

REPLY TO REVIEWERS

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We would like to thank the referee for taking the time to carefully read the article. The suggestions and corrections have been helpful and increase the quality of this article. The document has been edited under consideration of the remarks made in the review.

The attached documents contain a version of the paper where all changes are highlighted in blue. In the following, we address each of the remarks individually:

- (1) The paper is globally well written, with all the key arguments and algorithms clearly explained. The reference list is comprehensive.
- (2) Some notations are non fully rigorous. Typical is line 182 page 5 where $\hat{\mathbf{u}}_i$ is the i -th moment (scalar function of \mathbf{x} and t), while $\hat{\mathbf{u}}_j$ in line 225 page 6 is something else (a vector, the index is the cell number). I recommend to check all these kinds of formal inaccuracies.

Response: We added a time index to $\hat{\mathbf{u}}_j$ which was missing in the original version. Furthermore, we added a comment when introducing the moments at the spacial cells in line 189 in order to avoid confusion.

- (3) Is algorithm 1 (page 7) similar to the one by Poette? Can you explain the differences?

Response: Algorithm 1 is very similar to the original IPM algorithm by Poette. The only difference is that we are not using the old moments $\hat{\mathbf{u}}_j^{n+1}$ in the finite volume update (line 11 in algorithm 1), but we are using the term $\langle \mathbf{u}_s((\boldsymbol{\lambda}_j^n)^T \boldsymbol{\varphi}) \boldsymbol{\varphi}^T \rangle_Q^T$. This step ensures that the error from Newton's method does not yield non-physical moments which can lead to a failure of the method, see Section 4.3 in [8].

- (4) The formulation of Th. 1 is incomplete because $\boldsymbol{\lambda}_*$ is not specified. Are conditions on $\boldsymbol{\lambda}_*$? It must be explained.

Response: We added an explanation for $\boldsymbol{\lambda}^*$, see line 312. Furthermore, to underline that $\boldsymbol{\lambda}^*$ collects the dual variables at all spatial cells, we now denote $\boldsymbol{\lambda}^*$ as $\boldsymbol{\lambda}_\Delta^* = (\boldsymbol{\lambda}_1^*, \dots, \boldsymbol{\lambda}_{N_x}^*)^T$.

- (5) The super-attractive property of the Newton-Raphson method is well known, and in my mind, this is exactly what means $\partial_\lambda d = 0$. Can you confirm?

Response: Yes, this is correct. At the fixed point $\partial_\lambda d = 0$ which comes from choosing the Newton-Raphson method to iterate the dual variables.

- (6) Line 357 page 10 is impossible to understand since the right hand side seems independent of j and n . The same for line 364 same page. Actually it is mandatory

that the dependence with respect to j and n are the same throughout the whole proof.

Response: You are correct. We forgot to include indices and the manuscript lacked a clear notation at some points. We now added indices and define our notation more clearly.

- (7) And also my own calculations from (11) yield that $\partial_u d = -H(\lambda)^{-1}\lambda$. Can you check?

Response: We have that

$$\begin{aligned} d(\lambda, \hat{u}) &= \lambda - H^{-1}(\lambda) (\langle u_s(\lambda^T \varphi) \varphi^T \rangle_Q^T - \hat{u}) \\ &= \lambda - H^{-1}(\lambda) \langle u_s(\lambda^T \varphi) \varphi^T \rangle_Q^T + H^{-1}(\lambda) \hat{u} \end{aligned}$$

Since only the last term $H^{-1}(\lambda) \hat{u}$ depends on the moment vector \hat{u} , we have

$$\partial_{\hat{u}} d(\lambda, \hat{u}) = H^{-1}(\lambda).$$

Note that for the One-Shot method, the moments \hat{u} and dual variables λ do not depend on each other.

- (8) page 12 in Algo. 3, it would be better to write "Determine Refinement Level " instead of "DetermineRefinementLevel ". Please check all similar notations.

Response: We have changed this in the manuscript.

- (9) Line 496 page 14, I do not understand the need to specify the temperature T since it does not show up in line 487.

Response: We added a short description on how the defined quantities relate to the conserved variables in line 517.

- (10) I missed the definition of the reference solution u_Δ . It must be explicit (or more visible).

Response: We compute the reference solution with a finely resolved collocation method. The one-dimensional testcase uses 100 Gauss-Legendre quadrature points (as mentioned in line 506 in the original manuscript), the two-dimensional testcase uses 50^2 quadrature points (as mentioned in line 620 in the original manuscript) and the three-dimensional testcase uses 50^3 quadrature points (as mentioned in line 661 in the original manuscript). You are correct that this information gets lost in the text which is why we added a short description whenever we plot reference solutions (which we do in Figures 1, 7 and 8).

- (11) Figure 3: the threshold $\varepsilon = 10^{-6}$ is not visible on the plots, as it should be. Also the intrusive and non-intrusives methods must be more explicit: which ones (number of equations, . . .).

Response: We added the names of intrusive and non-intrusive methods and include references to the corresponding algorithms. The residual cannot be depicted in the plot itself, since the y-axis is the relative error, not the residual. We added references to the relative error as well as to the residual equations to avoid confusion.

- (12) Lines 553/554 page 17 seem weird. What do you call positivity of mass? Is it total mass? If yes, there is not problem with a conservative scheme.

Response: We meant to say density, which we changed in the manuscript. The total mass (here the expected value of the total mass) is conserved by the scheme, however the solution at individual values for ξ might yield negative densities, since we are using polynomials to span the solution in ξ .

- (13) Line 691 in the conclusion (Additionally, the ability to adaptively change 692 the truncation order helps intrusive methods to compete with SC in terms of computational runtime) are an interesting asset of the methods presented in this work.