Week 5 Thursday L-10

Mean intensity
$$J_{\lambda}=rac{1}{4\pi}\int I_{\lambda}d\Omega$$

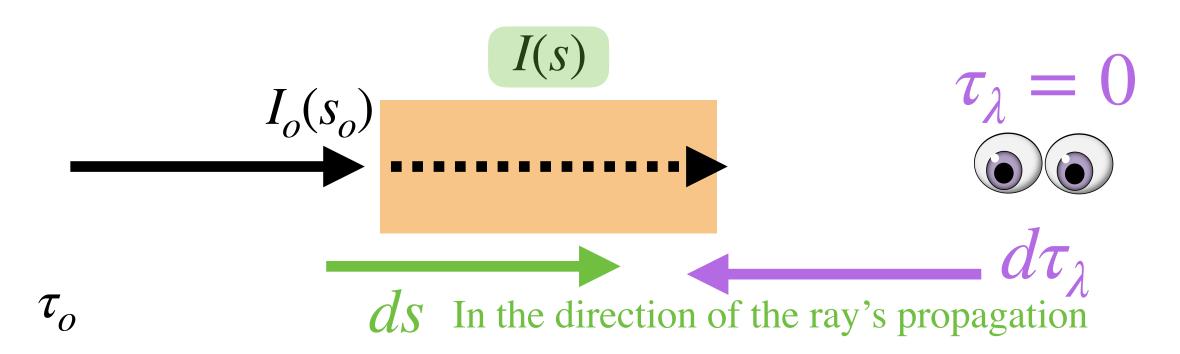
$$F_{\lambda} = \int I_{\lambda} \cos \theta d\Omega$$

Spherical, with azimuthal symmetry $(u = \cos \theta)$

$$J_{\lambda} = \frac{1}{2} \int_{-1}^{+1} I_{\lambda}(u) \ du$$

$$\frac{F_{\lambda}}{4\pi} = \frac{1}{2} \int_{-1}^{+1} I_{\lambda}(u) \ u \ du$$

Change in intensity



Optical depth

$$d\tau_{\lambda} = -\kappa_{\lambda} \rho ds$$

$$\tau_{\lambda}(s) - 0 = -\int_{S_{\text{obs}}}^{S} \kappa_{\lambda}(s) \ \rho(s) \ ds$$

Absorption

$$dI_{\lambda} = -\kappa_{\lambda}(s)\rho(s)dsI_{\lambda}$$

General Solution

$$\tau(s) = \int_{s'=s}^{s'=\infty} \kappa(s')\rho(s')ds' \qquad \tau(s) = \kappa\rho((s_o+d)-s)$$

$$dI_{\lambda} = -\kappa_{\lambda}(s)\rho(s)dsI_{\lambda} \quad I(s) = I_{o}e^{-\int_{s'=s_{o}}^{s'=s} \kappa(s')\rho(s')ds}$$

$$I(s) = I_o e^{\tau(s) - \tau(s = s_o)}$$

Constant properties in material

$$\tau(s) = \kappa \rho((s_o + d) - s)$$

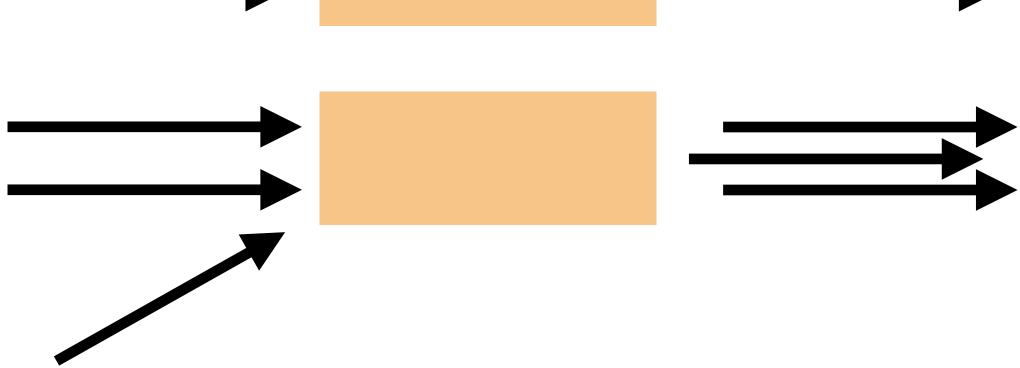
$$I(s) = I_o e^{-\kappa \rho (s - s_o)}$$

Emission: how can we gain intensity?

1. Pure emission: create a photon



2. Scattering: deflect a photon



(Q: what are the units of κ_{λ} ?)

Absorption:
$$dI_{\lambda} = - \rho ds \kappa_{\lambda} I_{\lambda}$$

(Q: what is the meaning of ρ ds κ_{λ} ?)

Emission:

Emission: how can we gain intensity?

1. Pure emission: create a photon



2. Scattering: deflect a photon

Absorption:

$$\boldsymbol{\lambda}$$

(Q: why no I_{λ} ?)

Emission:

$$II_{\lambda}$$
 =

$$\rho$$

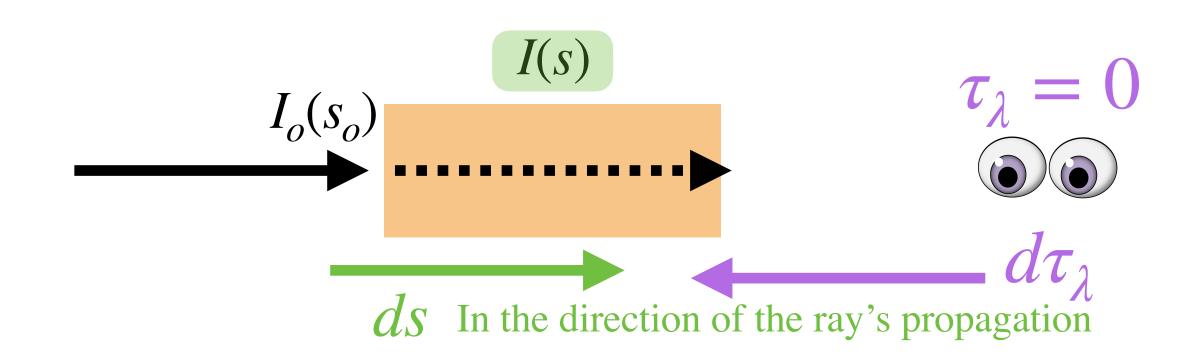
(Q: what are the units of j_{λ} ?)

Let's solve this equation, if we know $I(s_o) = I_o$

$$dI_{\lambda} = +j_{\lambda}(s) \rho(s) ds$$

Together on the board...

$$I(s) = I_o + \int_{s_o}^{s} j(s)\rho(s)ds$$



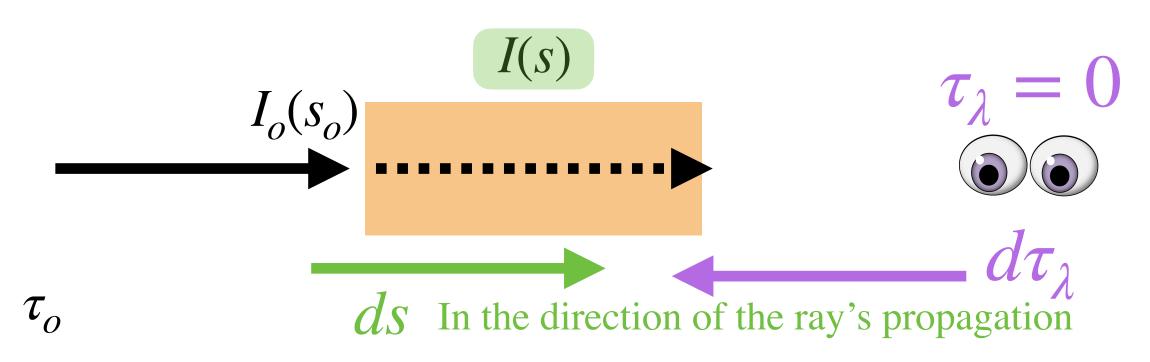
If the material has constant properties:

$$I(s) = I_o + j\rho[s - s_o]$$

In notebook:

- graph I(s) for constant values
- graph I(s) for linearly decreasing density

Change in intensity



Optical depth

$$d\tau_{\lambda} = -\kappa_{\lambda} \rho ds$$

$$\tau_{\lambda}(s) - 0 = -\int_{S_{\text{obs}}}^{s} \kappa_{\lambda}(s) \ \rho(s) \ ds$$

Absorption

$$dI_{\lambda} = -\kappa_{\lambda}(s)\rho(s)dsI_{\lambda}$$

Emission

$$dI_{\lambda} = + j_{\lambda}(s)\rho(s)ds$$

General Solution

$$\tau(s) = \int_{s'=s}^{s'=\infty} \kappa(s')\rho(s')ds' \qquad \tau(s) = \kappa\rho((s_o + d) - s)$$

$$dI_{\lambda} = -\kappa_{\lambda}(s)\rho(s)dsI_{\lambda} \quad I(s) = I_{o}e^{-\int_{s'=s_{o}}^{s'=s} \kappa(s')\rho(s')ds}$$

$$I(s) = I_o e^{\tau(s) - \tau(s = s_o)}$$

$$dI_{\lambda} = + j_{\lambda}(s)\rho(s)ds$$
 $I(s) = I_o + \int_{s'=s_o}^{s'=s} j(s')\rho(s')ds'$ $I(s) = I_o + j\rho(s - s_o)$

Constant properties in material

$$\tau(s) = \kappa \rho((s_o + d) - s)$$

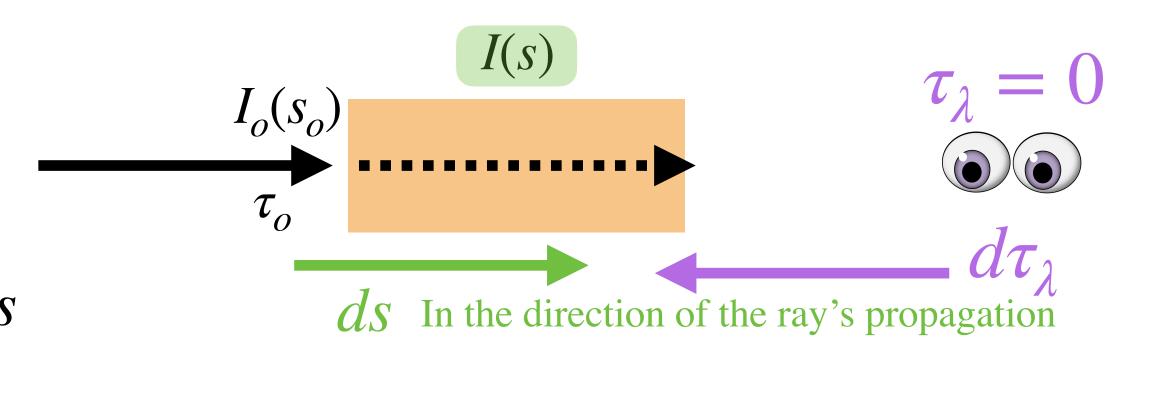
$$I(s) = I_o e^{-\kappa \rho (s - s_o)}$$

$$I(s) = I_o + j\rho(s - s_o)$$

Now, what if we have both absorption and emission: The "formal solution" of RT

(Q: why can't we integrate this directly like earlier?)

$$\frac{dI_{\lambda}}{-\kappa_{\lambda}(s) \rho(s) ds} = -\kappa_{\lambda}(s) \rho(s) ds I_{\lambda}(s) + j_{\lambda}(s) \rho(s) ds -\kappa_{\lambda}(s) \rho(s) ds -\kappa_{\lambda}(s) \rho(s) ds$$



$$\frac{dI_{\lambda}}{d\tau_{\lambda}} = I_{\lambda}(s) - \frac{j_{\lambda}(s)}{\kappa_{\lambda}(s)}$$

The "source function" $S_{\lambda}(s)$ (Q: what is the meaning of $S(\lambda)$?)

$$\frac{dI_{\lambda}}{d\tau_{\lambda}} = I_{\lambda}(\tau_{\lambda}) - S_{\lambda}(\tau_{\lambda})$$

Here, I can use the optical depth as my "coordinate system" instead of the physical distance

Now, what if we have both absorption and emission: The "formal solution" of RT

$$\frac{dI_{\lambda}}{d\tau_{\lambda}} = I_{\lambda}(\tau_{\lambda}) - S_{\lambda}(\tau_{\lambda})$$

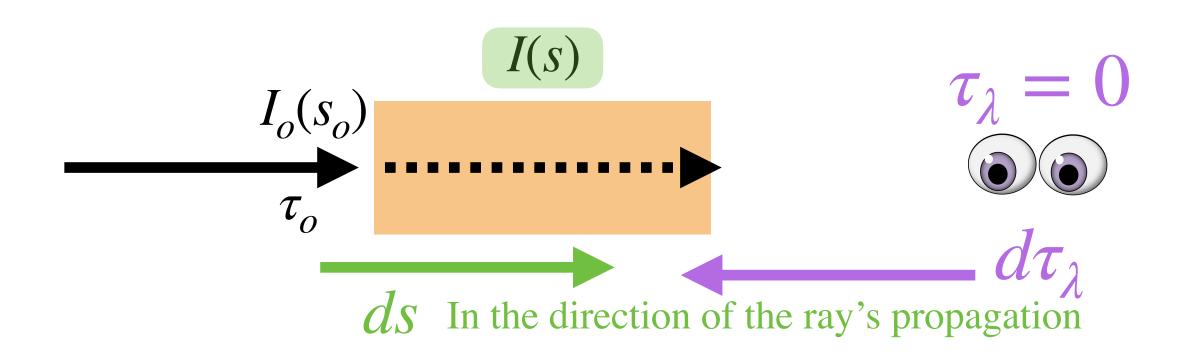
(Just dropping the λ s for more clarity)

$$\frac{dI}{d\tau} e^{-\tau} - I(\tau)e^{-\tau} = -S(\tau) e^{-\tau}$$

$$\frac{d}{d\tau} \left[I(\tau) \quad e^{-\tau} \right] = - S(\tau) e^{-\tau}$$

$$\int_{\text{known position }(I_{o}, \tau_{o})}^{\text{unknown position }(I, \tau)} = -\int_{\tau_{o}}^{\tau} S(\tau) e^{-\tau} d\tau$$

$$I(\tau)e^{-\tau} - I_o e^{-\tau_o} = \int_{\tau'=\tau_o}^{\tau'=\tau} S(\tau') e^{-\tau'} d\tau'$$



Here, I added explicit 'to the integrants, to distinguish it from the τ in the upper bound.

(This is the same as our integrals with a bound of *r* when we integrated the continuity and hydrostatic equations, but here it is a bit harder to visualize)

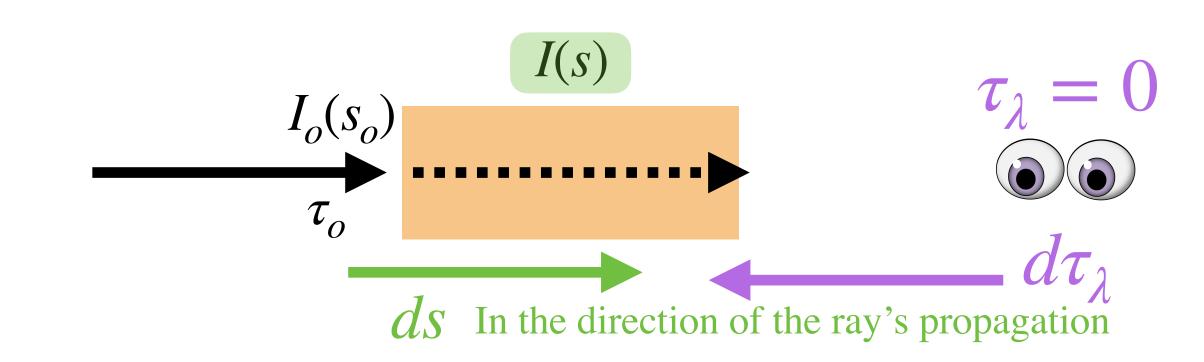
Now, what if we have both absorption and emission: The "formal solution" of RT

$$I(\tau)e^{-\tau} - I_o e^{-\tau_o} = \int_{\tau'=\tau_o}^{\tau'=\tau} S(\tau') e^{-\tau'} d\tau'$$

$$I(\tau)e^{-\tau} = I_o e^{-\tau_o} + \int_{\tau'=\tau_o}^{\tau'=\tau} S(\tau') e^{-\tau'} d\tau'$$

$$e^{-\tau}$$

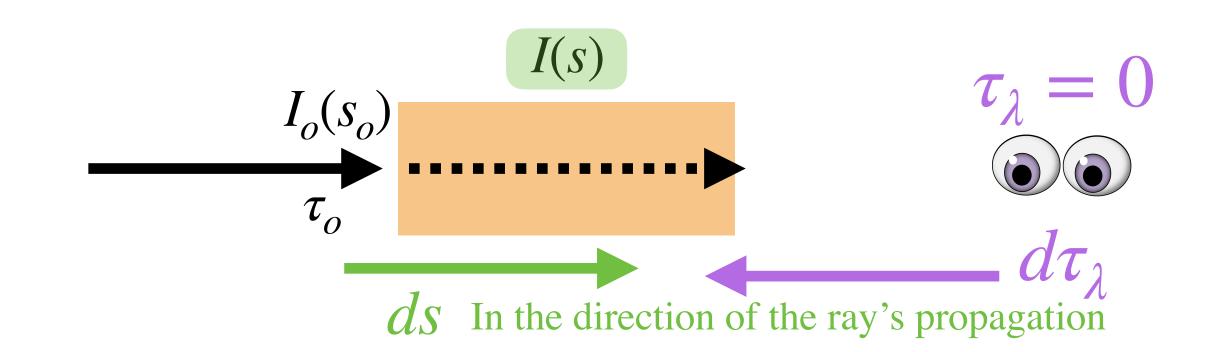
$$I(\tau) = I_o e^{\tau - \tau_o} + \int_{\tau' = \tau_o}^{\tau' = \tau} S(\tau') e^{\tau - \tau'} d\tau'$$
 That is the "formal solution"



Now, what if we have both absorption and emission: The "formal solution" of RT

$$I(\tau) = I_o e^{\tau - \tau_o} + \int_{\tau' = \tau_o}^{\tau' = \tau} S(\tau') e^{\tau - \tau'} d\tau'$$

$$I(\tau = 1) = I_o e^{1-\frac{5}{5}} + \int_{\tau'=5}^{\tau'=1} S(\tau') e^{1-\tau'} d\tau'$$







Replacing the integral by a sum, for illustration:

$$I(\tau = 1) = I_o e^{1-5} + S(\tau' = 4) e^{1-4} \Delta \tau' + S(\tau' = 3) e^{1-3} \Delta \tau' + \dots$$

1. What if the source function is constant $S(\tau) = S_o$?

$$I(\tau) = I_o e^{\tau - \tau_o} + \int_{\tau' = \tau_o}^{\tau' = \tau} S(\tau') e^{\tau - \tau'} d\tau'$$

- 1. Get the source function out of the integral
- 2. Note that $e^{\tau \tau'} = e^{\tau}e^{-\tau'}$ and e^{τ} can also be taken out of the integral (not a function of τ')

 $I_o(s_o)$

3. We can integrate

$$I(\tau) = I_o e^{\tau - \tau_o} + S_o \left[1 - e^{\tau - \tau_o} \right]$$

And we also would like to be able to plot I(s) in our notebook,

ds In the direction of the ray's propagation

2. What if the other material properties are also constant (κ_0, ρ_0) ? so we need to go back to that coordinate system!

$$I(s) = I_o e^{\tau(s) - \tau_o} + S_o \left[1 - e^{\tau(s) - \tau_o} \right]$$

$$\tau(s) - 0 = -\int_{s_{\text{obs}}}^{s} \kappa(s) \, \rho(s) \, ds = -\kappa_o \rho_o \, [s - s_{\text{obs}}]$$

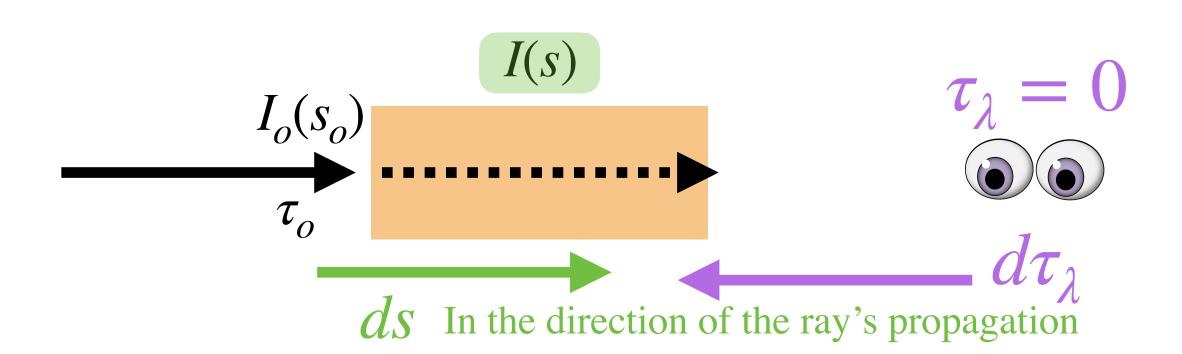
$$\tau(s_o) = -\kappa_o \rho_o \, [s_o - s_{\text{obs}}]$$

$$I(s) = I_o e^{-\kappa_o \rho_o [s-s_o]} + S_o \left[1 - e^{-\kappa_o \rho_o [s-s_o]}\right]$$

$$\tau(s) - \tau_o = -\int_{s_o}^{s} \kappa(s) \ \rho(s) \ ds = -\kappa_o \rho_o[s - s_o]$$

 $\tau(s) - \tau(s_o) = -\kappa_o \rho_o [s - s_o]$

$$\frac{dI_{\lambda}}{d\tau_{\lambda}} = I_{\lambda}(\tau_{\lambda}) - S_{\lambda}(\tau_{\lambda})$$

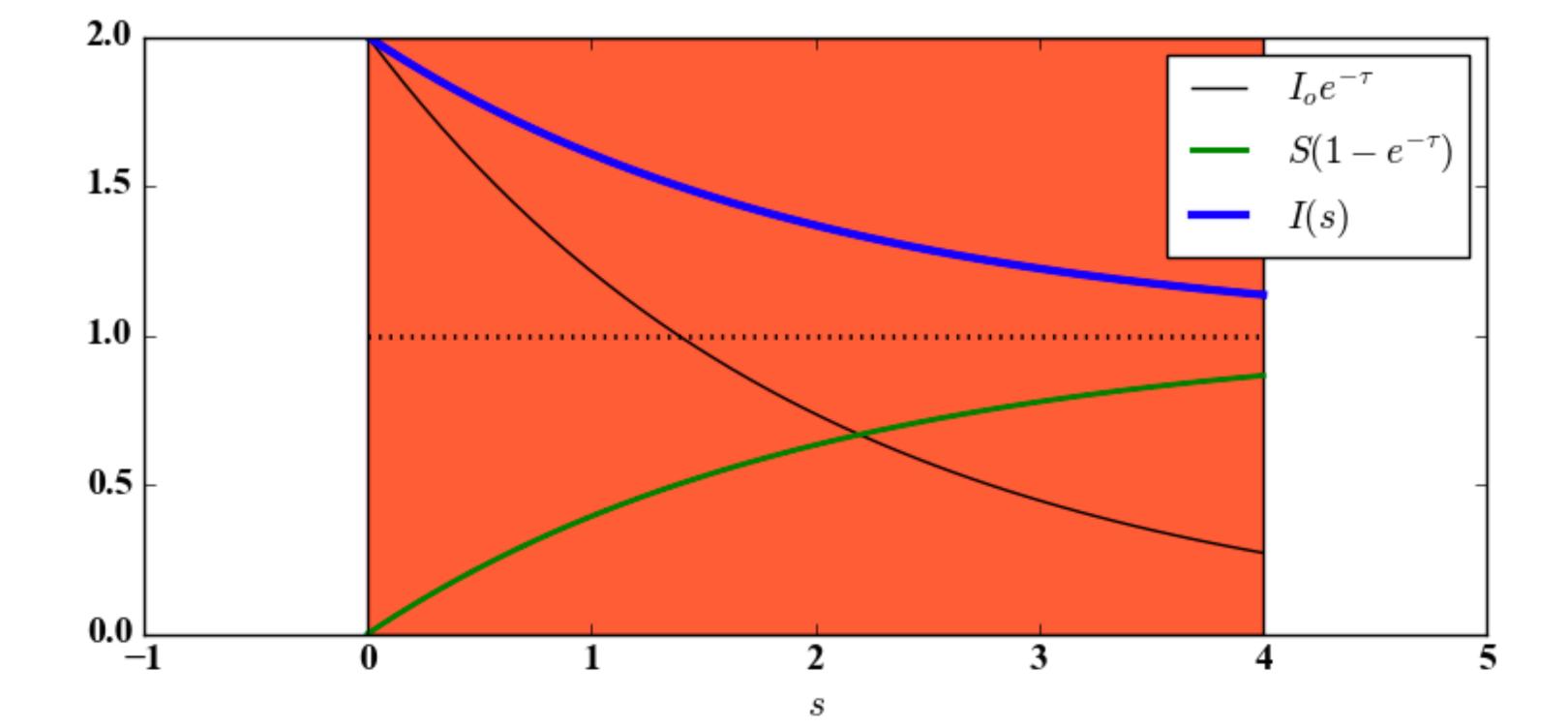


1. What if the source function is constant $S(\tau) = S_o$?

$$I(s) = I_o e^{\tau(s) - \tau_o} + S_o \left[1 - e^{\tau(s) - \tau_o} \right]$$

2. What if the other material properties are also constant (κ_o, ρ_o) ?

$$I(s) = I_o e^{-\kappa_o \rho_o [s-s_o]} + S_o [1 - e^{-\kappa_o \rho_o [s-s_o]}]$$

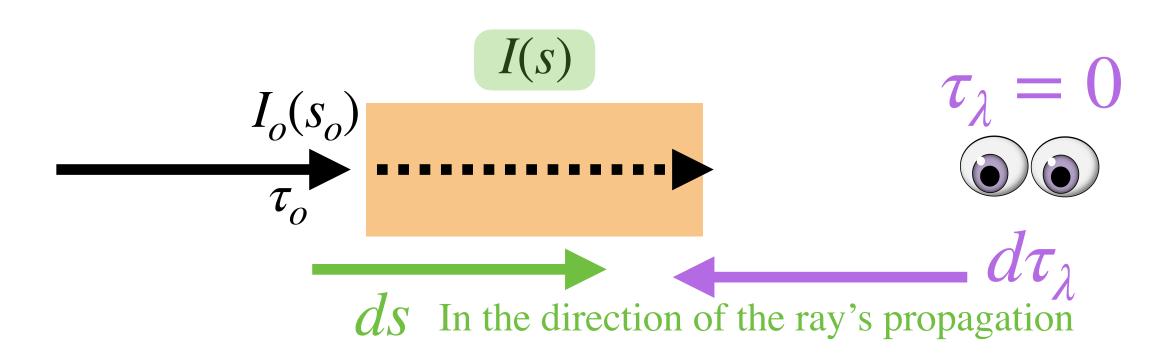


Q: What if $S_o = I_o$?

Q: What would happen to *I*(*s*) if the slab was very long?

-> The source function is what the intensity will become, if given enough optical depth to reach it!

Change in intensity



$$\frac{dI_{\lambda}(\tau_{\lambda})}{d\tau_{\lambda}} = I_{\lambda}(\tau_{\lambda}) - S_{\lambda}(\tau_{\lambda})$$

General Solution

$$I(\tau(s)) = I_o e^{\tau(s) - \tau_o} + \int_{\tau' = \tau(s)}^{\tau' = \tau_o} S(\tau') e^{\tau(s) - \tau'} d\tau'$$

Constant source

function

$$I(\tau(s)) = I_o e^{\tau(s) - \tau_o} + S \left[1 - e^{\tau(s) - \tau_o} \right]$$