SALEM SETS AVOIDING ROUGH CONFIGURATIONS

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Recall that a set $X \subset \mathbf{R}^d$ is a *Salem set* of dimension s if it has Hausdorff dimension s, and for every $\varepsilon > 0$, there exists a probability measure μ_{ε} supported on X such that for all $\xi \in \mathbf{R}^d$,

$$\sup_{\xi \in \mathbf{R}^d} |\xi|^{s-\varepsilon} |\widehat{\mu}_{\varepsilon}(\xi)| < \infty.$$

Our goal in these notes is to obtain, for each set $Z \subset \mathbf{R}^{dn}$ with Minkowski dimension s, a Salem set $X \subset \mathbf{R}^d$ with dimension

$$\frac{nd-s}{s}$$

such that for each set of n distinct elements $x_1, \ldots, x_n \in X$, $(x_1, \ldots, x_n) \notin Z$. We hope that we can rely on the random selection approach of our paper on rough configurations to obtain such a result.

1. Orlicz Norm

We define a convex function $\psi_2:[0,\infty)\to[0,\infty)$ by $\psi_2(t)=e^{t^2}-1$, and define a corresponding Orlicz norm on the family of scalar valued random variables X over a probability space by setting

$$||X||_{\psi_2(L)} = \inf \{ A \in (0, \infty) : \mathbf{E}(\psi_2(|X|/A)) \le 1 \}.$$

The family of random variables $\psi_2(L)$ are known as subgaussian random variables. Here are some important properties:

• (Gaussian Tails): There exists a universal constant c > 0 such that for any random variable X, $||X||_{\psi_2(L)} \le A$ if and only if for each $t \ge 0$,

$$\mathbf{P}(|X| \ge t) \le 2\exp(-ct^2/A^2).$$

• (Bounded Variables are Subgaussian): For any random X,

$$||X||_{\psi_2(L)} \lesssim ||X||_{L^{\infty}}.$$

 \bullet (Centering) For any random variable X,

$$||X - \mathbf{E}(X)||_{\psi_2(L)} \lesssim ||X||_{\psi_2(L)}.$$

• (Union Bound) If X_1, \ldots, X_N are random variables, then

$$||X_1 + \dots + X_N||_{\psi_2(L)} \le ||X_1||_{\psi_2(L)} + \dots + ||X_N||_{\psi_2(L)}.$$

• (Hoeffding's Inequality): If X_1, \ldots, X_N are independent random variables, then

$$||X_1 + \dots + X_N||_{\psi_2(L)} \lesssim (||X_1||_{\psi_2(L)}^2 + \dots + ||X_N||_{\psi_2(L)}^2)^{1/2}.$$

The Orlicz norm is a convenient notation to summarize calculations of subgaussian concentration inequalities. Roughly speaking, we can think of a random variable X with $\|X\|_{\psi_2(L)} \leq A$ as having the large majority of it's support in the interval [-A, A], and having mass rapidly decaying outside of this interval.

2. A Family of Cubes

We fix sequences of integers $\{N_m: m \geq 1\}$ and $\{M_m: m \geq 1\}$. We then define two sequences of real numbers $\{l_m: m \geq 0\}$ and $\{r_m: m \geq 0\}$, by initially setting $l_0 = r_0 = 1$, and then, for each $m \geq 1$, setting $r_m = l_{m-1}/M_k$, and $l_m = r_m/N_m$. For each $m \geq 0$ and d, we define two collections of strings

$$\Sigma_m^d = \mathbf{Z}^d \times [M_1]^d \times \dots \times [N_1]^d \times \dots \times [M_m]^d \times [N_m]^d$$

and

$$\Pi_m^d = \mathbf{Z}^d \times [M_1]^d \times [N_1]^d \times \dots \times [N_{m-1}]^d \times [M_m]^d.$$

For each string $i \in \Sigma_m^d$, we define a vector $a_i \in (l_m \mathbf{Z})^d$ by setting

$$a_i = i_0 + \sum_{k=1}^{m} i_{2k-1} \cdot r_k + i_{2k} \cdot l_k$$

Then each string $i \in \Sigma_m^d$ can be identified with a sidelength l_m cube

$$Q_i = \prod_{j=1}^d \left[a_{ij}, a_{ij} + l_m \right].$$

centered at a_i . Similarly, for each string $i \in \Pi_m^d$, we define a vector $a \in (r_m \mathbf{Z})^d$ by setting, for each $1 \le j \le d$,

$$a_i = i_0 + \left(\sum_{k=1}^{m-1} i_{2k-1} \cdot r_k + i_{2k} \cdot l_k\right) + i_{2m-1} \cdot r_m,$$

and then define a sidelength r_m cube

$$R_i = \prod_{j=1}^d \left[a_{ij}, a_{ij} + r_m \right].$$

We let $\mathcal{Q}_m^d = \{Q_i : i \in \Sigma_m^d\}$, and $\mathcal{R}_m^d = \{R_i : i \in \Pi_m^d\}$. Then

- For each m, \mathcal{Q}_m^d and \mathcal{R}_m^d are covers of \mathbf{R}^d .
- If $Q_1, Q_2 \in \bigcup_{m=0}^{\infty} \mathcal{Q}_m^d$, then either Q_1 and Q_2 have disjoint interiors, or one cube is contained in the other. Similarly, if $R_1, R_2 \in \bigcup_{m=1}^{\infty} \mathcal{R}_m^d$, then either R_1 and R_2 have disjoint interiors, or one cube is contained in the other.

• For each cube $Q \in \mathcal{Q}_m$, there is a unique cube $Q^* \in \mathcal{R}_m$ with $Q \subset Q^*$. We refer to Q^* as the *parent cube* of Q. Similarly, if $R \in \mathcal{R}_m$, there is a unique cube in $R^* \in \mathcal{Q}_{m-1}$ with $R \subset R^*$, and we refer to R^* as the *parent cube* of R.

We say a set $E \subset \mathbf{R}^d$ is \mathcal{Q}_m discretized if it is a union of cubes in \mathcal{Q}_m^d , and we then let $\mathcal{Q}_m(E) = \{Q \in \mathcal{Q}_m^d : Q \subset E\}$. Similarly, we say a set $E \subset \mathbf{R}^d$ is \mathcal{R}_m discretized if it is a union of cubes in \mathcal{R}_m^d , and we then let $\mathcal{R}_m(E) = \{R \in \mathcal{R}_m^d : R \subset E\}$. We set $\Sigma_m(E) = \{i \in \Sigma_m^d : Q_i \in \mathcal{Q}_m(E)\}$, and $\Pi_m(E) = \{i \in \Pi_m^d : R_i \in \mathcal{R}_m(E)\}$. We say a cube $Q_1 \times \cdots \times Q_n \in \mathcal{Q}_m^{dn}$ is strongly non diagonal if $Q_i^* \cap Q_j^* = \emptyset$ for each $i \neq j$.

3. A Family of Mollifiers

We now consider a family of mollifiers, which we will use to smooth out the Fourier transform of the measures we study. Begin by choosing a non-negative C^{∞} function ψ supported on $[-1,1]^d$ such that

$$\int_{\mathbf{R}^d} \psi(x) \, dx = 1,\tag{1}$$

and for each $x \in \mathbf{R}^d$,

$$\sum_{n \in \mathbf{Z}^d} \psi(x+n) = 1. \tag{2}$$

Since ψ is C^{∞} and compactly supported, then for each $t \in [0, \infty)$, we conclude

$$\sup_{\xi \in \mathbf{R}^d} |\xi|^t |\widehat{\psi}(\xi)| < \infty. \tag{3}$$

Now we rescale to obtain a mollifier at each scale of the argument. For each k > 0, we let

$$\psi_k(x) = l_k^{-d} \psi(l_k \cdot x).$$

Then ψ_k is supported on $[-l_k, l_k]^d$. Equation (1) implies that for each $x \in \mathbf{R}^d$,

$$\int_{\mathbf{R}^d} \psi_k = 1. \tag{4}$$

Equation (2) implies

$$\sum_{n \in \mathbf{Z}^d} \psi(x + l_k \cdot n) = l_k^{-d}. \tag{5}$$

For each $\xi \in \mathbf{R}^d$, $\widehat{\psi}_k(\xi) = \widehat{\psi}(l_k \xi)$, and so in particular, (3) implies that for each $t \in [0, \infty)$,

$$|\widehat{\psi}_k(\xi)| \lesssim_t l_k^{-t} |\xi|^{-t}. \tag{6}$$

These properties are sufficient to mollify the functions we consider.

4. Discrete Lemma

Lemma 1. Fix $s \in [1, dn)$ and $\varepsilon \in [0, (n-s)/4)$. Let $T \subset [0, 1]^d$ be a non-empty, Q_m discretized set, and let μ_T be a smooth probability measure compactly supported on T, together with a constant C such that for each $m \in \mathbb{Z}^d$,

$$|\widehat{\mu_T}(m)| \le C|m|^{-s/2}$$
.

Let $B \subset \mathbf{R}^{dn}$ be a non-empty, \mathcal{Q}_{m+1} discretized set such that

$$\#(\mathcal{Q}_{m+1}(B)) \le (1/l_{m+1})^{s+\varepsilon}.$$

Then there exists a constant $C(\mu_T, n, s)$ such that if

$$M_{m+1} \ge C(T, \mu_T, n, d, s), \tag{7}$$

$$N_{m+1} \ge C(T, \mu_T, n, d, s) \cdot M_{m+1}^{\frac{s+\varepsilon}{dn-s-2\varepsilon}}, \tag{8}$$

and

$$N_{m+1} \ge 10^{1/\varepsilon},\tag{9}$$

then there exists a Q_{m+1} discretized set $S \subset T$ together with a smooth probability measure supported on S such that

• For any strongly non-diagonal cube

$$Q = Q_1 \times \cdots \times Q_n \in \mathcal{Q}_{m+1}(B),$$

There exists i such that $Q_i \notin \mathcal{Q}_{m+1}(S)$.

• For any $m \in \mathbf{Z}^d$,

$$|\widehat{\mu}(m)| \le (C + l_{m+1}^{-1})|m|^{-s/2}$$

Proof. For each $i \in \Pi_{m+1}^d$, let j_i be a random integer vector chosen from $[N_{m+1}]^d$, such that the family $\{j_i : i \in \Pi_{m+1}^d\}$ is independent. We define a measure ν_S by setting, for each $x \in \mathbf{R}^d$,

$$d\nu_S(x) = r_{m+1}^d \sum_{i \in \Pi_{m+1}^d} \psi_{m+1}(x - a_{ij_i}) d\mu_T(x).$$

We then normalize, defining

$$\mu_S = \frac{\nu_S}{\nu_S(\mathbf{R}^d)}.$$

If we set

$$S = \bigcup \{Q \in \mathcal{Q}_{m+1}^d : \mu_S(Q) > 0\},\$$

then by definition, S is a \mathcal{Q}_{m+1} discretized set, μ_S is supported on S, and $S \subset T$. It now suffices to show that with nonzero probability, S and μ_S satisfy the properties guaranteed by the lemma.

For each $i \in \Pi_{m+1}(T)$, define a random measure ν_i by setting

$$d\nu_i(x) = r_{m+1}^d \psi_{m+1}(x - a_{ij_i}) d\mu_T(x).$$

Then $\nu_S = \sum_{i \in \Pi_{m+1}(T)} \nu_i$. Note that if $j, j' \in [N_{m+1}]^d$, then

$$|a_{ij} - a_{ij'}| = |j - j'| \cdot l_{m+1} \lesssim_d N_{m+1} l_{m+1} = r_{m+1},$$

which implies

$$\left| r_{m+1}^{d} \int_{\mathbf{R}^{d}} \psi_{m+1}(x - a_{ij}) \mu_{T}(x) - r_{m+1}^{d} \int_{\mathbf{R}^{d}} \psi_{m+1}(x - a_{ij'}) \mu_{T}(x) \right|$$

$$\leq r_{m+1}^{d} \int_{\mathbf{R}^{d}} \psi_{m+1}(x) \left| \mu_{T}(x + a_{ij}) - \mu_{T}(x + a_{ij'}) \right|$$

$$\lesssim_{d} r_{m+1}^{d+1} \int_{\mathbf{R}^{d}} \psi_{m+1}(x) \| \nabla \mu_{T} \|_{L^{\infty}(\mathbf{R}^{d})} = r_{m+1}^{d+1} \| \nabla \mu \|_{L^{\infty}(\mathbf{R}^{d})}.$$

$$(10)$$

Thus (10) implies that for each i,

$$\|\nu_i(\mathbf{R}^d) - \mathbf{E}(\nu_i(\mathbf{R}^d))\|_{L^{\infty}} \lesssim_d r_{m+1}^{d+1} \|\nabla \mu\|_{L^{\infty}(\mathbf{R}^d)}. \tag{11}$$

Furthermore, (5) implies

$$\sum_{i \in \Pi_{m+1}^{d}} \mathbf{E}(\nu_{i}(\mathbf{R}^{d})) = r_{m+1}^{d} \sum_{(i,j) \in \Sigma_{m+1}^{d}} \mathbf{P}(j_{i} = j) \int_{\mathbf{R}^{d}} \psi_{m+1}(x - a_{ij}) \mu_{T}(x) dx$$

$$= \frac{r_{m+1}^{d}}{N_{m+1}^{d}} \int_{\mathbf{R}^{d}} \left(\sum_{(i,j) \in \Sigma_{m+1}^{d}} \psi_{m+1}(x - a_{ij}) \right) \mu_{T}(x) dx \qquad (12)$$

$$= \frac{r_{m+1}^{d} l_{m+1}^{-d}}{N_{m+1}^{d}} = 1.$$

For all but $O_d(r_{m+1}^{-d})$ indices i, $\nu_i = 0$ almost surely. Thus we can apply the triangle inequality together with (11), and (12), we conclude

$$\|\nu_{S}(\mathbf{R}^{d}) - 1\|_{L^{\infty}} = \|\sum_{i \in \Pi_{m+1}^{d}} \left[\nu_{i}(\mathbf{R}^{d}) - \mathbf{E}(\nu_{i}(\mathbf{R}^{d}))\right]\|_{L^{\infty}}$$

$$\leq \sum_{i \in \Pi_{m+1}^{d}} \|\nu_{i}(\mathbf{R}^{d}) - \mathbf{E}(\nu_{i}(\mathbf{R}^{d}))\|_{L^{\infty}}$$

$$\lesssim_{d} r_{m+1}^{-d} r_{m+1}^{d+1} \|\nabla \mu\|_{L^{\infty}(\mathbf{R}^{d})}$$

$$= r_{m+1} \|\nabla \mu\|_{L^{\infty}(\mathbf{R}^{d})}.$$

$$(13)$$

Thus if (7) holds for an appropriately chosen constant depending on d and $\|\nabla \mu\|_{L^{\infty}(\mathbf{R}^d)}$, we can apply (13) to conclude

$$\|\nu_S(\mathbf{R}^d) - 1\|_{L^{\infty}} \le 1/2.$$
 (14)

Thus normalizing by $\nu_S(\mathbf{R}^d)$ only introduces a neglible constant.

For each $i \in \Pi_{m+1}^d$, let

$$S_i = \bigcup \{ Q \in \mathcal{Q}_{m+1}^d : \nu_i(Q) > 0 \}.$$

Then $S = \bigcup_{i \in \Pi_{m+1}^d} S_i$. Because j_i is selected uniformly from $[N_{m+1}]^d$ for each i, and ψ_{m+1} is supported on $[-l_{m+1}, l_{m+1}]^d$,

$$S_i \subset \bigcup \{R_{i_0} : R_{i_0} \cap R_i \neq \emptyset\}.$$

For any cube $Q_{ij} \in \Sigma_{m+1}^d$, there are $O_d(1)$ pairs (i_0, j_0) such that $Q_{i_0j_0} \cap Q_{ij} \neq \emptyset$, and so a union bound gives

$$\mathbf{P}(Q_{ij} \in \mathcal{Q}_{m+1}(S)) \le \sum_{Q_{i_0,i_0} \cap Q_{ij} \neq \emptyset} \mathbf{P}(j_{i'} = j') \lesssim_d N_{m+1}^{-d}.$$

Without loss of generality, removing cubes from B if necessary, we may assume all cubes in B are strongly non-diagonal. Let $Q = Q_{i_1j_1} \times \cdots \times Q_{i_nj_n}$ be a strongly non-diagonal cube in $Q_{m+1}(B)$. Since Q is strongly diagonal, the events $\{Q_{i_kj_k} \in S\}$ are independent from one another, which implies that

$$\mathbf{P}(Q \in \mathcal{Q}_{m+1}(S^n)) = \mathbf{P}(Q_{i_1j_1} \in S) \dots \mathbf{P}(Q_{i_nj_n} \in S) \lesssim_{d,n} N_{m+1}^{-dn}. \tag{15}$$

Taking expectations over all cubes in B, and applying (15) gives

$$\mathbf{E}(\#(\mathcal{Q}_{m+1}(B) \cap \mathcal{Q}_{m+1}(S^n))) \lesssim_{d,n} \#(\mathcal{Q}_{m+1}(B)) \cdot N_{m+1}^{-dn}$$

$$\leq l_{m+1}^{-(s+\varepsilon)} N_{m+1}^{-dn}$$

$$= \frac{M_{m+1}^{s+\varepsilon} l_m^{-(s+\varepsilon)}}{N_{m+1}^{dn-s-\varepsilon}}.$$
(16)

If (8) holds, for an appropriately chosen constant depending only on l_m, d, n , and s, we can apply Markov's inequality together with (9) and (16) to conclude

$$\mathbf{P}(\mathcal{Q}_{k+1}(B) \cap \mathcal{Q}_{k+1}(S^n) \neq \varnothing) = \mathbf{P}(\#(\mathcal{Q}_{k+1}(B) \cap \mathcal{Q}_{k+1}(S^n)) \geq 1)$$

$$\leq \mathbf{E}(\#(\mathcal{Q}_{m+1}(B) \cap \mathcal{Q}_{m+1}(S^n)))$$

$$\leq 1/N_{m+1}^{\varepsilon} \leq 1/10.$$
(17)

Thus $Q_{k+1}(S^n)$ is disjoint from $Q_{k+1}(B)$ with high probability.

Now we analyze the Fourier transform of the measure ν . For each $i \in \Pi_{m+1}^d$, and $m \in \mathbb{Z}$, define $X_{im} = \widehat{\nu}_i(m) - \widehat{\mathbf{E}(\nu_i)}(m)$. Note that

$$\sum_{i \in \Pi_{m+1}^{d}} \widehat{\mathbf{E}(\nu_{i})}(m) = \sum_{i \in \Pi_{m+1}^{d}} l_{m+1}^{d} \sum_{j \in [N_{m+1}]^{d}} \int_{\mathbf{R}^{d}} e^{-2\pi i m \cdot x} \psi_{m+1}(x - a_{ij}) d\mu_{T}(x)$$

$$= \int_{\mathbf{R}^{d}} e^{-2\pi i m \cdot x} d\mu_{T}(x) = \widehat{\mu_{T}}(m).$$
(18)

For each i and m, the standard (L^1, L^{∞}) bound on the Fourier transform, combined with (11), shows

$$||X_{im}||_{\psi_2(L)} \leq ||X_{im}||_{L^{\infty}} \leq ||\nu_i(\mathbf{R}^d)||_{L^{\infty}} + \mathbf{E}(\nu_i)(\mathbf{R}^d)$$

$$\lesssim_d \mathbf{E}(\nu_i)(\mathbf{R}^d) + r_{m+1}^{d+1} ||\nabla \mu_T||_{L^{\infty}(\mathbf{R}^d)}.$$
(19)

For a fixed m, the family of random variables $\{X_{im}\}$ are independent. Furthermore, $\sum X_{im} = \widehat{\nu}(m) - \widehat{\mathbf{E}(\nu)}(m)$, and

$$\mathbf{E}(\widehat{\nu_{S}}(m)) = \frac{r_{m+1}^{d}}{N_{m+1}^{d}} \int_{\mathbf{R}^{d}} e^{-2\pi i m \cdot x} \left(\sum_{(i,j) \in \Sigma_{m+1}^{d}} \psi_{m+1}(x - a_{ij}) \right) d\mu_{T}(x)$$

$$= \frac{r_{m+1}^{d} l_{m+1}^{-d}}{N_{m+1}^{d}} \int_{\mathbf{R}^{d}} e^{-2\pi i m \cdot x} d\mu_{T}(x)$$

$$= \frac{r_{m+1}^{d} l_{m+1}^{-d}}{N_{m+1}^{d}} \widehat{\mu_{T}}(m) = \widehat{\mu_{T}}(m).$$

Thus we may apply Hoeffding's inequality to (19) to conclude that

$$\|\widehat{\nu}(m) - \widehat{\mu_T}(m)\|_{\psi_2(L)} \lesssim_d \left(\sum \mathbf{E}(\nu_i)(\mathbf{R}^d)^2\right)^{1/2} + r_{m+1}^{d/2+1} \|\nabla \mu_T\|_{L^{\infty}(\mathbf{R}^d)}.$$

Now taking in absolute values gives the inequality

$$\mathbf{E}(\nu_{i}(\mathbf{R}^{d})) = l_{m+1}^{d} \sum_{j \in [N_{m+1}]^{d}} \int \psi_{m+1}(x - a_{ij}) d\mu_{T}(x)$$

$$\leq r_{m+1}^{d} \|\mu_{T}\|_{L^{\infty}(\mathbf{R}^{d})},$$

and so

$$\|\widehat{\nu}(m) - \widehat{\mu_T}(m)\|_{\psi_2(L)} \lesssim_d [\|\mu_T\|_{L^{\infty}(\mathbf{R}^d)} + \|\nabla \mu_T\|_{L^{\infty}(\mathbf{R}^d)}] r_{m+1}^{d/2}.$$
 (20)

If (8) holds, for an appropriately chosen constant depending only on l_m, d, n , and s, then we can apply a union bound over $D = \{m \in \mathbf{Z}^d : |m| \le 10l_{m+1}^{-1}\}$ to conclude that there exists a constant $c(\mu_T, d)$ such that

$$\mathbf{P}\left(\|\widehat{\nu} - \widehat{\mu_T}\|_{L^{\infty}(D)} \ge r_{m+1}^{d/2} \log(M_{m+1})\right) \lesssim_d l_{m+1}^{-d} \exp\left(-c(\mu_T, d) \log(M_{m+1})^2\right) < 1/10.$$

Thus $\widehat{\nu}$ and $\widehat{\mu_T}$ are highly like to differ only by a miniscule amount over small frequencies.

Finally, it suffices to analyze the values of $\widehat{\nu}_S(m)$ when $|m| \ge 10l_{m+1}^{-1}$. We note that if we define a random measure

$$\alpha = r_{m+1}^d \sum_{\substack{i \in \Pi_{m+1}^d \\ d(a_i, T) \leq 2r_{m+1}^{-1}}} \delta_{a_{ij_i}},$$

then $\nu_S = (\alpha * \psi_{m+1})\mu_T$. Thus we have $\widehat{\nu}_S = \widehat{\mu}_T * (\widehat{\alpha} \cdot \widehat{\psi}_{m+1})$. Since μ_T is compactly supported, we can define, for each t > 0,

$$A(t) = \sup |\widehat{\mu_T}(\xi)| |\xi|^t < \infty.$$

Similarity, since

$$\widehat{\psi_{m+1}}(\xi) = \widehat{\psi}\left(l_{m+1}^{-1}\xi\right),\,$$

if we set, for each t > 0,

$$B(t) = \sup |\widehat{\psi}(\xi)| |\xi|^t,$$

then

$$\sup |\widehat{\psi_{m+1}}(\xi)||\xi|^t = l_{m+1}^{-t}B(t).$$

It now suffices to bound

$$\sup_{|\eta| \ge 10l_{m+1}^{-1}} |\eta|^{s/2} \int \widehat{\mu_T} (\eta - \xi) \widehat{\alpha}(\xi) \widehat{\psi_{m+1}}(\xi) d\xi.$$

Since $\alpha(\mathbf{R}^d) \leq 2^d$, $\|\widehat{\alpha}\|_{L^{\infty}(\mathbf{R}^d)} \leq 2^d$, so it suffices to understand

$$\int |\widehat{\mu_T}(\eta - \xi)| |\widehat{\psi_{m+1}}(\xi)| d\xi.$$

If $|\xi| \le |\eta|/2$, $|\eta - \xi| \ge |\xi|/2$, so for all t > 0,

$$\int_{0 \le |\xi| \le |\eta|/2} |\widehat{\mu_T}(\eta - \xi)| |\widehat{\psi_{m+1}}(\xi)| d\xi \le \frac{A(t)2^{t-d}}{|\eta|^{t-d}}.$$

If we set t = d + 1 + s/2 and apply (8) for an appropriate chosen constant depending only on d and μ_T , we conclude

$$\int_{0 \le |\xi| \le |\eta|/2} |\widehat{\mu_T}(\eta - \xi)| |\widehat{\psi_{m+1}}(\xi)| \ d\xi \le \frac{A(d+1+s/2)2^{1+s/2}|\eta|^{-1}}{|\eta|^{s/2}} \le \frac{|\eta|^{-s/2}}{10 \cdot 2^d}.$$

Conversely, if $|\xi| \ge 2|\eta|$, then $|\eta - \xi| \ge |\xi|/2$, so for each t > 0,

$$\int_{|\xi| \ge 2|\eta|} |\widehat{\mu_T}(\eta - \xi)| |\widehat{\psi}_{m+1}(\xi)| d\xi \le \int_{|\xi| \ge 2|\eta|} \frac{A(t)2^t}{|\xi|^t} \\
\lesssim_d \int_{2|\eta|}^{\infty} r^{d-t} A(t) l_{m+1}^{-t'} 2^t.$$

Provided t > d + 1, this integral is finite, and is

$$\lesssim_{d,t} A(t) 2^{d+1} |\eta|^{d+1-t}$$
.

Setting t = d + 2 + s/2, and applying (8), we conclue

$$\int_{|\xi| \ge 2|\eta|} |\widehat{\mu_T}(\eta - \xi)||\widehat{\psi}_{m+1}(\xi)|| d\xi \le \frac{|\eta|^{-s/2}}{10 \cdot 2^d}.$$

Finally, we conclude that for each t > 0,

$$\int_{|\eta|/2 \le |\xi| \le 2|\eta|} |\widehat{\mu_T}(\eta - \xi)||\widehat{\psi_{m+1}}(\xi)| d\xi \le \frac{2^d |\eta|^d B(t) 2^t}{|\xi|^t}.$$

If we set t=d+s/2+1 and apply (8), we conclude

$$\int_{|\eta|/2 \le |\xi| \le 2|\eta|} |\widehat{\mu_T}(\eta - \xi)||\widehat{\psi_{m+1}}(\xi)| d\xi \le \frac{|\xi|^{-s/2}}{10 \cdot 2^d}.$$

Summing up the three bounds, we conclude that if $|\eta| \ge 10 l_{m+1}^{-1}$, then $|\widehat{\nu_S}(\eta)| \le |\eta|^{-s/2}$.