

1. Let the wavefunction $\Psi(x, t)$ be a solution to the time-dependent Schrödinger equation when the potential energy is given by $V(x)$. What is the solution to the Schrödinger equation if we now consider a potential of $V(x) + V_0$ where V_0 is a real positive constant.

2. A particle is observed in a quantum state described by the wavefunction

$$\Psi(x, t) = A \exp \left(-a \left(\frac{mx^2}{\hbar} + it \right) \right),$$

where A and a are real positive constants.

(a) Normalize Ψ .

As shown in Griffiths, if Ψ is normalized at any time t it is normalized at all times t . We will normalize Ψ at $t = 0$.

$$\begin{aligned} 1 &= \int_{-\infty}^{\infty} \Psi(x, t = 0) \, dx \\ &= \int_{-\infty}^{\infty} A \exp \left(\frac{-amx^2}{\hbar} \right) \, dx \\ &= A \int_{-\infty}^{\infty} \exp \left(\frac{-amx^2}{\hbar} \right) \, dx \\ &= A \sqrt{\frac{\pi}{\left(\frac{am}{\hbar}\right)}} \\ &= A \sqrt{\frac{\pi \hbar}{am}}. \end{aligned}$$

Thus, $A = \sqrt{am/(\pi \hbar)}$.

(b) What is the potential $V(x)$ that this particle finds itself within?

Ψ is, by definition, a solution to the Schrödinger equation. Thus,

$$i\hbar \frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + V\Psi.$$

Arranging for the potential energy function V we get

$$V = \frac{1}{\Psi} \left(i\hbar \frac{\partial \Psi}{\partial t} + \frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} \right).$$

To make these partial derivatives less threatening we can begin by writing Ψ in the form $\Psi = A\psi(x)\phi(t)$.¹ The A component is simply A as solved for above during normalization. The exponential term in Ψ can be separated into an x -dependent and t -dependent component;

$$\exp\left(\frac{-amx^2}{\hbar} - ait\right) = \exp\left(-\frac{amx^2}{\hbar}\right) \exp(-ait),$$

which become ψ and ϕ respectively. We know

$$\phi(t) = \exp\left(-\frac{iEt}{\hbar}\right) = \exp(-ait),$$

thus $a = E/\hbar$. We can now write V in a more approachable way with ordinary derivatives:

$$V = \frac{1}{A\psi\phi} \left(i\hbar A\psi \frac{d\phi}{dt} + \frac{\hbar^2}{2m} A\phi \frac{d^2\psi}{dx^2} \right).$$

The ordinary derivatives are

$$\begin{aligned} \frac{d\psi}{dt} &= \frac{d}{dt} [\exp(-ait)] \\ &= -ai \exp(-ait) \\ &= -ai\psi, \end{aligned}$$

and

$$\begin{aligned} \frac{d^2\phi}{dx^2} &= \frac{d^2}{dx^2} \left[\exp\left(-\frac{amx^2}{\hbar}\right) \right] \\ &= -\frac{2amx}{\hbar} \exp\left(-\frac{amx^2}{\hbar}\right) \\ &= -\frac{2amx}{\hbar} \psi. \end{aligned}$$

¹Since separable equations are the only ones we know how to deal with Ψ must be separable!

(c) Determine the expectation values $\langle x \rangle$, $\langle x^2 \rangle$, $\langle p \rangle$, $\langle p^2 \rangle$.

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$$\langle x \rangle = \int_{-\infty}^{\infty} x |\psi| \, dx$$

(d) Determine the standard deviations for position, σ_x , and momentum, σ_p .

(e) Are your values for σ_x and σ_p consistent with the uncertainty principle?

3. An electron is trapped in a harmonic quadratic potential. Suppose the expectation value for its position is given by $\langle x \rangle = \frac{a}{2} \sin(\omega t)$. Here, a is a real constant with units of length and ω is an angular frequency. What, if anything, can be concluded about the electron's momentum?