

1. A particle is trapped in an infinite square well with a width a . At $t = 0$ is equally probable for the particle to be found anywhere on the left-side of the well and impossible for it to be found on the right-side of the well.

(a) What is the normalized wavefunction $\Psi(x, 0)$ that represents this initial state?

Note that Ψ breaks one of the fundamental rules about wavefunctions.

The wavefunction, disregarding the cases of x outside the range $[0, a]$, is given by

$$\Psi(x, 0) = \begin{cases} A & 0 \leq x \leq \frac{a}{2} \\ 0 & \frac{a}{2} < x \leq a \end{cases}$$

where A is the normalization constant. We see the issue with Ψ at $x = a/2$ where we have a discontinuity. The normalization constant can be found by the typical normalization process:

$$1 = \int_{-\infty}^{\infty} |\Psi|^2 dx.$$

We recognize $\Psi(x, t = 0)$ takes on non-zero values only on the interval $[0, \frac{a}{2}]$ and that $\Psi(x, t = 0)$ is real.

$$1 = A^2 \int_0^{a/2} dx = A^2 \frac{a}{2}.$$

Thus,

$$A = \sqrt{\frac{2}{a}}$$

and consequently

$$\Psi(x, 0) = \begin{cases} \sqrt{\frac{2}{a}} & 0 \leq x \leq \frac{a}{2} \\ 0 & \frac{a}{2} < x \leq a \end{cases}$$

(b) What is the probability that you would measure an energy of

$$E = \frac{4\pi^2\hbar^2}{2ma^2}$$

at $t = 0$.

The energy of any infinitely-square-well-ed particle is given by

$$E_n = \frac{n^2\pi^2\hbar^2}{2ma^2}.$$

By inspection, our particle is in the $n = 2$ state. The probability of finding a particle in the $n = 2$ state is given by $|C_2|^2$ where C_2 is found through linearization of Ψ . The n th linear coefficient is given as

$$C_n = \int_{-\infty}^{\infty} \psi_n^* \Psi \, dx = \frac{2}{a} \int_0^{a/2} \sin\left(\frac{n\pi}{a}x\right) \, dx = \frac{2}{n\pi} \left(1 - \cos\left(\frac{n\pi}{2}\right)\right).$$

Thus, C_2 is given as

$$C_2 = \frac{2}{\pi},$$

hence, the probability of measuring the 2nd energy state is given as

$$P(E_2) = |C_2|^2 = \left(\frac{2}{\pi}\right)^2 \approx 0.41.$$

2. Consider the standard infinite square well with width a . The stationary state solutions are $\psi_n(x)$.

(a) Compute $\langle x \rangle$ and $\langle x^2 \rangle$ for $\psi_n(x)$.

$$\begin{aligned}\langle x \rangle &= \int_{-\infty}^{\infty} \psi_n^*[x] \psi_n \, dx \\ &= \end{aligned}$$

(b) Compute $\langle p \rangle$ and $\langle p^2 \rangle$ for $\psi_n(x)$.

(c) Compute σ_x and σ_p and confirm that the uncertainty principle is satisfied for all allowed n .

3. A particle infinitely-square-well-ed to width a is initially observed in a quantum state described by the wavefunction

$$\Psi(x, 0) = A(\psi_1(x) + \psi_3(x)),$$

where A is a real positive constant and both $\psi_1(x)$ and $\psi_3(x)$ are solutions to the time-independent Schrödinger equation for $n = 1$ and $n = 3$ respectively.

- (a) Normalize $\Psi(x, 0)$ assuming ψ_1 and ψ_3 are separately normalized.

We begin with typical normalization.

$$1 = \int_{-\infty}^{\infty} |\Psi(x, 0)|^2 dx = A^2 \left(\int_{-\infty}^{\infty} |\psi_1|^2 dx + \int_{-\infty}^{\infty} |\psi_3|^2 dx \right) = 2A^2,$$

where we have used the assumption that ψ_1 and ψ_3 are separately normalized to evaluate the integrals. Thus,

$$A = 2^{-1/2}.$$

- (b) Compute $|\Psi|^2$ and simplify as much as possible.

The time-dependent Ψ is obtained by adding the typical $\phi(t)$ to each stationary state describing Ψ .

$$\Psi = \frac{1}{\sqrt{2}} \left(\psi_1 e^{-it \frac{E_1}{\hbar}} + \psi_3 e^{-it \frac{E_3}{\hbar}} \right).$$

Then,

$$\begin{aligned} |\Psi|^2 &= \Psi^* \Psi \\ &= \frac{1}{2} \left(\psi_1 e^{it \frac{E_1}{\hbar}} + \psi_3 e^{it \frac{E_3}{\hbar}} \right) \left(\psi_1 e^{-it \frac{E_1}{\hbar}} + \psi_3 e^{-it \frac{E_3}{\hbar}} \right) \\ &= \frac{1}{2} \left(\psi_1^2 + \psi_3^2 + \psi_1 \psi_3 \exp \left(\frac{1}{\hbar} (E_1 - E_3 - E_1 + E_3) \right) \right) \\ &= \frac{1}{2} (\psi_1^2 + \psi_3^2). \end{aligned}$$

- (c) If you measured the particle's energy, what value(s) might you possibly obtain and what is the probability of measuring them?

The factor $1/2$ can be interpreted as both $|C_1|^2$ and $|C_3|^2$. Since this is the representation of Ψ as a linear combination of stationary states we say that all other coefficients C_n are 0 for $n \neq 1, 3$. Thus, the energy states E_1 and E_3 are the only possible energy states each with an equal probability of being measured.