1. An electron is in the spin state

$$\chi = A \begin{bmatrix} 3i \\ 4 \end{bmatrix}.$$

(a) Normalize χ to determine A.

We begin by recognizing that

$$\chi = A \begin{bmatrix} 3i \\ 4 \end{bmatrix} = A (3i\langle \uparrow | + 4 | \downarrow \rangle).$$

Then,

$$1 = |\chi|^2 = A^* A ((3i)(-3i)\langle \uparrow | \downarrow \rangle + (4)(4)\langle \downarrow | \downarrow \rangle + 0 + 0) = A^2 (9 + 16),$$

which is satisfied for A = 1/5.

(b) Determine $\langle S_x \rangle$, $\langle S_y \rangle$, $\langle S_z \rangle$.

From Griffiths Eq. 4.145 and Eq. 4.147,

$$\widehat{S}_x = \frac{\hbar}{2} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \quad \widehat{S}_y = \frac{\hbar}{2} \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}, \quad \widehat{S}_z = \frac{\hbar}{2} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}.$$

Then, in general,

$$\langle S_a \rangle = \langle \chi | \widehat{S}_a | \chi \rangle = \frac{1}{25} \begin{bmatrix} -3i & 4 \end{bmatrix} \widehat{S}_a \begin{bmatrix} 3i \\ 4 \end{bmatrix},$$

where a represents x, y, or z. Thus,

$$\langle S_x \rangle = \frac{1}{25} \frac{\hbar}{2} \begin{bmatrix} -3i & 4 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} 3i \\ 4 \end{bmatrix} = 0,$$

$$\langle S_y \rangle = \frac{1}{25} \frac{\hbar}{2} \begin{bmatrix} -3i & 4 \end{bmatrix} \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} 3i \\ 4 \end{bmatrix} = -\frac{24}{50} \hbar,$$

$$\langle S_z \rangle = \frac{1}{25} \frac{\hbar}{2} \begin{bmatrix} -3i & 4 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 3i \\ 4 \end{bmatrix} = -\frac{7}{50} \hbar.$$

(c) Determine the uncertainties σ_{S_x} , σ_{S_y} , and σ_{S_x} . In general,

$$\sigma_a = \sqrt{\langle a^2 \rangle - \langle a \rangle^2},$$

where a represents some observable. Our square spin operators are then

$$\widehat{S_x^2} = \frac{\hbar^2}{4} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}, \quad \widehat{S_y^2} = \frac{\hbar^2}{4} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}, \quad \widehat{S_z^2} = \frac{\hbar^2}{4} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}.$$

Since the operators are equal, their expectation values on χ will be equal. We also notice that the matrix for each operator is the identity matrix. Thus,

$$\left\langle S_{x}^{2}\right\rangle =\left\langle S_{y}^{2}\right\rangle =\left\langle S_{z}^{2}\right\rangle =\frac{\hbar^{2}}{4}\langle\chi|\widehat{S_{x}^{2}}|\chi\rangle =\frac{\hbar^{2}}{4}\langle\chi\,|\,\chi\rangle =\frac{\hbar^{2}}{4}.$$

Therefore,

$$\sigma_{S_x} = \sqrt{\langle S_x^2 \rangle - \langle S_x \rangle^2} = \sqrt{\frac{\hbar^2}{4} - 0} = \frac{\hbar}{2},$$

$$\sigma_{S_y} = \sqrt{\langle S_y^2 \rangle - \langle S_y \rangle^2} = \sqrt{\frac{\hbar^2}{4} - \left(-\frac{12}{25}\hbar\right)^2} = \frac{7}{50}\hbar$$

$$\sigma_{S_z} = \sqrt{\langle S_z^2 \rangle - \langle S_y \rangle^2} = \sqrt{\frac{\hbar^2}{4} - \left(-\frac{7}{50}\hbar\right)^2} = \frac{12}{25}\hbar.$$

(d) Confirm that the results are consistent with all three uncertainty principles; that is,

$$\sigma_{S_x}\sigma_{S_y} \ge \frac{\hbar}{2} \left| \left\langle S_z \right\rangle \right|,$$

and the cyclic permutations. Notice from part (c),

$$\sigma_{S_x} = \frac{\hbar}{2}, \quad \sigma_{S_y} = \frac{7}{50}\hbar = |\langle S_z \rangle|, \quad \sigma_{S_z} = \frac{12}{50}\hbar = |\langle S_y \rangle|.$$

Then,

$$\sigma_{S_x}\sigma_{S_y} \ge \frac{\hbar}{2} \left| \langle S_z \rangle \right| \Rightarrow \frac{\hbar}{2} \left| \langle S_z \rangle \right| \ge \frac{\hbar}{2} \left| \langle S_z \rangle \right|,$$

$$\sigma_{S_y}\sigma_{S_z} \ge \frac{\hbar}{2} \left| \langle S_x \rangle \right| \Rightarrow \left(\frac{7}{50} \hbar \right) \left(\frac{12}{50} \hbar \right) \ge 0,$$

$$\sigma_{S_z}\sigma_{S_x} \ge \frac{\hbar}{2} \left| \langle S_y \rangle \right| \Rightarrow \left| \langle S_y \rangle \right| \frac{\hbar}{2} \ge \frac{\hbar}{2} \left| \langle S_y \rangle \right|.$$

2. For a generalized spinor,

$$\chi = a\chi_{+,z} + b\chi_{-,z} = a\langle\uparrow| + b|\downarrow\rangle,$$

where $\chi_{+,z}$ is spin up and $\chi_{-,z}$ is spin down, compute $\langle S_x \rangle$, $\langle S_y \rangle$, $\langle S_z \rangle$, $\langle S_x^2 \rangle$, $\langle S_y^2 \rangle$, $\langle S_z^2 \rangle$. Check that

$$\langle S_x^2 \rangle + \langle S_y^2 \rangle + \langle S_z^2 \rangle = \langle S^2 \rangle$$
.

From Griffiths Eq. 4.145 and Eq. 4.147,

$$\widehat{S_x} = \frac{\hbar}{2} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \quad \widehat{S_y} = \frac{\hbar}{2} \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}, \quad \widehat{S_z} = \frac{\hbar}{2} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}.$$

Then,

$$\widehat{S}_x^2 = \widehat{S}_y^2 = \widehat{S}_z^2 = \frac{\hbar^2}{4} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}.$$

Consequently,

$$\widehat{S}^2 = \widehat{S}_x^2 + \widehat{S}_y^2 + \widehat{S}_z^2 = \hbar^2 \frac{3}{4} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}.$$

Then,

$$\langle S_x \rangle = \langle \chi | \widehat{S_x} | \chi \rangle = \frac{\hbar}{2} \begin{bmatrix} a^* & b^* \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} a \\ b \end{bmatrix} = \frac{\hbar}{2} (a^*b + b^*a),$$

$$\langle S_y \rangle = \langle \chi | \widehat{S_y} | \chi \rangle = \frac{\hbar}{2} \begin{bmatrix} a^* & b^* \end{bmatrix} \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} a \\ b \end{bmatrix} = i\frac{\hbar}{2} (ab^* - a^*b),$$

$$\langle S_z \rangle = \langle \chi | \widehat{S_z} | \chi \rangle = \frac{\hbar}{2} \begin{bmatrix} a^* & b^* \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} a \\ b \end{bmatrix} = \frac{\hbar}{2} (a^*a + b^*b).$$

Then,

$$\left\langle S_x^2 \right\rangle = \left\langle S_y^2 \right\rangle = \left\langle S_z^2 \right\rangle = \frac{\hbar^2}{4}.$$

Finally,

$$\left\langle S^{2}\right\rangle =\left\langle S_{x}^{2}\right\rangle +\left\langle S_{y}^{2}\right\rangle +\left\langle S_{z}^{2}\right\rangle =\hbar^{2}\frac{3}{4}.$$