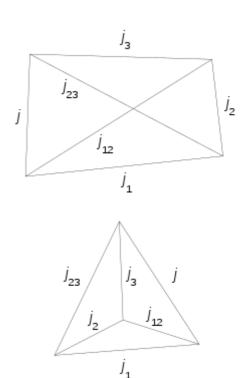
Racah W-coefficient

Racah's W-coefficients were introduced by <u>Giulio Racah</u> in 1942. These coefficients have a purely mathematical definition. In physics they are used in calculations involving the <u>quantum mechanical</u> description of angular momentum, for example in atomic theory.

The coefficients appear when there are three sources of angular momentum in the problem. For example, consider an atom with one electron in an <u>s orbital</u> and one electron in a <u>p orbital</u>. Each electron has <u>electron spin</u> angular momentum and in addition the p orbital has orbital angular momentum (an s orbital has zero orbital angular momentum). The atom may be described by *LS* coupling or by *jj* coupling as explained in the article on <u>angular momentum coupling</u>. The transformation between the wave functions that correspond to these two couplings involves a Racah W-coefficient.

Apart from a phase factor, Racah's W-coefficients are equal to Wigner's <u>6-j symbols</u>, so any equation involving Racah's W-coefficients may be rewritten using 6-*j* symbols. This is often advantageous because the symmetry properties of 6-*j* symbols are easier to remember.

Racah coefficients are related to recoupling coefficients by



Angular momenta in the Racah W coefficients. The top is a 2d plane projection as a quadrilateral, the bottom is a 3d tetrahedral arrangement.

$$W(j_1j_2Jj_3;J_{12}J_{23}) \equiv rac{\langle (j_1,(j_2j_3)J_{23})J|((j_1j_2)J_{12},j_3)J
angle}{\sqrt{(2J_{12}+1)(2J_{23}+1)}}.$$

Recoupling coefficients are elements of a <u>unitary transformation</u> and their definition is given in the next section. Racah coefficients have more convenient symmetry properties than the recoupling coefficients (but less convenient than the 6-j symbols). [2]

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Recoupling coefficients

Coupling of two angular momenta $\mathbf{j_1}$ and $\mathbf{j_2}$ is the construction of simultaneous eigenfunctions of $\mathbf{J^2}$ and J_z , where $\mathbf{J} = \mathbf{j_1} + \mathbf{j_2}$, as explained in the article on Clebsch–Gordan coefficients. The result is

$$|(j_1j_2)JM
angle = \sum_{m_1=-j_1}^{j_1} \sum_{m_2=-j_2}^{j_2} |j_1m_1
angle |j_2m_2
angle \langle j_1m_1j_2m_2|JM
angle,$$

where
$$J=|j_1-j_2|,\ldots,j_1+j_2$$
 and $M=-J,\ldots,J$.

Coupling of three angular momenta $\mathbf{j_1}$, $\mathbf{j_2}$, and $\mathbf{j_3}$, may be done by first coupling $\mathbf{j_1}$ and $\mathbf{j_2}$ to $\mathbf{J_{12}}$ and next coupling $\mathbf{J_{12}}$ and $\mathbf{j_3}$ to total angular momentum \mathbf{J} :

$$|((j_1j_2)J_{12}j_3)JM
angle = \sum_{M_{12}=-J_{12}}^{J_{12}}\sum_{m_3=-j_3}^{j_3}|(j_1j_2)J_{12}M_{12}
angle|j_3m_3
angle\langle J_{12}M_{12}j_3m_3|JM
angle$$

Alternatively, one may first couple $\mathbf{j_2}$ and $\mathbf{j_3}$ to $\mathbf{J_{23}}$ and next couple $\mathbf{j_1}$ and $\mathbf{J_{23}}$ to \mathbf{J} :

$$|(j_1,(j_2j_3)J_{23})JM
angle = \sum_{m_1=-j_1}^{j_1}\sum_{M_{23}=-J_{23}}^{J_{23}}|j_1m_1
angle|(j_2j_3)J_{23}M_{23}
angle\langle j_1m_1J_{23}M_{23}|JM
angle$$

Both coupling schemes result in complete orthonormal bases for the $(2j_1+1)(2j_2+1)(2j_3+1)$ dimensional space spanned by

$$|j_1m_1
angle|j_2m_2
angle|j_3m_3
angle, \;\; m_1=-j_1,\ldots,j_1; \;\; m_2=-j_2,\ldots,j_2; \;\; m_3=-j_3,\ldots,j_3.$$

Hence, the two total angular momentum bases are related by a unitary transformation. The matrix elements of this unitary transformation are given by a $\underline{\text{scalar product}}$ and are known as recoupling coefficients. The coefficients are independent of M and so we have

$$|((j_1j_2)J_{12}j_3)JM
angle = \sum_{J_{23}} |(j_1,(j_2j_3)J_{23})JM
angle \langle (j_1,(j_2j_3)J_{23})J|((j_1j_2)J_{12}j_3)J
angle.$$

The independence of M follows readily by writing this equation for M = J and applying the <u>lowering</u> operator J_{-} to both sides of the equation.

Algebra

Let

$$\Delta(a,b,c) = [(a+b-c)!(a-b+c)!(-a+b+c)!/(a+b+c+1)!]^{1/2}$$

be the usual triangular factor, then the Racah coefficient is a product of four of these by a sum over factorials,

$$W(abcd;ef) = \Delta(a,b,e)\Delta(c,d,e)\Delta(a,c,f)\Delta(b,d,f)w(abcd;ef)$$

where

$$w(abcd;ef) \equiv \sum_z rac{(-1)^{z+eta_1}(z+1)!}{(z-lpha_1)!(z-lpha_2)!(z-lpha_3)!(z-lpha_4)!(eta_1-z)!(eta_2-z)!(eta_3-z)!}$$

and

$$egin{array}{ll} lpha_1 = a + b + e; & eta_1 = a + b + c + d; \ lpha_2 = c + d + e; & eta_2 = a + d + e + f; \ lpha_3 = a + c + f; & eta_3 = b + c + e + f; \ lpha_4 = b + d + f. \end{array}$$

The sum over z is finite over the range^[3]

$$\max(\alpha_1, \alpha_2, \alpha_3, \alpha_4) \leq z \leq \min(\beta_1, \beta_2, \beta_3).$$

Relation to Wigner's 6-j symbol

Racah's W-coefficients are related to Wigner's <u>6-j symbols</u>, which have even more convenient symmetry properties

$$W(abcd;ef)(-1)^{a+b+c+d}=egin{cases} a & b & e \ d & c & f \end{pmatrix}\!.$$

 $Cf.^{[4]}$ or

$$W(j_1j_2Jj_3;J_{12}J_{23})=(-1)^{j_1+j_2+j_3+J}igg\{egin{array}{ccc} j_1 & j_2 & J_{12} \ j_3 & J & J_{23} \ \end{array}igg\}.$$

See also

- Clebsch–Gordan coefficients
- 3-j symbol
- 6-j symbol
- Pandya theorem

Notes

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- 4. Brink, D M & Satchler, G R (1968). *Angular Momentum* (Oxford University Press) 3rd ed., p. 142. online (https://archive.org/details/AngularMomentum)

Further reading

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External links

■ "Racah-Wigner coefficients" (https://www.encyclopediaofmath.org/index.php?title=Racah-Wigner coefficients), *Encyclopedia of Mathematics*, EMS Press, 2001 [1994]

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