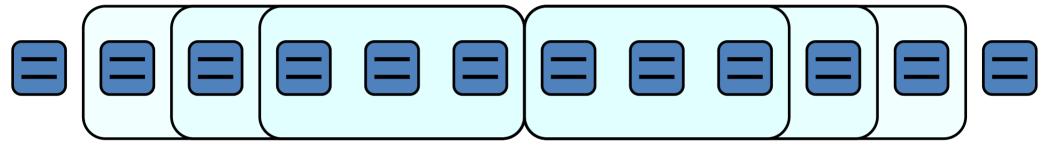
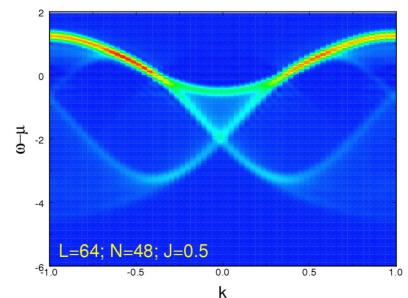
Understanding strongly correlated quantum matter

Adrian Feiguin





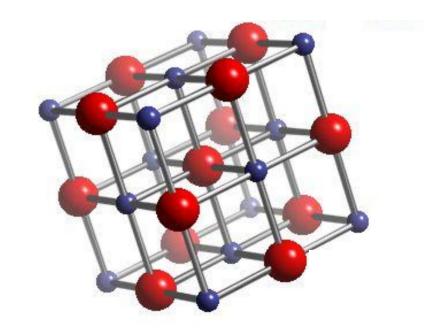




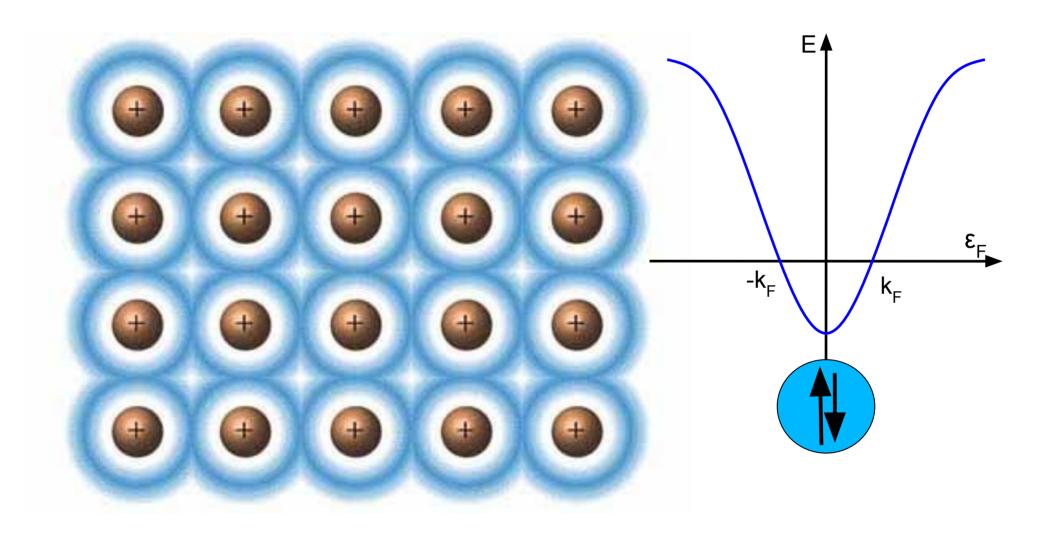
Condensed matter

- -Phenomena emerging from the collective behavior of a large number of degrees of freedom inside a solid.
- -We start with putting together single constituents (particles) with well defined properties (mass, spin, charge), and turning on interactions.

-We observe different phases of matter, and transitions between them.



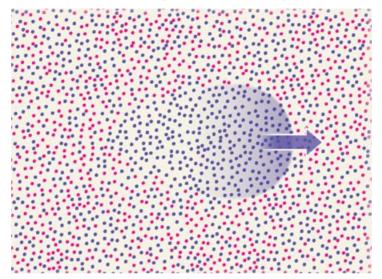
The Fermi Sea



Weakly interacting electrons

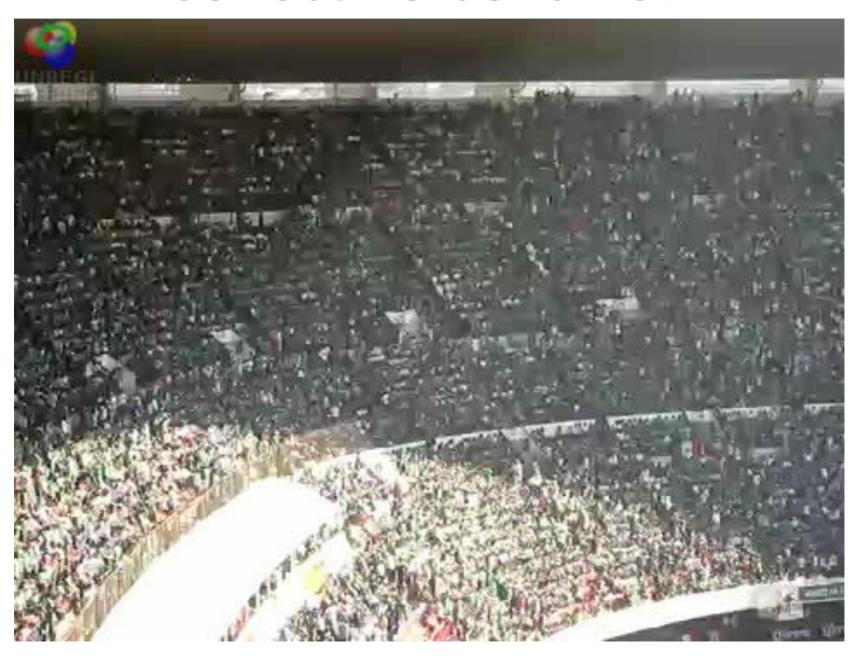


Electrons in solids: We think "quasi-particles"

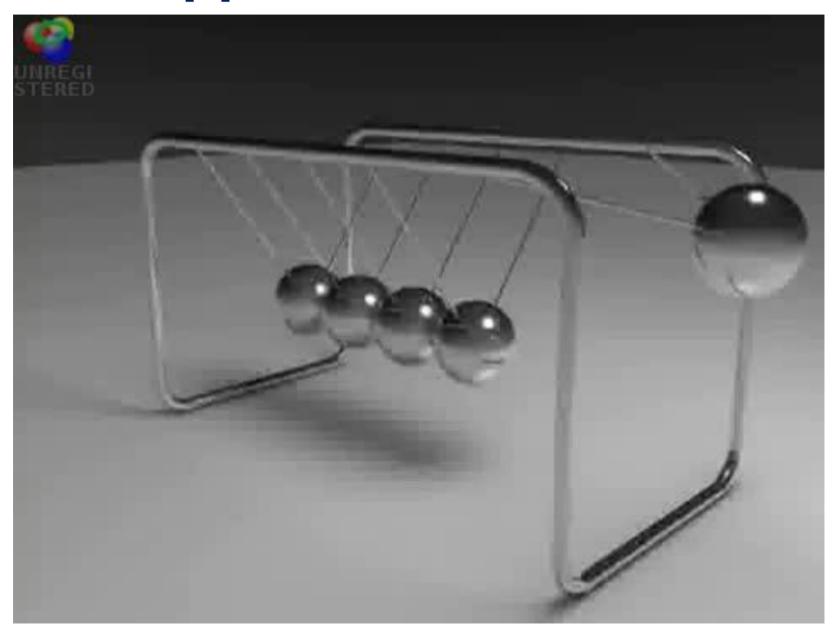


In a normal Fermi liquid, each constituent particle (dots) is surrounded by a polarization cloud of particles (blue shading) that is dragged along by its motion (arrow). The constituent particle and 'dressing cloud', which together form what is known as a quasi-particle, act like weakly interacting particles.

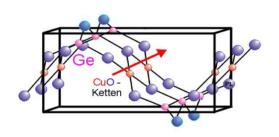
Collective behavior

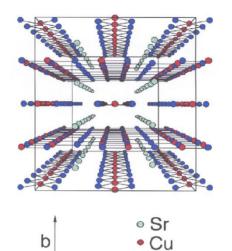


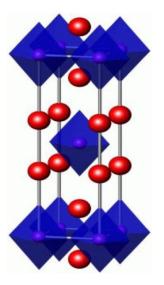
What happens in one dimension?



Low dimensional compounds



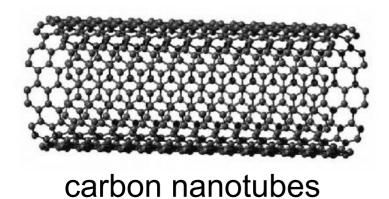


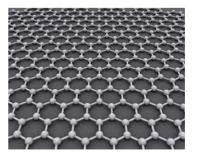


CuGeO₃ (chains)

Sr Cu O 41 (ladders)

La₂CuO₄ (planes)

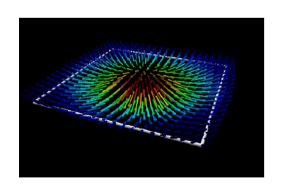


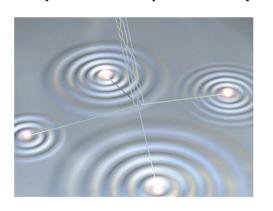


graphene

Strongly correlated problems

- High temperature superconductivity
- Colossal magnetoresistance (Manganites)
- Mott-Hubbard transition
- Non Fermi liquid behaviour
- Anderson and Kondo impurity problems
- Transport in low dimensions (quantum dots)
- Fractional Quantum Hall Effect
- Spin liquids; topological phases (novel quantum phases)







Emergent behavior defies intuition!



Areas of interest

- Exotic phases of matter of quantum origin: quantum magnetism, superconductivity...
- Quantum transport and non-equilibrium physics. "Mottronics" and molecular devices. Materials for light harvesting applications.
- Spintronics, using spin for logical gates and computing. NV centers in diamond.
- Basic understanding of quantum phenomena using cold atoms.
- New computational algorithms for quantum chemistry and condensed matter.

Why we need numerical methods?

- Analytical methods are no longer reliable!
- We need numerical methods to:
 - Validate predictions from theory
 - Understand the physics when theory fails
 - Compare competing orders (phases)
 - Find new phases
 - Make quantitative predictions for experiments

"Hydrogen" molecule

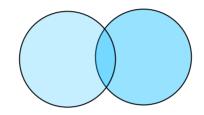
Case 1: atoms are far apart



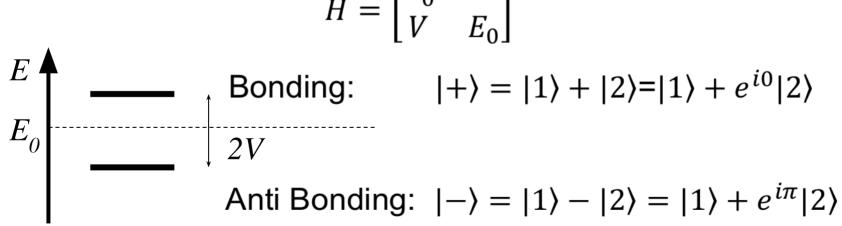
$$H = \begin{bmatrix} E_0 & 0 \\ 0 & E_0 \end{bmatrix}$$

"Hydrogen" molecule

Case 2: we bring atoms closer together

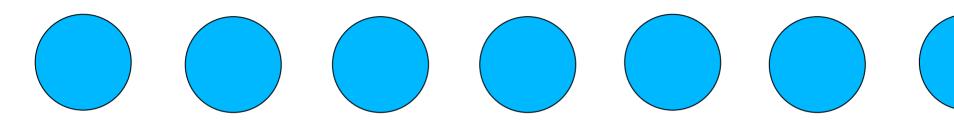


$$H = \begin{bmatrix} E_0 & V \\ V & E_0 \end{bmatrix}$$



We can classify these states as $|k\rangle = |1\rangle + e^{ik}|2\rangle$ with $k=0,\pi$

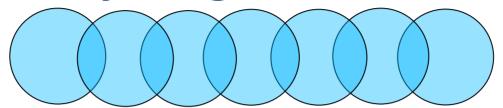
"Hydrogen" chain



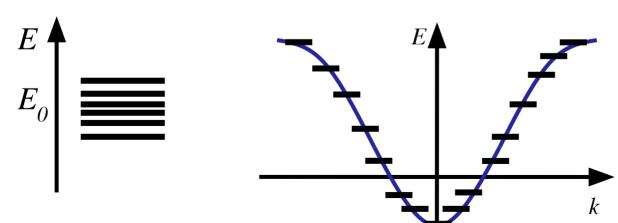
$$H = \begin{pmatrix} E_0 & 0 & \cdots & 0 \\ 0 & E_0 & & 0 \\ \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & E_0 \end{pmatrix}$$

$$E
brack \\
E_0
brack N-fold degenerate$$

"Hydrogen" chain

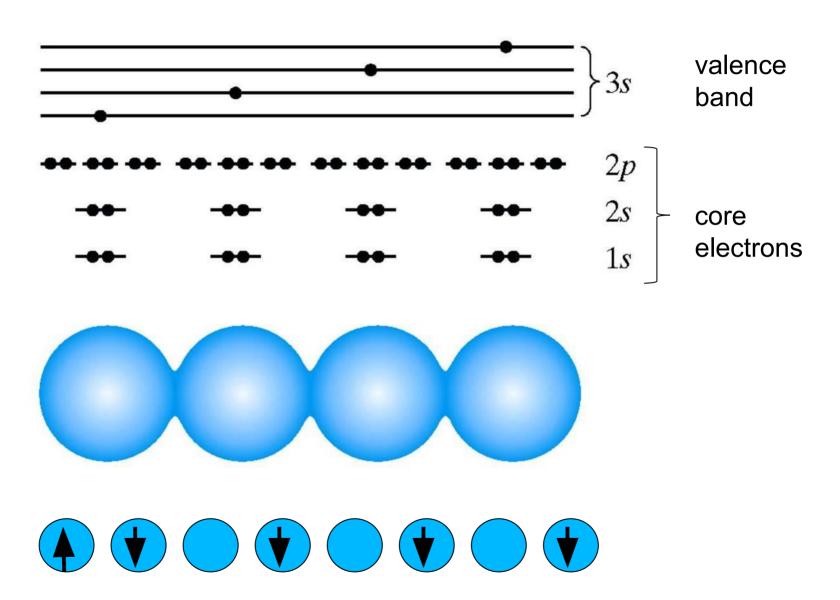


$$H = \begin{bmatrix} E_0 & V & 0 & \dots & 0 & 0 \\ V & E_0 & V & & & 0 \\ 0 & V & E_0 & & & 0 \\ \vdots & & \ddots & & \vdots \\ 0 & & & E_0 & V \\ 0 & 0 & \dots & V & E_0 \end{bmatrix}$$



State now lie on an "energy band"

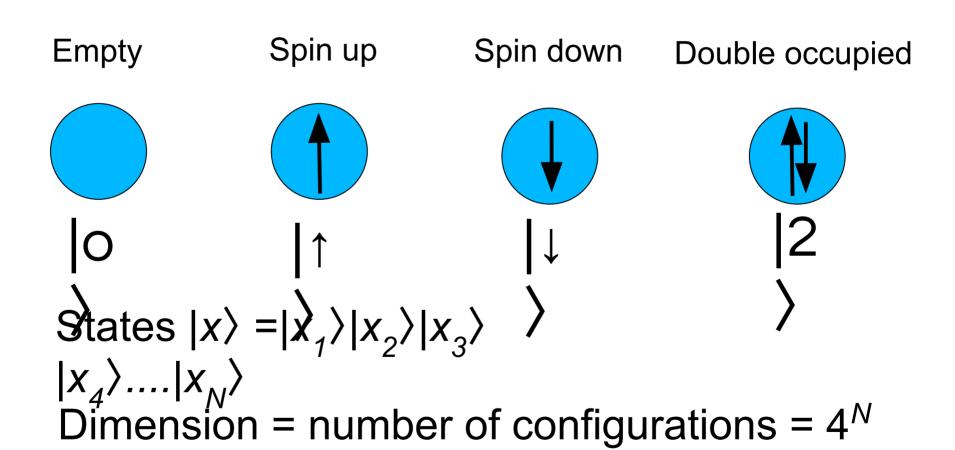
Effective models



Basis of states

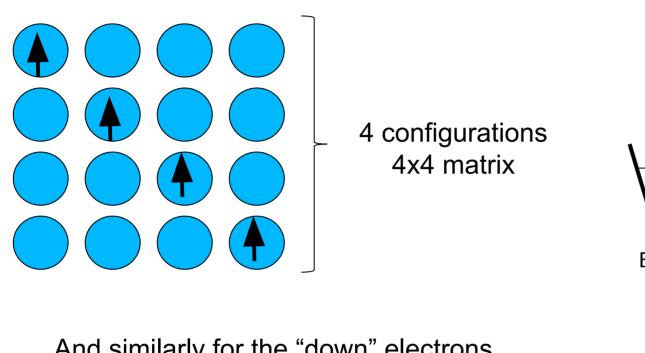
Occupation number representation

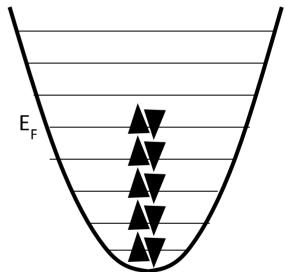
(1 orbital per site, spin ½)



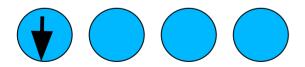
N: number of lattice sites

Single particle vs. many body picture



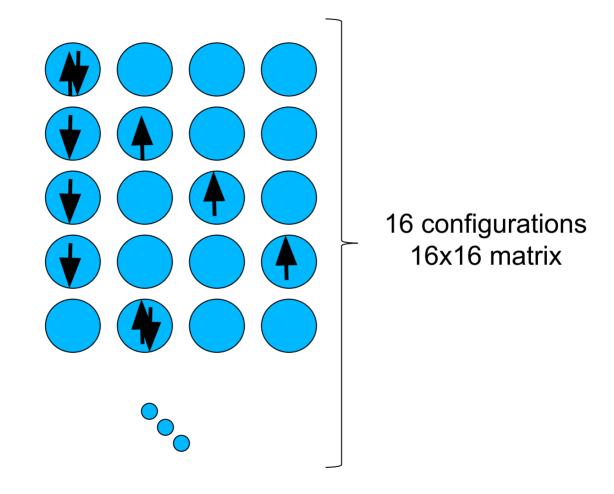


And similarly for the "down" electrons



The state of the system is a "product state" of single particle states. We only need to solve the one-particle problem.

Single particle vs. many body picture



The state of the system is "highly entangled". It cannot be written as a "product state", and the behavior of each electron is dictated by the behavior of the rest.

Numerical methods

- Quantum Monte Carlo
- Exact Diagonalization (Lanczos)
- Density Matrix Renormalization Group

Exact diagonalization

"brute force" diagonalization of the Hamiltonian matrix.

Schrödinger's Equation:

$$H|x\rangle = E|x\rangle$$
 |x\rangle : eigenstate

H: Hamiltonian operator

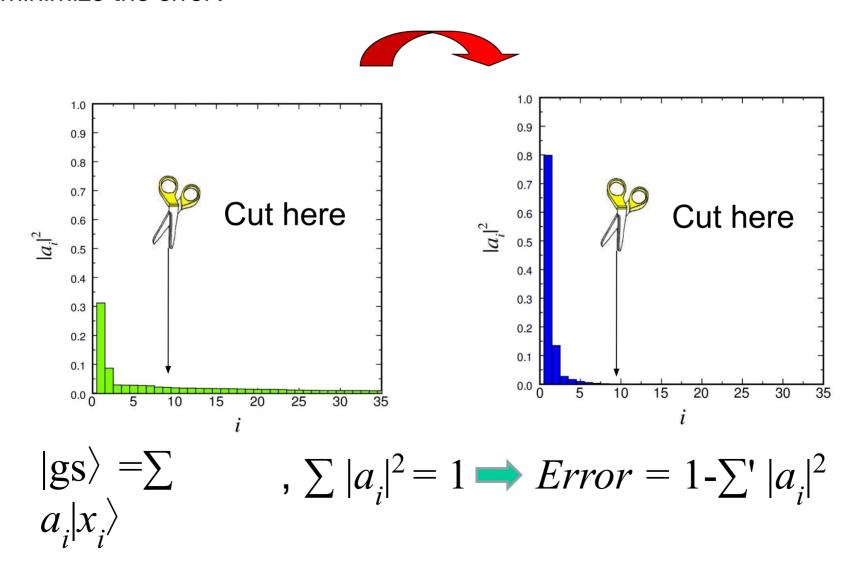
E: eigenvalue (ENERGY)



- Direct evaluation of the ground state of the full Hamiltonian...
- Excitation energies (gaps).
- Correlation functions.
- ... anything you want to know about the state... but... only small systems

Truncating the Hilbert space

Can we rotate our basis to one where the weights are more concentrated, to minimize the error?



"Classical" analogy Image compression algorithms (e.g.

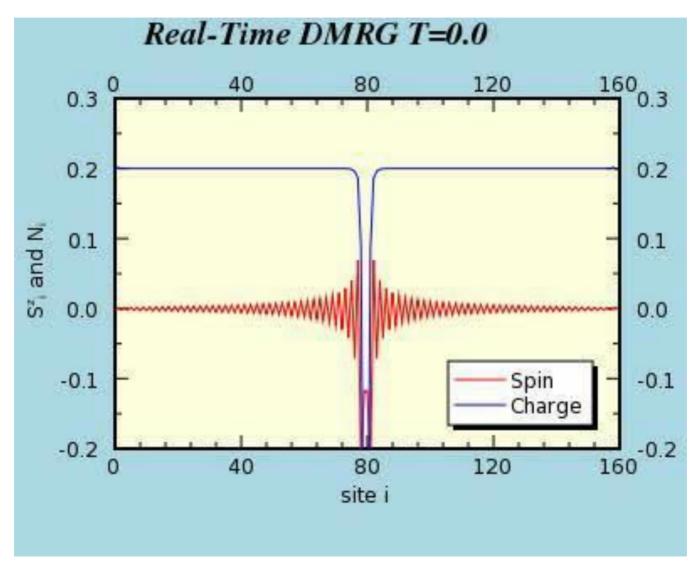
Jpeg)





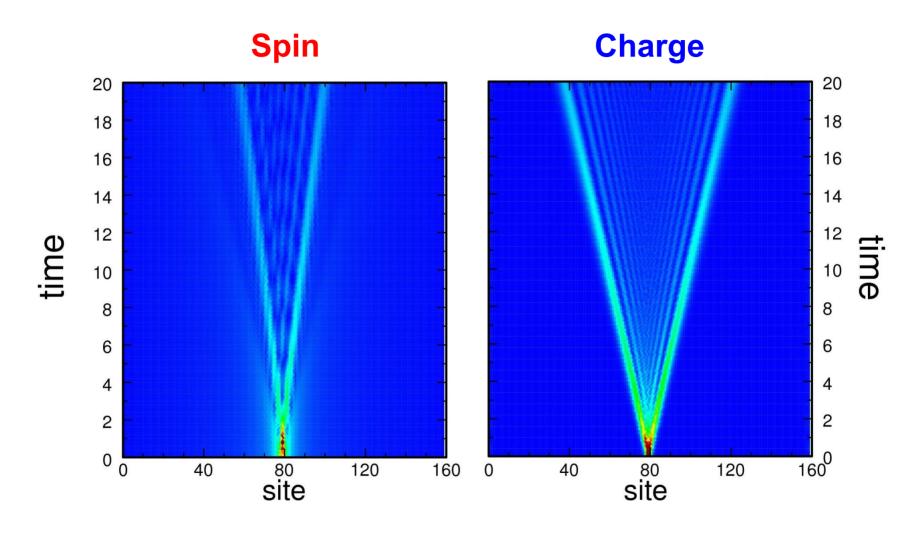
We want to achieve "lossless compression" ... or at least minimize *the loss of information*

Real-time simulation Half-filled Hubbard model (L=160, U=4)



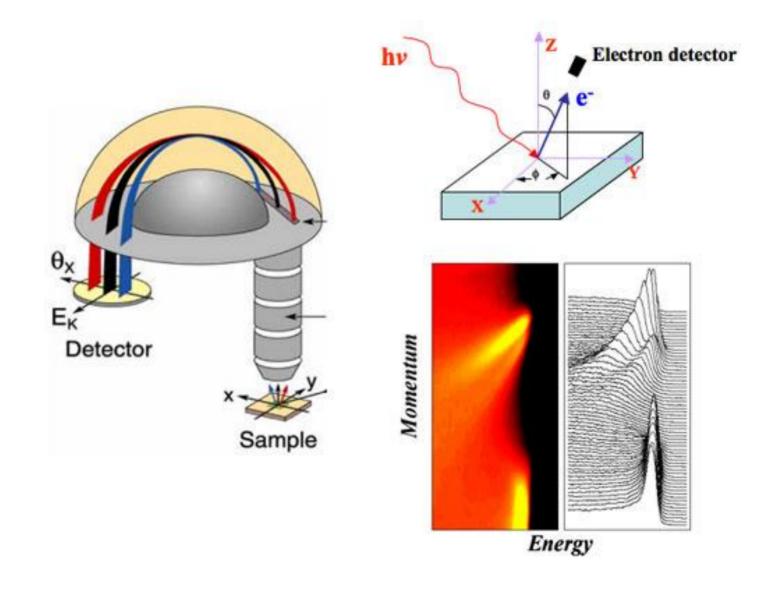
S.R.White and AEF, PRL (2004), Daley et al, J. Stat. Mech.: Theor. Exp. (2004); AEF and S.R.White, PRB (2005), Rapid Comm.

Lightcones

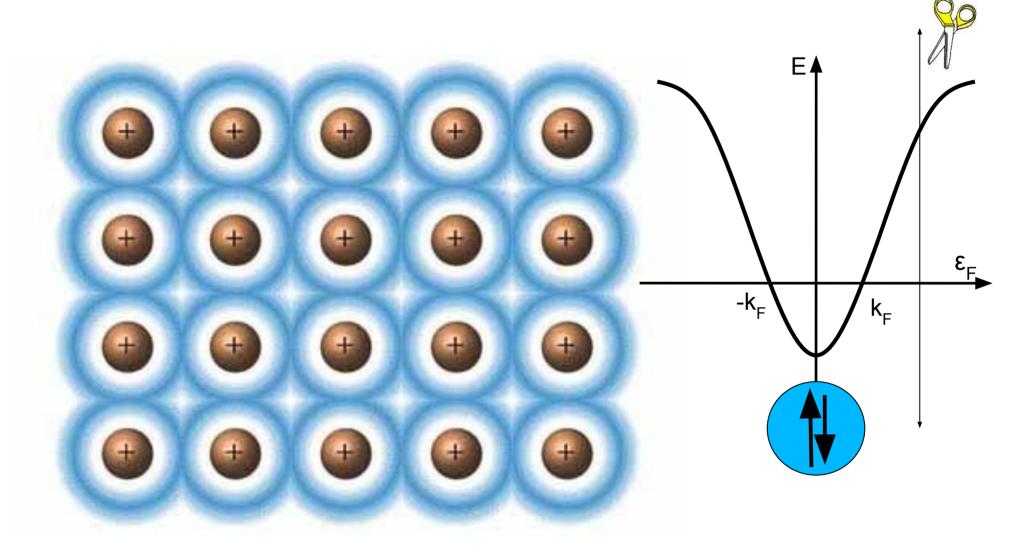


Spin and charge propagate at different velocities

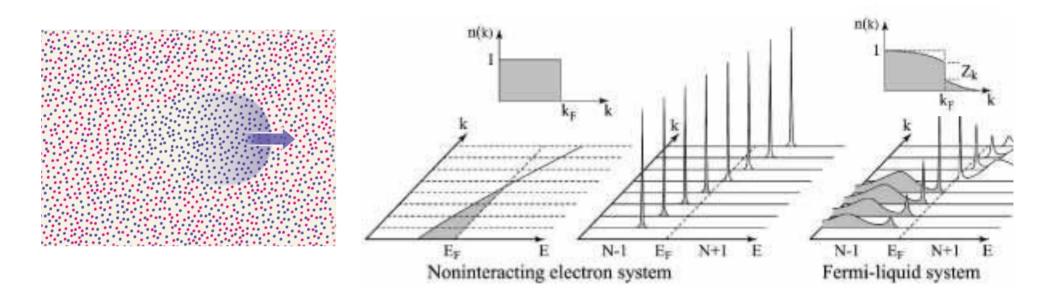
ARPES



The Fermi Sea



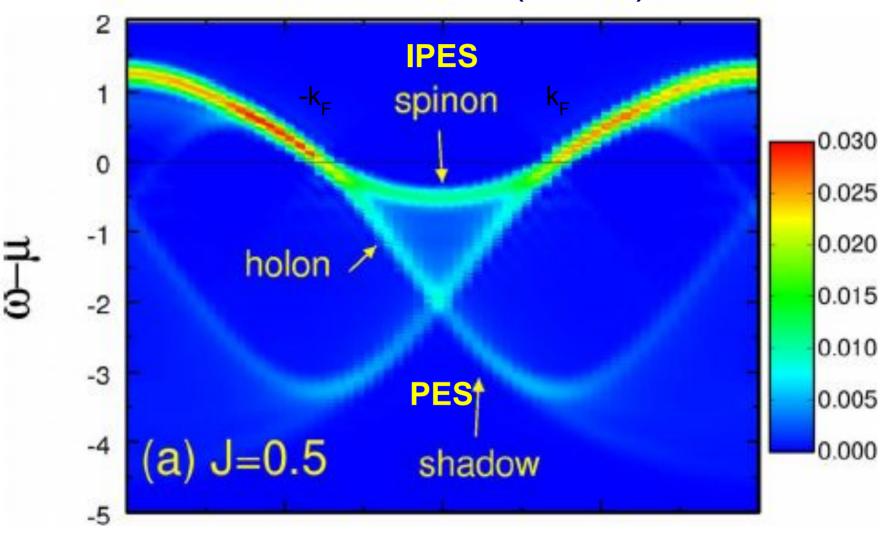
Electrons in solids: We think "quasi-particles"



In a normal Fermi liquid, each constituent particle (dots) is surrounded by a polarization cloud of particles (blue shading) that is dragged along by its motion (arrow). The constituent particle and 'dressing cloud', which together form what is known as a quasi-particle, act like weakly interacting particles.

ARPES at T=0

1D t-J model (J=0.5)



Variational Neural Networks

$$|\psi\rangle = \sum_{\{\sigma_1 \dots \sigma_N\}} \psi_{\sigma_1 \dots \sigma_N} |\sigma_1 \dots \sigma_N\rangle$$

$$\psi(\vec{\sigma}^z, \vec{a}, \vec{b}, W) = \sum_{h_1, h_2, \dots, h_M} e^{-E(\vec{\sigma}, \vec{h})},$$

with -

$$-E(\vec{\sigma}, \vec{h}) = \sum_{i=1}^{N} a_i \sigma_i^z + \sum_{i=1}^{N} b_i h_i + \sum_{i=1}^{N} \sum_{j=1}^{M} W_{ij} \sigma_i^z h_j,$$

 $h_i = \pm 1$ Visible Hidden

Figure 1: A restricted Boltzmann machine. Physical degrees of freedom are represented by the visible layer.

where $h_i = \{-1,1\}$ are "hidden variables", W's are weights and a,b are biases

"Training" the NN corresponds a "variational minimization": finding the weights and biases *a*,*b* and *W* that minimize the "energy functional"

$$E = \frac{\langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle}$$
 We use "Variational Monte Carlo"

Excited states (spectrum)

We want to calculate the "Green's function"

$$G_{ij}(z) = \langle \psi | A_i^{\dagger} \frac{1}{z - \hat{H}} A_j | \psi \rangle$$

where A is some operator of interest and $z=E_0+\omega+i\eta$ We introduce auxiliary states:

$$|A_i\rangle = \hat{A}_i|\psi\rangle$$

 $|\chi_j(z)\rangle = \frac{1}{z - \hat{H}}|A_j\rangle,$

where $|\chi_j(z)\rangle$ is called the "correction vector". Explicitly, $|\chi_j(z)\rangle$ can be obtained by solving the equation:

$$(z - \hat{H})|\chi_j(z)\rangle = \hat{A}_j|\psi\rangle = |A_j\rangle.$$

The spectral function is defined as the imaginary part of the Green's function, $A_{ij}(\omega) = -\frac{1}{\pi} \text{Im} G_{ij}(z)$, or:

$$A_{ij}(\omega) = -\frac{1}{\pi} \operatorname{Im} \langle A_i | \chi_j(z) \rangle.$$

The calculation has to be carried out using Monte Carlo for each value of frequencies *z*!!!

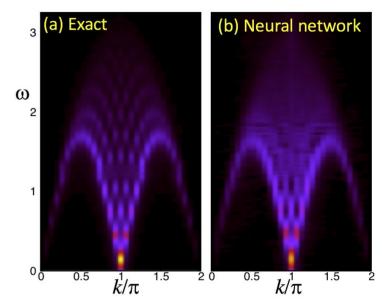


Figure 2: Excitation spectrum of a quantum spin chain using quantum ML (Ref.[5]), compared to the exact result.

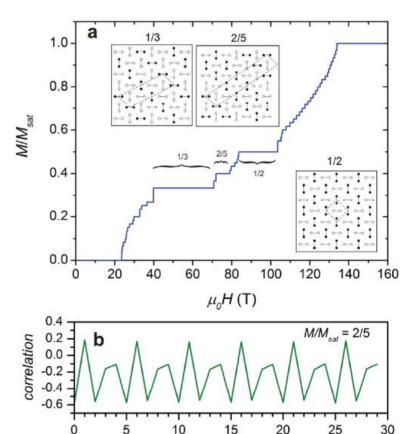
Applications (experiments)

Quantum magnetism

On-going collaborations with experimental groups in LANL (M. Jaime), ETH Zurich (H. Ott and F. Casola), PSI (C. Ruegg)



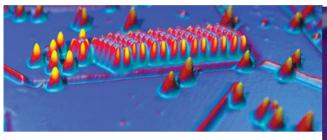
Magnetostriction and magnetic texture to 100.75
Tesla in frustrated SrCu₂(BO₃)₂
Marcelo Jaime^{a,b,1}, Ramzy Daou^{c,d}, Scott A. Crooker^{a,b},
Franziska Weickert^b, Atsuko Uchida^{a,b}, Adrian E.
Feiguin^e, Cristian D. Batista^f, Hanna A. Dabkowska^g, and
Bruce D. Gaulin^{g,b}
PNAS (2012)



site

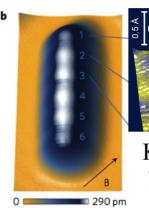
Atomic scale magnetic structures

Ladders and magnetic clusters



S. Loth et al Science (2013)

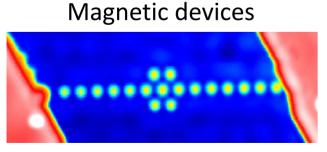
Spin chains



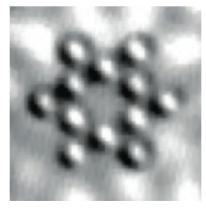
Khajetoorians et al Nat. Phys. (2012)

Spinelli et al Nat. Mat. (2014)

Frustration

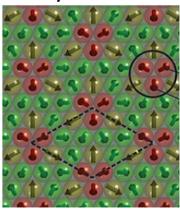


Khajetoorians et al Science (2011)



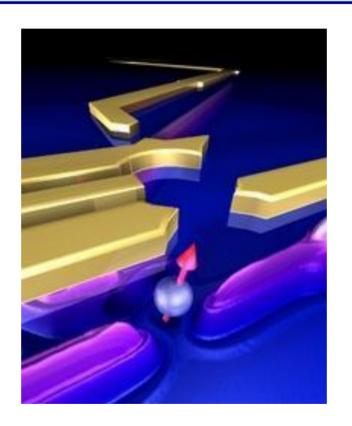
Khajetoorians et al Science (2011)

Skyrmions



Von Bergmann Nano. Lett. (2015)

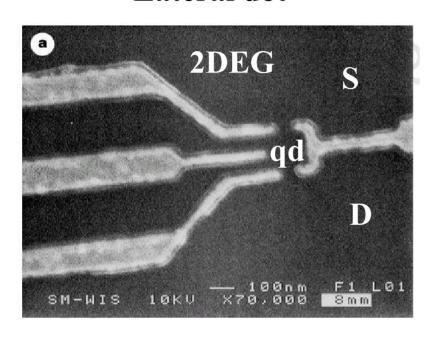
Quantum Transport



References: PRB **78**, 195317 (2008); PRA **78**, 013620 (2008); PRL **100**, 166403 (2008); PRB **73**, 195304 (2006); New. J. of Phys (2010)

Quantum dots

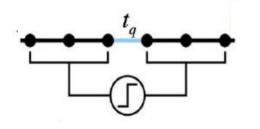
Lateral dot



D. Goldhaber-Gordon, Hadas Shtrikman, D. Mahalu, David Abusch-Magder, U. Meirav and M. A. Kastner, Nature **391** (1998) 156

- Semiconductor heterostructure
- 0-d system
- Droplet of few confined electrons
- Very precise control and tuning
- Single electron transistors
- Size quantization: "Artifical atoms"

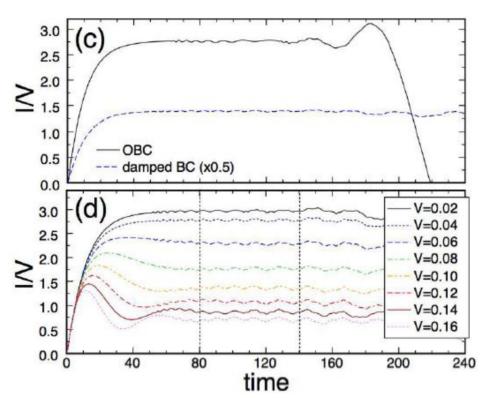
Example: transport in 1d



Interacting spinless fermions

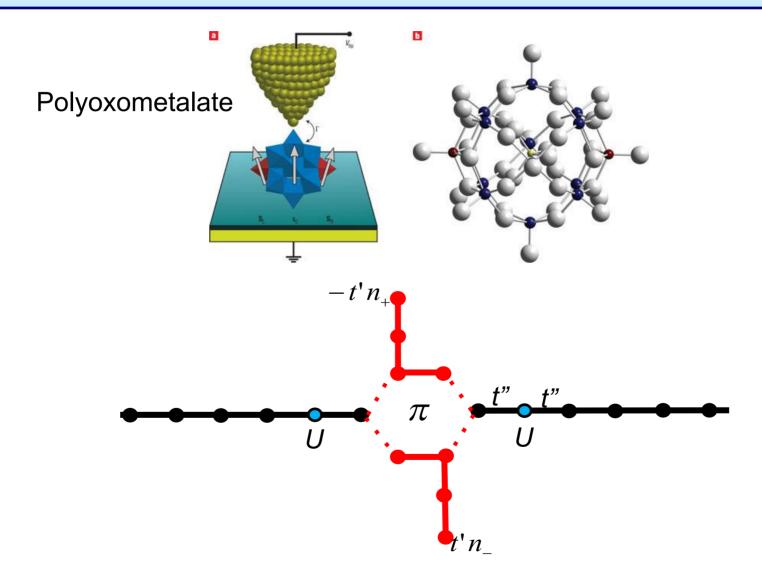
Typical behavior:

- 1) Short time transient
- 2) Plateau (we measure!)
- 3) Reflection at the boundaries. Current changes sign.



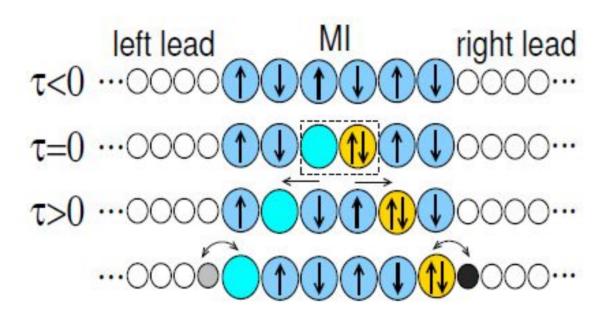
AEF, P. Fendley, MPA Fisher, C. Nayak, PRL0

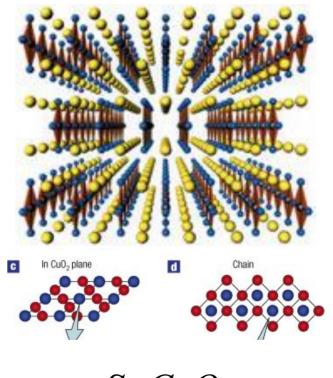
Molecular devices: multi-terminal devices and logical gates



J. Lehmann et al, Nature, Nanotech. 2, 312 (2007), C. A. Busser and AEF (2012)

Exciton creation and recombination: many-body effects in light harvesting materials





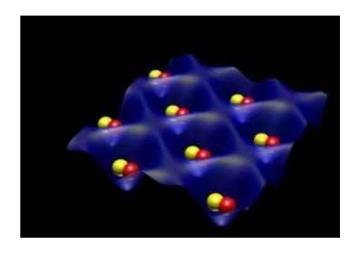
 Sr_2CuO_3

CM phenomena in cold atom systems



Condensed Matter

- •Bulk
- Disordered
- Unknown interactions



Cold Atoms

- Finite, trapped
- Fine tuned interactions
- Quantum coherence
- •Engineered Hamiltonians
- •"Artificial" quantum phases
- •Free of dislocations or defects
- No phonons

Areas of interest

- Exotic phases of matter of quantum origin: quantum magnetism, superconductivity...
- Quantum transport and non-equilibrium physics.
 "Mottronics" and molecular devices. Materials for light harvesting applications.
- Spintronics, using spin for logical gates and computing.
 NV centers in diamond.
- Basic understanding of quantum phenomena using cold atoms.
- New computational algorithms for quantum chemistry and condensed matter.