



Psalter: Particle Spectrum for Any Tensor Lagrangian

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GRTL meeting

Claire Rigouzzo 26/04/2024



Plan:

- **1.** What?
- 2. How?
- 3. Why?
- 4. Where?
- 5. Let's give it a go!

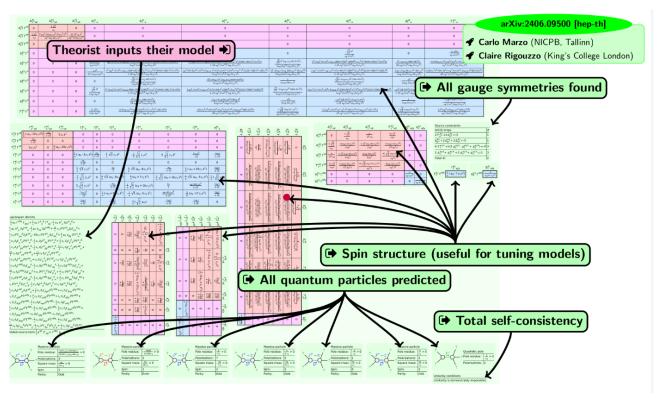


A new Mathematica package that allows you to explore particle content of Lagrangians.

(part of the xAct package)

[1] PSALTer: Particle Spectrum for Any Tensor Lagrangian, W. Barker (Cambridge U), C. Marzo (NICPB, Tallinn), C. Rigouzzo (King's Coll. London), 2406.09500





Slide courtesy of Will Barker

Any Lagrangian?

Yes and No



- ✓ Abstract indices.
- ✓ Any tensor (up to rank 3).
- ✓ Any symmetry of the tensor.
- ✓ Agnostic about formulation of gravity.

Any Lagrangian?

Yes and No



- × No parity breaking terms (as of now)
- Maximum quadratic in tensors (as of now)
- **×** Expansion around flat background (as of now)

 Based on an interesting piece of physics: Spin Projection Operators (SPOs).

Remember group theory:

The Lorentz algebra is the direct sum of independent SU(2) algebras

$$\mathfrak{so}(1,3)\simeq\mathfrak{su}(2)_-\oplus\mathfrak{su}(2)_+$$

(j,j_+)	dim	Type	Example
(0,0)	1	Scalar	$\pi^0, \pi^{\pm}, \text{Higgs}$
(1/2,0)	2	Left-handed spinor	Neutrinos
(0, 1/2)	2	Right-handed spinor	Anti-neutrinos
$(1/2,0) \oplus (0,1/2)$	4	Dirac spinor	e^{\pm}, p, n
(1/2, 1/2)	4	Vector	γ, W^{\pm}, Z^0, g
(1,1)	9	Traceless metric tensor	"Gravity"

Clebsch-Gordan decomposition of a vector:

$$\left(rac{1}{2},rac{1}{2}
ight)
ightarrow \left(1^-
ight) \oplus \left(0^+
ight)$$

Clebsch-Gordan decomposition of a rank two tensor:

$$\left(rac{1}{2},rac{1}{2}
ight)\otimes\left(rac{1}{2},rac{1}{2}
ight)
ightarrow\left(2^{+}
ight)\oplus\left(1^{+}
ight)\oplus2ig(1^{-}ig)\oplus2ig(0^{+}ig)$$

Clebsch-Gordan decomposition of a rank two tensor :

$$\left(rac{1}{2},rac{1}{2}
ight)\otimes\left(rac{1}{2},rac{1}{2}
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ightarrow\left(2^{+}
ight)\oplus\left(1^{+}
ight)\oplus2ig(1^{-}ig)\oplus2ig(0^{+}ig)$$

Conclusion: possible to have all sorts of propagating spinorial sectors, and we want to keep track of them!

Clebsch-Gordan decomposition of a rank two tensor :

$$\left(rac{1}{2},rac{1}{2}
ight)\otimes\left(rac{1}{2},rac{1}{2}
ight)
ightarrow\left(2^{+}
ight)\oplus\left(1^{+}
ight)\oplus2\left(1^{-}
ight)\oplus2\left(0^{+}
ight)$$

Idea 1: First decompose into each spin subsector

→ Use Spin Projection Operators (SPOs).

- Idea 2: massage the equation to obtain propagator
- → All the information (mass? Tachyon? Ghost?) is encoded in it.

Particles	Propagator
Spin 0 (scalar fields (Higgs, pions ,)	$\frac{i}{q^2 - m^2}$
Spin 1/2	$\frac{i}{p^2 - m^2} = i \frac{p^2 + m}{p^2 + m^2}$
Spin 1 massive (W,Z weak boson)	$\frac{-i\left(g_{\mu\nu} - \frac{q_{\mu}q_{\nu}}{m^2}\right)}{q^2 - m^2}$
Spin 1 massless(photon)	$\frac{-ig_{\nu\mu}}{q^2}$

Example: massive spin 0 field

$$Propagator \propto \frac{A}{k^2 - m^2}$$

- No ghosts: *A* < 0</p>
- No tachyons: $m^2 > 0$

(break causality/unitarity)

(subluminal speed)

III. Why?

 Was first motivated to explore the zoo of modified theory of gravity.

• In Palatini formulation, the metric and affine connection are independent: $(h_{ij}, \Gamma_{l[mn]})$

$$egin{aligned} \left(rac{1}{2},rac{1}{2}
ight)\otimes\left(rac{1}{2},rac{1}{2}
ight)\oplus\left(rac{1}{2},rac{1}{2}
ight)\otimes\left((1,0)\oplus(0,1)
ight)\ &
ightarrow2ig(2^+ig)\oplusig(2^-ig)\oplus3ig(1^+ig)\oplus4ig(1^-ig)\oplus3ig(0^+ig)\oplusig(0^-ig) \end{aligned}$$

III. Why?

Real need for computers to do this analysis...

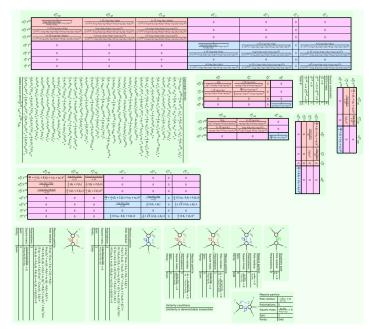


FIG. 21. The particle spectrum of the most general parity-preserving PGT. The fields f_{xy}^{\dagger} and θ_{xy}^{\dagger} contain respectively 16 and 24 d.o.f. The Poincaré symmetry eliminates 2×10 d.o.f, and two are accounted for by the graviton polarisations. The remaining 18 d.o.f are partitioned amongst the six massive species shown above. As is well known, only for special cases of the PGT action in Eq. (8) do the masses and pole residues of these species allow for unitarity: the general case shown here is sick. All quantities are defined in Figs. 18 and 19.

Proof by intimidation

 It is also an useful tool for scalarisation/ vectorisation.

 <u>Definition:</u> Scalarisation is a mechanism that endows self-gravitating bodies, such as BH and neutron stars, with a scalar-field configuration [1]

III. Why?

 At the perturbative level, spontaneous scalarisation is signalled by a tachyon.

→ Can use Psalter to find tachyons.

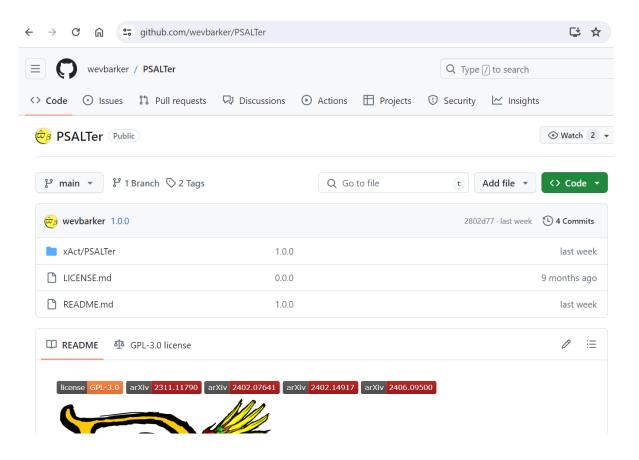
III. Why?

• Can also do R^2 , $R_{\mu\nu}$ $R^{\mu\nu}$... in **any** formulation of gravity



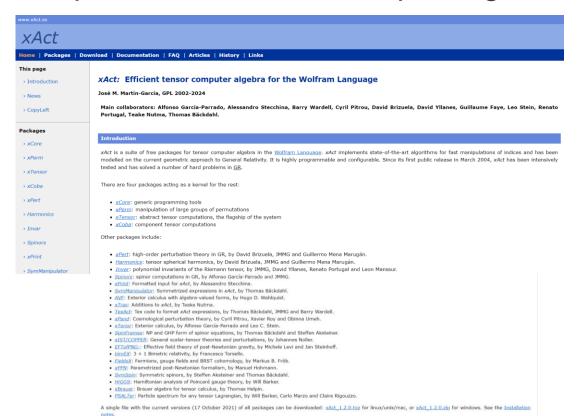








A part of the official xAct packages :



V. Let's give it a go!

- Let's work out a quick example together ©
 - Start with the usual graviton:

$$L \propto \int d^4x \sqrt{-g}R$$

Massive gravity:

$$L \propto \int d^4x \sqrt{-g} \left[\alpha R + \beta \left(h_{ab}h^{ab} - hh\right)\right]$$

Sick massive gravity:

$$L \propto \int d^4x \sqrt{-g} \left[\alpha R + \beta h_{ab} h^{ab} - \gamma hh \right]$$

V. Let's give it a go!

• Just R^2?:

$$L \propto \int d^4x \sqrt{-g} \ \alpha R^2$$

Conclusion and outlook:

Application:

- Deals with alternative formulation of gravity: you give a Lagrangian, you get the particle content (and more).
- Scalarisation

Outlook:

- Extend the code to deal with parity odd Lagrangian (c.f Karananas).
- Extend the code to higher interaction terms/ deal with loop computations.
- Extend the code to deal with non-trivial background (dS, AdS, torsion condensate...)



Questions? Comments?