Stability and detection performance of a GEM-based Optical Readout TPC with He/CF₄ gas mixtures

E. Baracchini, a,b L. Benussi, c S. Bianco, c C. Capoccia, c M. Caponero, c,d G. Cavoto, e,f A. Cortez, a,b I. A. Costa, e E. Di Marco, e G. D'Imperio, e G. Dho, a,b F. Iacoangeli, e G. Maccarrone, e M. Marafini, e,h G. Mazzitelli, e A. Messina, e,f R. A. Nobrega, e A. Orlandi, e E. Paoletti, e L. Passamonti, e F. Petrucci, e D. Piccolo, e D. Pierluigi, e D. Pinci, e F. Renga, e F. Rosatelli, e A. Russo, e G. Saviano, e and S. Tomassini e

Laboratori Nazionali del Gran Sasso, Assergi, Italy

Laboratori Nazionali di Frascati, Italy

E-mail: davide.pinci@roma1.infn.it

ABSTRACT: The performance and long term stability of an optically readout Time Projection Chamber with an electron amplification structure based on three Gas Electron Multipliers was studied. He/CF₄ based gas mixtures were used in two different proportions (60/40 and 70/30) in a CYGNO prototype with 7 litres sensitive volume. With electrical configurations providing very similar electron gains, an almost full detection efficiency in the whole detector volume was found with both mixtures, while a light yield about 20% larger for the 60/40 was found. The electrostatic stability was tested by monitoring voltages and currents during 25 days. The detector worked in very stable and safe condition for the whole period. In the presence of less CF₄, a larger probability of unstable events was clearly detected.

KEYWORDS: Dark Matter detectors; Gaseous imaging and tracking detectors; Micropattern gaseous detectors; Time projection Chambers.

^aGran Sasso Science Institute,

L'Aquila, Italy

^bIstituto Nazionale di Fisica Nucleare,

^cIstituto Nazionale di Fisica Nucleare ,

^dENEA Centro Ricerche Frascati, Frascati, Italy

^eIstituto Nazionale di Fisica Nucleare,

Sezione di Roma, Italy

^f Dipartimento di Fisica Sapienza Università di Roma, Italy

^g Universidade Federal de Juiz de Fora, Juiz de Fora, Brazil

^hMuseo Storico della Fisica e Centro Studi e Ricerche "Enrico Fermi",

Piazza del Viminale 1, Roma, Italy

ⁱDipartimento di Matematica e Fisica, Università Roma TRE, Roma, Italy

^j Istituto Nazionale di Fisica Nucleare, Sezione di Roma TRE, Roma, Italy

 $[^]k$ Dipartimento di Ingegneria Chimica, Materiali e Ambiente, Sapienza Università di Roma, Roma, Italy

¹Corresponding author.

Contents

1	Exp	erimen	ntal setup	2
	1.1	LEMO	ON detector	2
	1.2	Opera	ating conditions	2
	1.3	Data A	Acquisition	3
2	Ope	eration	stability	4
3	Sim	ulation	of gas mixtures	7
4	Dete	ector pe	erformance	8
	4.1	Data a	analysis	8
	4.2	Light	yield and energy resolution	10
		4.2.1	Measurements with sCMOS	10
		4.2.2	Measurements with PMT	10
		4.2.3	Energy resolution	11
		4.2.4	Measurement of diffusion effect	11
	4.3	Detec	ction efficiency	12
5	Con	clusion	n	12

Introduction

Liquid and gaseous Time Projection Chambers have been successfully proposed and exploited in last decades for very different applications, from High Energy Physics experiments on colliders to the searches of Dark Matter (DM) massive particles. In this latter field, a very promising technique involves the optical reading of gas electro-luminescence produced during the processes of electron multiplication. [1–6].

Thanks to the great progresses achieved in recent years in both the performance of Micro Pattern Gas Detectors and CMOS-based light sensors, optical readout provides several crucial advantages:

- the very good performance of optical sensors allows detection and reconstruction of very low energy releases;
- sensors can be installed outside of the sensitive volume reducing the interference with the detector operation and possible sources of gas contamination;
- the use of suitable lenses allows to acquire large surfaces with small sensors.

The CYGNO collaboration is developing the optical technique on Gas Electron Multipliers (GEM) [7] working with He/CF₄ based gas mixtures, with the aim of realising a cubic meter demonstrator to study its performance for low mass Dark Matter directional search. For such application, high detection efficiency and energy resolution are needed together with a very good detector stability ensuring safe operation for long data takings.

In this paper, the performance of a CYGNO prototype is presented and discussed for the first time together with a detailed study of long term reliability in operating conditions with two different He/CF₄ mixtures (60/40 and 70/30).

1 Experimental setup

1.1 LEMON detector

All studies presented in this paper have been carried out with the *LEMON* prototype [8–10]. A sketch of this detector (described in more details in [11]) is shown in Fig 1. It has a 7 litres sensitive volume (A) with a 20 cm long drift gap surrounded by an elliptical field cage closed on one side by a semitransparent cathode and on the other side by a 20 \times 24 cm² triple GEM structure. Light produced in multiplication channels, is acquired by:

- a photo-multiplier¹ with 5 ns rise-time, a maximum QE of 12% for 420 nm and a 76 × 76 mm² square-window (trough the cathode);
- a scientific-CMOS based camera²) with 2048×2048 pixels with an active area of $6.5 \times 6.5 \,\mu\text{m}^2$ each, equipped with a Schneider lens with 25 mm focal length and 0.95 aperture at a distance of 52.5 cm. The sCMOS sensor provides a quantum efficiency of about 70% at 600 nm.

The drift volume was filled with He/CF₄ based gas mixtures. The operating performance with the two different gas mixtures (60/40 and 70/30) was studied. According to previous studies, electro-luminesce spectra of He/CF₄ based mixtures show two main peaks: one around 300 nm and one around 620 nm. The relative light production depends on the relative amount of the two components in the gas mixture: the height of second peak increases with respect to the first one for a larger CF₄ percentage in the mixture [12].

1.2 Operating conditions

The typical working configuration of the detector is:

- a gas flux of 200 cc/min;
- an electric field within the sensitive volume $E_D = 0.5 \text{ kV/cm}$;
- an electric field in the 2 mm wide gaps between the GEMs $E_{Transf} = 2.5 \text{ kV/cm}$;
- a voltage difference across each GEM V_{GEM} = 460 V while operating with 60/40 mixture and V_{GEM} = 425 V while operating with 70/30 mixture.

¹Photonics XP3392

²Hamammatsu ORCA Flash 4

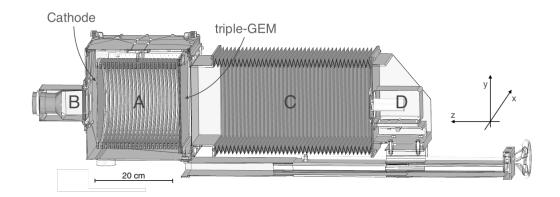


Figure 1. Drawing of the experimental setup. In particular, the elliptical field cage closed on one side by the triple-GEM structure and on the other side by the semitransparent cathode (A), the PMT (B), the adaptable bellow (C) and the CMOS camera with its lens (D).

According to results presented in [13], the behavior of electron gain with the two gas mixtures as a function of V_{GEM} can be described as:

$$G_{60/40}(V_{\text{Gem}}) = 28.36 \times e^{0.0235 \cdot V_{\text{Gem}}(V)}$$
 (1.1)

$$G_{70/30}(V_{\text{Gem}}) = 22.82 \times e^{0.0258 \cdot V_{\text{Gem}}(V)}$$
 (1.2)

Therefore, since:

$$G_{60/40}(460 \text{ V})/G_{70/30}(425 \text{ V}) = 1.065$$
 (1.3)

the two values chosen for V_{GEM} should allow to operate with very similar charge gains.

After two days of conditioning, the detector was closed in a 5 cm thick lead box, in order to partially shield it from soft cosmic rays and external natural radioactivity background. A slit on the box side, allowed to irradiate the sensitive volume with 5.9 keV photons produced by a 740 MBq 55 Fe source at a y position corresponding to half height of the field cage and free to move along z axis.

1.3 Data Acquisition

In the events presented in this paper the sCMOS sensor was operated in auto-trigger mode with an exposure time of 40 ms. The PMT waveform was sent into a 4 GS/s digitizer. Only in the case that during the CMOS exposure window PMT waveform exhibited an amplitude larger than 80 mV it was acquired for a time interval of 25 μ s and the image was stored.

In all tested configurations, runs with 1000 events were acquired.

In order to monitor and study the working conditions of the detector, voltage outputs and currents drawn by all high voltage channels supplying the different electrodes of the detectors, were logged with a sampling rate of 1 Hz. The used high voltage supply module is able to measure only negative currents. For this reason, it was possible to acquire only the currents provided to GEM and field cage cathodes.

2 Operation stability

Detector operational stability was evaluated during a 25 days long test. Figure 2 shows the behavior of current drawn by the high voltage channel supplying the upper electrode (the one which faces the second GEM) of the third GEM in the stack (I_{G3U}) averaged day by day. On this electrode, a fraction of ions produced during the avalanche is collected and thus I_{G3U} is proportional to the charge produced by ionization in the sensitive volume amplified by the triple-GEM gain. For $E_D \cong 0$, so that with no ionization charge collected on the GEM, it was found to be compatible with the sensitivity of current measurement (20 nA).

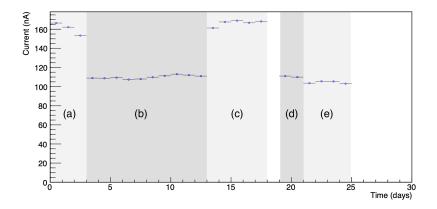


Figure 2. Average current drawn by the upper electrode of the third GEM I_{G3U} as a function of time. Details about the experimental set-up in different periods are summarised in Table 1.

The data were collected for five consecutive periods, whose main characteristics are summarised in Tab 1

Period	Gas Proportion	Pb	⁵⁵ Fe	Collimator	Avg. Current
	(He/CF ₄)	Shielding	Source		(nA)
(a)	60/40	No	No	No	164 ± 2
(b)	60/40	Yes	No	No	110 ± 1
(c)	60/40	Yes	Yes	No	168 ± 2
(d)	60/40	Yes	Yes	Yes	110 ± 1
(e)	70/30	Yes	Yes	Yes	104 ± 2

Table 1. Description of main settings in different periods of test

The 5 cm thick lead shield reduced the level of ionization due to natural background by about 30% ((a) and (b) in Fig. 2). The ratio between the current drawn in periods (d) and (e) results to be 1.06 ± 0.02 in good agreement with the value 1.065 expected for these operating conditions (see Sect. 1.1).

During normal operation, even if for most of the time the currents drawn by the detector electrodes were very stable, from time to time spikes occurred indicating detector instabilities.

Two different kinds of electrostatic instabilities have been observed:

• **Hot spots.** Appearance of small luminous spots on the GEM surface as the one shown in Fig.3, usually accompanied by negligible current increase.

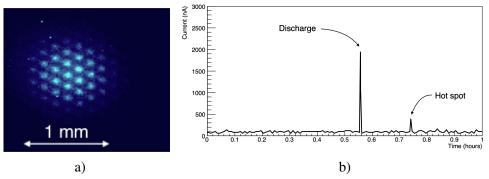


Figure 3. Zoomed image of a luminous spot on the GEM surface (*hot spot*) very likely due to self-sustaining micro-discharge occurring in one or more GEM channels (a). The structure of GEM channels is clearly visible. Example of the behavior of current drawn by the upper electrode of the third GEM I_{G3U} as a function of time (b). The occurrence of a *discharge* and a *hot spot* are indicated (see text for details).

Some of these spots fades out with the time but in some case they start to slowly grow up (on a time scale of minutes). At some point they could even involve a current drawn by the GEM as large as to be measurable by the power supply (tens of nA, as shown on the right of Fig. 3). These are probably due to self-sustaining micro-discharges happening in one or few GEM channels. It was found that a decrease of the voltage across all GEMs of 100 V is enough to dump this self-sustaining process and that, when V_{GEM} is restored the hot-spot do not re-appears. An automatic hot-spot dump procedure was implemented to decrease all V_{GEM} if the drawn current increases above some threshold. The V_{GEM} is thus restored in 5 steps of 20 V with a 30 sec pause between each step. This procedure lasts about 3 minutes and introduce a *dead time* in the detector operation.

• **Discharges** High charge density due to very high ionizing particles or charge accumulation on electrode imperfections can suddenly discharge across GEM channels. In these events, a sudden increase in the drawn current is recorded with a voltage restoring on the electrodes through protection resistors on a few seconds time basis (an example is shown on the left of Fig. 3). Also these events trigger the recovery procedure. Even if these events are less frequent than hot spots they can be dangerous for the GEM structure and the energy released in the discharge can, in principle, damage it.

To study the occurrence of above events for the two gas mixtures, the detector currents were acquired while all other setting parameters were kept in stable and operational mode. For the 40/60 mixture the LEMON behavior has been acquired for 130 hours (5.4 days). In the whole period a total of 66 hot spots and 31 discharges were recorded giving an occurrence probability of 12.2 hotspots/day and 5.7 discharges/day. While operated with the 70/30 gas mixture, the detector showed a less stable conditions and therefore a shorted test time was needed to reach statistically significant

results. In 60 hours of test (2.5 days), 121 hot-spots and 39 discharges occurred, corresponding to a probability of 48.4 hot-spots/day and 15.6 discharges/day.

Figure 4 shows the distributions of the time interval between the appearance of two subsequent hot-spots for the two gas mixtures.

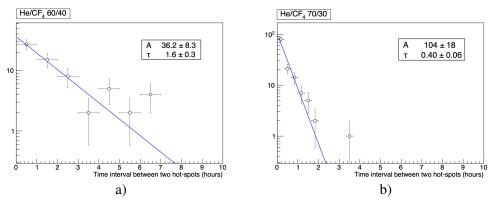


Figure 4. Distributions of the interval between the appearance of two subsequent hot-spots (a) for He/CF₄ (60/40) and (b) for He/CF₄ (70/30)) with superimposed fits to exponential decrease.

The superimposed exponential fit to a function:

$$y = A e^{-\frac{\Delta t}{\tau}} \tag{2.1}$$

demonstrates that this distribution is the one expected for events occurring at completely random moments without any evident correlation [14]. In particular, the results of the fit confirm that the rate of hot-spots in He/CF_4 (70/30) mixture is 4 times larger than He/CF_4 (60/40).

As it was done for the hot-spots, the distributions of the time intervals between the occurrence of two subsequent discharges was studied for both the mixtures (Fig. 5).

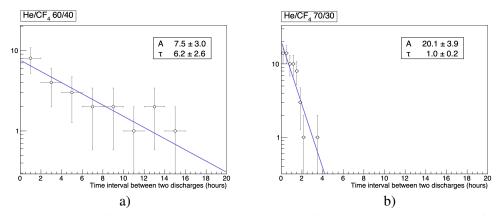


Figure 5. Distributions of the interval between the occurrence of two subsequent discharges (a) for He/CF₄ (60/40) and (b) for He/CF₄ (70/30)) with superimposed fits to exponential decrease.

The behavior is well described by the function in 2.1 showing that also the occurrence of these events is random in time without any evident correlation. The results of the fit confirm that the probability of having a discharge in He/CF₄ (70/30) mixture is 6 times larger than He/CF₄ (60/40).

It results evident that a lower amount of CF_4 , able to quench and keep under control possible production of large amount of charge during the multiplication processes, resulted in a less stable electrostatic configuration. The instability events gave rise to a detection inefficiency due to dead time introduced by recovering procedures of 3.8% (60/40) and 13.3% (70/30).

Nevertheless, it is important to outline that, in both cases, the detector behavior seemed not to be critical and the provided performance was satisfactory.

3 Simulation of gas mixtures

The parameters of the two gas mixtures relevant to study the electron transport in the field cage were calculated by means of Garfield [15, 16].

Because electrons diffuses along their drift in the field cage, when they arrive on the GEM, they trigger avalanche multiplications of an area larger than the ionization region. The photons emitted by gas electro-luminescence in the avalanches will thus create the light-spots on the sensor (see for example on the right of Fig. 8).

After a drift over a distance z, transverse and longitudinal profiles of electron clouds produced in these spots can be described by Gaussian curves with standard deviations that can be calculated as:

$$\sigma_{\rm T} = \sqrt{\sigma_{\rm T0}^2 \oplus D_{\rm T}^2 \cdot z} \tag{3.1}$$

$$\sigma_{\rm L} = \sqrt{\sigma_{\rm L0}^2 \oplus D_{\rm L}^2 \cdot z} \tag{3.2}$$

where σ_{T0} and σ_{L0} are constant contributions due diffusion in the GEM structure and to channel pitch and $D_{\rm L}$ and $D_{\rm L}$ are transverse and longitudinal diffusion coefficients that depend on the gas mixture and the electric field.

The behavior of the diffusion coefficients for different electric fields is reported in Fig. 6 In

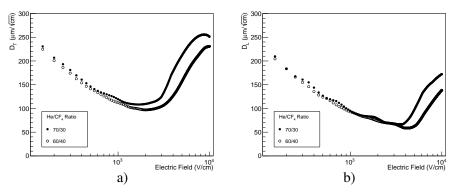


Figure 6. Transverse (a) and longitudinal (b) diffusion coefficients for the two mixtures as a function of the electric field.

particular, for an electric field of 0.5 kV/cm (i.e. the value set in the sensitive volume for the measurements presented in this paper) they were evaluated to be:

$$D_{\rm T}^{60/40} = 140 \frac{\mu \rm m}{\sqrt{\rm cm}} \text{ and } D_{\rm T}^{70/30} = 145 \frac{\mu \rm m}{\sqrt{\rm cm}}$$

$$D_{\rm L}^{60/40} = 120 \frac{\mu \rm m}{\sqrt{\rm cm}}$$
 and $D_{\rm L}^{70/30} = 125 \frac{\mu \rm m}{\sqrt{\rm cm}}$

The electron drift velocity as a function of the electric field is shown on Fig. 7 (a).

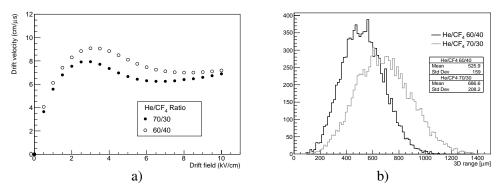


Figure 7. Electron drift velocities as evaluated with Garfield for the two gas mixtures (a) and distribution of ranges of 5.9 keV electrons evaluated with GEANT4 (b) for the two gas mixtures

It is clearly visible that a larger fraction of CF₄ allows to have slightly larger drift velocities. In particular, for an electric field of 500 V/cm:

$$v_{\rm drift}^{60/40} = 40.6 \ \mu {\rm m/ns} \ {\rm and} \ v_{\rm drift}^{70/30} = 36.5 \ \mu {\rm m/ns}$$

Mean free path of the 5.9 keV electrons produced by photo-electric effect in gas by the ⁵⁵Fe photons was evaluated with GEANT4 [17] for the two gas mixtures. The obtained distributions are shown in Fig. 7 (b).

Even if a larger fraction of Helium increases the average path, in both cases, average ranges of several hundreds of micrometers were found. Therefore, interaction of ⁵⁵Fe photons are expected to produce light spots with dimensions mainly due to the diffusion of electrons in the gas.

4 Detector performance

The detection performance provided by the two gas mixtures were evaluated by studying signals produced in the sensitive volume by 5.9 keV photons. As described in [11], the interactions of low energy photons with atoms in the gas mixtures create photo-electrons that release their whole energy in few hundreds of micrometers. The subsequent diffusion of ionized electrons, makes these events produce round-like spots on the CMOS sensor with diameters of 2-3 mm as the ones shown in Fig. 8 due to two photons independently interacting in the gas.

4.1 Data analysis

Data acquired with CMOS camera were analysed by using the DBSCAN-based algorithm described in details in [18]. In each image, clusters of illuminated pixels are individuated and used to reconstruct the signal spots. For subsequent analysis, the position, the size and the total light of each spot are recorded.

For all events, a very simple analysis was performed on waveform provided by the PMT:

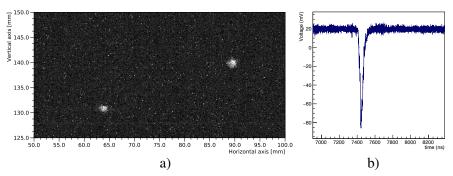


Figure 8. Detection of interactions of ⁵⁵Fe photons in the sensitive volume. (a) sCMOS image with two light spots due to two photons interacting in the gas. (b) an example of PMT signal from a photon interaction.

- signal is integrated to evaluate the charge provided;
- the FWHM of peak shape is measured to evaluate the signal duration.

The source was placed at a distance of 14 cm from border of the sensitive volume in a metal collimator that made it possible to limit the "illuminated" region to a cone. For most of following measurements the source was kept at a distance of 10.5 cm from the GEM stack. The maps of the positions of all reconstructed clusters in a typical run are shown in Fig. 9.

The shape of the collimator hole was changed during the data taking. As a consequence, as it is well visible from the maps, the number of events due to interactions of 5.9 keV photons in the runs with the 70/30 (right panel in Fig. 9) is found to be almost five times larger with respect to the ones 60/40 (left panel in Fig. 9).

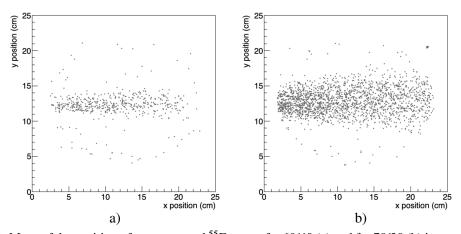


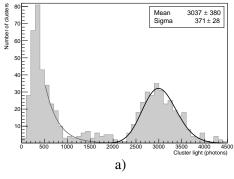
Figure 9. Maps of the position of reconstructed ⁵⁵Fe spots for 60/40 (a) and for 70/30 (b) in two typical runs.

Due to the source position, a large number of clusters due to ⁵⁵Fe photon interactions is reconstructed in the central region of the sensitive volume (at an height between 10 cm and 15 cm).

4.2 Light yield and energy resolution

4.2.1 Measurements with sCMOS

Figure 10 shows the spectra of the amount of light detected in spots reconstructed on sCMOS sensor. The distributions are fitted to a function obtained by the sum of an exponential decay (to describe the background behavior) and a Polya [19] for the 5.9 keV peak.



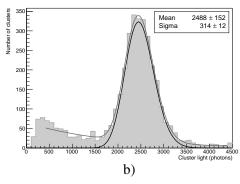


Figure 10. Distribution of the light content in spots reconstructed for the 60/40 (a) and the 70/30 (b) gas mixtures with superimposed fit to an exponential plus a Polya function. The Polya contribution is also indicated.

As described in Sect 1.2, in the chosen working conditions slightly different electron gains are expected, with a value 6.5% larger for the 60/40. Average light yields for the two mixtures were evaluated by dividing the average value of the Polya fit to the two distributions by 5.9 keV:

- 60/40 provides an average value of 514 ± 63 detected photons per keV released in the gas (in agreement with results obtained with lower V_{GEM} and E_{Transf} [11]) with a relative fluctuation of 12.2%;
- 70/30 provides an average value of 420 ± 53 detected photons per keV released in the gas with a relative fluctuation of 12.6%;

A light production 18% larger for 60/40 was measured with respect to 70/30. The slightly lower light detected by the sCMOS with less CF_4 is in good agreement with expectations (see Sect. 1.3) and confirms the component around 600 nm to be due to CF_3^* dis-excitation.

4.2.2 Measurements with PMT

The spectra of the charge integrated in Photo-multiplier waveforms are shown in Figure 11 with a superimposed Polya fit.

The two distributions were fitted to a Polya function to evaluate the average light detected by the PMT with the two mixtures and light yield are obtained by dividing for 5.9 keV:

- 60/40 provides an average value of (12.0 ± 0.2) pC per keV released in the gas with a relative fluctuation of 15.5%;
- 70/30 provides an average value of (13.4 ± 0.1) pC per keV released in the gas with a relative fluctuation of 16.3%;

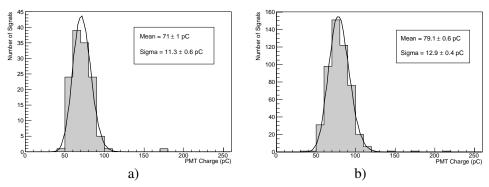


Figure 11. Distribution of the charge provided by the PMT for the 60/40 (a) and the 70/30 (b) gas mixtures.

The PMT collects 12% more light for 70/30 than 60/40. This can be explained by the larger amount of UV light produced by a more He rich mixture to whom PMT photo-cathode is sensitive (Sect. 1.3).

4.2.3 Energy resolution

Since the energy release in gas is practically constant, the fluctuations of the response can be used to evaluate the resolution on the energy measurements at 5.9 keV. The results above, obtained with sCMOS and PMT indicate that with the sCMOS sensor a resolution of about 12% is achieved with both mixtures, slightly better than what obtained with the PMT (between 15% and 16%) probably because of the very low noise level of the sCMOS sensor. With both methods, a similar energy resolution was measured for the two gas mixtures confirming that the main contribution to this parameter is due to the statistical fluctuations of gas ionization and electron multiplication processes.

4.2.4 Measurement of diffusion effect

The distributions of σ_T for spots reconstructed with the source at a distance z of 10.5 cm from the GEM were studied to evaluate the effect of the diffusion and are reported in Fig. 12 with a superimposed gaussian fit.

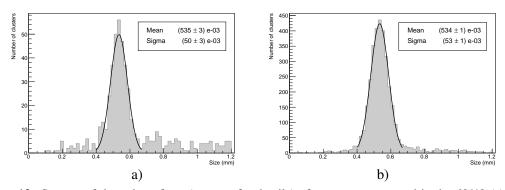


Figure 12. Spectra of the value of σ_T (see text for details) of spots reconstructed in the 60/40 (a) and the 70/30 (b) gas mixture.

The values of σ_{T0} were evaluated from Eq. 3.1 by using the diffusion coefficient given from the simulation (see Sect. 3). For the two mixtures they were found to be:

$$\sigma_{\rm T0}^{60/40}$$
 = (280 ± 60) $\mu{\rm m}$ and $\sigma_{\rm T0}^{70/30}$ = (260 ± 60) $\mu{\rm m}$

They are comparable within the measurement uncertainties confirming that they are mainly due to diffusion in the GEM stack.

4.3 Detection efficiency

The number of spots reconstructed as a function of the source position along z axis was studied to evaluate the behavior of detection efficiency for interactions at different distances from the GEM.

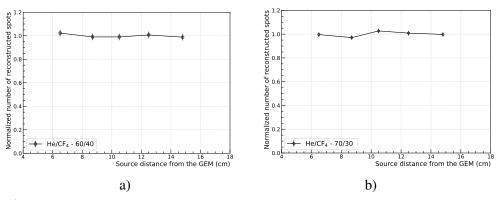


Figure 13. Number of spots reconstructed (normalized to their average value) as a function of the position of the 55 Fe source along the z axis (a) for the 60/40 and (b) and 70/30).

Figure 13 shows the behavior of the ratio between the number of reconstructed spots in the runs at different z and and the average value in the whole scan. No evidence of a dependence of the detection efficiency on z was found, allowing to conclude that, a constant detection efficiency is provided at all studied depths. This result is good in agreement with the very low electron absorption probability due to electron attachment estimated for He/CF₄ mixtures with Garfield for electric field values of E_D used (see Sect. 1.2).

5 Conclusion

The performance of an optical readout TPC with a sensitive volume of 7 litres was studied with two He/CF₄ based mixtures in different proportions (60/40 and 70/30). The chosen detector electrical configurations, allowed to operate in very similar electron gas gain conditions. Performed studies indicate that a constant detection efficiency was found in the sensitive volume together with a very good energy resolution around 13% in both cases, (even if the light yield for the 70/30 mixture resulted to be 18% lesser) indicating that this is mainly due to ionization and multiplication statistics. A sub-keV resolution is also very promising for application in Dark Matter search being the maximum energy released to an He nucleus by 1 GeV mass particle about 1 keV. Detector operation was monitored for a 25 days period. A detailed study of the behavior of currents and voltages provided by the supply system has shown the presence of two different kind instability

events in GEM channels: discharges (with a sudden and fast current increase) and hot-spots (self sustaining events involving less current and creating small light spots on GEM surface). Thanks to its quenching properties a larger presence of CF₄, showed to ensure a quite better stability with lower rate for both kinds of events. Their occurrence was anyway found to be random in time with frequencies always lesser than few per hour and no evidence were found of correlation between two subsequent events.

Results presented in this paper demonstrated the possibility of operating the CYGNO prototype safe and stable conditions while providing promising performance (light yield, energy resolution and detection efficiency) in view of a larger TPC for dark matter search.

Acknowledgments

This work was supported by the European Research Council (ERC) under the European Union's Horizon 2020 research and innovation program (grant agreement No 818744)".

References

- [1] L. M. S. Margato, F. A. F. Fraga, S. T. G. Fetal, M. M. F. R. Fraga, E. F. S. Balau, A. Blanco et al., Performance of an optical readout GEM-based TPC, Nucl. Instrum. Meth. A535 (2004) 231.
- [2] C. M. B. Monteiro, A. S. Conceicao, F. D. Amaro, J. M. Maia, A. C. S. S. M. Bento, L. F. R. Ferreira et al., Secondary scintillation yield from gaseous micropattern electron multipliers in direct dark matter detection, Phys. Lett. **B677** (2009) 133.
- [3] C. M. B. Monteiro, L. M. P. Fernandes, J. F. C. A. Veloso, C. A. B. Oliveira and J. M. F. dos Santos, Secondary scintillation yield from GEM and THGEM gaseous electron multipliers for direct dark matter search, Phys. Lett. B714 (2012) 18.
- [4] A. Bondar, A. Buzulutskov, A. Grebenuk, A. Sokolov, D. Akimov, I. Alexandrov et al., *Direct observation of avalanche scintillations in a THGEM-based two-phase Ar avalanche detector using Geiger-mode APD*, *JINST* 5 (2010) P08002 [1005.5216].
- [5] M. Marafini, V. Patera, D. Pinci, A. Sarti, A. Sciubba and E. Spiriti, *ORANGE: A high sensitivity particle tracker based on optically read out GEM, Nucl. Instrum. Meth.* **A845** (2017) 285.
- [6] V. C. Antochi, E. Baracchini, G. Cavoto, E. D. Marco, M. Marafini, G. Mazzitelli et al., *Combined readout of a triple-GEM detector*, *JINST* 13 (2018) P05001 [1803.06860].
- [7] F. Sauli, GEM: A new concept for electron amplification in gas detectors, Nucl. Instrum. Meth. A 386 (1997) 531.
- [8] D. Pinci, E. Di Marco, F. Renga, C. Voena, E. Baracchini, G. Mazzitelli et al., *Cygnus: development of a high resolution TPC for rare events*, *PoS* **EPS-HEP2017** (2017) 077.
- [9] G. Mazzitelli et al., A high resolution TPC based on GEM optical readout, in 2017 IEEE Nuclear Science Symposium and Medical Imaging Conference, 2017, DOI.
- [10] G. Mazzitelli, V. C. Antochi, E. Baracchini, G. Cavoto, A. De Stena, E. Di Marco et al., A high resolution tpc based on gem optical readout, in 2018 IEEE Nuclear Science Symposium and Medical Imaging Conference (NSS/MIC), vol. Under publication in IEEE Nuclear Science Symposium Medical Imaging Conference, 2018, 2018.

- [11] I. A. Costa, E. Baracchini, F. Bellini, L. Benussi, S. Bianco, M. Caponero et al., *Performance of optically readout GEM-based TPC with a 55fe source, Journal of Instrumentation* **14** (2019) P07011.
- [12] M. M. F. R. Fraga, F. A. F. Fraga, S. T. G. Fetal, L. M. S. Margato, R. Ferreira-Marques and A. J. P. L. Policarpo, *The GEM scintillation in He CF*₄, *Ar CF*₄, *Ar TEA and Xe TEA mixtures*, *Nucl. Instrum. Meth.* **A504** (2003) 88.
- [13] R. Campagnola, Study and optimization of the light-yield of a triple-GEM detector, Feb, 2018.
- [14] G. Knoll, Radiation Detection and Measurement (4th ed.). John Wiley, Hoboken, NJ, 2010.
- [15] R. Veenhof, Garfield, a drift chamber simulation program, Conf. Proc. C 9306149 (1993) 66.
- [16] R. Veenhof, GARFIELD, recent developments, Nucl. Instrum. Meth. A 419 (1998) 726.
- [17] GEANT4 collaboration, GEANT4: A Simulation toolkit, Nucl. Instrum. Meth. A 506 (2003) 250.
- [18] E. Baracchini et al., A density-based clustering algorithm for the CYGNO data analysis, 2007.01763.
- [19] W. Blum, L. Rolandi and W. Riegler, *Particle detection with drift chambers*, Particle Acceleration and Detection, ISBN = 9783540766834. 2008, 10.1007/978-3-540-76684-1.