. Evidence for dark matter interactions in cosmological precision data? *JCAP*, 1602(02):037, 2016.
- [62] Y. Ali-Haïmoud, J. Chluba, and M. Kamionkowski. Constraints on Dark Matter Interactions with Standard Model Particles from Cosmic Microwave Background Spectral Distortions. *Physical Review Letters*, 115(7):071304, August 2015.

- [63] R. Essig, A. Manalaysay, J. Mardon, P. Sorensen, and T. Volansky. First Direct Detection Limits on Sub-GeV Dark Matter from XENON10. *Physical Review Letters*, 109(2):021301, July 2012.
- [64] C. Bøhm, J. A. Schewtschenko, R. J. Wilkinson, C. M. Baugh, and S. Pascoli. Using the Milky Way satellites to study interactions between cold dark matter and radiation. *MNRAS*, 445:L31–L35, November 2014.
- [65] K. Jedamzik, V. Katalinić, and A. V. Olinto. Limit on Primordial Small-Scale Magnetic Fields from Cosmic Microwave Background Distortions. *PRL*, 85:700–703, July 2000.
- [66] H. Tashiro, E. Sabancilar, and T. Vachaspati. CMB distortions from superconducting cosmic strings. *Phys. Rev. D.*, 85(10):103522, May 2012.
- [67] A. D. Dolgov and D. Ejlli. Resonant high energy graviton to photon conversion at the post-recombination epoch. *Phys. Rev. D.*, 87(10):104007, May 2013.
- [68] H. Tashiro, J. Silk, and D. J. E. Marsh. Constraints on primordial magnetic fields from CMB distortions in the axiverse. *Phys. Rev. D.*, 88(12):125024, December 2013.
- [69] R. R. Caldwell and N. A. Maksimova. Spectral distortion in a radially inhomogeneous cosmology. *Phys. Rev. D.*, 88(10):103502, November 2013.
- [70] Manoj Kaplinghat, Lloyd Knox, and Yong-Seon Song. Determining neutrino mass from the CMB alone. *Phys. Rev. Lett.*, 91:241301, 2003.
- [71] Michael Levi et al. The DESI Experiment, a whitepaper for Snowmass 2013. 2013.
- [72] J. S. Dunlop. The Cosmic History of Star Formation. *Science*, 333:178, July 2011.
- [73] Planck Collaboration, R. Adam, P. A. R. Ade, N. Aghanim, M. Arnaud, J. Aumont, C. Baccigalupi, A. J. Banday, R. B. Barreiro, J. G. Bartlett, and et al. Planck intermediate results. XXX. The angular power spectrum of polarized dust emission at intermediate and high Galactic latitudes. *ArXiv e-prints*, September 2014.
- [74] Planck Collaboration, P. A. R. Ade, N. Aghanim, C. Armitage-Caplan, M. Arnaud, M. Ashdown, F. Atrio-Barandela, J. Aumont, C. Baccigalupi, A. J. Banday, and et al. Planck 2013 results. XVIII. The gravitational lensing-infrared background correlation. *Astron. Astrophys.*, 571:A18, November 2014.
- [75] P. Madau and M. Dickinson. Cosmic Star-Formation History. *ARA&A*, 52:415–486, August 2014.
- [76] J. C. Hill, N. Battaglia, J. Chluba, S. Ferraro, E. Schaan, and D. N. Spergel. Taking the Universe’s Temperature with Spectral Distortions of the Cosmic Microwave Background. *Physical Review Letters*, 115(26):261301, December 2015.
- [77] Planck Collaboration, P. A. R. Ade, N. Aghanim, C. Armitage-Caplan, M. Arnaud, M. Ashdown, F. Atrio-Barandela, J. Aumont, C. Baccigalupi, A. J. Banday, and et al. Planck 2013 results. XX. Cosmology from Sunyaev-Zeldovich cluster counts. *Astron. Astrophys.*, 571:A20, November 2014.

- [78] S. Y. Sazonov and R. A. Sunyaev. Cosmic Microwave Background Radiation in the Direction of a Moving Cluster of Galaxies with Hot Gas: Relativistic Corrections. *Ap. J.*, 508:1–5, November 1998.
- [79] N. Itoh, Y. Kohyama, and S. Nozawa. Relativistic corrections to the sunyaev-zeldovich effect for clusters of galaxies. *Ap. J.*, 502:7, July 1998.
- [80] A. Challinor and A. Lasenby. Relativistic corrections to the sunyaev-zeldovich effect. *Ap. J.*, 499:1, May 1998.
- [81] J. A. Rubiño-Martín, C. Hernández-Monteagudo, and R. A. Sunyaev. The imprint of cosmological hydrogen recombination lines on the power spectrum of the CMB. *Astron. Astrophys.*, 438:461–473, August 2005.
- [82] C. Hernández-Monteagudo, J. A. Rubiño-Martín, and R. A. Sunyaev. On the influence of resonant scattering on cosmic microwave background polarization anisotropies. *MNRAS*, 380:1656–1668, October 2007.
- [83] A. Lewis. Rayleigh scattering: blue sky thinking for future CMB observations. *JCAP*, 8:053, August 2013.
- [84] K. Basu, C. Hernández-Monteagudo, and R. A. Sunyaev. CMB observations and the production of chemical elements at the end of the dark ages. *Astron. Astrophys.*, 416:447–466, March 2004.
- [85] D. R. G. Schleicher, D. Galli, F. Palla, M. Camenzind, R. S. Klessen, M. Bartelmann, and S. C. O. Glover. Effects of primordial chemistry on the cosmic microwave background. *Astron. Astrophys.*, 490:521–535, November 2008.

## **2 Curriculum Vitae**

### **3 Summary of Work Effort**

### **4 Current and Pending Support**

## **5 Letters of Support**

## **6 Budget Details - Narrative**

### **6.1 Team, and Work Effort**

#### **6.1.1 Funded Team Members**

#### **6.1.2 Non-Funded Team Members**

### **6.2 Costing Principles**

- **Summer Salaries:**
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### **6.3 University of Minnesota Budget**

#### **6.3.1 Direct Labor**

#### **6.3.2 Supplies**

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#### **6.3.4 Other Direct Costs**

#### **Publications and Teleconferencing**

#### **Other Subcontracts**

#### **6.3.5 Facilities and Administrative Costs**



## 7 Budget Sheets

**ACS** attitude control system

**ADC** analog-to-digital converters

**ADS** attitude determination software

**AHWP** achromatic half-wave plate

**AMC** Advanced Motion Controls

**ARC** anti-reflection coatings

**ATA** advanced technology attachment

**BRC** bolometer readout crates

**BLAST** Balloon-borne Large-Aperture Submillimeter Telescope

**CANbus** controller area network bus

**CIB** cosmic infrared background

**CMB** cosmic microwave background

**CMM** coordinate measurement machine

**CSBF** Columbia Scientific Balloon Facility

**CCD** charge coupled device

**DAC** digital-to-analog converters

**DASI** Degree Angular Scale Interferometer

**dGPS** differential global positioning system

**DfMUX** digital frequency domain multiplexer

**DLFOV** diffraction limited field of view

**DSP** digital signal processing

**EBEX** E and B Experiment

**EBEX2013** EBEX2013

**ELIS** EBEX low inductance striplines

**ETC** EBEX test cryostat

**FDM** frequency domain multiplexing

**FPGA** field programmable gate array

**FCP** flight control program

**FOV** field of view

**FWHM** full width half maximum

**GPS** global positioning system

**HDPE** high density polyethylene

**HIM** high index materials

**HWP** half-wave plate

**IA** integrated attitude

**IP** instrumental polarization

**JSON** JavaScript Object Notation

**LDB** long duration balloon

**LED** light emitting diode

**LCS** liquid cooling system

**LC** inductor and capacitor

**LZH** Lazer Zentrum Hannover

**MCP** multi-color pixel

**MSM** millimeter and sub-millimeter

**MLR** multilayer reflective

**MAXIMA** Millimeter Anisotropy eXperiment IMaging Array

**NASA** National Aeronautics and Space Administration

**NDF** neutral density filter

**PCB** printed circuit board

**PE** polyethylene

**PME** polarization modulation efficiency

**PSF** point spread function

**PV** pressure vessel

**PWM** pulse width modulation

**RMS** root mean square

**SLR** single layer reflective

**SMB** superconducting magnetic bearing

**SQUID** superconducting quantum interference device

**SQL** structured query language

**STARS** star tracking attitude reconstruction software

**SWS** sub-wavelength structures

**TES** transition edge sensor

**TDRSS** tracking and data relay satellites

**TM** transformation matrix

**UHMWPE** ultra high molecular weight polyethylene

**UMN** University of Minnesota