# Radial Multipliers

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March 17, 2022

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#### Research Outline

The question of the  $L^p$  boundedness of translation-invariant operators on **R**<sup>n</sup> has proved central to the development of modern harmonic analysis. Indeed, answers to these questions underpin any subtle understanding of the Fourier transform, since with essentially any such operator T, we can associate a function  $m: \mathbb{R}^n \to \mathbb{C}$ , known as the symbol of T, such that for any function f, the Fourier transform of Tf obeys the relation  $\widehat{Tf} = m\widehat{f}$ ; thus translation invariant operators are also called Fourier multiplier operators. Initial questions about Fourier multipliers emerged from classical questions concerning the convergence properties of Fourier series, and in the study of the classical equations of physics, like the heat and wave equation. Such operators often have rotational symmetry, so it is natural to restrict our attention to multiplier operators which are also rotationinvariant. These operators are called radial Fourier multipliers, since the associated symbol is then a radial function. This research project proposes the study of necessary and sufficient conditions to guarantee  $L^p$  boundedness of radial multiplier operators, stimulated by recent developments which indicate lines of attack for three related problems in the field.

The general study of the boundedness of Fourier multipliers was intiated in the 1960s. It was quickly realized that the most fundamental estimates for a translation-invariant operator T were  $L^p$  estimates of the form  $||Tf||_{L^p(\mathbb{R}^n)} \lesssim ||f||_{L^p(\mathbb{R}^n)}$  in the range  $1 \leqslant p \leqslant 2$ . For p=1 and p=2, mathematicians found simple necessary and sufficient conditions to ensure boundedness [7]. But the problem of finding necessary and sufficient conditions for boundedness in the range 1 proved impenetrable. Indeed, many interesting problems about the boundedness of*spe*-

*cific* Fourier multiplier operators for values of *p* in the range, such as the Bochner-Riesz conjecture, remain largely unsolved today.

Thus it came as a surprise when recent results indicated necessary and sufficient conditions for radial Fourier multiplier operators to be bounded for values of p in this range. First came the result of [3], which gave a simple sufficient and necessary condition for bounds of the form  $||Tf||_{L^p(\mathbb{R}^n)} \lesssim$  $||f||_{L^p(\mathbf{R}^n)}$  to hold for a particular radial multiplier operator T, uniformly over radial functions f, precisely in the range 1 . It isnatural to conjecture that the same criterion, applied to the same range of p, gives the bound  $||Tf||_{L^p(\mathbb{R}^n)} \lesssim ||f||_{L^p(\mathbb{R}^n)}$  for general functions f. In this outline we call this statement the radial multiplier conjecture. We now know, by the results of [6] and [2], that the radial multiplier conjecture is true when n > 4 and 1 , and when <math>n = 4 and 1 . We also know [2] the criterion in the conjectureis sufficient to obtain a restricted weak type bound  $||Tf||_{L^p(\mathbb{R}^n)} \lesssim ||f||_{L^{p,1}(\mathbb{R}^n)}$ when n = 3 and 1 . But the radial multiplier conjecture has not yet been completely resolved in any dimension n, we do not have any strong type  $L^p$  bounds when n = 3, and we don't have any bounds whatsoever when n = 2.

If fully proved, the radial multiplier conjecture would imply the Bochner-Riesz conjecture, and thus the Kakeya and restriction conjectures as a result. All three consequences are major unsolved problems in harmonic analysis, so a complete resolution of the conjecture is far beyond the scope of current research techniques. On the other hand, the Bochner Riesz conjecture is completely resolved when n = 2, while in contrast, no results related to the radial multiplier conjecture are known in this dimension at all. And in any dimension n > 2, the range under which the Bochner-Riesz multiplier is known to hold [4] is strictly larger than the range under which the radial multiplier conjecture is known to hold, even for the restricted weak-type bounds obtained in [2]. Thus it still seems within hope that the techniques recently applied to improve results for Bochner-Riesz problem, such as broad-narrow analysis [1], the polynomial Wolff axioms [8], and methods of incidence geometry and polynomial partitioning [12] can be applied to give improvements to current results characterizing the boundedness of general radial Fourier multipliers.

Our hopes are further emboldened when we consult the proofs in [6] and [2], which reduce the radial multiplier conjecture to the study of upper bounds of quantities of the form  $\|\sum_{(y,r)\in\mathcal{E}} F_{y,r}\|_{L^p(\mathbb{R}^n)}$ , where  $\mathcal{E} \subset$ 

 $\mathbf{R}^n \times (0,\infty)$  is a finite collection of pairs, and  $F_{y,r}$  is an oscillating function supported on a O(1) neighborhood of a sphere of radius r centered at a point y. The  $L^p$  norm of this sum is closely related to the study of the tangential intersections of these spheres, a problem successfully studied in more combinatorial settings using incidence geometry and polynomial partitioning methods [13], which provides further estimates that these methods might yield further estimates on the radial multiplier conjecture.

When n=3, [2] is only able to obtain bounds on the  $L^p$  sums in the last paragraph when  $\mathcal{E}$  is a Cartesian product of two subsets of  $(0,\infty)$  and  $\mathbf{R}^n$ . This is why only restricted weak-type bounds have been obtained in this dimension. It is therefore an interesting question whether different techniques enable one to extend the  $L^p$  bounds of these sums when the set  $\mathcal{E}$  is not a Cartesian product, which would allow us to upgrade the result of [2] in n=3 to give strong  $L^p$  bounds. This question also has independent interest, because it would imply new results for the 'endpoint' local smoothing conjecture, which concerns the regularity of solutions to the wave equation in  $\mathbf{R}^n$ . Incidence geometry has been recently applied to yield results on the 'non-endpoint' local smoothing conjecture [5], which again suggests these techniques might be applied to yield the estimates needed to upgrade the result of [2] to give strong  $L^p$ -type bounds.

A third line of questioning about the radial multiplier conjecture is obtained by studying natural analogues of Fourier multiplier operators on Riemannian manifolds. Using functional calculi, for any function  $m:[0,\infty)\to \mathbb{C}$ , we can associate an operator  $m(\sqrt{-\Delta})$  on a compact Riemannian manifold M, where  $\Delta$  is the Laplace-Beltrami on M. Just like multiplier operators on  $\mathbb{R}^n$  are crucial to an understanding of the Fourier transform, the operators  $m(\sqrt{-\Delta})$  are crucial to understand the behaviour of eigenfunctions of the Laplace-Beltrami operator on M.

The direct analogue of the boundedness problems for multipliers is trivial in this setting, since any compactly supported function m will induce an operator  $T=m(\sqrt{-\Delta})$  satisfying estimates of the form  $\|Tf\|_{L^p(M)}\lesssim \|f\|_{L^p(M)}$  for all  $1\leqslant p\leqslant 2$ . The correct formulation of the problem is instead to consider the alternate bound  $\sup_{R>0}\|m(\sqrt{-\Delta}/R)f\|_{L^p(M)}\lesssim \|f\|_{L^p(M)}$ . In fact, a transference principle of Mitjagin [10] implies that if the latter bound holds for a multiplier  $m(\sqrt{-\Delta})$ , then a bound of the form  $\|Tf\|_{L^p(\mathbf{R}^n)}\lesssim \|f\|_{L^p(\mathbf{R}^n)}$  holds where T is the Fourier multiplier operator associated with

the symbol  $m(|\xi|)$ . Thus boundedness in this alternate sense is at least as hard on a compact manifold M as it is in  $\mathbb{R}^n$ . The research project here is to try and extend the radial multiplier conjecture to this setting.

For general compact manifolds difficulties arise in generalizing the conjecture, connected to the fact that analogues of the Kakeya / Nikodym conjecture are false in this general setting [9]. But these problems do not arise for constant curvature manifolds, like the sphere. The sphere also has over special properties which make it especially amenable to analysis, such as the fact that solutions to the wave equation on spheres are periodic. Best of all, there are already results which achieve the analogue of [3] on the sphere. Thus it seems reasonable that current research techniques can obtain interesting results for radial multipliers on the sphere, at least in the ranges established in [6] or even [2].

In conclusion, the results of [6] and [2] indicate three lines of questioning about radial Fourier multiplier operators, which current research techniques place us in reach of resolving. The first question is whether we can extend the range of exponents upon which the conjecture of [3] is true, at least in the case n = 2 where Bochner-Riesz has been solved. The second is whether we can use more sophisticated arguments to prove the  $L^p$  sum bounds obtained in [2] when n = 3 when the sums are no longer cartesian products, thus obtaining strong  $L^p$  characterizations in this settiong, as well as new results about the endpoint local smoothing conjecture. The third question is whether we can generalize these bounds obtained in these two papers to study radial Fourier multipliers on the sphere.

# Sogge: $L^p$ Estimates for the Wave Equation and Applications

When solving an oscillatory integral,

# Seeger: Singular Convolution Operators in $L^p$ Spaces

Let  $m(\xi)$  be a function acting as a Fourier multiplier. If the resulting operator m(D) was bounded from  $L^p(\mathbf{R}^d)$  to  $L^p(\mathbf{R}^d)$  with operator norm A, then the operator would also be bounded 'at all scale'. That is, if we consider a littlewood Paley decomposition  $f = P_0 f + P_1 f + \ldots$ , where  $\widehat{P_i f} = \eta_i \cdot \widehat{f}$  is supported on  $|\xi| \sim 2^i$  for  $i \ge 1$ , and  $|\xi| \le 1$  for i = 0, then for each i, we would have estimates of the form

$$||m(D)P_if||_{L^p(\mathbf{R}^d)} = ||m_i(D)f||_{L^p(\mathbf{R}^d)} \le A||f||_{L^p(\mathbf{R}^d)},$$
 (3.1)

where the implicit constant is uniform in i. The main focus of the paper in question is to determine whether a uniform bound of the form (3.1) for each i implies m(D) is bounded. More precisely, if (3.1) holds for all i, is it necessarily true that  $\|m(D)f\|_{L^p(\mathbf{R}^d)} \lesssim_p A\|f\|_{L^p(\mathbf{R}^d)}$ ?

The condition certainly works for p=2, since (3.1) would then imply that  $\|m_i\|_{L^\infty(\mathbf{R}^d)} \leq A$  for all i, and since  $\sum \eta_i(\xi)=1$  for all i, and at most O(1) of the functions are supported at each point, this implies  $\|m\|_{L^\infty(\mathbf{R}^d)} \lesssim A$ , and so  $\|m(D)f\|_{L^2(\mathbf{R}^d)} \leq \|m\|_{L^\infty(\mathbf{R}^d)} \|f\|_{L^2(\mathbf{R}^d)} \lesssim \|f\|_{L^2(\mathbf{R}^d)}$ . On the other hand, for p=1 or  $p=\infty$ , this condition doesn't work at all. Intuitively we shouldn't expect the condition to work, since Littlewood-Paley theory runs into all kinds of problems for these values of p. More precisely, the condition would be equivalent to showing that if  $\|K*\widehat{\eta_i}\|_{L^1(\mathbf{R}^d)} \leq A$  for all

*i*, then  $||K||_{L^1(\mathbf{R}^d)} \lesssim A$ . If

$$K_N(x) = \int_{-2^N}^{2^N} e^{2\pi i \xi \cdot x} d\xi$$

is the Dirichlet kernel, then for  $i \leq N$ ,  $\|K * \widehat{\eta_i}\|_{L^1(\mathbf{R})} = \|\widehat{\eta_i}\|_{L^1(\mathbf{R})} \sim 1$ , and for  $i \geq N$ ,  $K * \widehat{\eta_i} = 0$ . On the other hand,  $\|K_N\|_{L^1(\mathbf{R})} \sim N$ , so we cannot have  $\|K\|_{L^1(\mathbf{R})} \lesssim 1$  uniformly in N. Baire category techniques can thus be used to produce a kernel K not in  $L^1(\mathbf{R})$ , but such that  $\|K * \eta_i\|_{L^1(\mathbf{R})} \lesssim 1$  uniformly in N.

The result actually fails for 2 , due to an examples of Triebel. For simplicity, let's work in**R** $. If we fix a bump function <math>\phi \in C_c^{\infty}(\mathbf{R})$  supported in [-1,1], and set

$$m_N(\xi) = \sum_{k=N}^{2N} e^{2\pi i (2^k \xi)} \phi(\xi - 2^k),$$

then  $m_N(\xi)\eta_i(\xi) = m_{N,i}(\xi)$ , where  $m_{N,i}(\xi) = e^{2\pi i(2^k\xi)}\phi(\xi - 2^k)$ , and so  $K_{N,i}(x) = \widehat{m_{N,i}}(x) = e^{2\pi i 2^k(x-2^k)}\widehat{\phi}(x-2^k)$ , hence

$$||m_{N,i}(D)f||_{L^p(\mathbf{R}^d)} = ||K_{N,i} * f||_{L^p(\mathbf{R}^d)} \le ||\widehat{\phi}||_{L^1(\mathbf{R})} ||f||_{L^p(\mathbf{R})} \lesssim ||f||_{L^p(\mathbf{R})}.$$

On the other hand, the operator norm of  $m_N(D)$  from  $L^p(\mathbf{R})$  to  $L^p(\mathbf{R})$  is actually  $\gtrsim_p N^{|1/p-1/2|}$ , and thus not bounded uniformly in N, so Baire category shows things don't work so well here.

This paper shows that one *can* get uniform bounds assuming an additional, very weak smoothness condition, which rules out the example  $m_N$  above. Under the most simple assumptions, if (3.1) holds, and  $\|m_i\|_{\Lambda^{\varepsilon}} \lesssim 2^{-ik}$ , where  $\Lambda^{\varepsilon}$  is the  $\varepsilon$ -Lipschitz norm, then  $\|m(D)f\|_{L^r(\mathbf{R}^d)} \lesssim \|f\|_{L^r(\mathbf{R}^d)}$  whenever |1/r - 1/2| < |1/p - 1/2|. Under slightly stronger smoothness assumptions, we can actually conclude  $\|m(D)f\|_{L^p(\mathbf{R}^d)} \lesssim \|f\|_{L^p(\mathbf{R}^d)}$ .

To prove the result, we rely on Littlewood-Paley theory and the Fefferman-Stein sharp maximal function. Without loss of generality we may assume that 2 . We will actually show that if for all <math>i and  $\omega \ge 0$ ,

$$\int_{|x| \geqslant \omega} |K_i(x)| \, dx \leqslant B(1 + 2^i \omega)^{-\varepsilon},$$

consistent with the fact that, if  $m_i$  was smooth, the uncertainty principle would say that  $K_i$  would live on a ball of radius  $1/2^i$ . We will then prove that  $||m(D)f||_{L^p(\mathbf{R}^d)} \leq A\widetilde{\log}(B/A)^{|1/2-1/p|}$ . Our goal is to show that if

$$S^{\#}f(x) = \sup_{x \in Q} \int_{Q} \left( \sum_{i=0}^{\infty} \left| m_{i}(D)f(y) - \int_{Q} m_{i}(D)f(z) dz \right|^{2} \right)^{1/2} dy,$$

then  $||S^{\#}f||_{L^p(\mathbf{R}^d)} \lesssim A\widetilde{\log}(B/A)^{1/2-1/p}||f||_{L^p(\mathbf{R}^d)}$ . It then follows by Littlewood-Paley theory implies

$$\|m(D)f\|_{L^{p}(\mathbf{R}^{d})} \lesssim_{p} \left\| \left( \sum_{k=0}^{\infty} |m_{i}(D)f|^{2} \right)^{1/2} \right\|_{L^{p}(\mathbf{R}^{d})}$$

$$\leq \left\| M \left[ \left( \sum_{k=0}^{\infty} |m_{i}(D)f|^{2} \right)^{1/2} \right] \right\|_{L^{p}(\mathbf{R}^{d})}$$

$$\lesssim \left\| S^{\#} \left( \sum_{k=0}^{\infty} |m_{i}(D)f|^{2} \right)^{1/2} \right\|_{L^{p}(\mathbf{R}^{d})}$$

$$\lesssim A\widetilde{\log}(B/A)^{1/2-1/p}.$$

To bound  $S^{\#}$ , we linearize using duality, picking  $Q_x$  for each x, and a family of functions  $\chi_i(x,y)$  such that  $\left(\sum |\chi_i(x,y)|^2\right)^{1/2} \le 1$ , such that

$$S^{\#}f(x) \approx \int_{Q_x} \sum_{i=0}^{\infty} \left( m_i(D) f(y) - \int_{Q_x} m_i(D) f(z) dz \right) \chi_i(x,y) dy.$$

Thus  $S^{\#}f = S_1f + S_2f$ , where if  $Q_x$  has sidelength  $2^{l(x)}$ ,

$$S_1 f(x) = \int_{Q_x} \sum_{|i+l(x)| \le \tilde{\log}(B/A)} \left( m_i(D) f(y) - \int_{Q_x} m_i(D) f(z) \, dz \right) \chi_i(x, y) \, dy$$

and

$$S_2 f(x) = \int_{Q_x} \sum_{|i+l(x)| \geqslant \tilde{\log}(B/A)} \left( m_i(D) f(y) - \int_{Q_x} m_i(D) f(z) \, dz \right) \chi_i(x, y) \, dy.$$

If  $|i + l(x)| \leq 1$ , then the uncertainty principle tells us that  $m_i(D)f$  is roughly constant on squares on radius  $Q_x$ , up to some small error, so that we should expect

$$\left| m_i(D)f(y) - \oint_{Q_x} m_i(D)f(z) \ dz \right| \lesssim \left| \oint_{Q_x} m_i(D)f(z) \ dz \right|.$$

Thus it is natural to use the bound,  $|S_1f(x)| \lesssim M(\sum_{i=0}^{\infty} |m_i(D)f|^2)^{1/2}$ , which implies

$$||S_1 f||_{L^2(\mathbf{R}^d)} \lesssim ||M(\sum_{i=0}^{\infty} |m_i(D)f|^2)^{1/2}||_{L^2(\mathbf{R}^d)}$$

$$\lesssim ||(\sum_{i=0}^{\infty} |m_i(D)f|^2)^{1/2}||_{L^2(\mathbf{R}^d)}$$

$$= \left(\sum_{i=0}^{\infty} ||m_i(D)f||_{L^2(\mathbf{R}^d)}^2\right)^{1/2}$$

and

$$||S_{1}f||_{L^{\infty}(\mathbf{R}^{d})} \leq ||M(\sum_{|i+l(x)| \leq |\tilde{\log}(B/A)}^{\infty} |m_{i}(D)f|^{2})^{1/2}||_{L^{\infty}(\mathbf{R}^{d})}$$

$$\leq ||\left(\sum_{|i+l(x)| \leq |\tilde{\log}(B/A)} |m_{i}(D)f|^{2}\right)^{1/2}||_{L^{\infty}(\mathbf{R}^{d})}$$

$$\lesssim ||\tilde{\log}(B/A)^{1/2} \sup_{i} ||m_{i}(D)f||_{L^{\infty}(\mathbf{R}^{d})}$$

Interpolation gives  $\|S_1 f\|_{L^p(\mathbf{R}^d)} \lesssim \tilde{\log}(B/A)^{1/2-1/p} \|m_i(D)f\|_{L^p_x(l^p_i)}$ . But now Littlewood-Paley theory shows that

$$\|m_i(D)f\|_{L^p_x(l^p_i)} \leqslant A \left( \sum_{i=0}^{\infty} \|P_i f\|_{L^p(\mathbf{R}^d)} \right)^{1/p} \leqslant A \left( \sum_{i=0}^{\infty} \|P_i f\|_{L^p(\mathbf{R}^d)}^2 \right)^{1/2} \lesssim A \|f\|_{L^p}.$$

Thus 
$$||S_1 f||_{L^p(\mathbf{R}^d)} \lesssim A\tilde{\log}(B/A)^{1/2-1/p}||f||_{L^p(\mathbf{R}^d)}$$
.

On the other hand, if i is much smaller than l(x), we should expect the error between  $m_i(D)f(y)$  and  $\oint_{Q_x} m_i(D)f(z)\,dz$  to be even smaller, and if i is much bigger, then  $m_i(D)f$  is no longer constant at this scale, and so the averages should be small, so  $m_i(D)f(x)$  should dominate  $\oint_{Q_x} m_i(D)f(z)$ . Now since our assumption implues that  $\|m(D)f\|_{L^2(\mathbb{R}^d)} \lesssim \|f\|_{L^2(\mathbb{R}^d)}$ , it is not so difficult to prove that

$$||S_2 f||_{L^2(\mathbf{R}^d)} \lesssim A ||f||_{L^2(\mathbf{R}^d)} \sim A ||\left(\sum |P_i f|^2\right)^{1/2}||_{L^2(\mathbf{R}^d)}.$$

The difficulty is proving  $\|S_2 f\|_{L^\infty(\mathbf{R}^d)} \lesssim A \|\left(\sum |P_i f|^{1/2}\right)\|_{L^\infty(\mathbf{R}^d)}$ , which we can interpolate into an inequality like above where we can apply Littlewood-Paley theory. To do this we perform another decomposition, writing

$$S_2 f = If + IIf$$

where

$$If(x) = \int_{Q_x} \sum_{|i+l(x)| \geqslant \tilde{\log}(B/A)} \left( m_i(D) (\mathbf{I}_{2Q_x} f)(y) - \int_{Q_x} m_i(D) (\mathbf{I}_{2Q_x} f)(z) \ dz \right) \chi_i(x,y) \ dy.$$

and

$$If(x) = \int_{Q_x} \sum_{|i+l(x)| \geqslant \log(B/A)} \left( m_i(D) (\mathbf{I}_{(2Q_x)^c} f)(y) - \int_{Q_x} m_i(D) (\mathbf{I}_{(2Q_x)^c} f)(z) \, dz \right) \chi_i(x,y) \, dy.$$

Now

$$||If||_{L^{\infty}} \leq \sup_{x} \int_{Q_{x}} \left( \sum |m_{i}(D)(\mathbf{I}_{2Q_{x}}f)|^{2} \right)^{1/2} dy \leq \sup_{x} |Q_{x}|^{-1/2} \left( \sum ||m_{i}(D)(\mathbf{I}_{2Q_{x}}f)||_{L^{2}(\mathbf{R}^{d})}^{2} \right)^{1/2} \leq A|Q_{x}|^{2} dy$$

#### Heo, Nazarov, and Seeger

In this chapter we give a description of the techniques of Heo, Nazarov, and Seeger's paper 2011 *Radial Fourier Multipliers in High Dimensions* [11]. Recall that if  $m \in L^{\infty}(\mathbf{R}^d)$  is the symbol of a Fourier multiplier operator  $T_m: L^2(\mathbf{R}^d) \to L^2(\mathbf{R}^d)$ . We let  $\|m\|_{M^p(\mathbf{R}^d)}$  denote the operator norm of  $T_m$  from  $L^p(\mathbf{R}^d)$  to  $L^p(\mathbf{R}^d)$ . The goal of this paper is to show that if  $m \in L^{\infty}(\mathbf{Z})$  is a radial function,  $d \ge 4$ , and  $\eta \in \mathcal{S}(\mathbf{R}^d)$  is nonzero, then

$$||m||_{M^p(\mathbf{R}^d)} \sim \sup_{t>0} t^{d/p} ||T_m(\mathrm{Dil}_t \eta)||_{L^p(\mathbf{R}^d)} \quad \text{for } p \in \left(1, \frac{2(d-1)}{d+1}\right),$$

where the implicit constant depends on p and  $\eta$ . Since

$$\sup_{t>0} t^{d/p} \|T_m(\mathrm{Dil}_t \eta)\|_{L^p(\mathbf{R}^d)} \sim \sup_{t>0} \frac{\|T_m(\mathrm{Dil}_t \eta)\|_{L^p(\mathbf{R}^d)}}{\|\mathrm{Dil}_t \eta\|_{L^p(\mathbf{R}^d)}}$$

we find that the boundedness of  $T_m$  on  $\mathcal{S}(\mathbf{R}^d)$  is equivalent to it's boundedness on the family  $\{\mathrm{Dil}_t \eta\}$ .

Note that the assumption of this result, if true, for m is compactly supported is equivalent to the assumption that  $\widehat{m}$  is in  $L^p(\mathbf{R}^n)$  (See Theorem 9.3 of this paper).

Note that, applying Littlewood-Paley theory, the assumption of this theorem is equivalent to the fact that

$$F_{p,2}^0 = t^d \| \left( \sum_{n=0}^{\infty} |m \widehat{\text{Dil}_{1/t}} \psi_n|^2 \right)^{1/2} \|_{L^p(\mathbb{R}^n)}$$

In Garrigós and Seeger's 2007 paper *Characterizations of Hankel Multi*pliers, it was proved that if  $\eta \in \mathcal{S}(\mathbf{R}^d)$  is a nonzero, radial Schwartz function, then

$$||m||_{M^p_{\text{rad}}(\mathbf{R}^d)} \sim \sup_{t>0} t^{d/p} ||T_m(\text{Dil}_t \eta)||_{L^p(\mathbf{R}^d)} \quad \text{for } p \in \left(1, \frac{2d}{d+1}\right),$$

where  $M_{\text{rad}}^p(\mathbf{R}^d)$  is the operator norm of  $T_m$  from  $L_{\text{rad}}^p(\mathbf{R}^d)$  to  $L^p(\mathbf{R}^d)$ . Thus, at least in the range  $p \in \left(1, 2\frac{d-1}{d+1}\right)$ , boundedness of  $T_m$  on radial functions is equivalent to boundedness on all functions.

Another consequence of the techniques of this paper is that an 'end-point' result for local smoothing is proved for the wave equation.

#### 4.1 Discretized Reduction

It is obvious that

$$\|m\|_{M^p(\mathbf{R}^d)} \gtrsim_{\eta} \sup_{t>0} t^{d/p} \|T_m(\mathrm{Dil}_t \eta)\|_{L^p(\mathbf{R}^d)},$$

so it suffices to show that

$$\|m\|_{M^p(\mathbf{R}^d)} \lesssim_{\eta} \sup_{t>0} t^{d/p} \|T_m(\mathrm{Dil}_t \eta)\|_{L^p(\mathbf{R}^d)},$$

We will show this via a convolution inequality, which can also be used to prove local smoothing results for the wave equation.

Let  $\sigma_r$  be the surface measure for the sphere of radius r centered at the origin in  $\mathbf{R}^d$ . Also fix a nonzero, radial Schwartz function  $\psi \in \mathcal{S}(\mathbf{R}^d)$ . Given  $x \in \mathbf{R}^d$  and  $r \ge 1$ , define  $f_{xr} = \operatorname{Trans}_x(\sigma_r * \psi)$ , which we view as a smoothened indicator function on a thickness  $\approx 1$  annulus of radius r centered at x. Our goal is to prove the following inequality.

**Lemma 4.1.** *For any*  $a : \mathbb{R}^d \times [1, \infty) \to \mathbb{C}$ *, and*  $1 \le p < 2(d-1)/(d+1)$ *,* 

$$\left\|\int_{\mathbf{R}^d}\int_1^\infty a_r(x)f_{xr}\,dx\,dr\right\|_{L^p(\mathbf{R}^d)}\lesssim \left(\int_{\mathbf{R}^d}\int_1^\infty |a_r(x)|^pr^{d-1}drdx\right)^{1/p}.$$

The implicit constant here depends on p, d, and  $\psi$ .

Why is Lemma 4.1 useful? Suppose  $m: \mathbf{R}^d \to \mathbf{C}$  is a radial multiplier given by some function  $\tilde{m}: [1, \infty) \to \mathbf{C}$ , and we set  $a_r(x) = \tilde{m}(r)f(x)$  for some  $f: \mathbf{R}^d \to \mathbf{C}$ . Then it is simple to check that

$$\int_{\mathbf{R}^d} \int_1^\infty a_r(x) f_{xr} \, dx \, dr = K * \psi * f$$

where  $K(x) = |x|^{1-d} m(x)$ . In this setting, Lemma 4.1 says that

$$||K * \psi * f||_{L^p(\mathbf{R}^d)} \lesssim ||m||_{L^p(\mathbf{R}^d)} ||f||_{L^p(\mathbf{R}^d)},$$

which is clearly related to the convolution bound we want to show if  $\psi = \hat{\eta}$ , provided that we are dealing with a multiplier supported away from the origin. To understand Lemma 4.1 it suffices to prove the following discretized estimate.

**Theorem 4.2.** Fix a finite family of pairs  $\mathcal{E} \subset \mathbf{R}^d \times [1, \infty)$ , which is discretized in the sense that  $|(x_1, r_1) - (x_2, r_2)| \ge 1$  for each distinct pair  $(x_1, r_1)$  and  $(x_2, r_2)$  in  $\mathcal{E}$ . Then for any  $a : \mathcal{E} \to \mathbf{C}$  and  $1 \le p < 2(d-1)/(d+1)$ ,

$$\left\| \sum_{(x,r)\in\mathcal{E}} a_r(x) f_{xr} \right\|_{L^p(\mathbf{R}^d)} \lesssim \left( \sum_{(x,r)\in\mathcal{E}} |a_r(x)|^p r^{p-1} \right)^{1/p},$$

where the implicit constant depends on p, d, and  $\psi$ , but most importantly, is independent of  $\mathcal{E}$ .

*Proof of Lemma 4.1 from Lemma 4.2.* For any  $a : \mathbf{R}^d \times [1, \infty) \to \mathbf{C}$ ,

$$\int_{\mathbf{R}^d} \int_1^\infty a_r(x) f_{xr} = \int_{[0,1)^d} \int_0^1 \sum_{n \in \mathbf{Z}^d} \sum_{m \in \mathbf{Z}} \operatorname{Trans}_{n,m}(a f_{rx}) dr dx$$

Minkowski's inequality thus implies that

$$\left\| \int_{\mathbf{R}^{d}} \int_{1}^{\infty} a_{r}(x) f_{xr} \right\|_{L^{p}(\mathbf{R}^{d})} \leq \int_{[0,1)^{d}} \int_{0}^{1} \left\| \sum_{n \in \mathbf{Z}^{d}} \sum_{m \in \mathbf{Z}} \operatorname{Trans}_{n,m}(a f_{rx}) \right\|_{L^{p}(\mathbf{R}^{d})} dr dx$$

$$\leq \int_{[0,1)^{d}} \int_{0}^{1} \left( \sum_{n \in \mathbf{Z}^{d}} \sum_{m \in \mathbf{Z}} |a_{r}(x)|^{p} r^{p-1} \right)^{1/p} dr dx$$

$$\leq \left( \int_{[0,1)^{d}} \int_{0}^{1} \sum_{n \in \mathbf{Z}^{d}} \sum_{m \in \mathbf{Z}} |a_{r}(x)|^{p} r^{p-1} dr dx \right)^{1/p}$$

$$= \left( \int_{\mathbf{R}^{d}} \int_{1}^{\infty} |a_{r}(x)|^{p} r^{d-1} dr dx \right)^{1/p}.$$

Lemma 4.2 is further reduced by considering it as a bound on the operator  $a \mapsto \sum_{(x,r) \in \mathcal{E}} a_r(x) f_{xr}$ . In particular, applying real interpolation, it suffices for us to prove a restricted strong type bound. Given any discretized set  $\mathcal{E}$ , let  $\mathcal{E}_k$  be the set of  $(x,r) \in \mathcal{E}$  with  $2^k \le r < 2^{k+1}$ . Then Lemma 4.2 is implied by the following Lemma.

**Lemma 4.3.** For any  $1 \le p < 2(d-1)/(d+1)$  and  $k \ge 1$ ,

$$\left\| \sum_{(x,r)\in\mathcal{E}_k} f_{xr} \right\|_{L^p(\mathbf{R}^d)} \lesssim 2^{k(d-1)} \#(\mathcal{E}_k)^{1/p} = 2^k \cdot (2^{k(d-p-1)} \#(\mathcal{E}_k))^{1/p}.$$

*Proof of Lemma 4.2 from Lemma 4.3.* Applying a dyadic interpolation result (Lemma 2.2 of the paper), Lemma 4.3 implies that

$$\|\sum_{(x,r)\in\mathcal{E}} f_{xr}\|_{L^p(\mathbf{R}^d)} \lesssim \left(\sum 2^{kp} 2^{k(d-p-1)} \#(\mathcal{E}_k)\right)^{1/p} = \left(\sum 2^{k(d-1)} \#(\mathcal{E}_k)\right)^{1/p}$$

This is a restricted strong type bound for Lemma 4.2, which we can then interpolate.  $\Box$ 

If  $\psi$  is compactly supported, and r is sufficiently large depending on the size of this support, then  $f_{xr}$  is supported on an annulus with centre

x, radius r, and thickness O(1). Thus  $||f_{xr}||_{L^p(\mathbf{R}^d)} \sim r^{(d-1)/p}$ , which implies that

$$\|\sum_{(x,r)\in\mathcal{E}_k} f_{xr}\|_{L^p(\mathbf{R}^d)} \gtrsim 2^{k(d-1)/p} \#(\mathcal{E}_k)^{1/p}.$$

Thus this bound can only be true if  $p \ge 1$ , and becomes tight when p = 1, where we actually have

$$\|\sum_{(x,r)\in\mathcal{E}_k} f_{xr}\|_{L^1(\mathbf{R}^d)} \sim 2^{k(d-1)} \#(\mathcal{E}_k)$$

because there can be no constructive interference in the  $L^1$  norm. Understanding the sum in Lemma 4.3 for 1 will require an understanding of the interference patterns of annuli with comparable radius. We will use almost orthogonality principles to understand these interference patterns.

**Lemma 4.4.** For any N > 0,  $x_1, x_2 \in \mathbb{R}^d$  and  $r_1, r_2 \ge 1$ ,

$$\begin{split} |\langle f_{x_1r_1}, f_{x_2r_2}\rangle| \lesssim_N (r_1r_2)^{(d-1)/2} (1+|r_1-r_2|+|x_1-x_2|)^{-(d-1)/2} \\ \sum_{\pm,\pm} (1+||x_1-x_2|\pm r_1\pm r_2|)^{-N}. \end{split}$$

In particular,

$$|\langle f_{x_1r_1}, f_{x_2r_2} \rangle| \lesssim \left(\frac{r_1r_2}{|(x_1, r_1) - (x_2, r_2)|}\right)^{(d-1)/2}$$

*Remark.* Suppose  $r_1 \le r_2$ . Then Lemma 4.4 implies that  $f_{x_1r_1}$  and  $f_{x_2r_2}$  are roughly uncorrelated, except when  $|x_1 - x_2|$  and  $|r_1 - r_2|$  is small, and in addition, one of the following two properties hold:

- $r_1 + r_2 \approx |x_1 x_2|$ , which holds when the two annuli are 'approximately' externally tangent to one another.
- $r_2 r_1 \approx |x_1 x_2|$ , which holds when the two annuli are 'approximately' internally tangent to one another.

Heo, Nazarov, and Seeger do not exploit the tangency information, though utilizing the tangencies seems important to improve the results they obtain. In particular, Laura Cladek's paper exploits this tangency information.

Proof. We write

$$\langle f_{x_1 r_1}, f_{x_2 r_2} \rangle = \left\langle \widehat{f}_{x_1 r_1}, \widehat{f}_{x_2 r_2} \right\rangle$$

$$= \int_{\mathbf{R}^d} \widehat{\sigma_{r_1} * \psi}(\xi) \cdot \widehat{\overline{\sigma_{r_2} * \psi}(\xi)} e^{2\pi i (x_2 - x_1) \cdot \xi} d\xi$$

$$= (r_1 r_2)^{d-1} \int_{\mathbf{R}^d} \widehat{\sigma}(r_1 \xi) \widehat{\overline{\sigma}(r_2 \xi)} |\widehat{\psi}(\xi)|^2 e^{2\pi i (x_2 - x_1) \cdot \xi} d\xi.$$

Define functions A and B such that  $B(|\xi|) = \hat{\sigma}(\xi)$ , and  $A(|\xi|) = |\hat{\psi}(\xi)|^2$ . Then

$$\langle f_{x_1r_1}, f_{x_2r_2} \rangle = C_d(r_1r_2)^{d-1} \int_0^\infty s^{d-1} A(s) B(r_1s) B(r_2s) B(|x_2 - x_1|s) \ ds.$$

Using well known asymptotics for the Fourier transform for the spherical measure, we have

$$B(s) = s^{-(d-1)/2} \sum_{n=0}^{N-1} (c_{n,+} e^{2\pi i s} + c_{n,-} e^{-2\pi i s}) s^{-n} + O_N(s^{-N}).$$

But now substituting in, assuming A(s) vanishes to order 100N at the origin, we conclude that

$$\begin{split} \langle f_{x_1r_1}, f_{x_2r_2} \rangle &= C_d \left( \frac{r_1r_2}{|x_1 - x_2|} \right)^{(d-1)/2} \sum_{n,\tau} c_{n,\tau} r_1^{-n_1} r_2^{-n_2} |x_2 - x_1|^{-n_3} \\ & \left\{ \int_0^\infty A(s) s^{-(d-1)/2} s^{-n_1 - n_2 - n_3} e^{2\pi i (\tau_1 r_1 + \tau_2 r_2 + \tau_3 |x_2 - x_1|) s} \, ds \right\} \\ &\lesssim_N \left( \frac{r_1 r_2}{|x_1 - x_2|} \right)^{\frac{d-1}{2}} \left( 1 + \frac{1}{|x_1 - x_2|^N} \right) \sum_{\tau} (1 + |\tau_1 r_1 + \tau_2 r_2 + \tau_3 |x_2 - x_1||)^{-5N} \\ &\lesssim_N \left( \frac{r_1 r_2}{|x_1 - x_2|} \right)^{\frac{d-1}{2}} \left( 1 + \frac{1}{|x_1 - x_2|^N} \right) \sum_{\tau} (1 + |\tau_1 \tau_3 r_1 + \tau_2 \tau_3 r_2 + |x_2 - x_1||)^{-5N} \, . \end{split}$$

This gives the result provided that  $1+|x_1-x_2|\geqslant |r_1-r_2|/10$  and  $|x_1-x_2|\geqslant 1$ . If  $1+|x_1-x_2|\leqslant |r_1-r_2|/10$ , then the supports of  $f_{x_1r_1}$  and  $f_{x_2r_2}$  are disjoint, so the inequality is trivial. On the other hand, if  $|x_1-x_2|\leqslant 1$ ,

then the bound is trivial by the last sentence unless  $|r_1 - r_2| \le 10$ , and in this case the inequality reduces to the simple inequality

$$\langle f_{x_1r_1}, f_{x_2r_2} \rangle \lesssim_N (r_1r_2)^{(d-1)/2}.$$

But this follows immediately from the Cauchy-Schwartz inequality.  $\Box$ 

The exponent (d-1)/2 in Lemma 4.4 is too weak to apply almost orthogonality directly to obtain  $L^2$  bounds on  $\sum_{(x,r)\in\mathcal{E}_k} f_{xr}$ . To fix this, we apply a 'density decomposition', somewhat analogous to a Calderon Zygmund decomposition, which will enable us to obtain  $L^2$  bounds. We say a 1-separated set  $\mathcal{E}$  in  $\mathbf{R}^d \times [R, 2R)$  is of *density type* (u, R) if

$$\#(B \cap \mathcal{E}) \leq u \cdot \operatorname{diam}(B)$$

for each ball B in  $\mathbb{R}^{d+1}$  with diameter  $\leq R$ . A covering argument then shows that for any ball B,

$$\#(B \cap \mathcal{E}) \lesssim_d u \cdot \left(1 + \frac{\operatorname{diam}(B)}{R}\right)^d \cdot \operatorname{diam}(B).$$

(NOTE: WE MIGHT BE ABLE TO DO BETTER USING THE FACT THAT  $\mathcal{E} \subset \mathbf{R}^d \times [R, 2R)$ , USING THE VALUE R).

**Theorem 4.5.** For any 1-separated set  $\mathcal{E} \subset \mathbf{R}^d \times [R, 2R)$ , we can consider a disjoint union  $\mathcal{E} = \bigcup_{k=1}^{\infty} \bigcup_{m=1}^{\infty} \mathcal{E}_k(2^m)$  with the following properties:

- For each k and m,  $\mathcal{E}_k(2^m)$  has density type  $(2^m, 2^k)$ .
- If B is a ball of radius  $\leq 2^m$  containing at least  $2^m rad(B)$  points of  $\mathcal{E}_k$ , then

$$B \cap \mathcal{E}_k \subset \bigcup_{m' \geqslant m} \mathcal{E}_k(2^{m'}).$$

• For each k and m, there are disjoint balls  $\{B_i\}$  of radius at most  $2^k$ , such that

$$\sum_{i} rad(B) \leqslant \frac{\#(\mathcal{E}_k)}{u}$$

such that  $\bigcup B_i^*$  covers  $\bigcup_{m' \geqslant m} \mathcal{E}_k(2^{m'})$ , where  $B_i^*$  denotes the ball with the same center as  $B_i$  but 5 times the radius.

Proof. Vitali Covering.

Given a sum  $F = \sum_{(x,r)\in\mathcal{E}} f_{xr}$ , decompose  $\mathcal{E}$  as  $\mathcal{E}_k(2^m)$ , and define  $F_{km}$  to be the sum over  $\mathcal{E}_k(2^m)$ . It follows from the convering argument above that measure of the support of  $F_{km}$  is  $O(2^{k(d-1)-m}\#(\mathcal{E}_k))$ . We define  $F_m = \sum_k F_{km}$ . To Prove Lemma 4.3, it will suffice to prove the following  $L^2$  estimate on  $F_m$ .

**Lemma 4.6.** Suppose  $\mathcal{E}$  is a set with density type  $(2^m, 2^k)$ . Then

$$\left\| \sum_{(x,r)\in\mathcal{E}} f_{x,r} \right\|_{L^2(\mathbf{R}^d)} \lesssim 2^{m/(d-1)} \sqrt{\log(2+2^m)} 2^{k(d-1)/2} \cdot \#(\mathcal{E}_k)^{1/2}.$$

Proof of Lemma 4.3 from Lemma 4.6. We have

$$||F_{km}||_{L^2(\mathbf{R}^d)} \lesssim 2^{m/(d-1)} \sqrt{\log(2+2^m)} 2^{k(d-1)/2} \#(\mathcal{E}_k)^{1/2}.$$

If we interpolate this bound with the support bound for  $F_{km}$ , we conclude that for 0 ,

$$||F_{km}||_{L^{p}(\mathbf{R}^{d})} \leq |\operatorname{Supp}(F_{km})|^{1/p-1/2} ||F_{km}||_{L^{2}(\mathbf{R}^{d})}$$

$$\lesssim (2^{(k(d-1)-m)})^{1/p-1/2} 2^{m/(d-1)} \sqrt{\log(2+2^{m})} 2^{k(d-1)/2} \#(\mathcal{E}_{k})^{1/2}$$

$$\lesssim 2^{m(1/p_{d}-1/p)} \sqrt{\log(2+2^{m})} \cdot 2^{k(d-1)/p} \#(\mathcal{E}_{k})^{1/2}.$$

where  $p_d = 2(d-1)/(d+1)$ . This bound is summable in m for  $p < p_d$ , which enables us to conclude that

$$||F_k||_{L^p(\mathbf{R}^d)} \lesssim_p 2^{k(d-1)/p} \#(\mathcal{E}_k)^{1/2}.$$

NOTE: THIS SEEMS LIKE A TYPO. Thus for  $1 \le p < p_d$ , we obtain the bound stated in Lemma 4.3.

Proving 4.6 is where the weak-orthogonality bounds from Lemma 4.4 come into play.

*Proof of Lemma 4.6.* Split the interval  $[2^k, 2^{k+1}]$  into  $\lesssim 2^{(1-\alpha)k}$  intervals of length  $2^{\alpha k}$ , for some  $\alpha$  to be optimized later. For appropriate integers a, let  $I_a = [2^k + (a-1)2^{\alpha k}, 2^k + a2^{\alpha k}]$ . Let  $\mathcal{E}_a = \{(x,r) \in \mathcal{E} : r \in I_a\}$ , and write

 $F = \sum f_{xr}$ , and  $F_a = \sum_{(x,r) \in \mathcal{E}_a} f_{xr}$ . Without loss of generality, splitting up the sum appropriately, we may assume that the set of a such that  $\mathcal{E}_a$  is nonempty is 10-separated. We calculate that

$$||F||_{L^{2}(\mathbf{R}^{d})}^{2} = \sum_{a} ||F_{a}||_{L^{2}(\mathbf{R}^{d})}^{2} + 2 \sum_{a_{1} < a_{2}} |\langle F_{a_{1}}, F_{a_{2}} \rangle|$$

Given  $a_1 < a_2$ ,  $(x_1, r_1) \in \mathcal{E}_{a_1}$ , and  $(x_2, r_2)$  such that  $\langle f_{x_1 r_1}, f_{x_2 r_2} \rangle \neq 0$ , then  $|x_1 - x_2| \leq 2^{k+2}$ . Since  $|r_1 - r_2| \leq 2^{k+1}$  follows because  $r_1, r_2 \in [2^k, 2^{k+1}]$ , it follows that  $|(x_1, r_1) - (x_2, r_2)| \leq 3 \cdot 2^{k+1}$ . For each such pair, since we may assume that  $a_2 - a_1 \geq 10$  without loss of generality, it follows that  $|r_1 - r_2| \geq 2^{\alpha k}$ , and so applying Lemma 4.4 together with the density property, we conclude that for  $d \geq 4$ ,

$$\begin{split} |\langle f_{x_1r_1}, F_{a_2} \rangle| &\leqslant \sum_{l=1}^{(1-\alpha)k+1} \sum_{2^l 2^{\alpha k} \leqslant |(x_1, r_1) - (x_2, r_2)| \leqslant 2^{l+1} 2^{\alpha k}} \langle f_{x_1r_1}, f_{x_2r_2}| \\ &\lesssim \sum_{l=1}^{(1-\alpha)k+1} (2^m 2^l 2^{\alpha k}) \left(\frac{2^{2k}}{2^l 2^{\alpha k}}\right)^{(d-1)/2} \\ &\lesssim \sum_{l=1}^{(1-\alpha)k+1} 2^m (2^k)^{(d-1) - (d-3)/2\alpha} 2^{-(d-3)/2 \cdot l} \\ &\lesssim 2^m (2^k)^{(d-1) - (d-3)/2\alpha}. \end{split}$$

Summing over all choices of  $x_1$  and  $r_1$ , we conclude that

$$2\sum_{a_1 < a_2} |\langle F_{a_1}, F_{a_2} \rangle| \lesssim 2^m (2^k)^{(d-1) - (d-3)/2\alpha} \#(\mathcal{E}).$$

On the other hand, TODO

Cladek: Improvements to Radial Multiplier Problem Using Incidence Geometry

# Mockenhaupt, Seeger, and Sogge: Exploiting Wave-Equation Periodicity

The main goal of the paper *Local Smoothing of Fourier Integral Operators* and *Carleson-Sjölin Estimates* is to prove local regularity theorems for a class of Fourier integral operators in  $I^{\mu}(Z,Y;\mathcal{C})$ , where Y is a manifold of dimension  $n \geq 2$ , and Z is a manifold of dimension n+1, which naturally arise from the study of wave equations. A consequence of this result will be a local smoothing result for solutions to the wave equation, i.e. that if  $2 , then there is <math>\delta$  depending on p and n, such that if  $T: Y \to Y \times \mathbf{R}$  is the solution operator to the wave equation, and Y is a compact manifold whose geodesics are periodic, then T is continuous from from  $L_c^p(Y)$  to  $L_{\alpha,\text{loc}}^p(Y \times \mathbf{R})$  for  $\alpha \leq -(n-1)|1/2-1/p|+\delta$ . Such a result is called local smoothing, since if we define  $Tf(t,x) = T_t f(x)$ , then the operator  $T_t$  is, for each t, a Fourier integral operator of order zero, with canonical relation

$$C_t = \{(x, y; \xi, \xi) : x = y + t\widehat{\xi}\},\$$

where  $\hat{\xi} = \xi/|\xi|$  is the normalization of  $\xi$ . Standard results about the regularity of hyperbolic partial differential equations show that each of the operators  $T_t$  is continuous from  $L_c^p(Y)$  To  $L_{\alpha,\text{loc}}^p(Y \times \mathbf{R})$  for  $\alpha \le -(n-1)|1/2-1/p|$ , and that this bound is sharp. Thus T is *smoothing* in the t variable, so that for any  $f \in L^p$ , the functions  $T_t f$  'on average' gain a regularity of  $\delta$  over the worst case regularity at each time. The local smoothing conjecture states that this result is true for any  $\delta < 1/p$ .

The class of Fourier integral operators studied are those satisfying the following condition: as is standard, the canonical relation  $\mathcal C$  is a conic Lagrangian manifold of dimension 2n+1. The fact that  $\mathcal C$  is Lagrangian implies  $\mathcal C$  is locally parameterized by  $(\nabla_\zeta H(\zeta,\eta),\nabla_\eta H(\zeta,\eta),\zeta,\eta)$ , where H is a smooth, real homogeneous function of order one. If we assume  $\mathcal C \to T^*Y$  is a submersion, then  $D_\xi[\nabla_\eta H(\zeta,\eta)]$  has full rank, which implies  $D_\eta[\nabla \xi H(\zeta,\eta)] = (D_\xi[\nabla_\eta H(\zeta,\eta)])^T$  has full rank, and thus the projection  $\mathcal C \to T^*Z$  is an immersion. We make the further assumption that the projection  $\mathcal C \to Z$  is a submersion, from which it follows that for each z in the image of this projection, the projection of points in  $\mathcal C$  onto  $T_z^*Z$  is a conic hypersurface  $\Gamma_z$  of dimension n. The final assumption we make is that all principal curvatures of  $\Gamma_z$  are non-vanishing.

*Remark.* The projection properties of  $\mathcal{C}$  imply that, in  $T^*(Z \times Y)$ , there exists a smooth phase  $\phi$  defined on an open subset of  $Z \times T^*Y$ , homogeneous in  $T^*Y$ , such that locally we can write  $\mathcal{C}$  as  $(z, \nabla_z \phi(z, \eta), \nabla_\eta \phi(z, \eta), \eta)$  for  $\eta \neq 0$ . Then, working locally on conic sets,

$$\Gamma_z = \{(\nabla_z \phi(z, \eta))\},\$$

and the curvature condition becomes that the Hessian  $H_{\eta\eta}\langle\nabla_z\phi,\nu\rangle$  has constant rank n-1, where  $\nu$  is the normal vector to  $\Gamma_z$ . This is a natural homogeneous analogue of the Carleson-Sjölin condition for non-homogeneous oscillatory integral operators, i.e. the Carleson-Sjölin condition is allowed to assume  $H_{\eta\eta}\phi$  has rank n, which cannot be possible in our case, since  $\phi$  is homogeneous here. An approach using the analytic interpolation method of Stein or the Strichartz / Fractional Integral approach generalizes the Carleson-Sjölin theorem to show that for any smooth, non-homogeneous phase function  $\Phi: \mathbf{R}^{n+1} \times \mathbf{R}^n \to \mathbf{R}$ , and any compactly supported smooth amplitude a on  $\mathbf{R}^{n+1} \times \mathbf{R}^n$ . Consider the operators

$$T_{\lambda}f(z) = \int a(z,y)e^{2\pi i\lambda\Phi(z,y)}f(y) dy.$$

If the associated canonical relation C, if C projects submersively onto  $T^*\mathbf{R}^n$ , so that for each  $z \in \mathbf{R}^{n+1}$  in the image of the projection map C, the set  $S_z \subset \mathbf{R}^{n+1}$  obtained from the inverse image of the projection of  $C \to Z$  at z is a n dimensional hypersurface with k non-vanishing curvatures. Then for  $1 \le p \le 2$ ,

$$||T_{\lambda}f||_{L^{q}(\mathbf{R}^{n+1})} \lesssim \lambda^{-(n+1)/q} ||f||_{L^{p}(\mathbf{R}^{n})}.$$

where  $q = p^*(1 + 2/k)$ .

*Remark.* We can also see these assumptions as analogues in the framework of cinematic curvature, splitting the z coordinates into 'time-like' and 'space-like' parts. Working locally, because  $\mathcal{C} \to T^*Y$  is a submersion, we can consider coordinates z=(x,t) so that, with the phase  $\phi$  introduced above,  $D_x(\nabla_\eta \phi)$  has full rank n, and that  $\partial_t \phi(x,t,\eta) \neq 0$ . Then for each z=(x,t), we can locally write  $\partial_t \phi(x,t,\eta)=q(x,t,\nabla_x \phi(x,t,\eta))$ , homogeneous in  $\eta$ , and then

$$C = \{(x, t, y; \xi, \tau, \eta) : (x, \xi) = \chi_t(y, \eta), \tau = q(x, t, \xi)\},\$$

where  $\chi_t$  is a canonical transformation. Our curvature conditions becomes that  $H_{\xi\xi}q$  has full rank n-1. This is the cinematic curvature condition introduced by Sogge.

Under these assumptions, the paper proves that any Fourier integral operator T in  $I^{\mu-1/4}(Z,Y;\mathcal{C})$  maps  $L^2_c(Y)$  to  $L^q_{loc}(Z)$  if

$$2\left(\frac{n+1}{n-1}\right) \leqslant q < \infty$$
 and  $\mu \leqslant -n(1/2-1/q)+1/q$ .

and maps  $L_c^p(Y)$  to  $L_{loc}^p(Z)$  if

$$p > 2$$
 and  $\mu \le -(n-1)(1/2 - 1/p) + \delta(p, n)$ .

If we introduce time and space variables locally as in the remark above, any operator in  $I^{\mu-1/4}(Z,Y;\mathcal{C})$  can be written locally as a finite sum of operators of the form

$$Tf(x) = \int_{-\infty}^{\infty} T_t f(x),$$

where

$$T_t f(x) = \int a(t,x,\eta) e^{2\pi i \phi(x,t,y,\eta)} f(y) \, dy \, d\eta.$$

is a Fourier integral operator whose canonical relation is a locally a canonical graph, then the general theory implies that each of the maps  $T_t$  maps  $L_c^2(Y)$  to  $L_{loc}^q(X)$  if

$$2 \leqslant q \leqslant \infty$$
 and  $\mu \leqslant -n(1/2 - 1/q)$ 

so that here we get local smoothing of order 1/q, and also maps  $L_c^p(Y)$  to  $L_{loc}^p(X)$  if

$$1 and  $\mu \le -(n-1)|1/p - 1/2|$$$

so we get  $\delta(p,n)$  smoothing. A consequence of the smoothing, via Sobolev embedding, is a maximal theorem result for the operator  $T_t$ , i.e. that for any finite interval I, the operator

$$Mt = \sup_{t \in I} |T_t f|$$

maps  $L_c^p(Y)$  to  $L_{\mathrm{loc}}^p(X)$  if  $\mu < -(n-1)(1/2-1/p)-(1/p-\delta(p,n))$ . If the local smoothing conjecture held, we would conclude that, except at the endpoint  $T^*$  has the same  $L_c^p(Y)$  to  $L_{\mathrm{loc}}^p(X)$  mapping properties as each of the operators  $T_t$ . For the subsequence proof, we also get mixed norm estimates, such that for any finite interval I, the operator

$$Sf(x) = \left(\int_{I} |T_t f(x)|^2 dt\right)^{1/2}$$

if

$$2 \leqslant p \leqslant$$
 and  $q = 3p/(p-1)$  and  $\mu < -2(1/2 - 1/q)$ ,

is bounded from  $L_c^p(Y)$  to  $L_{loc}^q(X)$ .

Our main reason to focus on this paper is the results of the latter half of the paper applying these techniques to radial multipliers on compact manifolds with periodic geodesics. Thus we consider a compact Riemannian manifold M, such that the geodesic flow is periodic with minimal period  $2\pi\Pi$ . We consider  $m \in L^{\infty}(\mathbf{R})$ , such that  $\sup_{s>0} \|\beta \mathrm{Dil}_s m\|_{L^2_{\alpha}(\mathbf{R})} = A_{\alpha}$  is finite for some  $\alpha > 1/2$ . We define a 'radial multiplier' operator

$$Tf = \sum m(\lambda) P_{\lambda} f$$

where  $P_{\lambda}$  is the projection of f onto the space of eigenfunctions for the operator  $\sqrt{-\Delta}$  on M with eigenvalue  $\lambda$ . We can also write this operator as  $m(\sqrt{-\Delta})$ . Then the wave propogation operator  $e^{2\pi i t \sqrt{-\Delta}}$  is periodic of period  $\Pi$ . The Weyl formula tells us that the number of eigenvalues of  $\sqrt{-\Delta}$  which are smaller than  $\lambda$  is equal to  $V(M)\lambda^n + O(\lambda^{n-1})$ .

**Theorem 6.1.** Let  $m \in L^2_{\alpha}(\mathbf{R})$  be supported on (1,2), and assume  $\alpha > 1/2$ , then for  $2 \le p \le 4$ , and  $f \in L^p(M)$ ,

$$\left\| \sup_{2^i \leqslant \tau \leqslant 2^{i+1}} |Dil_{\tau} m(\sqrt{-\Delta}) f| \right\|_{L^p(M)} \lesssim_{\alpha} \|m\|_{L^2_{\alpha}(M)} \|f\|_{L^p(M)}.$$

*Proof.* To understand the radial multipliers we apply the Fourier transform, writing

$$T_{\tau}f = (\mathrm{Dil}_{\tau}m)(\sqrt{-\Delta})f = m(\sqrt{-\Delta}/\tau)f = \int_{-\infty}^{\infty} \tau \widehat{m}(t\tau)e^{2\pi it\sqrt{-\Delta}}f \ dt.$$

If we define  $\beta \in C_c^{\infty}((1/2,8))$  such that  $\beta(s) = 1$  for  $1 \le s \le 4$ , and set  $L_k f = \mathrm{Dil}_{2^k} \beta(\sqrt{-\Delta}) f$ , then for  $2^k \le \tau \le 2^{k+1}$ 

$$T_{\tau}f = (\mathrm{Dil}_{\tau}m)(\sqrt{-\Delta})f = (\mathrm{Dil}_{\tau}m \cdot \mathrm{Dil}_{2^k}\beta)(\sqrt{-\Delta}) = T_{\tau}L_kf.$$

so Cauchy-Schwartz implies that

$$\begin{aligned} |T_{\tau}f(x)| &= \left| \int_{-\infty}^{\infty} \tau \, \widehat{m}(\tau) e^{2\pi i t \sqrt{-\Delta}} L_{k}f(x) \, dt \right| \\ &\leq \|m\|_{L_{\alpha}^{2}(M)} \left( \int_{-\infty}^{\infty} \frac{\tau}{(1 + |t\tau|^{2})^{\alpha}} |e^{2\pi i t \sqrt{-\Delta}} L_{k}f(x)|^{2} \right)^{1/2} \\ &\leq \|m\|_{L_{\alpha}^{2}(M)} \left( \int_{-\infty}^{\infty} \frac{2^{k}}{(1 + |2^{k}t|^{2})^{\alpha}} |e^{2\pi i t \sqrt{-\Delta}} L_{k}f(x)|^{2} \right) \end{aligned}$$

Because of periodicity, it suffices to prove that for  $\alpha > 1/2$ ,

$$\left\| \left( \int_0^{\Pi} \frac{2^k}{(1+|2^k t|^2)^{\alpha}} |e^{2\pi i t \sqrt{-\Delta}} L_k f(x)|^2 dt \right)^{1/2} \right\|_{L^p(M)} \lesssim_{\alpha,p} \|f\|_{L^p(M)}.$$

This is a weighted combination of the wave propogators, roughly speaking, assigning weight  $2^i$  for  $t \leq 1/2^i$ , and assigning weight 1/t to values  $t \geq 1/2^k$ . For a fixed  $0 < \delta$ , we can split this using a partition of unity into a region where  $t \geq \delta$  and a region where  $t \leq \delta$ , where  $\delta$  is independent of i. The square function estimate above can be applied in the region where  $t \geq \delta$ , because the weighted square integral above has weight  $O_{\delta}(1)$  uniformly for all i. If we localize near  $t \leq 1/2^k$  by multiplying by  $\phi(2^k t)$  for

some compactly supported smooth  $\phi$  supported on  $|t|\lesssim 1$ , then the family of functions

$$\left\| \left( \int |\phi(t)e^{2\pi i(t/2^k)\sqrt{-\Delta}} L_k f(x)|^2 Dt \right)^{1/2} \right\|_{L^p_x} \lesssim \sup \|e^{2\pi i(t/2^k)\sqrt{-\Delta}} L_k f\|_{L^p_x}$$

$$a_k(t) = 2^{-k/2} \hat{\phi}(t/2^k) \beta(\tau/2^k)$$

it suffices to uniformly bound quantities of the form

$$\left\| \left( \int 2^k \phi(2^k t) |e^{2\pi i \sqrt{-\Delta}} L_k f(x)|^2 dt \right)^{1/2} \right\|_{L^p(M)} \lesssim_{\alpha, p} \|f\|_{L^p(M)}$$

For  $t \lesssim \delta$ , we localize around each x, so that we can use local coordinates, and apply the approximation to the wave equation to get more detailed information, i.e. finding a symbol  $b \in S^0$  such that for  $|t| \leq 2\delta$ , we have  $e^{2\pi i t \sqrt{-\Delta}} = A_t + R_t$ , where  $R_t$  has  $C^{\infty}$  kernel, uniformly for  $0 \leq t \leq 2\delta$ , and such that

$$A_t g(x) = \int b(x,t,\xi) e^{2\pi i (\xi \cdot x + t |\xi| + t^2 r(x,t,\eta))} \widehat{g}(\xi) d\xi,$$

where  $r \in S^1$  is homogeneous in  $\eta$ . We have  $||R_t g||_{L^p_x(M)L^2_t[0,\delta]} \lesssim 1$  because of the  $C^{\infty}$  kernel. Applying  $L_k$  gives a frequency cutoff, i.e. if  $b_i(x,t,\xi) = b(x,t,\xi)\beta(\xi/2^i)$ , then

$$A_t L_k f(x) = \int b_i(x, t, \xi) e^{2\pi i (\xi \cdot x + t \mid \xi \mid + t^2 r(x, t, \eta))} \widehat{f}(\eta) d\eta.$$

We now apply a dyadic decomposition to deal with the smaller values of t. Let us assume for simplicity of notation that  $\delta < 1$ , and then consider a partition of unity  $1 = \sum_{j=1}^{\infty} \phi(2^j t)$  for  $0 \le t \le 1$ , and such that  $\phi$  is localized near  $1/4 \le t \le 2$ , then our goal is to bound the quantities

$$\left\| \left( \int \phi(2^j t) \frac{2^i}{(1+|2^i t|^2)^{\alpha}} |A_t L_k f(x)|^2 dt \right)^{1/2} \right\|_{L^p(M)},$$

which are each proportional to

# Lee and Seeger: Decomposition Arguments For Estimating Fourier Integral Operators

In the paper Lebesgue Space Estimates For a Class of Fourier Integral Operators Associated With Wave Propogation, Lee and Seeger prove a variable coefficient version of the result of Heo, Nazarov, and Seeger, i.e. generalizing their result from proving results about the boundedness of radial Fourier multipliers on  $\mathbf{R}^n$  to certain Fourier integral operators satisfying the cinematic curvature condition. Let  $T = m(-\sqrt{\Delta})$  be a radial multiplier on  $\mathbf{R}^n$ , i.e. such that

$$Tf(x) = \int m(|\xi|)e^{2\pi i\xi\cdot(x-y)}f(y) d\xi dy.$$

If m is a symbol, then we can interpret T directly as a Pseudodifferential Operator. But Heo, Nazarov, and Seeger's result discuss families of multipliers m that are not even necessarily smooth, but do satisfy certain integrability conditions. To fix this, we assume a priori that we have applyied a decomposition argument, so we may assume m is compactly supported away from the origin. Then (by Paley-Wiener)  $\hat{m}$  is a smooth symbol of some finite order satisfying some integrability properties, which indicates how we might apply the theory of Fourier integral operators, i.e. by taking the Fourier transform of m, we get that

$$Tf(x) = \int \widehat{m}(\rho) e^{2\pi i [\rho|\xi| + \xi \cdot (x - y)]} f(y) \, d\rho \, d\xi \, dy.$$

This is 'almost' a Fourier integral operator, except the phase is not smooth unless  $\hat{m}$  is supported away from the origin (fixed by a decomposition argument), and the phase is non-homogeneous. To fix the non-homogeneity, we just isolate the operator in  $\rho$ , writing

$$Tf(x) = \int_{-\infty}^{\infty} \widehat{m}(\rho) T_{\rho} f(x) d\rho,$$

where

$$T_{\rho}f(x) = e^{2\pi i\rho\sqrt{-\Delta}}f(x) = \int e^{2\pi i[\rho|\xi| + \xi \cdot (x-y)]}f(y) d\xi dy$$

is the propogation operator for the half-wave equation  $\partial_t u = \sqrt{-\Delta} \cdot u$ . It has phase  $\phi(x,y,\xi) = \rho|\xi| + \xi \cdot (x-y)$ , and thus we have a stationary frequency value when  $x = y - \rho \hat{\xi}$ , where  $\hat{\xi} = \xi/|\xi|$  is the normalization of  $\xi$ . This has canonical relation

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