

A Discontinuous Petrov-Galerkin method for compressible flow problems

Jesse Chan

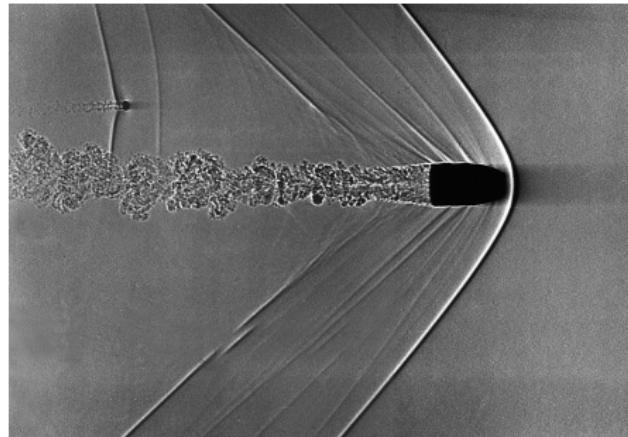
Supervisors: Leszek Demkowicz, Robert Moser

Committee: Todd Arbogast, Omar Ghattas, Venkat Raman

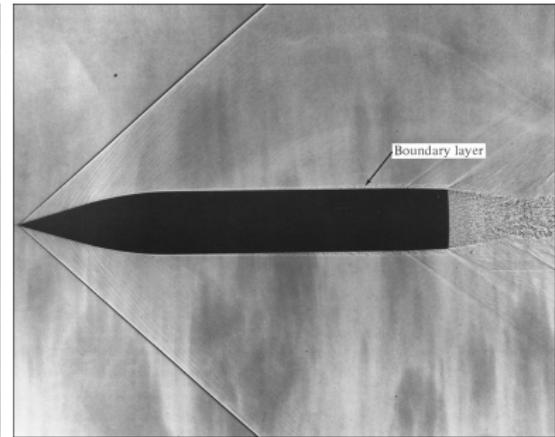
Institute for Computational Engineering and Sciences

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Phenomena in compressible flow



(a) Shock wave



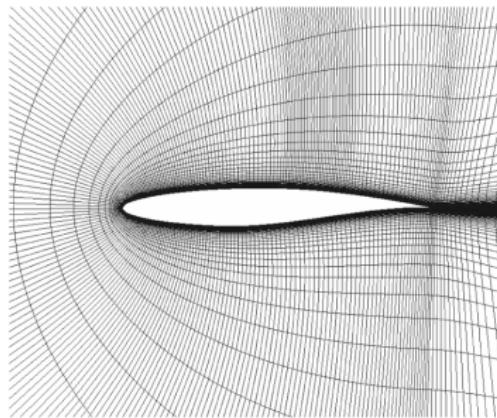
(b) Boundary layer

Compressible flow plays an important role in the aerospace and energy industries - transonic and supersonic aircraft, combustion engines, etc.

Compressible Navier-Stokes equations

Numerical difficulties:

- Resolving solution features
(sharp, localized viscous-scale phenomena)
 - Shocks
 - Boundary layers - resolution needed for drag/load
 - Turbulence (non-localized)
- Nonlinear convergence and uniqueness of solutions
- Stability of numerical schemes
 - Coarse/adaptive grids
 - Higher order



Idea: begin first with the model problem of convection-diffusion.

Robustness: convection-diffusion as a model problem

$$\nabla \cdot (\beta u) - \epsilon \Delta u = f, \quad \text{on } \Omega \in \mathbb{R}^3$$

In 1D: $\beta u' - \epsilon u'' = f$. Standard continuous Galerkin variational formulation: solve

$$b(u, v) = \ell(v), \quad u, v \in H_0^1([0, 1])$$

where

$$b(u, v) = \int_{\Omega} -\beta u v' + \epsilon u' v'$$

$$\ell(v) = \int_{\Omega} f v$$

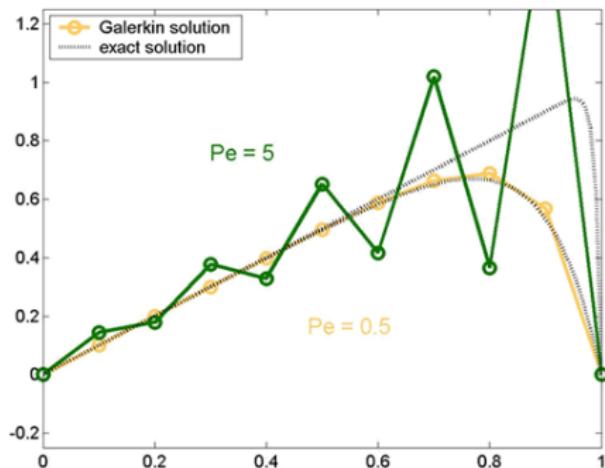


Figure: Solution for $f = 1$. Oscillations in the standard Galerkin method for large Peclet numbers $\text{Pe} := \frac{h}{\epsilon}$.

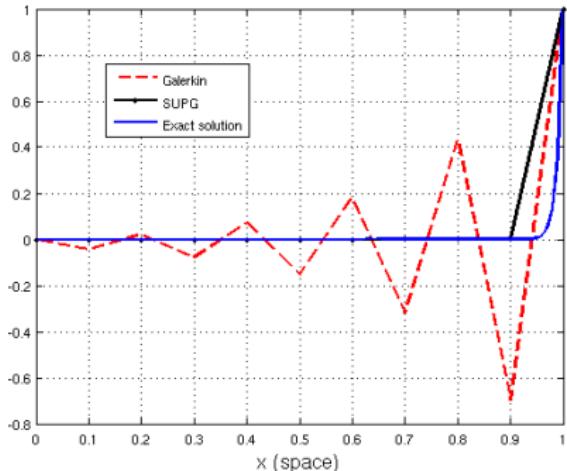
Streamline-upwind Petrov-Galerkin (SUPG)

SUPG solves $b_{\text{SUPG}}(u, v) = l_{\text{SUPG}}(v)$, where

$$b_{\text{SUPG}}(u, v) = b(u, v) + \sum_K \int_K \tau (L_{\text{adv}} v) L u$$

$$l_{\text{SUPG}}(v) = \ell(v) + \sum_K \int_K \tau (L_{\text{adv}} v) f.$$

- $L u = \nabla \cdot (\beta u) - \epsilon \Delta u$,
 $L_{\text{adv}} u = \nabla \cdot (\beta u)$, and τ is
 a parameter.
- Effective for $f \neq 0$.
- *Residual-based* stabilization.



Can be interpreted as a
Petrov-Galerkin method,

$$b(u, \tilde{v}_i) = \ell(\tilde{v}_i), \quad \forall i = 1, \dots, N-1,$$

where the SUPG test function \tilde{v}_i is
defined elementwise as¹

$$\tilde{v}_i = \phi_i(x) + \tau L_{\text{adv}} \phi_i.$$

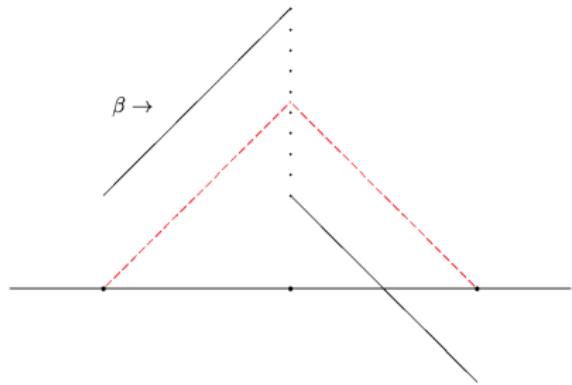


Figure: SUPG test function v_i .

¹A. Brooks and T. Hughes. Streamline upwind/Petrov-Galerkin formulations for convection dominated flows with particular emphasis on the incompressible Navier-Stokes equations. *Comp. Meth. Appl. Mech. Engr*, 32:199–259, 1982

DPG: a minimum residual method via optimal testing

Given a trial space U and Hilbert test space V ,

$$b(u, v) = \ell(v), \quad u \in U, \quad v \in V$$

This is equivalent to the operator equation posed in V'

$$Bu = \ell$$

if we identify $B : U \rightarrow V'$ and $\ell \in V'$ such that

$$\begin{aligned} \langle Bu, v \rangle_V &:= b(u, v), \quad u \in U, v \in V, \\ \langle \ell, v \rangle_V &:= \ell(v), \quad v \in V. \end{aligned}$$

We seek to minimize the dual functional over $U_h \subset U$

$$J(u_h) = \frac{1}{2} \|Bu_h - \ell\|_{V'}^2 := \frac{1}{2} \sup_{v \in V \setminus \{0\}} \frac{|b(u_h, v) - \ell(v)|^2}{\|v\|_V^2}.$$

Let $R_V : V \rightarrow V'$ be the isometric Riesz map st.

$$\langle R_V v, \delta v \rangle_V := (v, \delta v)_V, \quad \forall \delta v \in V.$$

Then, our functional $J(u_h)$ is equal to

$$\min_{u_h \in U_h} J(u_h) = \frac{1}{2} \|Bu_h - \ell\|_{V'}^2 = \frac{1}{2} \|R_V^{-1}(Bu_h - \ell)\|_V^2.$$

First order optimality: Gâteaux derivative is zero in all directions $\delta u \in U_h$

$$\begin{aligned} & (R_V^{-1}(Bu_h - \ell), R_V^{-1}B\delta u)_V = 0, \quad \forall \delta u \in U_h. \\ & \rightarrow \langle (Bu_h - \ell), R_V^{-1}B\delta u \rangle = 0, \\ & \rightarrow b(u_h, R_V^{-1}B\delta u) - \ell(R_V^{-1}B\delta u) = 0 \end{aligned}$$

Summary: select test functions to minimize residuals

For $\delta u \in U_h$, define the **optimal test function** $v_{\delta u}$.

$$v_{\delta u} := R_V^{-1} B \delta u.$$

Then, the residual $J(u_h) = \frac{1}{2} \|Bu_h - \ell\|_{V'}^2$ is minimized by the solution of

$$b(u_h, v_{\delta u}) = \ell(v_{\delta u}), \quad \forall \delta u \in U_h.$$

Practical details of DPG

Computation of $v_{\delta u} := R_V^{-1} B \delta u$ is **global** and **infinite-dimensional**.

- By choosing a **broken** test space V and **localizable** norm $\|v\|_V^2 = \sum_K \|v\|_{V(K)}^2$, test functions can be determined locally.
- In practice, we use an **enriched space** $V_h \subset V$, where $\dim(V_h) > \dim(U_h)$ elementwise, and **optimal test functions** are approximated by computing $v_{\delta u} := R_{V_h}^{-1} B \delta u \in V_h$ through²

$$(v_{\delta u}, \delta v)_V = b(\delta u, \delta v), \quad \delta u \in U_h, \quad \forall \delta v \in V_h$$

Typically, if $U_h = \mathcal{P}^p(\mathbb{R}^n)$, $V_h = \mathcal{P}^{p+\Delta p}(\mathbb{R}^n)$, where $\Delta p \geq n$.³

²L. Demkowicz and J. Gopalakrishnan. A class of discontinuous Petrov-Galerkin methods. II. Optimal test functions. *Num. Meth. for Partial Diff. Eq.*, 27:70–105, 2011

³J. Gopalakrishnan and W. Qiu. An analysis of the practical DPG method. Technical report, IMA, 2011. Submitted

Properties of DPG

DPG provides a symmetric positive-definite stiffness matrix. Let $\{\phi_j\}_{j=1}^m$ be a basis for U_h , and $\{v_i\}_{i=1}^n$ a basis for V_h , such that $n > m$. Then, for

$$\begin{aligned}B_{ji} &= b(\phi_j, v_i), \\l_i &= \ell(v_i),\end{aligned}$$

DPG solves the discrete system for degrees of freedom u

$$\left(B^T R_V^{-1} B \right) u = \left(B^T R_V^{-1} \right) l,$$

where, under a localizable norm and discontinuous test functions, R_V^{-1} is block diagonal.

Properties of DPG

Additional properties of DPG include⁴

- DPG provides the best approximation in the **energy norm**

$$\|u\|_E = \|Bu\|_{V'} = \sup_{\|v\|_V=1} |b(u, v)|.$$

- The energy error is computable through the residual

$$\|u - u_h\|_E = \|B(u - u_h)\|_{V'} = \|R_V^{-1}(I - Bu_h)\|_V = \|e\|_V$$

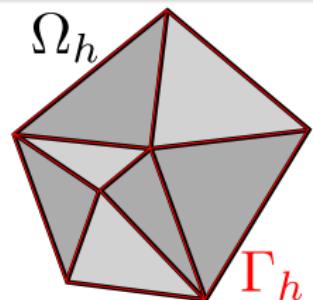
where the **error representation function** e is defined through
 $(e, \delta v)_V = \ell(v) - b(u_h, \delta v)$ for all $\delta v \in V$.

⁴L. Demkowicz, J. Gopalakrishnan, and A. Niemi. A class of discontinuous Petrov-Galerkin methods. Part III: Adaptivity. *Appl. Numer. Math.*, 62(4):396–427, April 2012

Ultra-weak formulation

Given a first order system $Au = f$, we identify the **partition** Ω_h and **mesh skeleton** Γ_h .

The ultra-weak formulation for $Au = f$ on Ω_h is



$$b((u, \hat{u}), v) := \sum_K \langle \hat{u}, v \rangle_{\partial K} + (u, A_h^* v)_{\Omega_h} = (f, v)_{\Omega_h}.$$

Under proper assumptions, $\sum_K \langle \hat{u}, v \rangle_{\partial K} = \langle \hat{u}, [v] \rangle_{\Gamma_h}$, with energy setting

$$u \in L^2(\Omega_h) \equiv L^2(\Omega), \quad v \in V = D(A_h^*), \quad \hat{u} \in \gamma(D(A)),$$

where $D(A_h^*)$ is the broken graph space of the formal adjoint A_h^* , and $\gamma(D(A))$ the trace space of the graph space of operator A .⁵

⁵L. Demkowicz and J. Gopalakrishnan. Analysis of the DPG method for the Poisson equation. *SIAM J. Numer. Anal.*, 49(5):1788–1809, September 2011

The canonical “graph” test norm

Recall $\|u\|_E := \sup_{v \in V \setminus \{0\}} \frac{b(u, v)}{\|v\|_V}$. Under the ultra-weak formulation, the trial norm

$$\|(u, \hat{u})\|_U := \|u\|_{L^2(\Omega)}^2 + \|\hat{u}\|^2$$

generates a test norm *equivalent* to

$$\|v\|_{V_{\text{opt}}} := \|A_h^* v\|_{L^2(\Omega)}^2 + \left(\sup_{\hat{u}} \frac{\langle \hat{u}, [v] \rangle_{\Gamma_h}}{\|\hat{u}\|} \right)^2.$$

This norm is not localizable, so we instead substitute the jump terms for an L^2 term, giving us the *graph norm*⁶

$$\|v\|_V := \|A_h^* v\|_{L^2(\Omega)}^2 + \|v\|_{L^2}.$$

⁶T. Bui-Thanh, L. Demkowicz, and O. Ghattas. A unified discontinuous Petrov-Galerkin method and its analysis for Friedrichs' systems. *Submitted to SIAM J. Numer. Anal.*, 2011. Also ICES report 11-34, November 2011

Goal: robust higher order adaptive methods

Our aim is to design an adaptive method for compressible laminar flows that is

- robust in Reynolds number
- stable for arbitrary elements

This proposal will outline

- 1 a robust DPG method for convection-diffusion,
- 2 extension to nonlinear problems and systems,
- 3 proposed work.

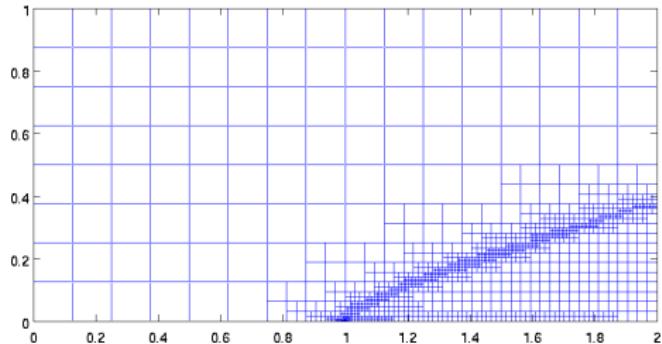
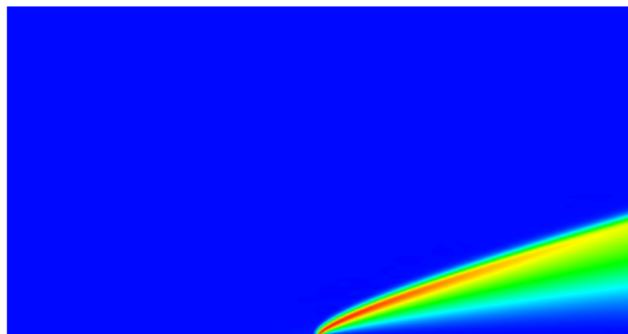


Figure: u_2 and $p = 2$ mesh for $Re = 1000$.

Ultra-weak formulation for convection-diffusion

In first order form, the convection-diffusion equation is

$$A(u, \sigma) := \begin{bmatrix} \nabla \cdot (\beta u - \sigma) \\ \frac{1}{\epsilon} \sigma - \nabla u \end{bmatrix} = \begin{bmatrix} f \\ 0 \end{bmatrix}.$$

The variational formulation is

$$\begin{aligned} b\left(\left(u, \sigma, \hat{u}, \hat{f}_n\right), (v, \tau)\right) &= (u, \nabla_h \cdot \tau - \beta \cdot \nabla_h v)_{\Omega_h} + (\sigma, \epsilon^{-1} \tau + \nabla_h v)_{\Omega_h} \\ &\quad - \langle [\![\tau \cdot n]\!], \hat{u} \rangle_{\Gamma_h} + \left\langle \hat{f}_n, [\![v]\!] \right\rangle_{\Gamma_h}, \end{aligned}$$

where $\hat{f}_n := \beta_n u - \sigma_n$ and $\left\langle \hat{f}_n, [\![v]\!] \right\rangle_{\Gamma_h}$ is defined

$$\left\langle \hat{f}_n, [\![v]\!] \right\rangle_{\Gamma_h} := \sum_K \int_{\partial K} \operatorname{sgn}(\vec{n}) \hat{f}_n v.$$

Graph norm under convection-diffusion

For convection-diffusion, the graph test norm is defined elementwise

$$\|(\boldsymbol{v}, \tau)\|_{V(K)}^2 = \|\nabla \cdot \tau - \beta \cdot \nabla \boldsymbol{v}\|_{L^2(K)}^2 + \|\epsilon^{-1} \tau + \nabla \boldsymbol{v}\|_{L^2(K)}^2 + \|\boldsymbol{v}\|_{L^2(K)}^2.$$

Problem with this test norm: approximability of test functions.

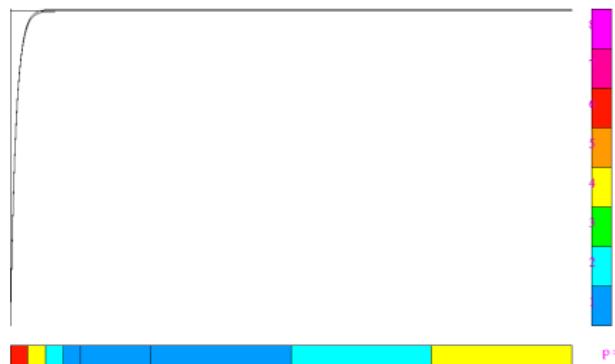


Figure: \boldsymbol{v} and τ components of the 1D optimal test functions for flux \hat{f}_n on the *right-hand* side of a unit element for $\epsilon = 0.01$.

Determining an alternative test norm

Recall the convection-diffusion bilinear form

$$\begin{aligned} b\left(\left(u, \sigma, \hat{u}, \hat{f}_n\right), (v, \tau)\right) = & (u, \nabla \cdot \tau - \beta \cdot \nabla v)_{\Omega_h} + (\sigma, \epsilon^{-1}\tau + \nabla v)_{\Omega_h} \\ & - \langle [\![\tau \cdot n]\!], \hat{u} \rangle_{\Gamma_h} + \left\langle \hat{f}_n, [\![v]\!] \right\rangle_{\Gamma_h}, \end{aligned}$$

We recover $\|u\|_{L^2(\Omega)}^2$ by choosing conforming (v, τ) satisfying the *adjoint equations*

$$\nabla \cdot \tau - \beta \cdot \nabla v = u$$

$$\frac{1}{\epsilon}\tau + \nabla v = 0$$

with boundary conditions s.t. $\langle [\![\tau \cdot n]\!], \hat{u} \rangle_{\Gamma}$ and $\left\langle \hat{f}_n, [\![v]\!] \right\rangle_{\Gamma}$ vanish.

A robust bound

“Necessary” conditions for robustness — let $\mathbf{U} = (u, \sigma, \hat{u}, \hat{f}_n)$. Then, by choosing specific (v, τ) satisfying the adjoint equations,

$$\|u\|_{L^2(\Omega)}^2 = b(\mathbf{U}, (v, \tau)) = \frac{b(\mathbf{U}, (v, \tau))}{\|(v, \tau)\|_V} \|(v, \tau)\|_V \leq \|\mathbf{U}\|_E \|(v, \tau)\|_V$$

Let \lesssim denote a robust bound - if $\|(v, \tau)\|_V \lesssim \|u\|_{L^2(\Omega)}$, then we have that

$$\|u\|_{L^2(\Omega)} \lesssim \|\mathbf{U}\|_E$$

Main idea: the test norm should measure adjoint solutions robustly.

Choice of inflow boundary condition

We impose the standard outflow wall boundary condition on u . For inflow boundary condition:

- The standard choice of inflow boundary condition: $u = u_0$.
- We impose the non-standard inflow condition:
 $\hat{f}_n := \beta_n u - \sigma_n \approx \beta_n u_0$ on Γ_{in} .

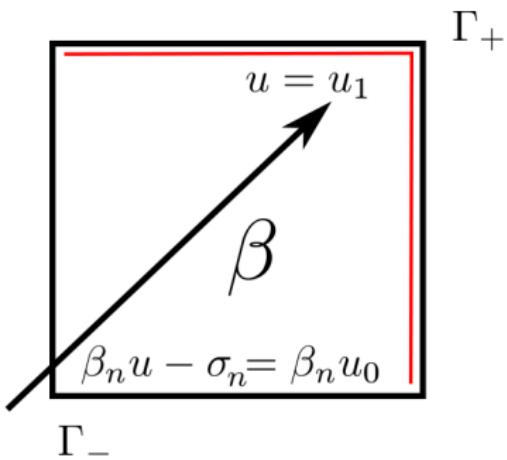


Figure: Non-standard inflow.

When $\sigma_n \approx 0$ near the inflow, condition on \hat{f}_n approximates condition on u .

Dirichlet inflow boundary condition

Standard choice of boundary condition: $u = u_0$ on inflow boundary Γ_{in} , induces boundary layers in adjoint problems, $\|\beta \cdot \nabla v\|_{L^2} = O(\epsilon^{-1})$.

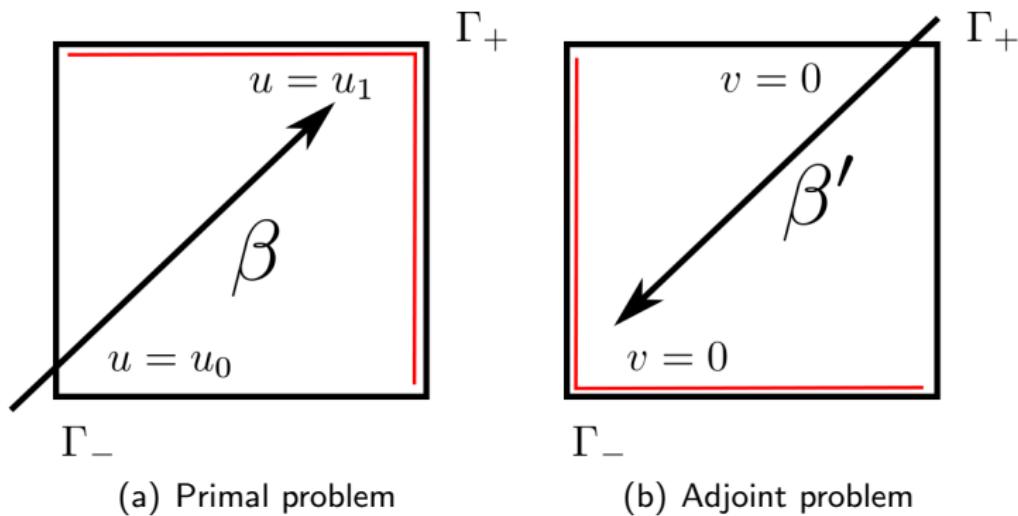


Figure: For the standard Dirichlet inflow condition, the solution to the adjoint problem can develop strong boundary layers at the outflow of the adjoint problem.

New inflow boundary condition on \hat{f}_n

Non-standard choice of boundary condition: $\hat{f}_n = \beta_n u_0$ on Γ_{in} , induces smoother adjoint problems, $\|\beta \cdot \nabla v\|_{L^2} = O(1)$.

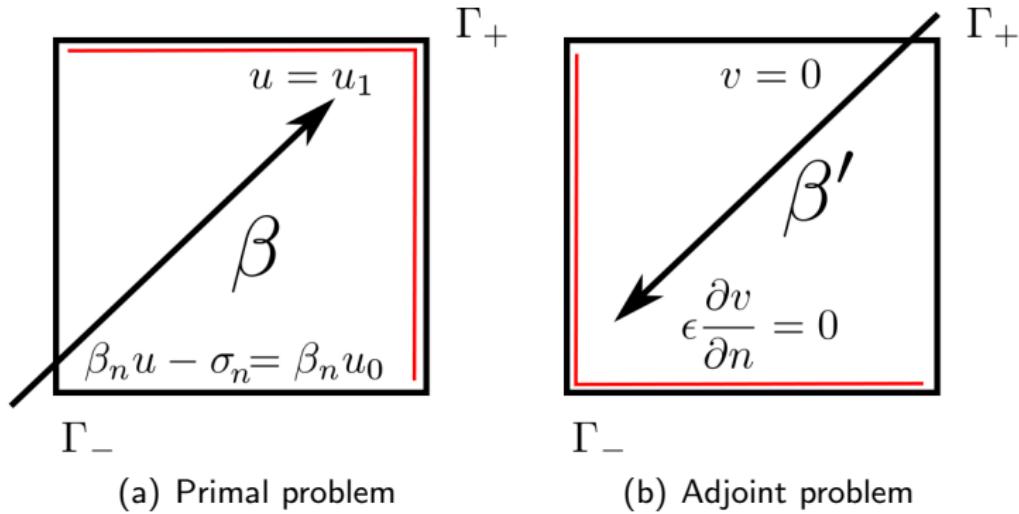
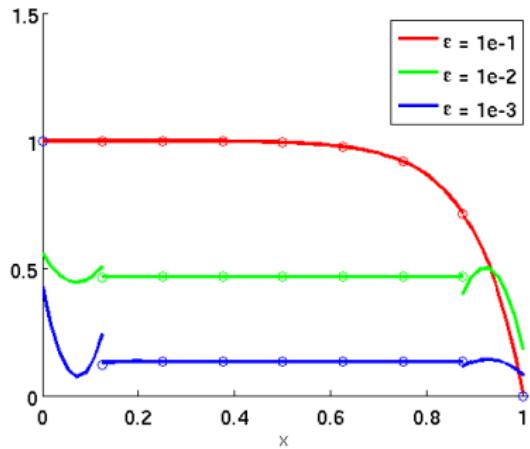


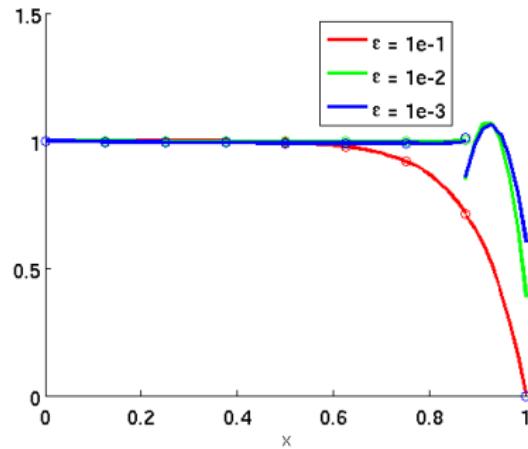
Figure: Under the new inflow condition, the wall-stop boundary condition is relaxed to a zero-stress condition at the outflow boundary of the adjoint problem.

Test norms and adjoint solutions

Intuition: the effectiveness of DPG under a test norm is governed by how a **specific test norm** measures the **solutions of the adjoint problem**.



(a) Dirichlet inflow



(b) "Convection" inflow

Figure: DPG solutions to convection-diffusion for both inflow conditions using an H^1 test norm.

Adjoint estimates, or “norm building blocks”

For solutions (v, τ) of the adjoint equations, the following quantities are robustly bounded from above by $\|u\|_{L^2(\Omega)}$.

$$\begin{aligned} & \|v\|, \sqrt{\epsilon} \|\nabla v\|, \|\beta \cdot \nabla v\| \\ & \|\nabla \cdot \tau\|, \frac{1}{\sqrt{\epsilon}} \|\tau\|. \end{aligned}$$

We will construct a test norm through a combination of the above quantities, such that

- v and τ decoupled (no systems).
- Coefficients are of equal order after transforming to the unit element (no boundary layers).

Mesh-scaled test norms

Our test norm, as defined over a single element K , is now

$$\|(\boldsymbol{v}, \tau)\|_{V,K}^2 = \min \left\{ \frac{\epsilon}{|K|}, 1 \right\} \|\boldsymbol{v}\|^2 + \epsilon \|\nabla \boldsymbol{v}\|^2 + \|\beta \cdot \nabla \boldsymbol{v}\|^2 + \|\nabla \cdot \tau\|^2 + \min \left\{ \frac{1}{\epsilon}, \frac{1}{|K|} \right\} \|\tau\|^2.$$

which induces the proven *robust* bound⁷

$$\|\boldsymbol{u}\|_{L^2(\Omega)} + \|\sigma\|_{L^2(\Omega)} + \epsilon \|\hat{\boldsymbol{u}}\| + \sqrt{\epsilon} \|\hat{\boldsymbol{f}}_n\| \lesssim \left\| (\boldsymbol{u}, \sigma, \hat{\boldsymbol{u}}, \hat{\boldsymbol{f}}_n) \right\|_E.$$

⁷ J. Chan, N. Heuer, T. Bui Thanh, and L. Demkowicz. Robust DPG method for convection-diffusion problems II: natural inflow conditions. Technical Report 12-21, ICES, June 2012. Submitted

Eriksson-Johnson model problem

On domain $\Omega = [0, 1]^2$, with $\beta = (1, 0)^T$, $f = 0$ and boundary conditions

$$\widehat{\beta_n u - \sigma_n} = \widehat{f}_n = u_0, \quad \beta_n \leq 0$$

$$\widehat{u} = 0, \quad \beta_n > 0$$

All numerical experiments done using Camellia codebase.⁸

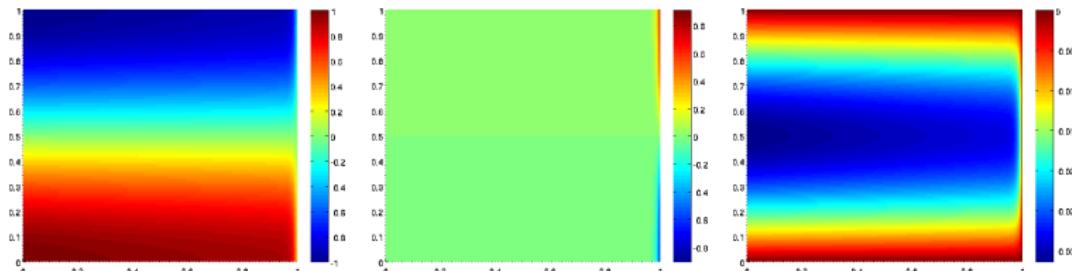


Figure: Exact solution for u , σ_x , and σ_y for $\epsilon = .01$, $C_1 = 1$, $C_n = 0$, $n \neq 1$

⁸N. Roberts, D. Ridzal, P. Bochev, and L. Demkowicz. A Toolbox for a Class of Discontinuous Petrov-Galerkin Methods Using Trilinos. Technical Report SAND2011-6678, Sandia National Laboratories, 2011

Error rates

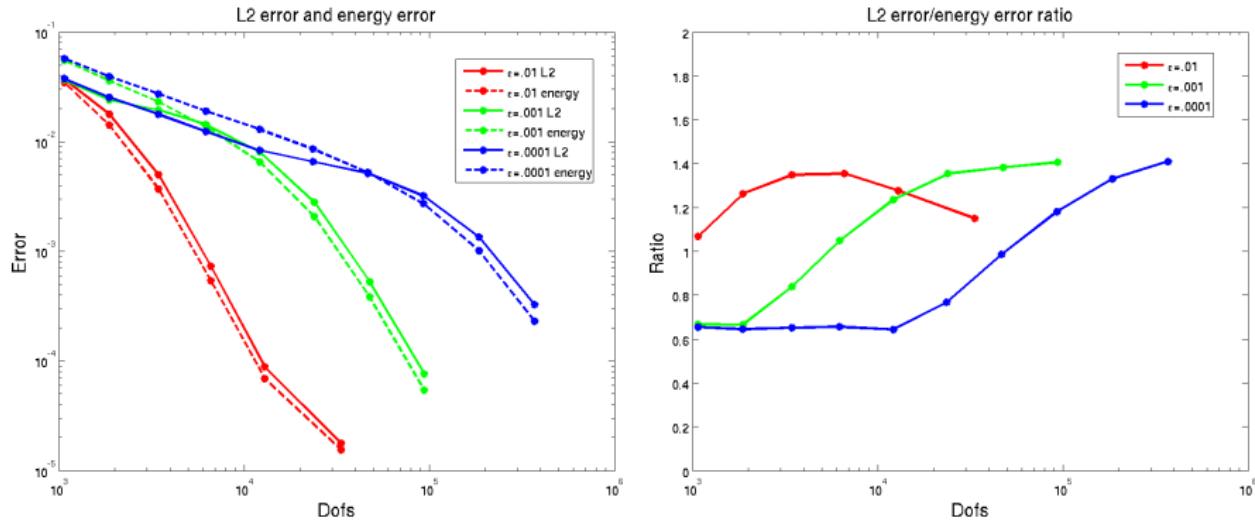


Figure: L^2 /energy errors and their ratio for $\epsilon = .01, .001, .0001$.

Regularization

For $\beta = (-y, x)^T$ on $\Omega = [-1, 1]^2$. Ill posed in the convection setting.
Similar tests have been done with discontinuous data.

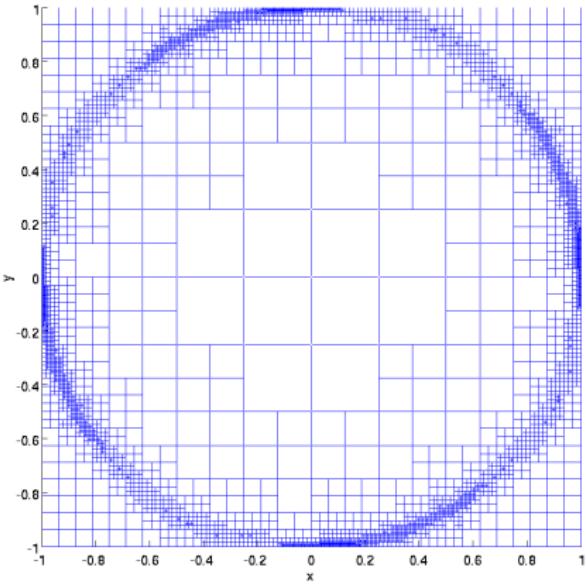
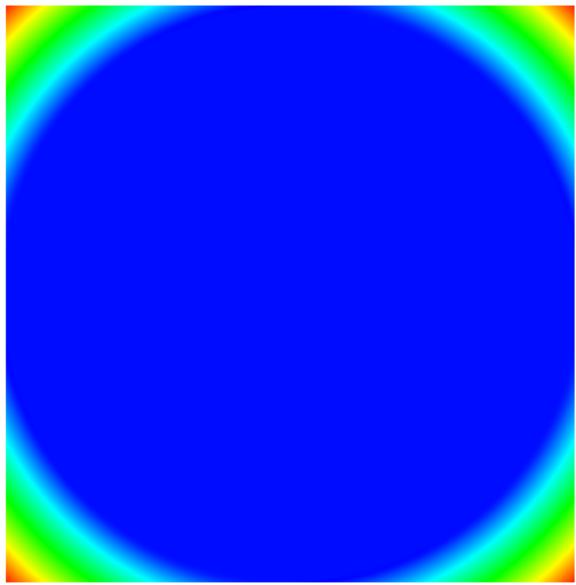


Figure: Steady vortex problem with $\epsilon = 1e - 4$, using a variant of the test norm.

DPG for nonlinear problems

Given some linearization technique (typically Newton-Raphson and pseudo-timestepping linearization), we measure

- size of the linearized update Δu

$$\|\Delta u\|_E := \|B_u \Delta u\|_{V'} = \|R_V^{-1} B_u \Delta u\|_V$$

- the nonlinear residual

$$\|R(u)\|_E := \|B(u) - \ell\|_{V'} = \|R_V^{-1} (B(u) - \ell)\|_V$$

Preliminary experiments were done in 1D⁹ and space-time.¹⁰

⁹J. Chan, L. Demkowicz, R. Moser, and N. Roberts. A New Discontinuous Petrov-Galerkin Method with Optimal Test Functions. Part V: Solution of 1D Burgers' and Navier-Stokes Equations. Technical Report 10-25, ICES, June 2010

¹⁰J. Chan, L. Demkowicz, and M. Shashkov. Space-time DPG for shock problems. Technical Report LA-UR 11-05511, LANL, September 2011

2D test case: Burgers equation

$$\frac{\partial (u^2/2)}{\partial x} + \frac{\partial u}{\partial y} + \epsilon \Delta u = f$$

Burgers equation can be written
with $\beta(u) = (u/2, 1)$

$$\nabla \cdot (\beta(u)u - \sigma) = f$$

$$\frac{1}{\epsilon} \sigma - \nabla u = 0.$$

i.e. nonlinear convection-diffusion
on domain $[0, 1]^2$.

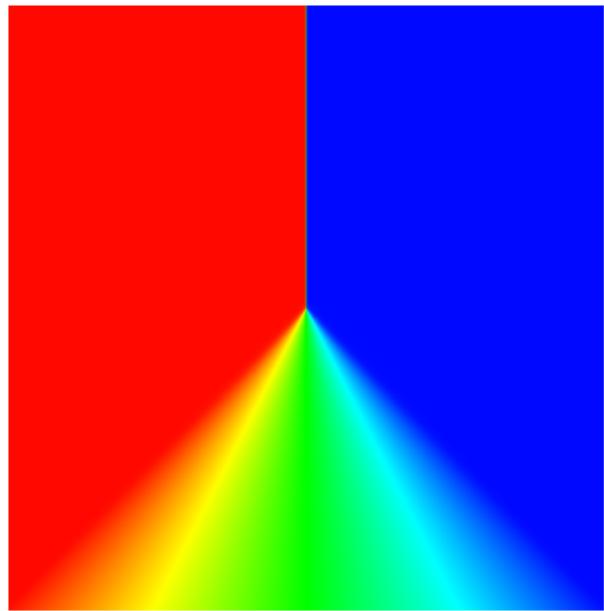


Figure: Shock solution for Burgers' equation, $\epsilon = 1e-4$, using Newton-Raphson.

Adaptivity begins with a cubic 4×4 mesh.

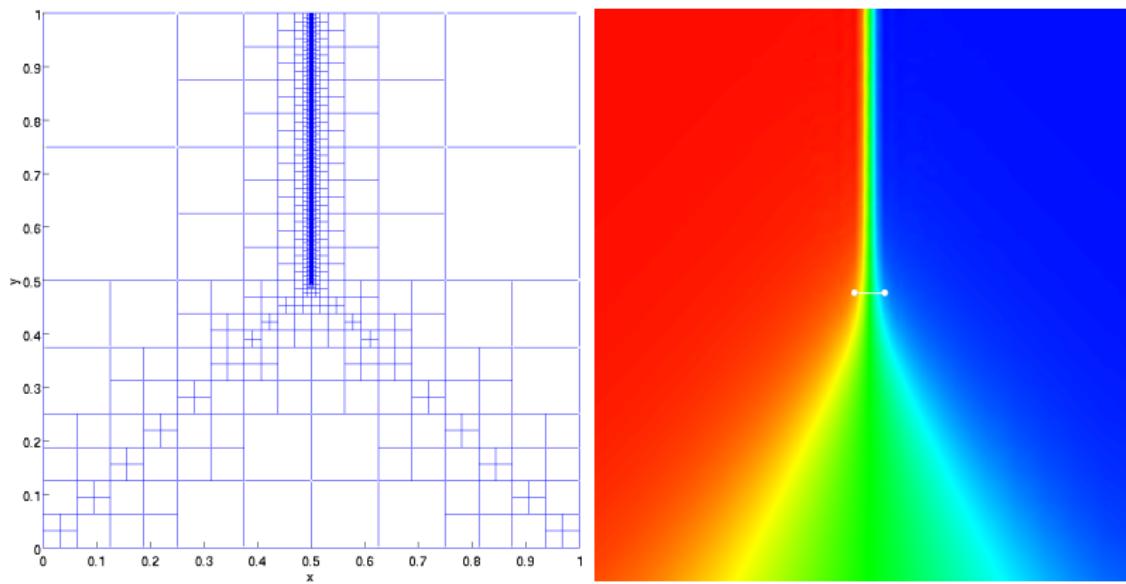


Figure: Adaptive mesh after 9 refinements, and zoom view at point $(.5,.5)$ with shock formation and $1e - 3$ width line for reference.

2D Compressible Navier-Stokes equations (ideal gas)

Given density ρ , velocities $\mathbf{u} = (u_1, u_2)$ and temperature T ,

$$\nabla \cdot \begin{bmatrix} \rho u_1 \\ \rho u_2 \end{bmatrix} = 0$$

$$\nabla \cdot \left(\begin{bmatrix} \rho u_1^2 + p \\ \rho u_1 u_2 \end{bmatrix} - \boldsymbol{\sigma}_1 \right) = 0$$

$$\nabla \cdot \left(\begin{bmatrix} \rho u_1 u_2 \\ \rho u_2^2 + p \end{bmatrix} - \boldsymbol{\sigma}_2 \right) = 0$$

$$\nabla \cdot \left(\begin{bmatrix} ((\rho e) + p) u_1 \\ ((\rho e) + p) u_2 \end{bmatrix} - \boldsymbol{\sigma} \mathbf{u} + \vec{q} \right) = 0$$

$$\frac{1}{2\mu} \boldsymbol{\sigma} - \frac{\lambda}{4\mu(\mu + \lambda)} \text{tr}(\boldsymbol{\sigma}) \mathbf{I} = \nabla \mathbf{u} - \text{Re } \boldsymbol{\omega}$$

$$\frac{1}{\kappa} \vec{q} = \nabla T$$

Stress law

σ for a Newtonian fluid: $\sigma_{ij} = \mu(u_{i,j} + u_{j,i}) + \lambda u_{k,k} \delta_{ij}$.

$$\frac{1}{2} (\nabla U + \nabla^T U) = \frac{1}{2\mu} \sigma_{ij} - \frac{\lambda}{4\mu(\mu + \lambda)} \sigma_{kk} \delta_{ij}$$

The symmetric part of the gradient is

$$\frac{1}{2} (\nabla U + \nabla^T U) = \nabla U - \omega$$

Our final form is

$$\nabla U - \omega = \frac{1}{2\mu} \sigma - \frac{\lambda}{4\mu(\mu + \lambda)} \text{tr}(\sigma) \mathbf{I}.$$

By enforcing strong symmetry of σ , taking the antisymmetric part implicitly defines ω to be the antisymmetric part of ∇u .

Extrapolation of test norms

Convection-diffusion:

$$\nabla \cdot (\beta u - \sigma) = f$$

$$\frac{1}{\epsilon} \sigma - \nabla u = 0.$$

Navier-Stokes: defining vector variables $U = \{\rho, u_1, u_2, T\}$ and $\Sigma = \{\sigma, \mathbf{q}, \omega\}$,

$$\nabla \cdot (A_{\text{invisc}} U - A_{\text{visc}} \Sigma) = R_{\text{conserv}}(U, \Sigma)$$

$$E_{\text{visc}} \Sigma - \nabla U = R_{\text{constit}}(U, \Sigma)$$

where $R_{\text{conserv}}(U, \Sigma)$ and $R_{\text{constit}}(U, \Sigma)$ are the conservation/constitutive residuals.

Test norms over one element

Convection-diffusion:

$$\begin{aligned} \|(\mathbf{v}, \tau)\|_{V,K}^2 &= \min \left\{ \frac{\epsilon}{|K|}, 1 \right\} \|\mathbf{v}\|^2 + \|\boldsymbol{\beta} \cdot \nabla \mathbf{v}\|^2 + \epsilon \|\nabla \mathbf{v}\|^2 \\ &\quad + \|\nabla \cdot \tau\|^2 + \min \left\{ \frac{1}{\epsilon}, \frac{1}{|K|} \right\} \|\tau\|^2. \end{aligned}$$

Navier-Stokes: let V and W be vectors of test functions v_i and τ_i .

$$\begin{aligned} \|(V, W)\|_{V,K}^2 &= \min \left\{ \frac{1}{\text{Re}|K|}, 1 \right\} \|V\|^2 + \|A_{\text{invisc}}^T \nabla V\|^2 + \frac{1}{\text{Re}} \|A_{\text{visc}}^T \nabla V\|^2 \\ &\quad + \|\nabla \cdot W\|^2 + \min \left\{ 1, \frac{1}{\text{Re}|K|} \right\} \|E_{\text{visc}}^T W\|^2. \end{aligned}$$

Carter's flat plate problem

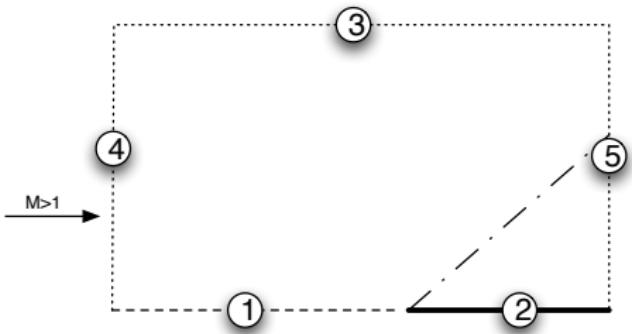
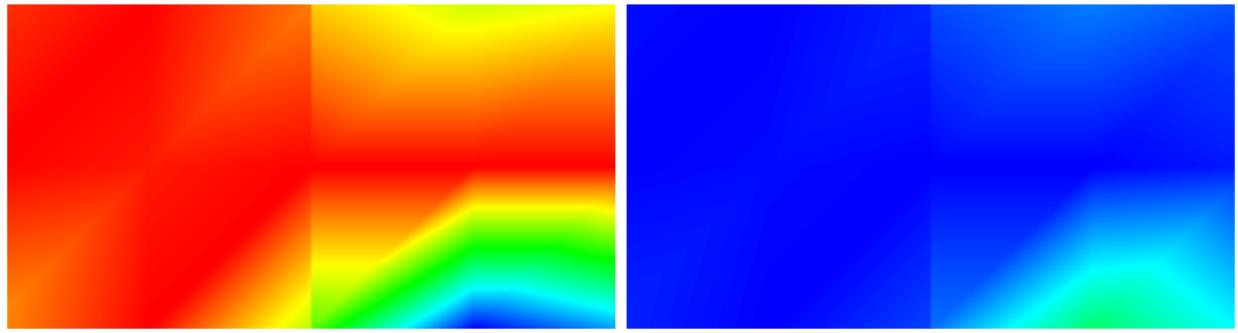
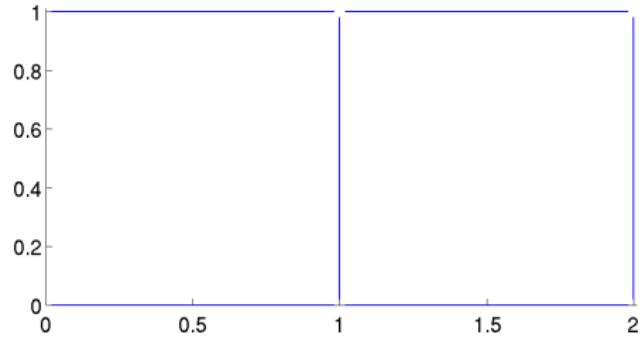


Figure: Carter flat plate problem on domain $[0, 2] \times [0, 1]$. Plate begins at $x = 1$, $\text{Re} = 1000$.

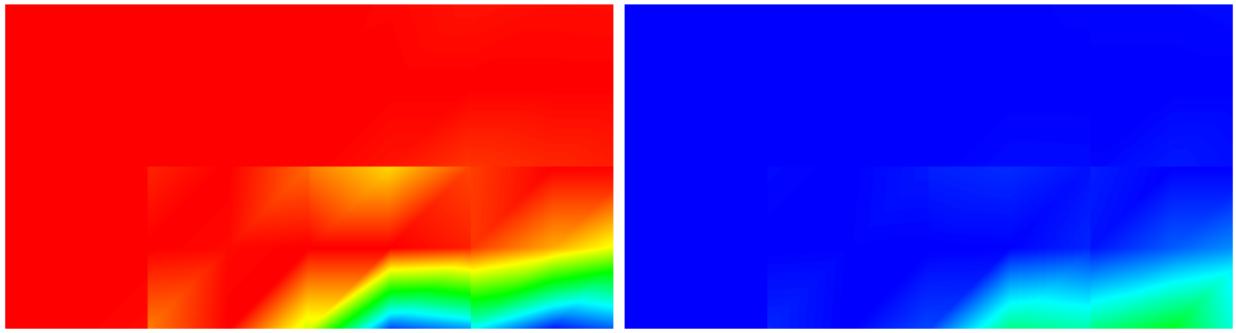
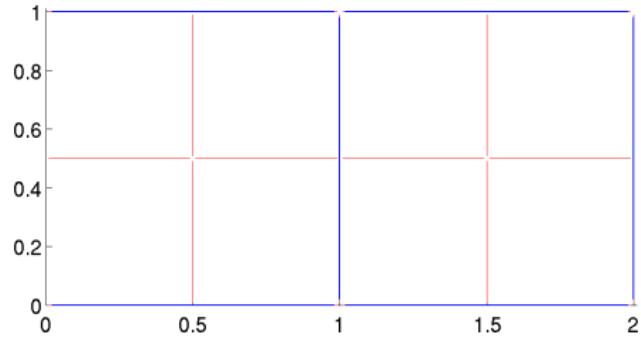
- 1 Symmetry boundary conditions.
- 2 Prescribed temperature and wall stagnation conditions.
- 3 Symmetry boundary conditions.
- 4 Inflow: conserved quantities specified using far-field values.
- 5 No outflow condition set.

Stress/heat flux boundary conditions are set in terms of the momentum and energy fluxes.

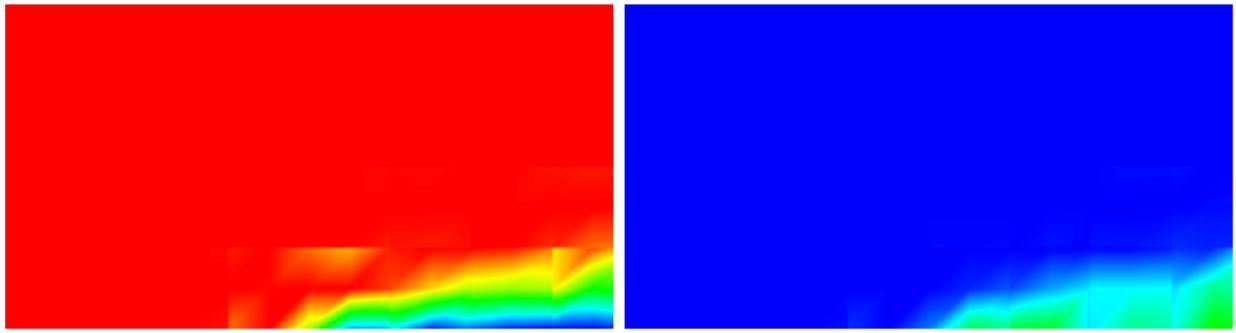
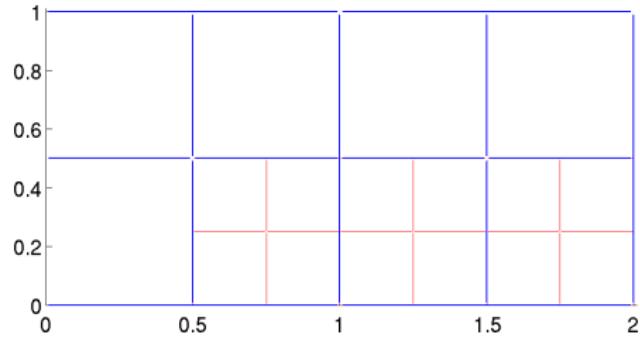
Refinement level 0

(a) u_1 (b) T 

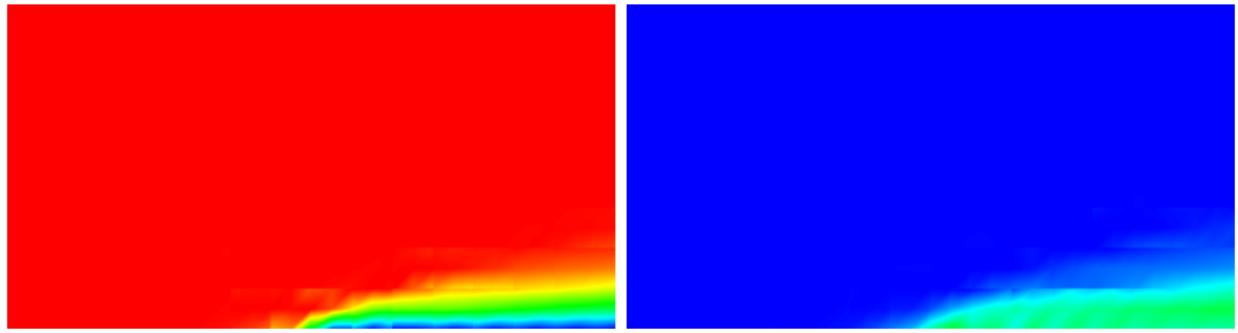
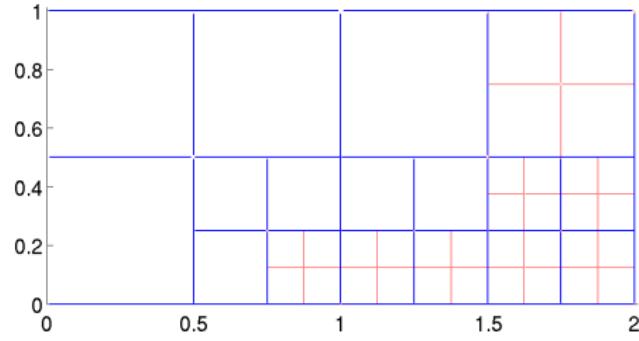
Refinement level 1

(a) u_1 (b) T 

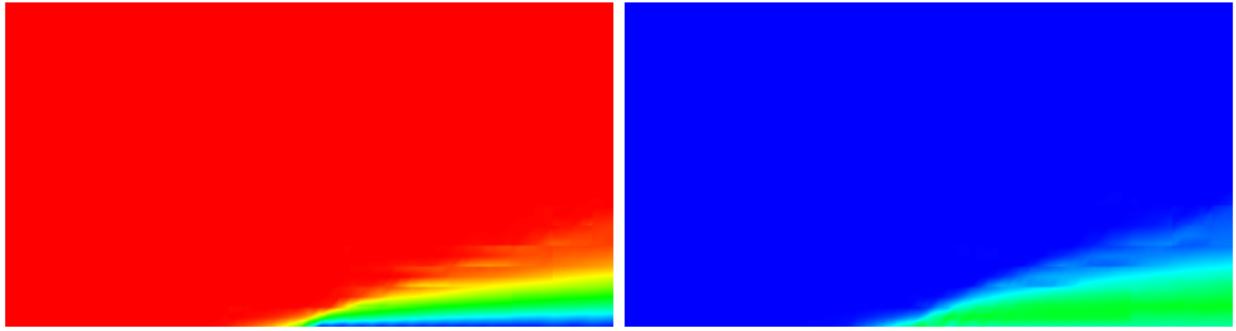
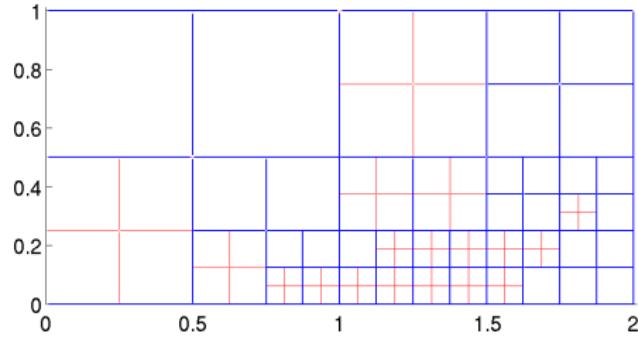
Refinement level 2

(a) u_1 (b) T 

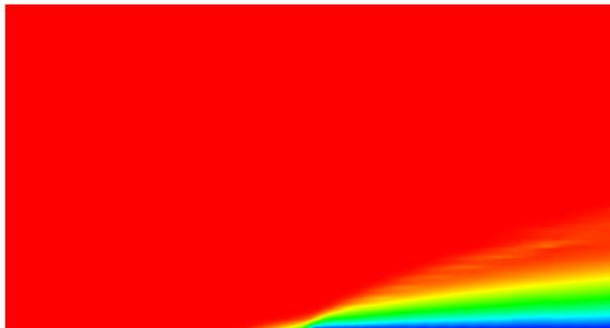
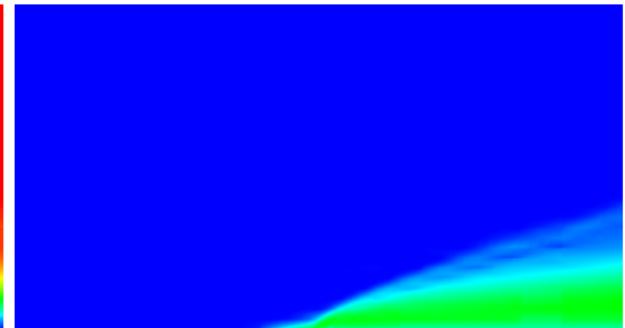
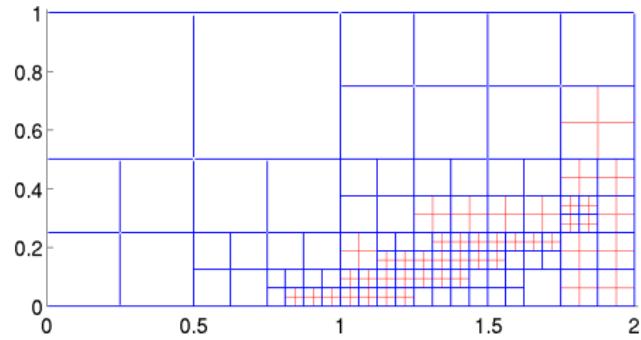
Refinement level 3

(a) u_1 (b) T 

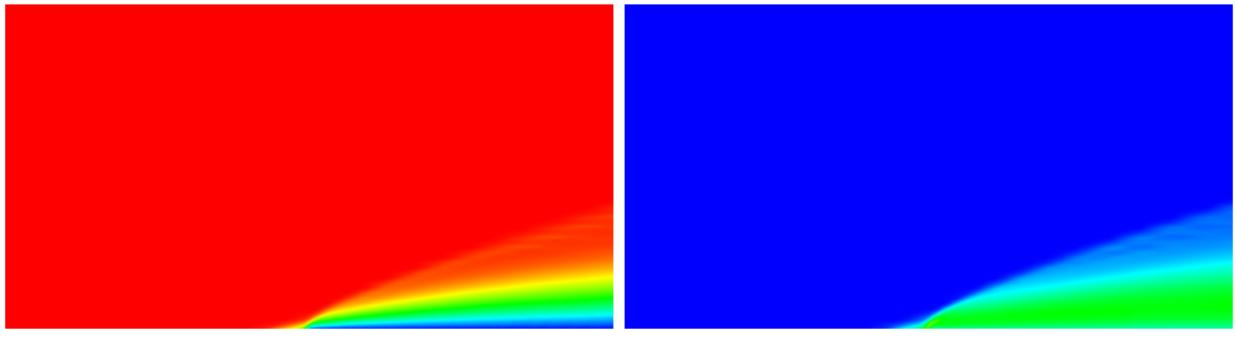
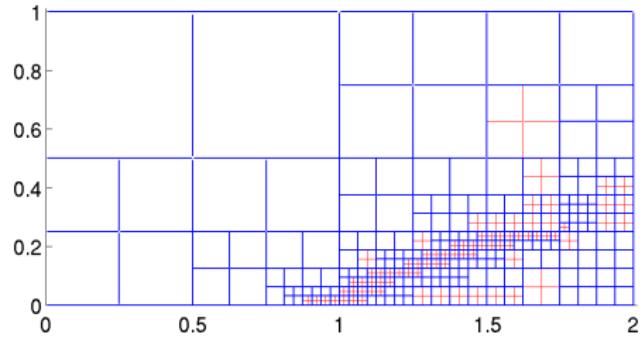
Refinement level 4

(a) u_1 (b) T 

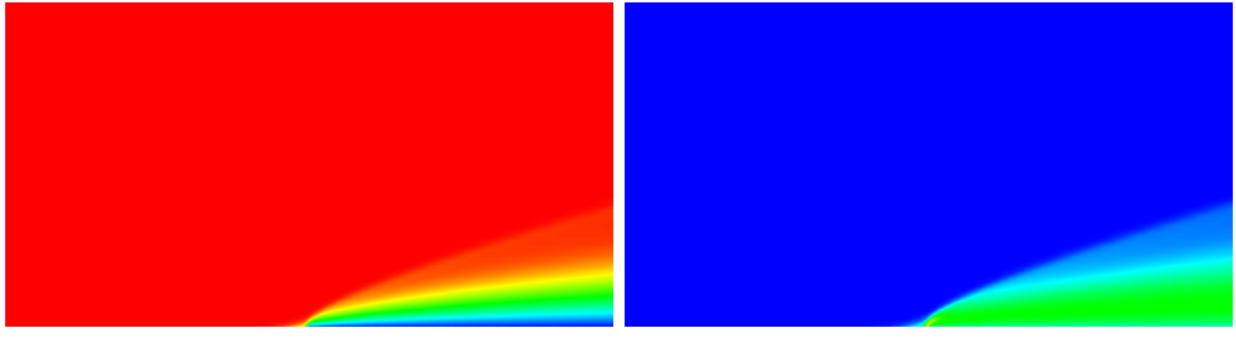
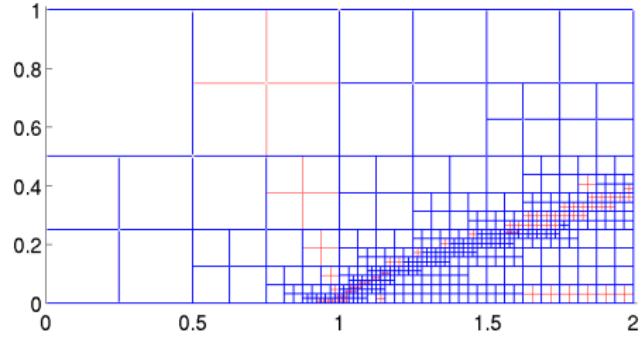
Refinement level 5

(a) u_1 (b) T 

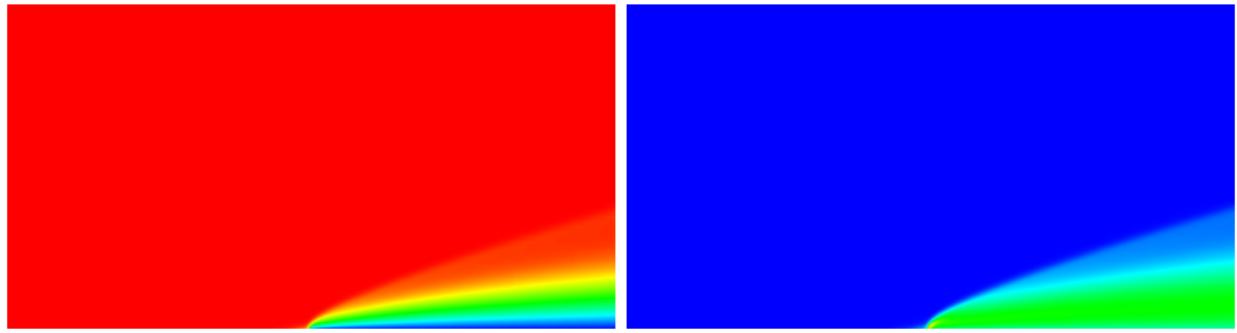
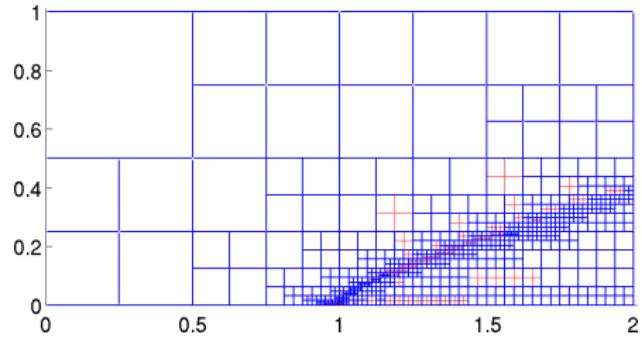
Refinement level 6

(a) u_1 (b) T 

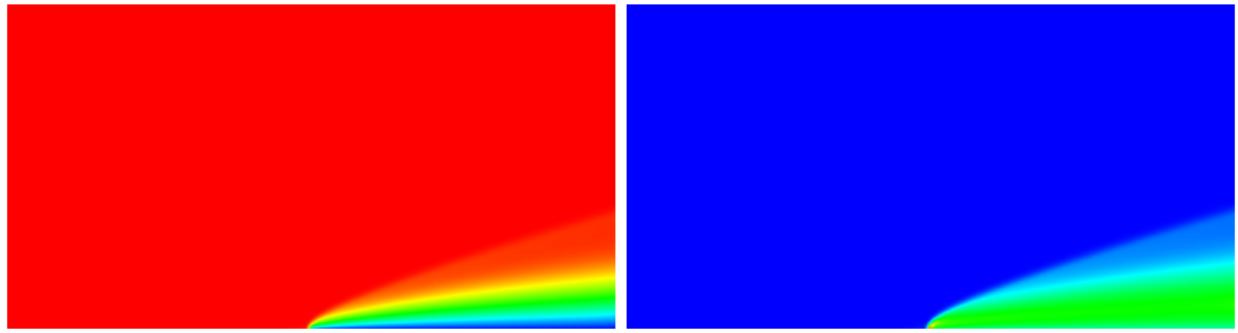
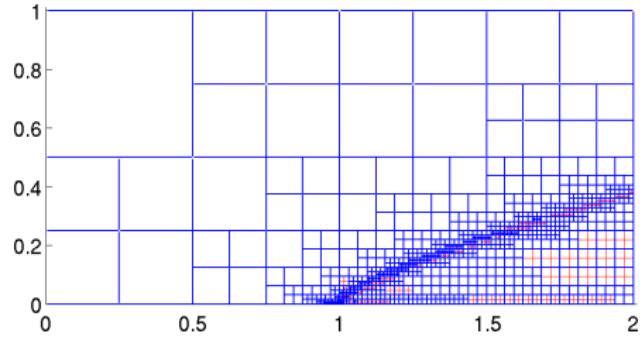
Refinement level 7

(a) u_1 (b) T 

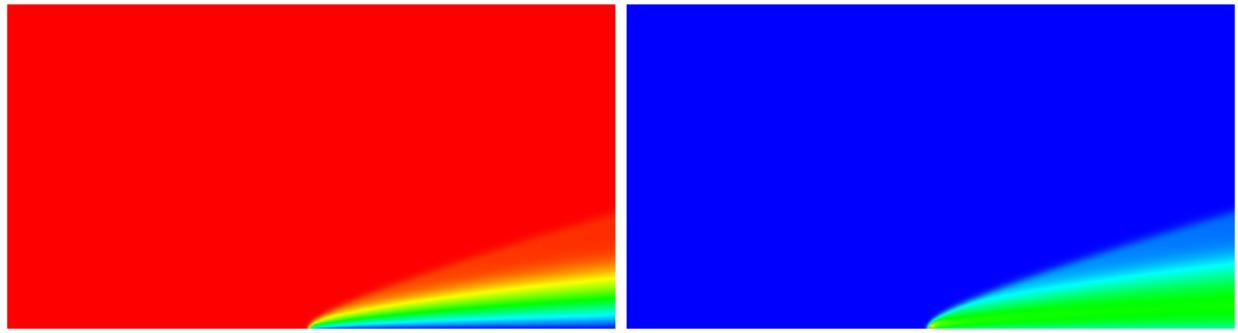
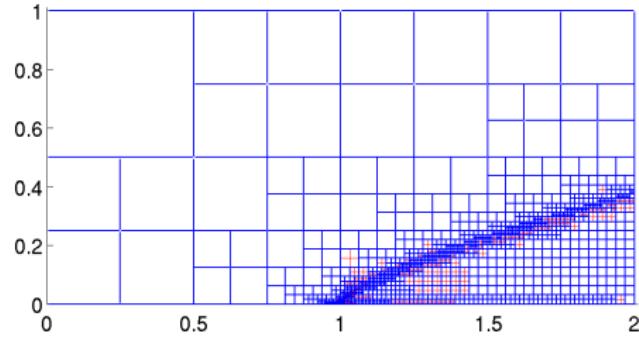
Refinement level 8

(a) u_1 (b) T 

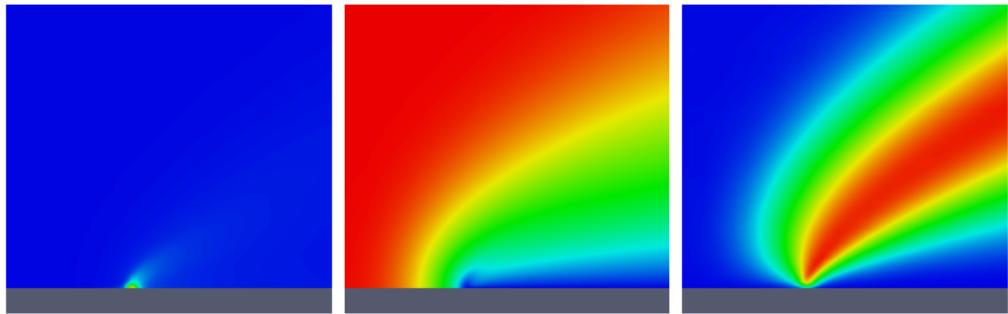
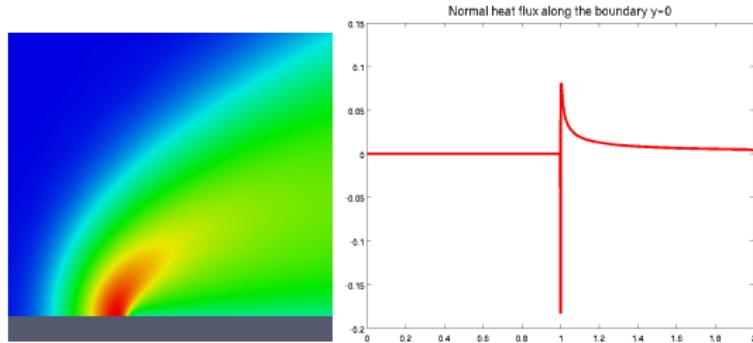
Refinement level 9

(a) u_1 (b) T 

Refinement level 10

(a) u_1 (b) T 

Zoomed solutions at plate/stagnation point

(a) ρ (b) u_1 (c) u_2 (d) T (e) q_n

Proposed work: Area A

- **Completed: Prove robustness of DPG method for the scalar convection-diffusion problem.**

We have introduced a test norm under which the DPG method robustly bounds the L^2 error in the field variables u and the scaled stress σ . Numerical results confirm the theoretical bounds given.

- **Proposed: Attempt analysis of the linearized Navier-Stokes system.**

We hope to analyze the linearized Navier-Stokes equations to determine an optimal extrapolation of the test norm for the scalar convection-diffusion problem to systems.

Proposed work: Area B

- **Completed: Collaborative work with Nathan Roberts on the higher order parallel adaptive DPG code Camellia.**

Numerical experiments use the higher-order adaptive codebase Camellia, built by Nathan Roberts upon the Trilinos library.

- **Proposed: Anisotropic refinements and hp -schemes.**

The error representation function drives refinement effectively, and we hope to generalize its use to anisotropic adaptive schemes.

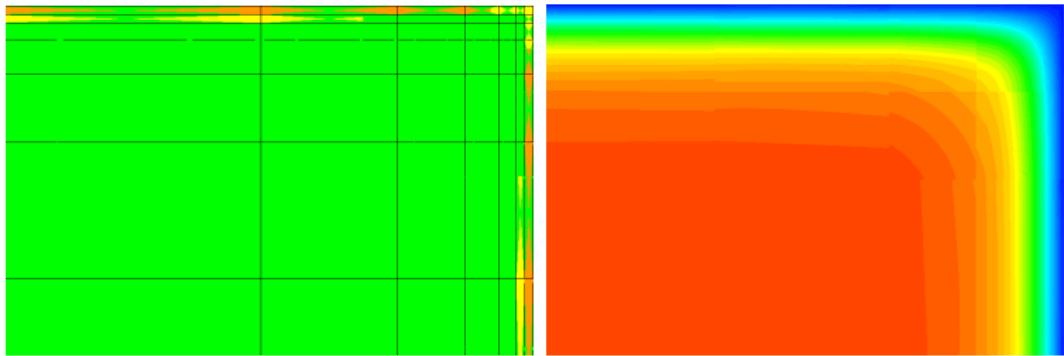


Figure: Anisotropic mesh for a convection-diffusion boundary layer.

Proposed: Distributed iterative static condensation.

$$Ku = \begin{bmatrix} A & B \\ B^T & D \end{bmatrix} \begin{bmatrix} u_{\text{flux}} \\ u_{\text{field}} \end{bmatrix} = \begin{bmatrix} f \\ g \end{bmatrix} = I$$

where D has a block-diagonal structure. The system can be reduced to yield the condensed system

$$(A - BD^{-1}B^T) u_{\text{flux}} = f - BD^{-1}g$$

where D^{-1} can be inverted block-wise. For FE stiffness matrices under the Laplace equation, the Schur complement has reduced condition number of $O(h^{-1})$ as opposed to $O(h^{-2})$.¹¹

¹¹L. Mansfield. On the Conjugate Gradient Solution of the Schur Complement System Obtained from Domain Decomposition. *SIAM Journal on Numerical Analysis*, 27(6):pp. 1612–1620, 1990

Proposed: a Nonlinear Hessian-based DPG method. Given a nonlinear variational problem $b(u, v) = \ell(v)$, linear in v but not in u , beginning with the *nonlinear* dual residual

$$J(u_h) = \frac{1}{2} \|B(u_h) - \ell\|_{V'}^2 := \frac{1}{2} \sup_{v \in V \setminus \{0\}} \frac{|b(u_h, v) - \ell(v)|^2}{\|v\|_V^2}.$$

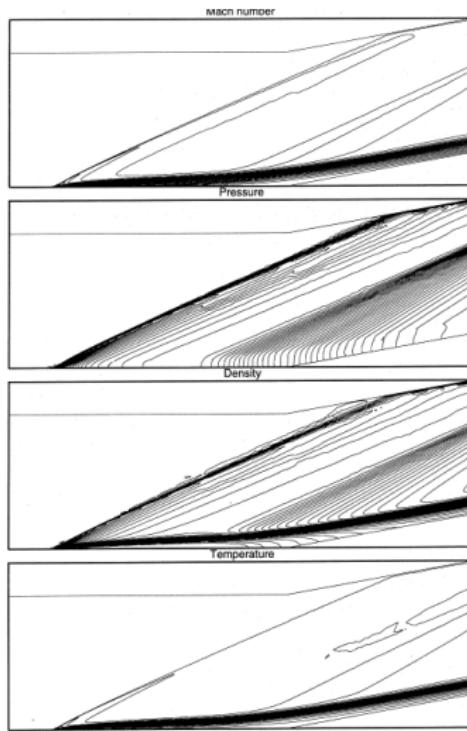
produces a Hessian-based DPG method, which solves

$$b_u(\Delta u, v) + b''(\Delta u, \delta u, v_{R(u)}) = \ell(v) - b(u, v) = r(u, v),$$

and aims to minimize the nonlinear dual residual instead of the linearized problem residual.

Proposed work: Area C

- Completed:
convection-dominated diffusion, Burgers, and a model problem for Navier-Stokes.
- Proposed: Range of Reynolds/Mach numbers.
- Proposed: Flow over a bump, Holden's ramp problem (shock-boundary layer interaction).
- Proposed (time permitting): airfoils, regularized Euler.



Thank you!

Questions?

Extended literature review

- Finite difference/volume methods, upwind stencils, monotonicity, ENO/WENO.
- Artificial diffusion, flux limiters/discontinuity-capturing operators.
- DG fluxes can be interpreted as residual-based stabilizations similar to SUPG.¹²
- Hybridized DG (HDG) methods introduce independent unknowns for numerical traces. Fluxes are defined in terms of DG fluxes and the numerical trace.¹³

¹²F. Brezzi, B. Cockburn, L.D. Marini, and E. Süli. Stabilization mechanisms in discontinuous Galerkin finite element methods. *Computer Methods in Applied Mechanics and Engineering*, 195(25–28):3293 – 3310, 2006

¹³B. Cockburn, J. Gopalakrishnan, and R. Lazarov. Unified hybridization of discontinuous Galerkin, mixed, and continuous Galerkin methods for second order elliptic problems. *SIAM J. Numer. Anal.*, 47(2):1319–1365, February 2009

Specifics of the ultra-weak formulation

Given a first order system $Au = f$, multiply by test function v and integrate

$$(Au, v) = \langle \gamma(Au), v \rangle + (u, A_h^* v) = (f, v)$$

If $u \in L^2(\Omega)$, the trace of u is undefined, so we identify boundary terms $\langle \gamma(Au), v \rangle_{\Gamma_h} = \langle \hat{u}, v \rangle_{\Gamma_h}$ as unknowns \hat{u} on Γ_h . For convection-diffusion,

$$\begin{aligned} b\left(\left(u, \sigma, \hat{u}, \hat{f}_n\right), (v, \tau)\right) &= (u, \nabla_h \cdot \tau - \beta \cdot \nabla_h v)_{\Omega_h} + (\sigma, \epsilon^{-1} \tau + \nabla_h v)_{\Omega_h} \\ &\quad - \langle [\![\tau \cdot n]\!], \hat{u} \rangle_{\Gamma_h} + \left\langle \hat{f}_n, [\![v]\!] \right\rangle_{\Gamma_h}, \end{aligned}$$

where

$$\hat{f}_n := \beta_n u - \sigma_n \in H^{-1/2}(\Gamma_h), \quad \hat{u} \in H^{1/2}(\Gamma_h)$$

with the minimum energy extension norm on Γ_h . We note that $H^{\pm 1/2}(\Gamma_h)$ are *closed* subspaces of $\prod_K H^{\pm 1/2}(\partial K)$.

Inviscid equations

Issue: consider pure convection, $\nabla \cdot \beta u = f$. The ultra-weak variational formulation is

$$\left\langle \hat{f}_n, v \right\rangle - (u, \beta \cdot \nabla v) = (f, v),$$

where $\hat{f}_n := \beta_n u$. When $\beta_n = 0$, v has only a streamline derivative, and \hat{f}_n becomes an ill-defined trace in the cross-stream direction. For hyperbolic systems, this issue manifests as *sonic lines*.

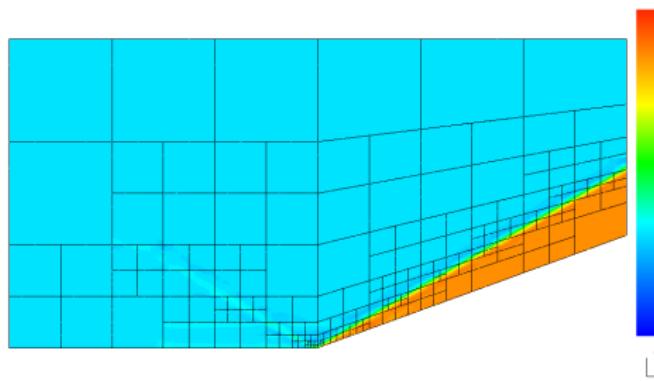


Figure: Sonic lines ($u_n - c = 0$) appear for linearized Euler.



F. Brezzi, B. Cockburn, L.D. Marini, and E. Süli.

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N. Roberts, D. Ridzal, P. Bochev, and L. Demkowicz.

A Toolbox for a Class of Discontinuous Petrov-Galerkin Methods Using Trilinos.
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