

Week 13 Tuesday

L-23

Homology

$$\frac{dM_r(r)}{dr} = 4\pi r^2 \rho(r)$$

$$\frac{dP(r)}{dr} = -\rho(r) \frac{GM_r(r)}{r^2}$$

$$\frac{dT(r)}{dr} = -\frac{3\kappa_R(r)\rho(r)}{64\pi\sigma} \frac{L_r(r)}{r^2 T(r)^3}$$

$$\frac{dL_r(r)}{dr} = 4\pi r^2 \rho(r) \epsilon(r)$$

$M_r(r)$	$P(r)$	$L_r(r)$	$T(r)$
$\rho(r)$	$\mu(r)$	$\epsilon_{\text{nuc}}(r)$	$\kappa_R(r)$

$$P(r) = \frac{\rho(r)}{\mu(r)} kT(r)$$

$$\mu(r) = f(\text{comp}, T(r), P(r))$$

$$\kappa_R(r) = f(\text{comp}, T(r), P(r))$$

$$\epsilon_{\text{nuc}}(r) = f(\text{comp}, T(r), P(r))$$

Other energy transport?

Ideal gas always valid?

Nuclear mechanism?

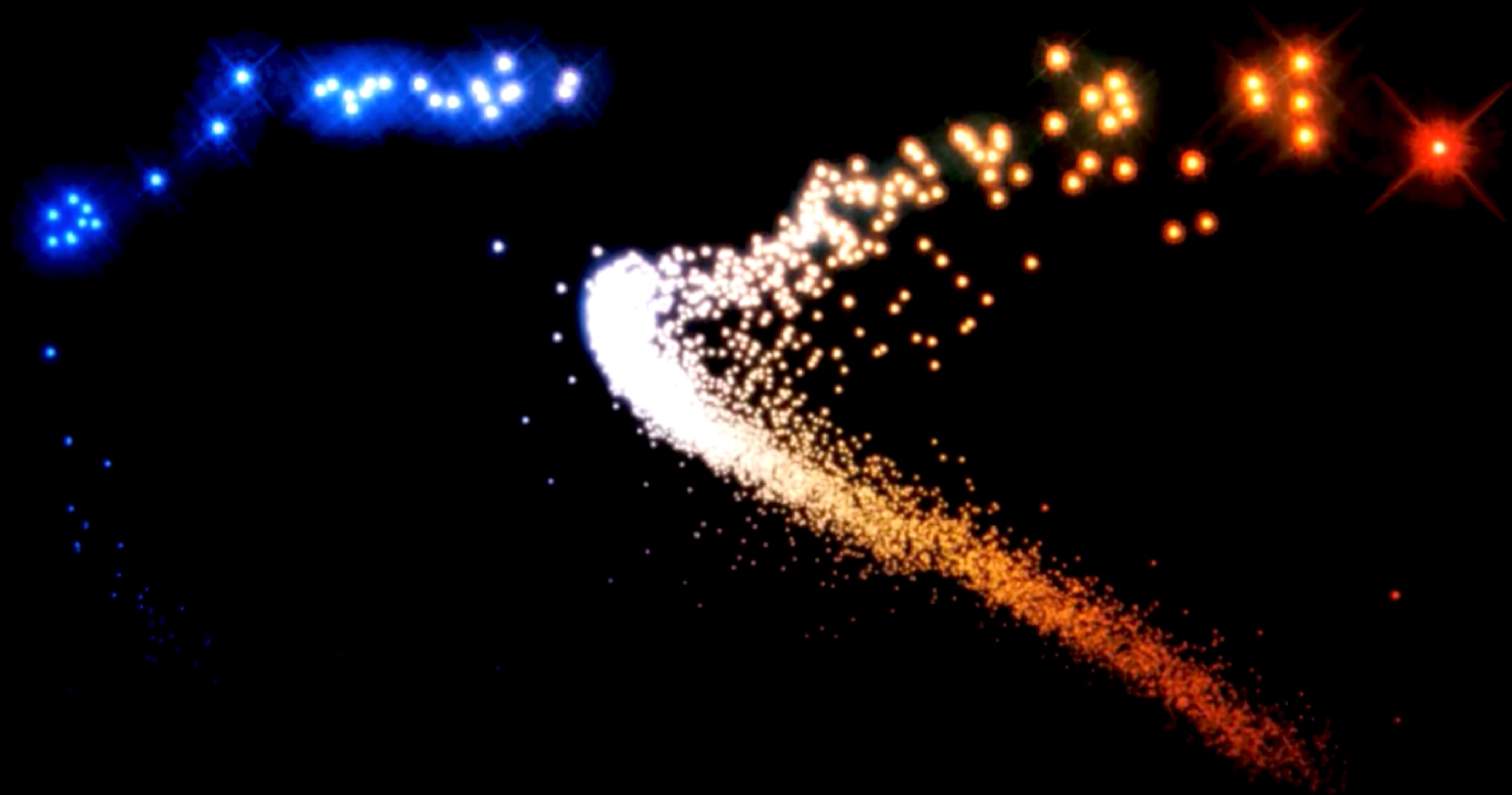
How to solve?

What changes with time?



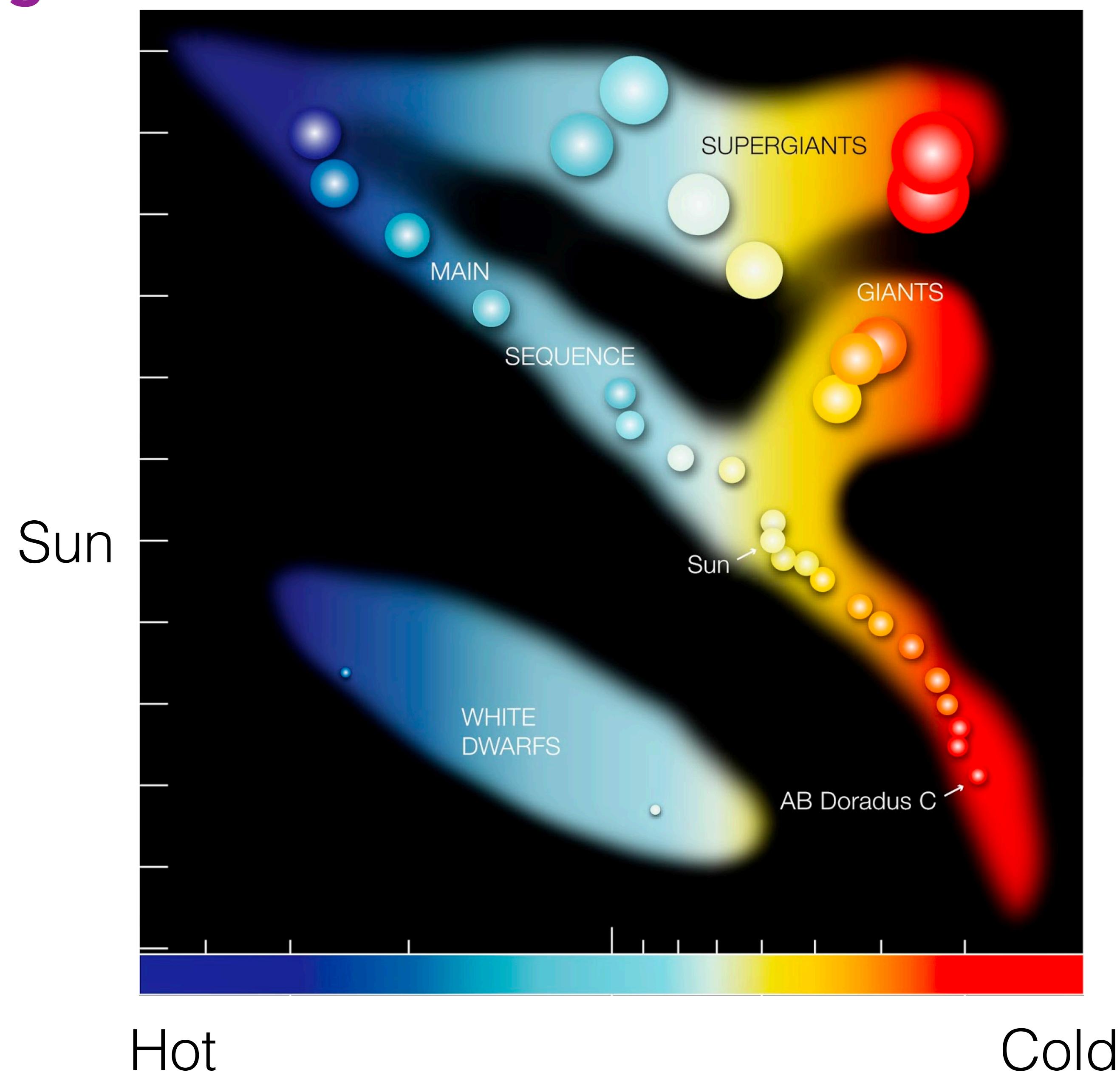
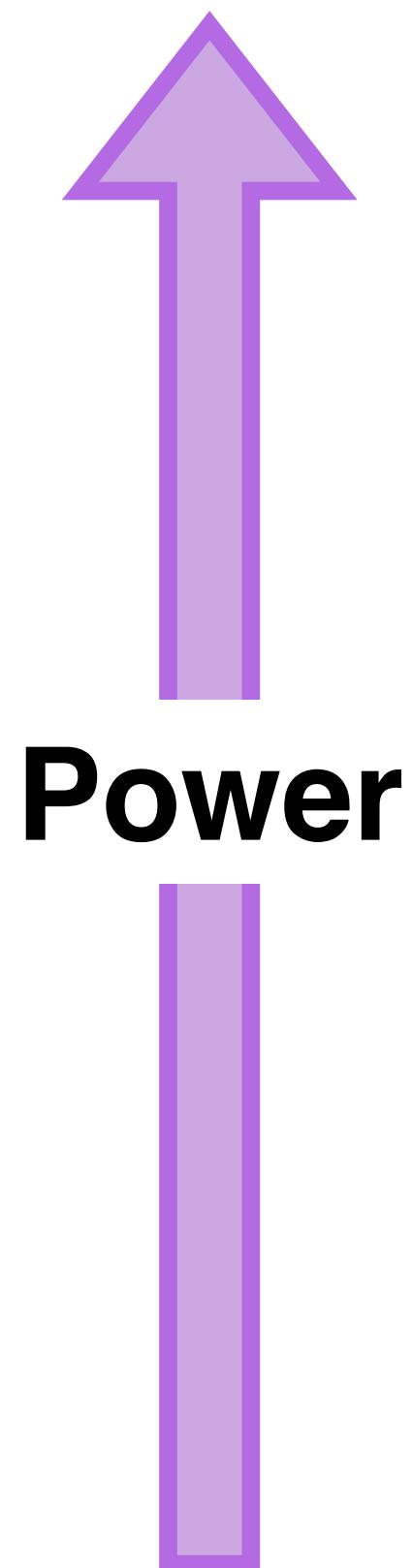
A wide-field photograph of the Omega Centauri Globular Cluster, one of the largest and most massive globular clusters known. The cluster is densely packed with stars of various colors, ranging from blue-white to yellow-orange, set against a dark, star-filled background. The central region is brighter and more concentrated, while the edges appear more diffuse. Several prominent foreground stars are visible as bright points with diffraction spikes.

Omega Centauri Globular Cluster

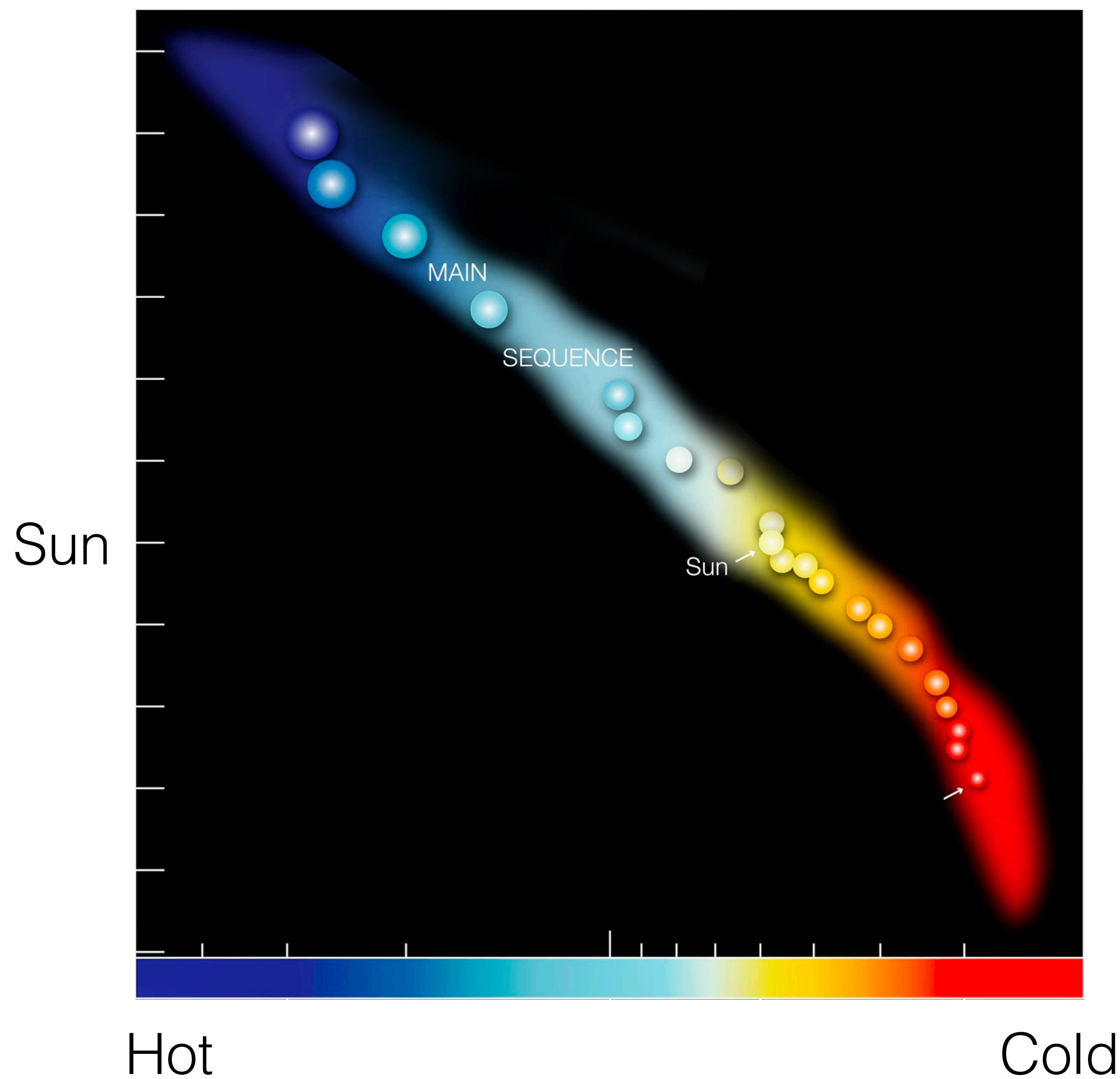
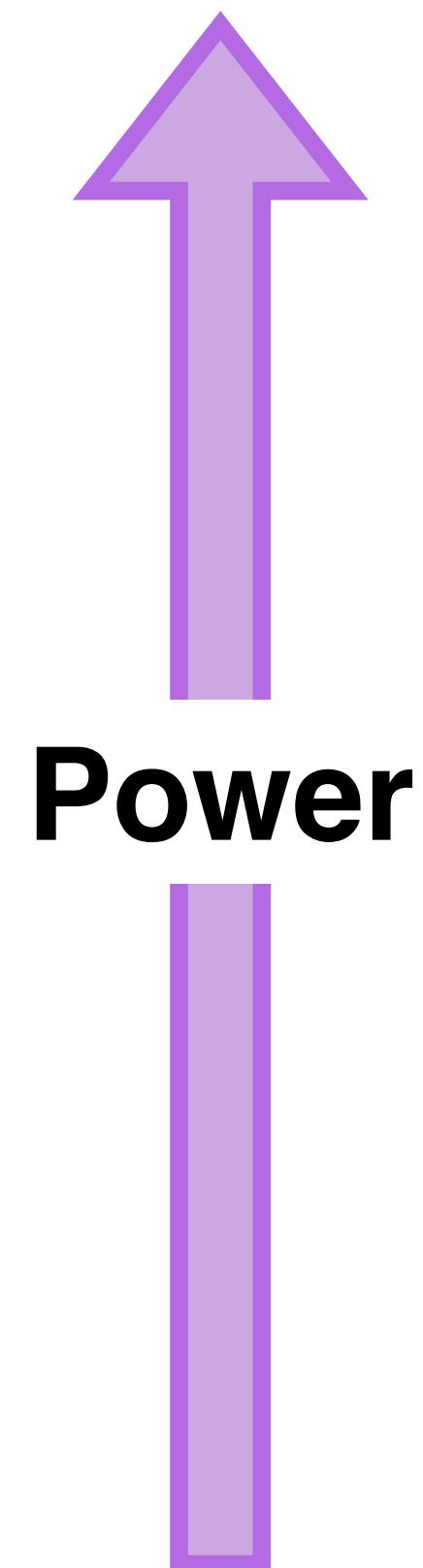


NASA, ESA, J. Anderson and R. van der Marel (STScI)

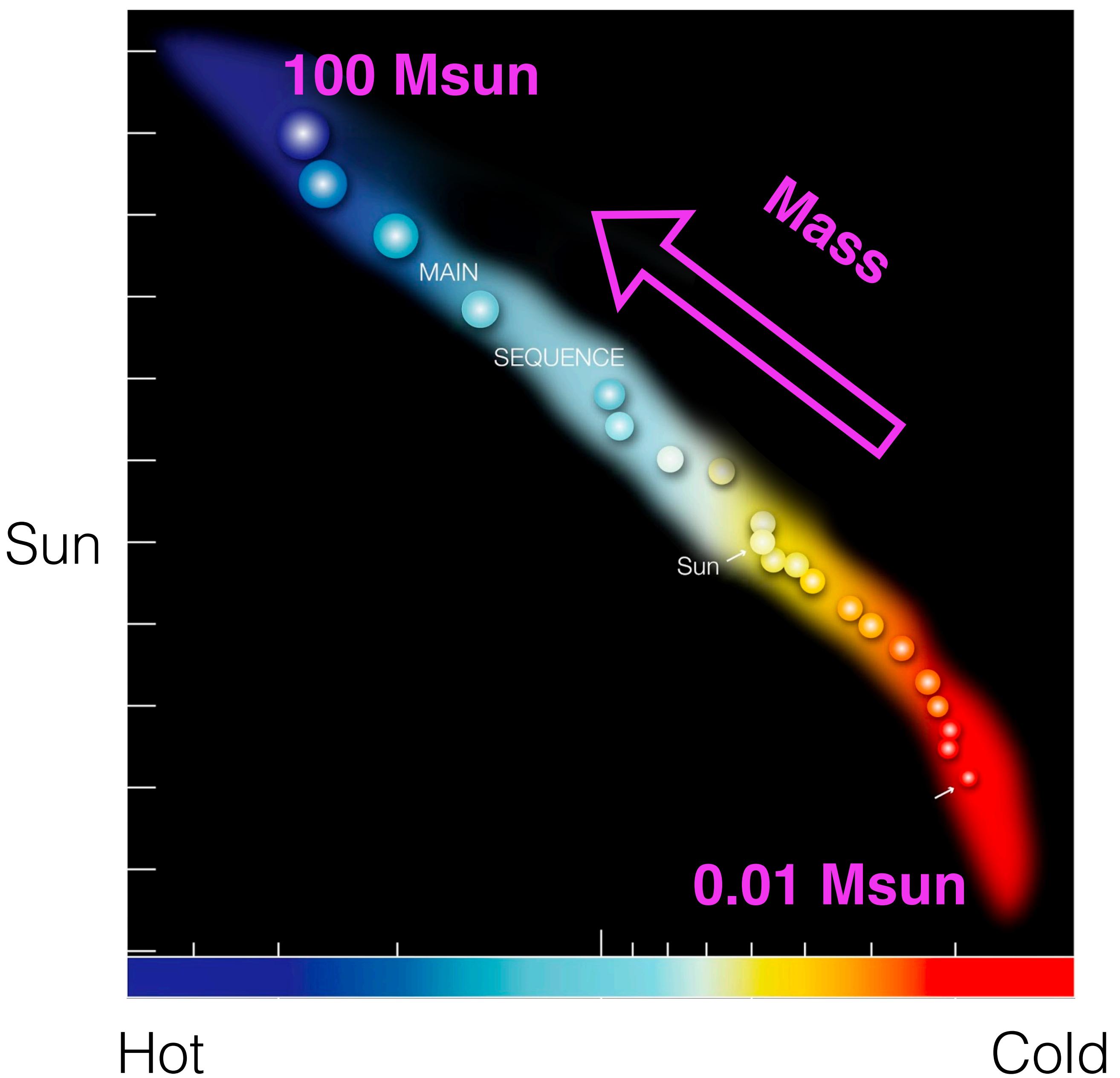
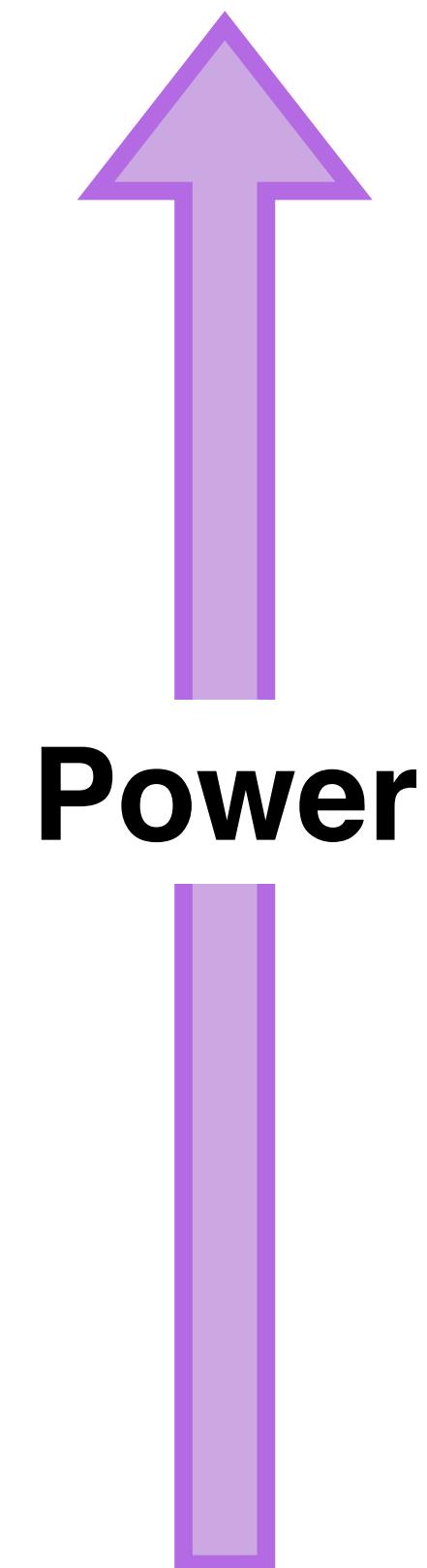
The Hertzsprung-Russell (HR) Diagram is a family portrait of stars

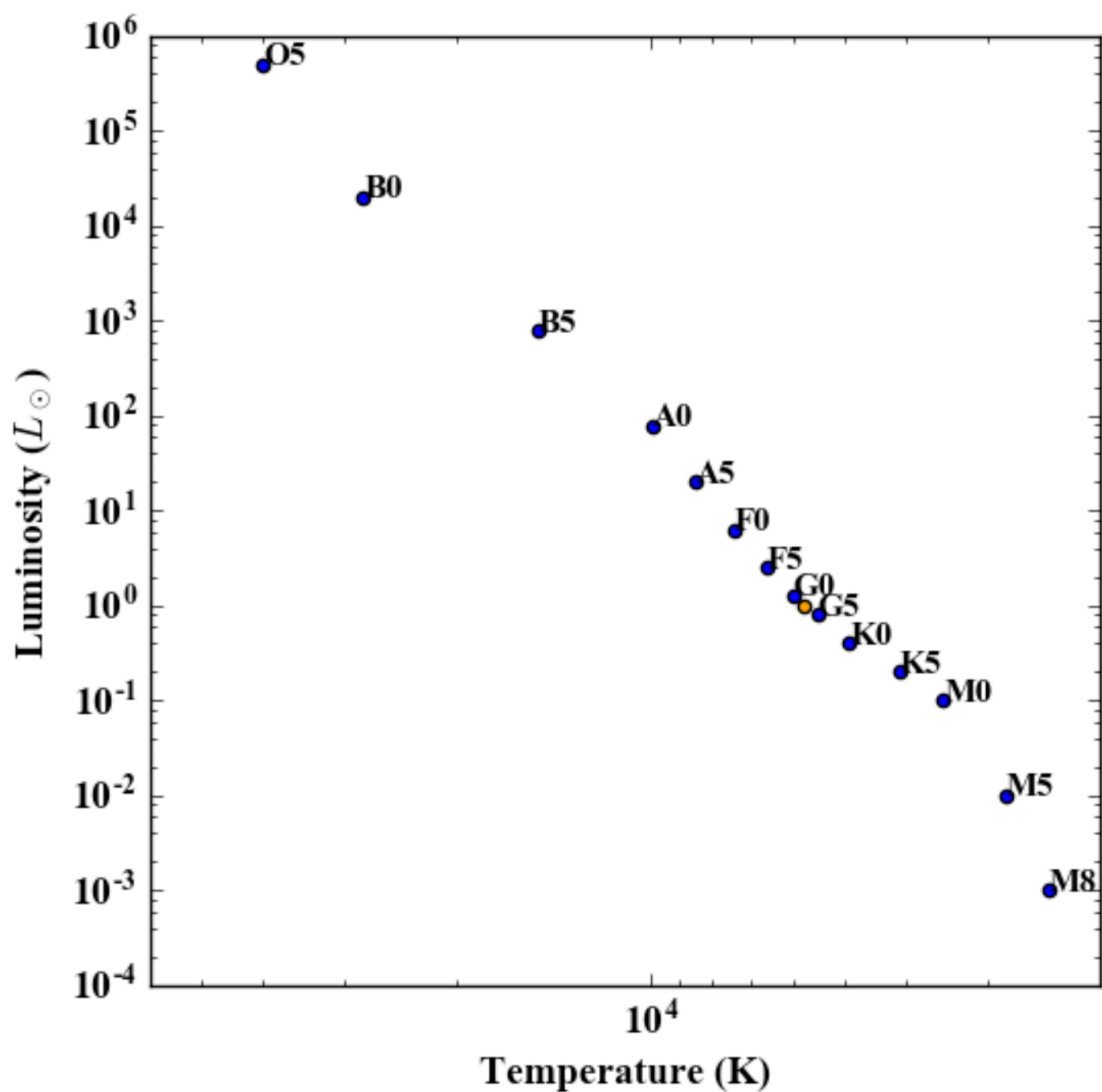


**Most stars in our galaxy are
on the “Main Sequence”**

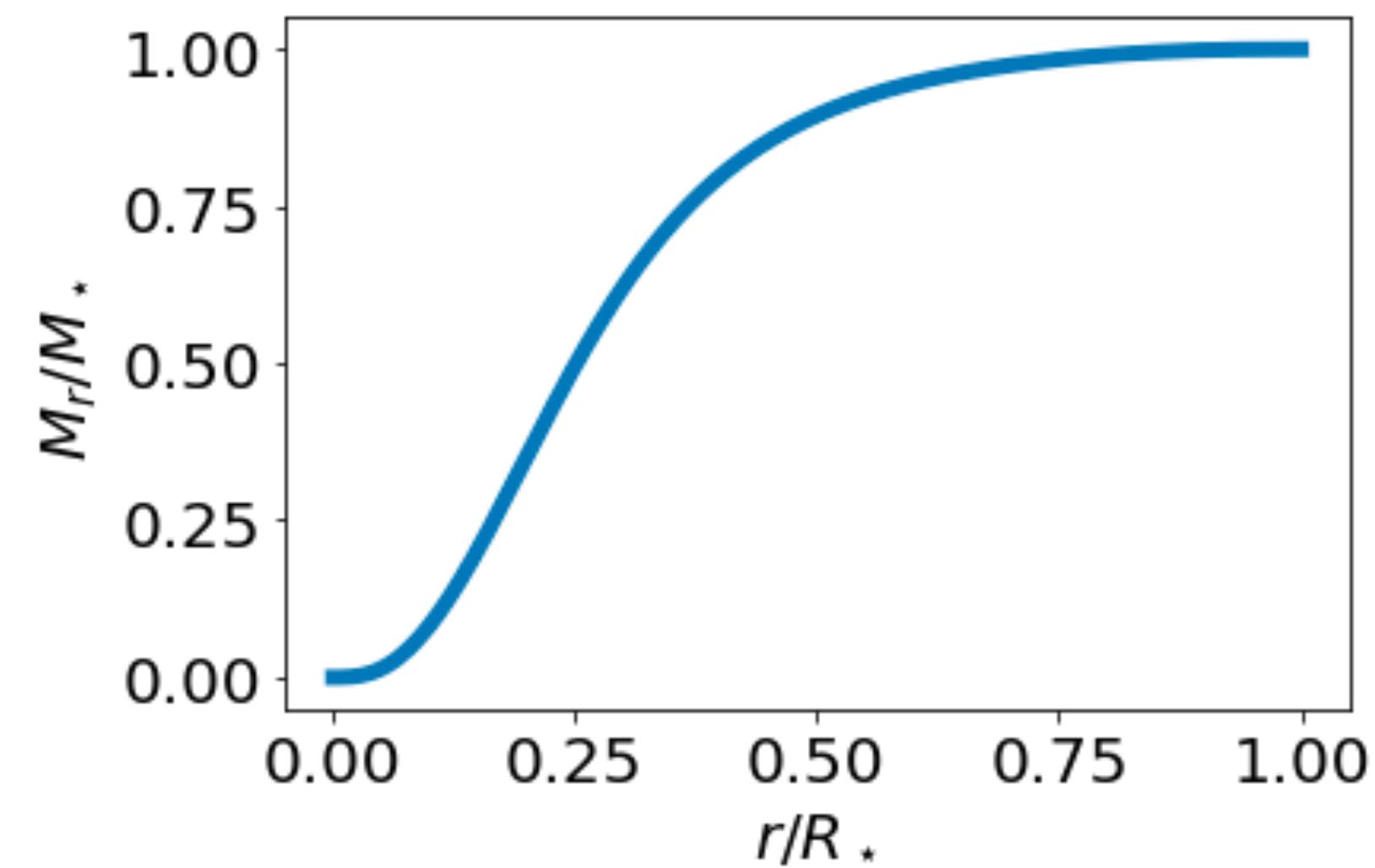


On this “Main Sequence”, stars are also ordered by their mass

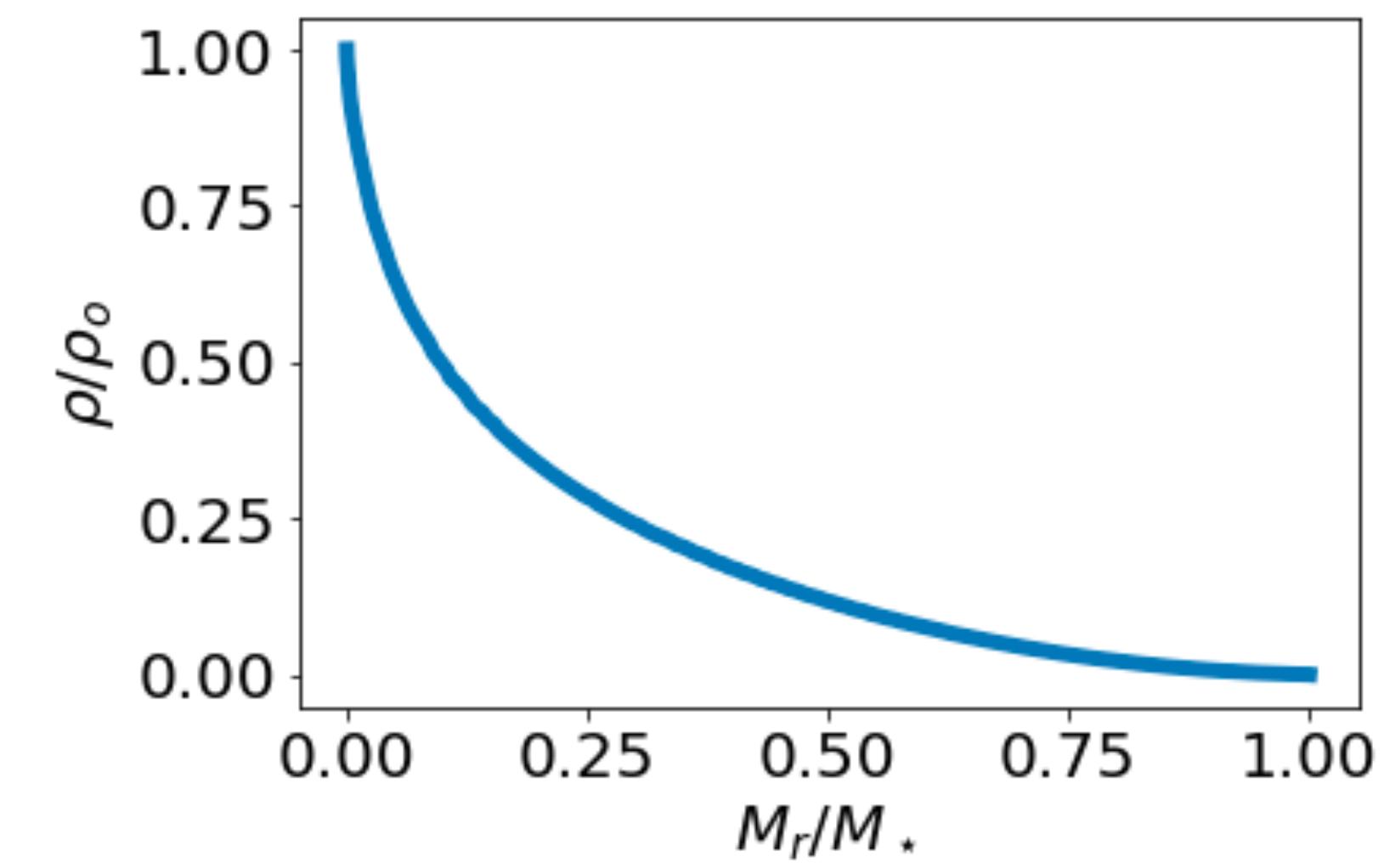
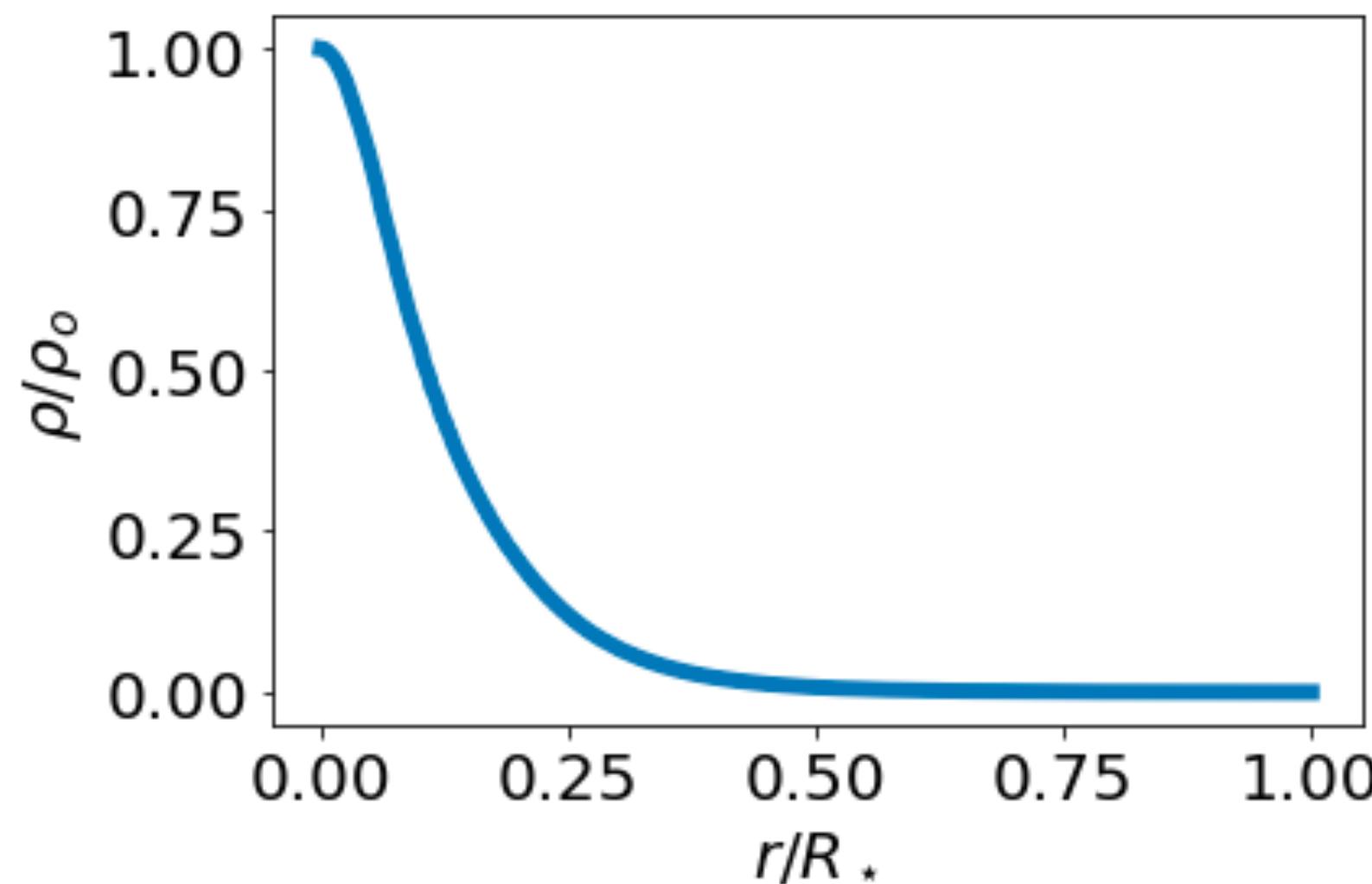
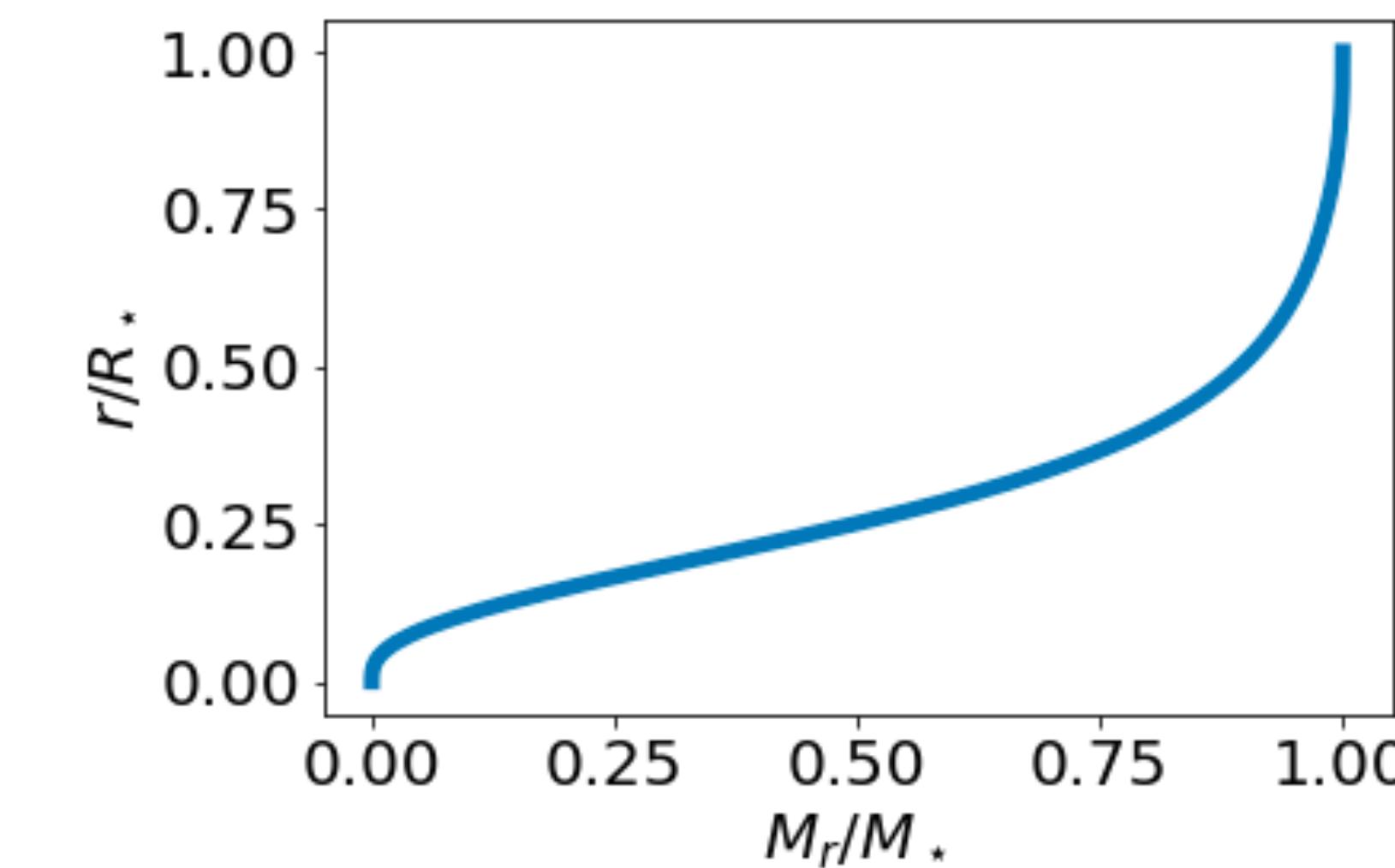




Radius coordinates



Enclosed mass coordinates



Radius coordinates

$$\frac{dM_r(r)}{dr} = 4\pi r^2 \rho(r)$$

$$\frac{dP(r)}{dr} = -\rho(r) \frac{GM_r(r)}{r^2}$$

$$\frac{dL_r(r)}{dr} = 4\pi r^2 \rho(r) \epsilon(r)$$

$$\frac{dT(r)}{dr} = -\frac{3\kappa_R(r)\rho(r)}{64\pi\sigma} \frac{L_r(r)}{r^2 T(r)^3}$$

Enclosed mass coordinates

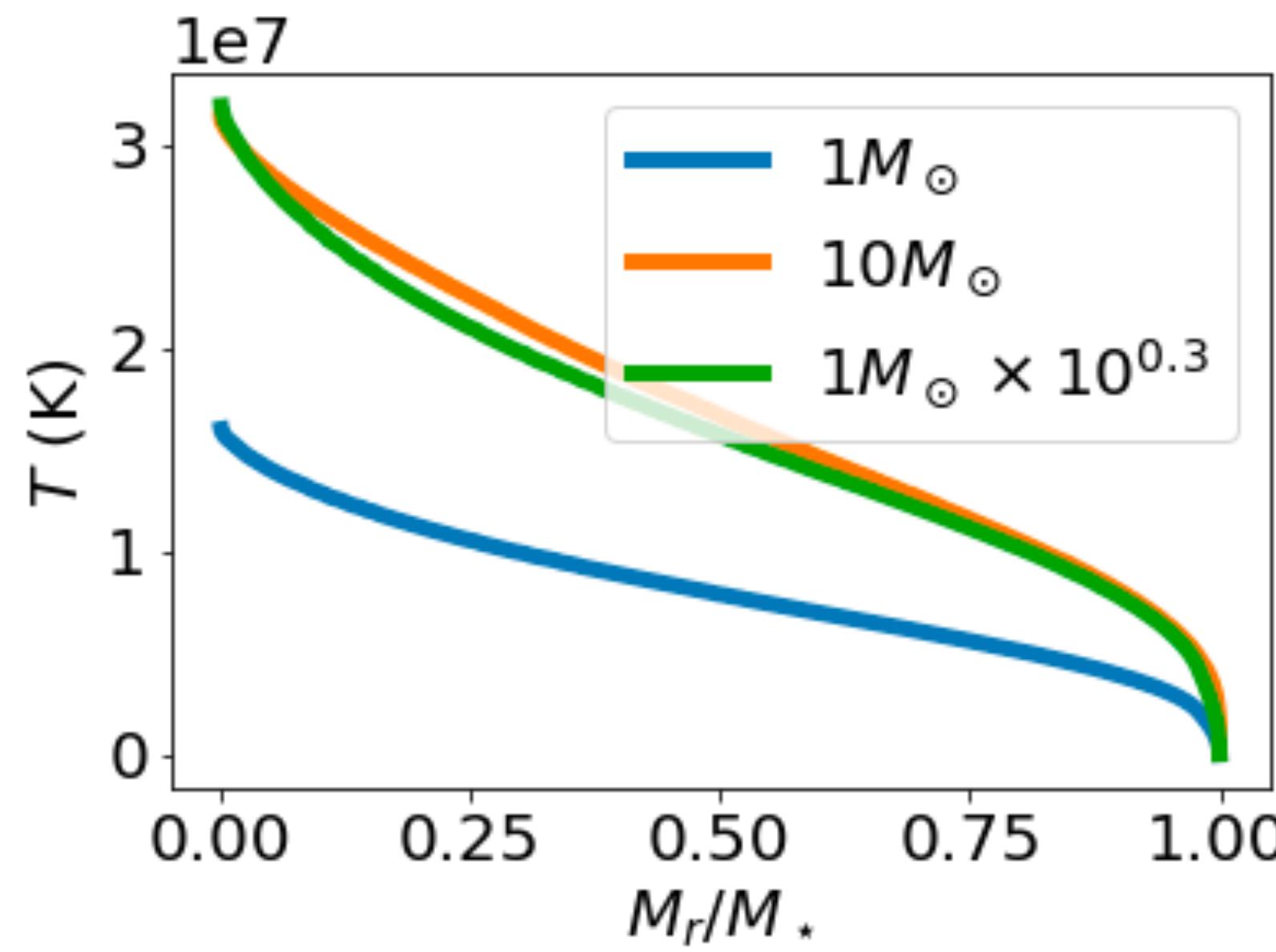
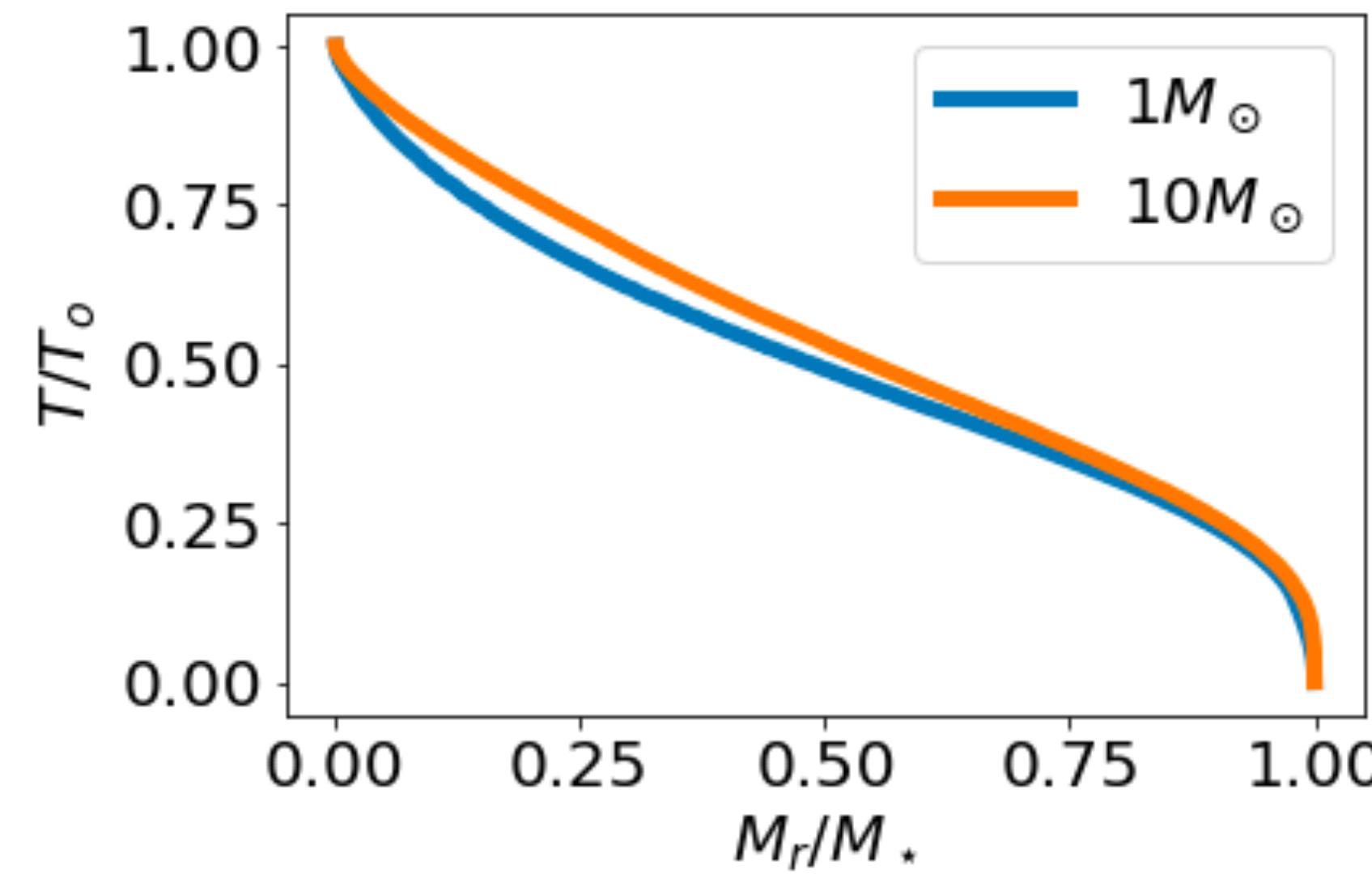
$$\frac{dr(M_r)}{dM_r} = \frac{1}{4\pi r^2(M_r) \rho(M_r)}$$

$$\frac{dP(M_r)}{dM_r} = -\frac{GM_r}{4\pi r^4(M_r)}$$

$$\frac{dL_r(M_r)}{dM_r} = \epsilon(M_r)$$

$$\frac{dT(M_r)}{dM_r} = -\frac{3\kappa(M_r) L_r(M_r)}{64\pi^2 ac r^4(M_r) T^3(M_r)}$$

Homology relations



$$q = \frac{M_r}{M_\star}$$

$$r = M_\star^{a_r} \tilde{r}(q)$$

$$P = M_\star^{a_P} \tilde{P}(q)$$

$$L_r = M_\star^{a_L} \tilde{L}_r(q)$$

$$T = M_\star^{a_T} \tilde{T}(q)$$

Enclosed mass coordinates

$$\frac{dr(M_r)}{dM_r} = \frac{1}{4\pi r^2(M_r) \rho(M_r)}$$

$$\frac{dP(M_r)}{dM_r} = -\frac{GM_r}{4\pi r^4(M_r)}$$

$$\frac{dL_r(M_r)}{dM_r} = \epsilon(M_r)$$

$$\frac{dT(M_r)}{dM_r} = -\frac{3\kappa(M_r) L_r(M_r)}{64\pi^2 ac r^4(M_r) T^3(M_r)}$$

Homology relations

$$q = \frac{M_r}{M_\star} \quad dM_r = M_\star dq$$

$$r = M_\star^{a_r} \tilde{r}(q)$$

$$P = M_\star^{a_P} \tilde{P}(q)$$

$$L_r = M_\star^{a_L} \tilde{L}_r(q)$$

$$T = M_\star^{a_T} \tilde{T}(q)$$

Enclosed mass coordinates

$$\frac{dP(M_r)}{dM_r} = -\frac{GM_r}{4\pi r^4(M_r)}$$

$$\frac{M_\star^{a_p} d\tilde{P}}{M_\star dq} = -\frac{G (qM_\star)}{4\pi \left(M_\star^{4a_r} \tilde{r}^4 \right)}$$

$$M_\star^{(4a_r + a_p - 2)} = -\frac{G q}{4\pi \tilde{r}^4 (d\tilde{P}/dq)}$$

$$(4a_r + a_p - 2) = 0$$

Homology relations

$$q = \frac{M_r}{M_\star} \quad dM_r = M_\star dq$$

$$r = M_\star^{a_r} \tilde{r}(q)$$

$$P = M_\star^{a_p} \tilde{P}(q)$$

$$L_r = M_\star^{a_L} \tilde{L}_r(q)$$

$$T = M_\star^{a_T} \tilde{T}(q)$$

Enclosed mass coordinates

$$\frac{dr(M_r)}{dM_r} = \frac{1}{4\pi r^2(M_r) \rho(M_r)}$$

$$\frac{M_\star^{a_r}}{M_\star} \frac{d\tilde{r}}{dq} = \frac{1}{4\pi M_\star^{2a_r} \tilde{r}^2(q) \rho}$$

$$\rho = \frac{1}{4\pi \tilde{r}^2(d\tilde{r}/dq)} M_\star^{1-3a_r}$$

$$\propto M_\star^{1-3a_r}$$

Homology relations

$$q = \frac{M_r}{M_\star} \quad dM_r = M_\star dq$$

$$r = M_\star^{a_r} \tilde{r}(q)$$

$$P = M_\star^{a_P} \tilde{P}(q)$$

$$L_r = M_\star^{a_L} \tilde{L}_r(q)$$

$$T = M_\star^{a_T} \tilde{T}(q)$$

Enclosed mass coordinates

$$\frac{dL_r(M_r)}{dM_r} = \epsilon(M_r)$$

$$\frac{M_\star^{a_L}}{M_\star} \frac{d\tilde{L}}{dq} = \boxed{\epsilon}$$

$$\boxed{\epsilon} = M_\star^{(a_L - 1)} \frac{d\tilde{L}_r}{dq}$$

$$\boxed{\epsilon \propto M_\star^{(a_L - 1)}}$$

Homology relations

$$q = \frac{M_r}{M_\star} \quad dM_r = M_\star dq$$

$$r = M_\star^{a_r} \tilde{r}(q)$$

$$P = M_\star^{a_P} \tilde{P}(q)$$

$$L_r = M_\star^{a_L} \tilde{L}_r(q)$$

$$T = M_\star^{a_T} \tilde{T}(q)$$

Enclosed mass coordinates

$$\frac{dT(M_r)}{dM_r} = -\frac{3\kappa(M_r) L_r(M_r)}{64\pi^2 ac r^4(M_r) T^3(M_r)}$$

$$\boxed{\kappa \propto M_\star^{(4a_T + 4a_r - a_L - 1)}}$$

Homology relations

$$q = \frac{M_r}{M_\star} \qquad dM_r = M_\star dq$$

$$r = M_\star^{a_r} \tilde{r}(q)$$

$$P = M_\star^{a_P} \tilde{P}(q)$$

$$L_r = M_\star^{a_L} \tilde{L}_r(q)$$

$$T = M_\star^{a_T} \tilde{T}(q)$$

Enclosed mass coordinates

$$\frac{dr(M_r)}{dM_r} = \frac{1}{4\pi r^2(M_r) \rho(M_r)}$$

$$\frac{dP(M_r)}{dM_r} = -\frac{GM_r}{4\pi r^4(M_r)}$$

$$\frac{dL_r(M_r)}{dM_r} = \epsilon(M_r)$$

$$\frac{dT(M_r)}{dM_r} = -\frac{3\kappa(M_r) L_r(M_r)}{64\pi^2 ac r^4(M_r) T^3(M_r)}$$

Homology relations

$$\rho \propto M_\star^{1-3a_r}$$

$$(4a_r + a_p - 2) = 0$$

$$\epsilon \propto M_\star^{(a_L - 1)}$$

$$\kappa \propto M_\star^{(4a_T + 4a_r - a_L - 1)}$$

But:

Ideal gas law says:

$$\rho \propto \frac{P}{T} \quad P = M_{\star}^{a_p} \tilde{P}(q)$$
$$T = M_{\star}^{a_T} \tilde{T}(q)$$

$$\propto M_{\star}^{(a_p - a_T)}$$

Homology relations

$$\rho \propto M_{\star}^{1-3a_r}$$

$$(4a_r + a_p - 2) = 0$$

$$\epsilon \propto M_{\star}^{(a_L - 1)}$$

$$\kappa \propto M_{\star}^{(4a_T + 4a_r - a_L - 1)}$$

$$(a_p - a_T) = 1 - 3a_r$$

$$3a_r + 1a_p - 1a_T = 1$$

But:

Ideal gas law says:

$$\rho \propto \frac{P}{T} \quad P = M_{\star}^{a_p} \tilde{P}(q) \quad T = M_{\star}^{a_T} \tilde{T}(q)$$
$$\propto M_{\star}^{(a_p - a_T)}$$

Homology relations

$$\rho \propto M_{\star}^{(1 - 3a_r)}$$

$$3a_r + 1a_p - 1a_T = 1$$

$$(4a_r + a_p - 2) = 0$$

$$4a_r + a_p = 2$$

The p-p chain says:

$$\epsilon \approx \rho T^4 \propto M_{\star}^{(1 - 3a_r)} M_{\star}^{4a_T}$$

$$\epsilon \propto M_{\star}^{(a_L - 1)}$$

$$3a_r + a_L - 4a_T = 2$$

Dominant opacity processes for (ρ, T) in the Sun:

$$\kappa \propto \rho T^{-7/2}$$

$$\kappa \propto M_{\star}^{(4a_T + 4a_r - a_L - 1)}$$

$$14a_r - 2a_L + 15a_T = 4$$

$$3a_r + 1a_p - 1a_T = 1$$

$4a_r$	a_p	a_L	a_T	=	2
$3a_r$	a_p	a_L	$-a_T$	=	1
$14a_r$	a_p	$-2a_L$	$15a_T$	=	4
$3a_r$	a_p	a_L	$-4a_T$	=	2

$$a_r = \frac{1}{13}$$

$$a_p = \frac{22}{13}$$

$$a_L = \frac{71}{13}$$

$$a_T = \frac{12}{13}$$

$$4a_r + a_p = 2$$

$$3a_r + a_L - 4a_T = 2$$

$$14a_r - 2a_L + 15a_T = 4$$

$$q = \frac{M_r}{M_\star}$$

$$a_r = \frac{1}{13}$$

$$r = M_\star^{a_r} \tilde{r}(q)$$

$$a_P = \frac{22}{13}$$

$$P = M_\star^{a_P} \tilde{P}(q)$$

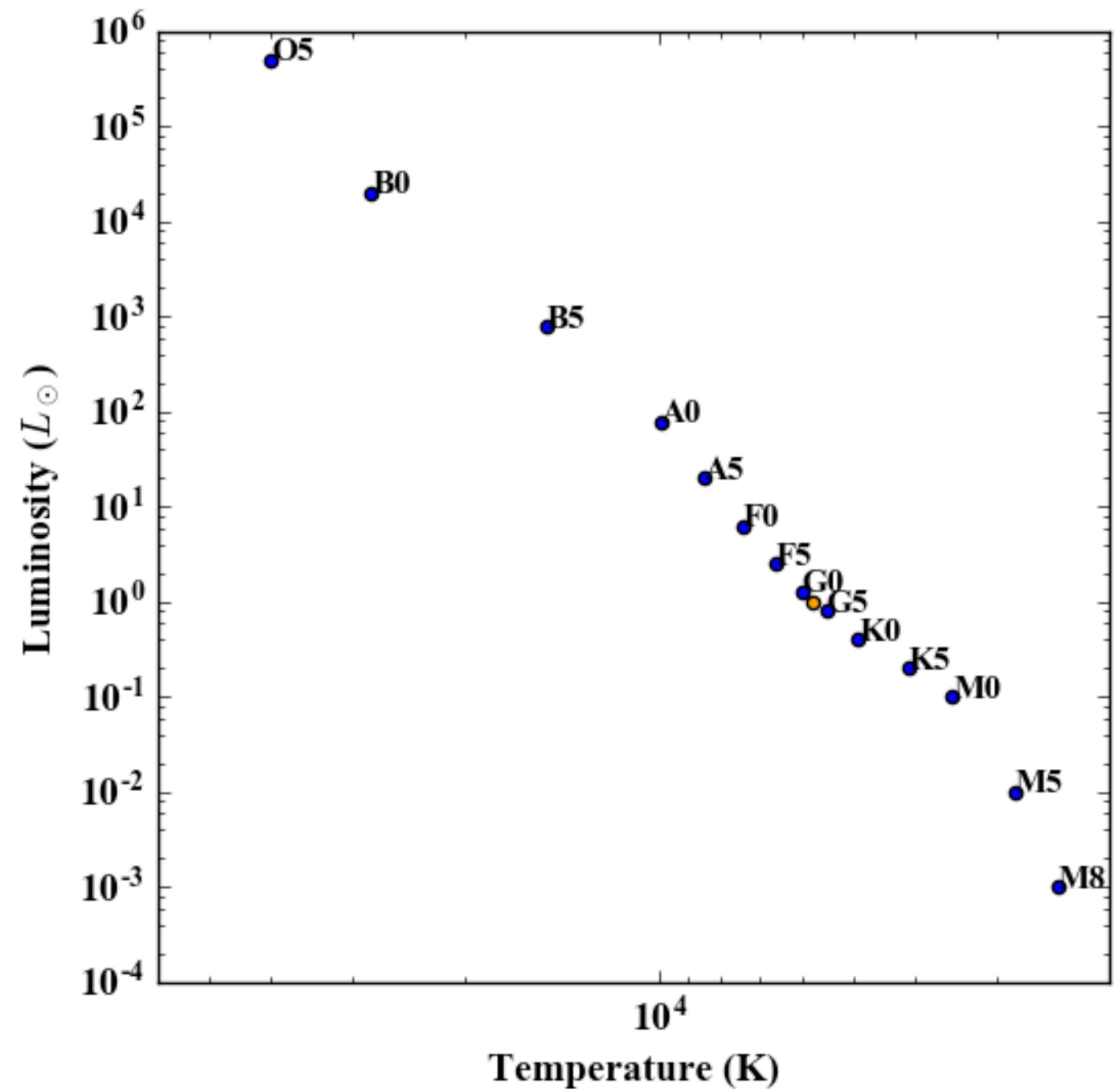
$$a_L = \frac{71}{13}$$

$$L_r = M_\star^{a_L} \tilde{L}_r(q)$$

$$a_T = \frac{12}{13}$$

$$T = M_\star^{a_T} \tilde{T}(q)$$

$$\frac{L_\star}{L_\odot} = \left(\frac{M_\star}{M_\odot} \right)^{71/13}$$



$$q = \frac{M_r}{M_\star}$$

$$a_r = \frac{1}{13}$$

$$r = M_\star^{a_r} \tilde{r}(q)$$

$$a_P = \frac{22}{13}$$

$$P = M_\star^{a_P} \tilde{P}(q)$$

$$a_L = \frac{71}{13}$$

$$L_r = M_\star^{a_L} \tilde{L}_r(q)$$

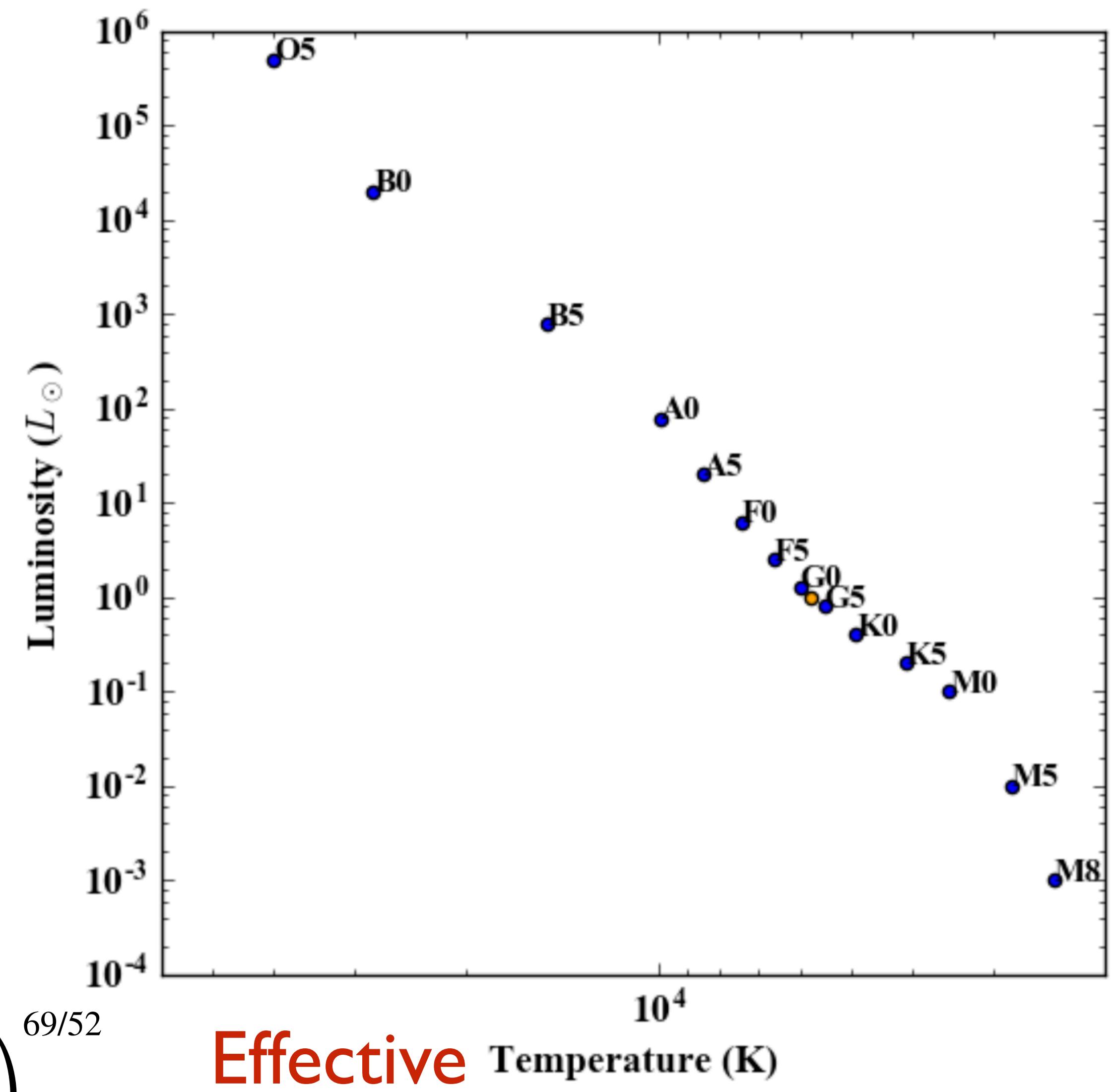
$$a_T = \frac{12}{13}$$

$$T = M_\star^{a_T} \tilde{T}(q)$$

The effective temperature of a star is the _____
of a _____ with the same _____

$$\frac{T_{\text{eff},\star}}{T_{\text{eff},\odot}} = \left(\frac{M_\star}{M_\odot}\right)^{(a_L - 2a_r)/4} = \left(\frac{M_\star}{M_\odot}\right)^{69/52}$$

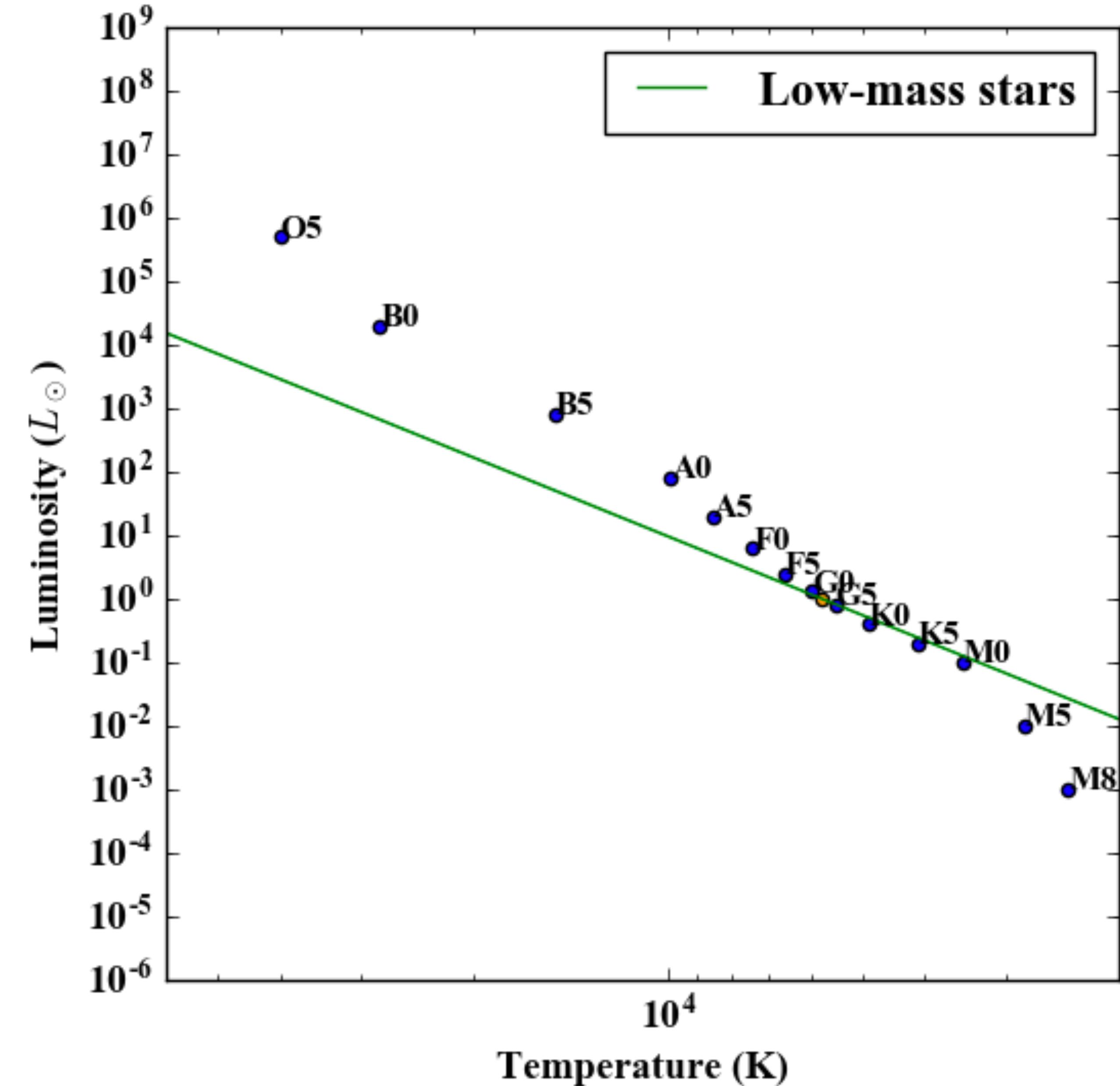
$$M_{\text{array}} = [1, 5, 10, 15, \dots] M_\odot$$



$$\frac{L_\star}{L_\odot} = \left(\frac{M_\star}{M_\odot} \right)^{71/13}$$

$$\frac{T_{\text{eff},\star}}{T_{\text{eff},\odot}} = \left(\frac{M_\star}{M_\odot} \right)^{(a_L - 2a_r)/4} = \left(\frac{M_\star}{M_\odot} \right)^{69/52}$$

$M_{\text{array}} = \text{np.linspace}(0.1, 40, 100)$



But:

Ideal gas law says:

$$\rho \propto \frac{P}{T} \propto M_{\star}^{(a_p - a_T)}$$

Homology relations

$$\rho \propto M_{\star}^{(1 - 3a_r)}$$

$$3a_r + 1a_p - 1a_T = 1$$

What needs to change for more massive stars?

$$(4a_r + a_p - 2) = 0$$

$$4a_r + a_p = 2$$

The p-p chain says:

$$\epsilon \approx \rho T^4 \propto M_{\star}^{(1 - 3a_r)} M_{\star}^{4a_T}$$

$$\epsilon \propto \rho T^{18} \propto M_{\star}^{(1 - 3a_r)} M_{\star}^{4a_T}$$

$$\epsilon \propto M_{\star}^{(a_L - 1)}$$

$$3a_r + a_L - 4a_T = 2$$

$$3a_r + a_L - 18a_T = 2$$

Dominant opacity processes for (ρ, T) in the Sun:

$$\kappa \propto \rho T^{-7/2}$$

$$\kappa \propto \text{constant}$$

$$\kappa \propto M_{\star}^{(4a_T + 4a_r - a_L - 1)}$$

$$14a_r - 2a_L + 15a_T = 4$$

$$4a_r - 1a_L + 4a_T = 1$$

$$3a_r + 1a_p - 1a_T = 1$$

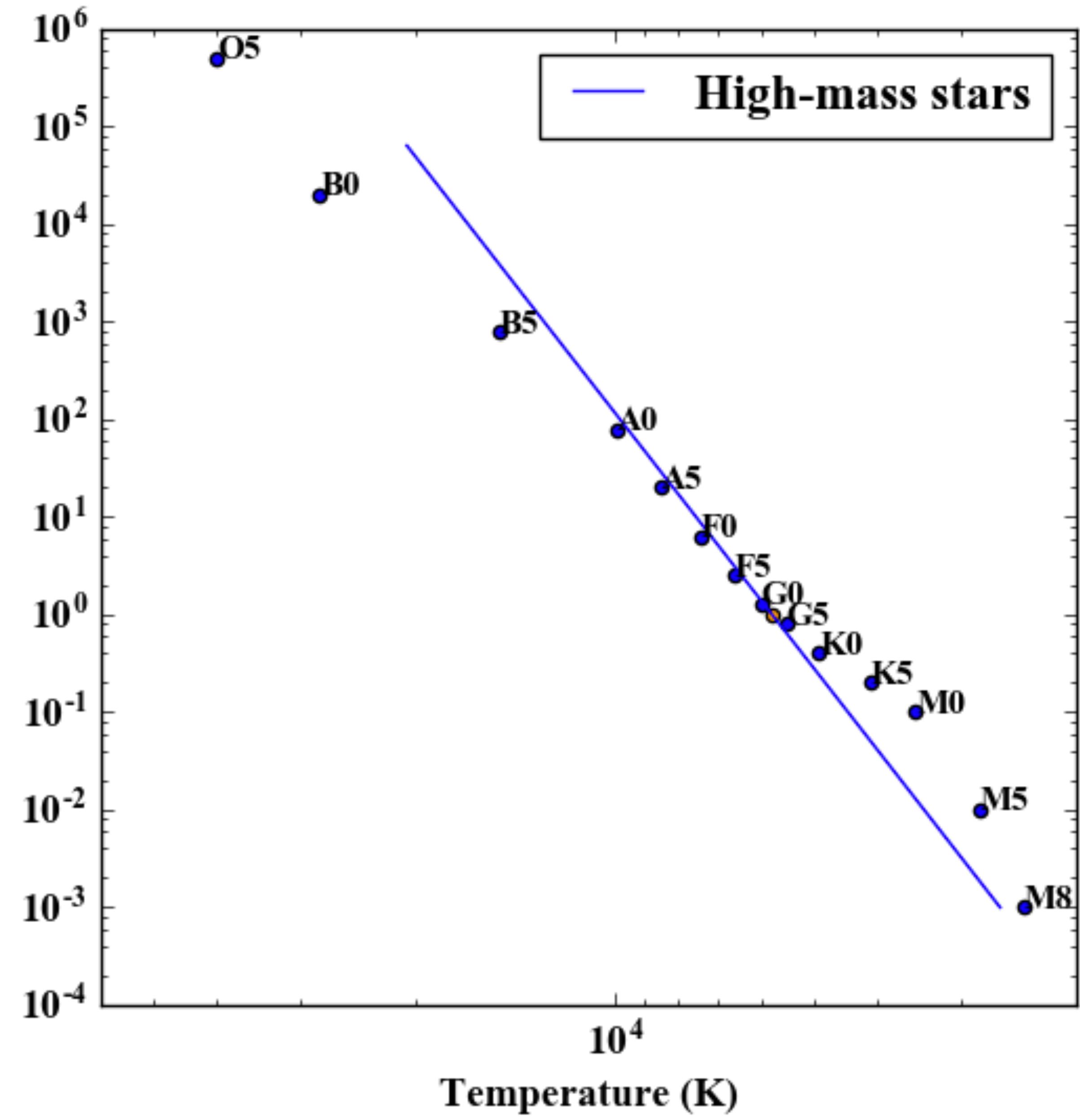
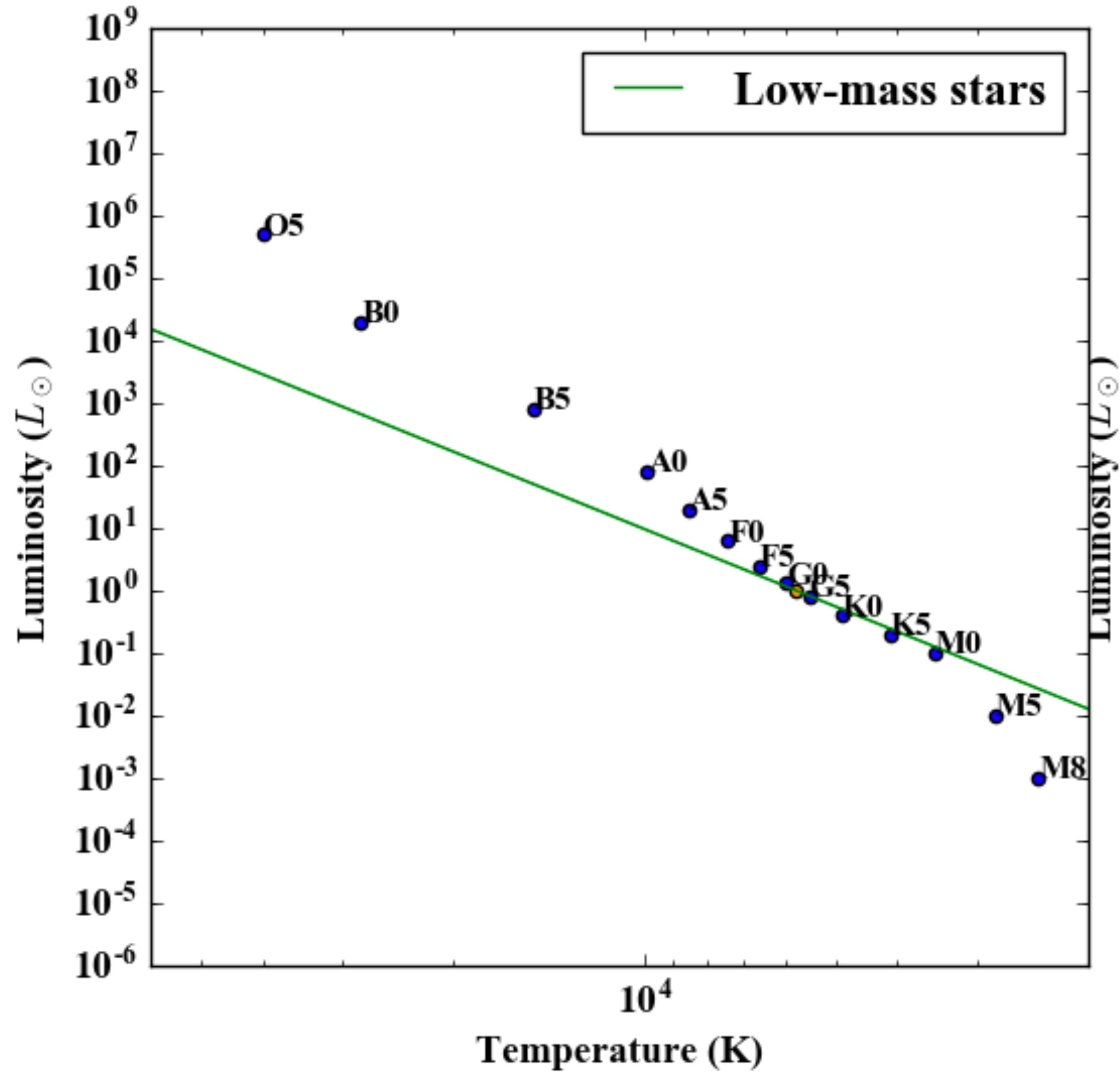
$$\begin{array}{rccccccccc} 4 & | & a_r & & 1 & | & a_p & 0 & | & a_L & 0 & | & a_T & = & 2 \\ 3 & | & a_r & & 1 & | & a_p & 0 & | & a_L & -1 & | & a_T & = & 1 \\ 4 & | & a_r & & 0 & | & a_p & -1 & | & a_L & 4 & | & a_T & = & 1 \\ 3 & | & a_r & & 0 & | & a_p & 1 & | & a_L & -18 & | & a_T & = & 2 \end{array}$$

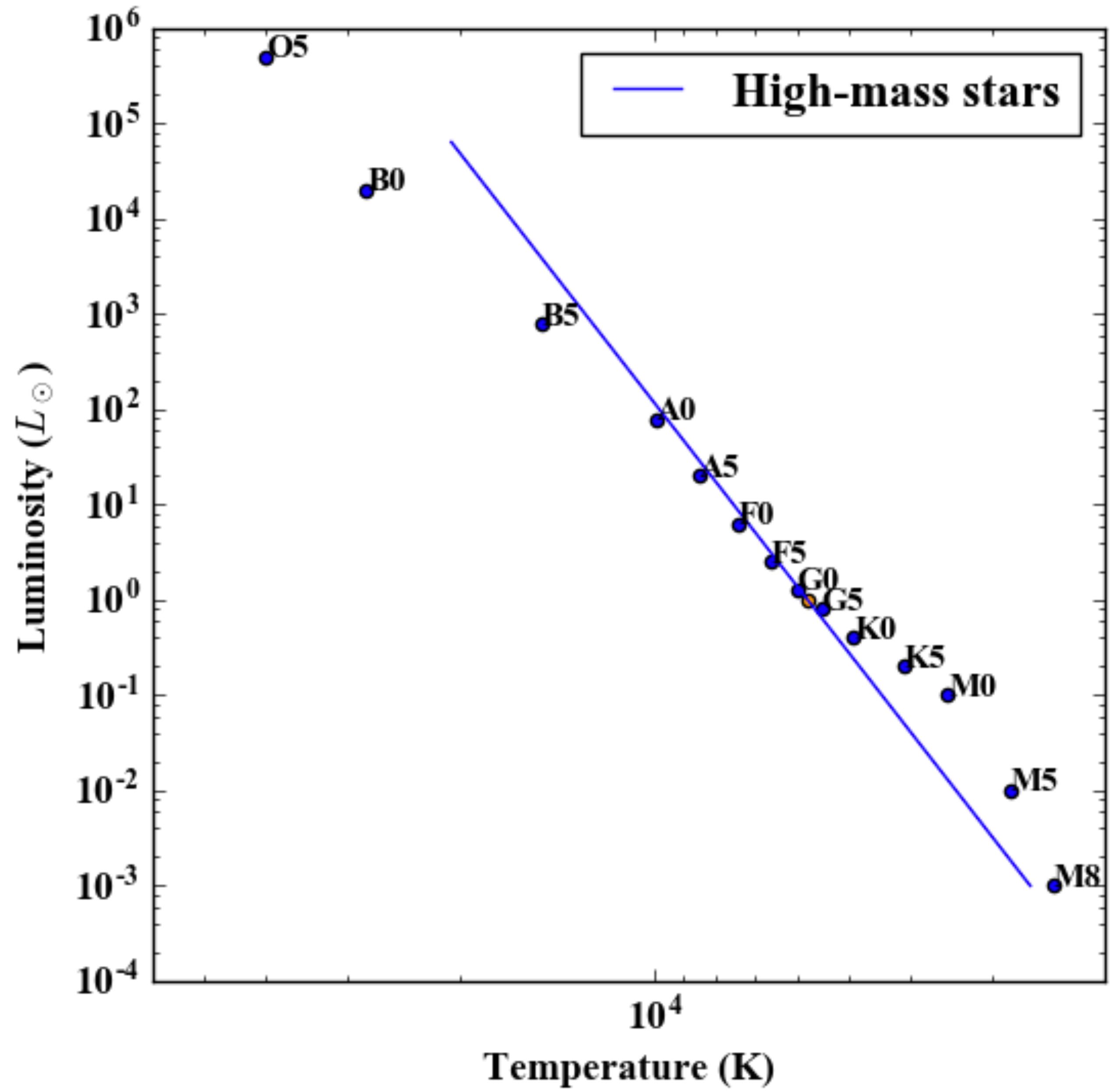
$$a_r = \frac{17}{21} \quad a_p = -\frac{26}{21} \quad a_L = 3 \quad a_T = \frac{4}{21}$$

$$4a_r + a_p = 2$$

$$3a_r + a_L - 18a_T = 2$$

$$4a_r - 1a_L + 4a_T = 1$$





What would happen if a massive stars was powered by PP chain instead of CNO cycle?

Let's define a generalized case:

$$\epsilon = \epsilon_0 \rho T^\eta$$

$$k = k_0(\alpha) (\rho T^{-7/2})^\alpha$$

$$\begin{bmatrix} 4 & 1 & 0 \\ 3 & 1 & 0 \\ (8+6\alpha) & 0 & -2 \\ 3 & 0 & 1 \end{bmatrix} \begin{bmatrix} a_r \\ a_p \\ a_L \\ q_T \end{bmatrix} = \begin{bmatrix} 2 \\ 1 \\ 2(\alpha+1) \\ 2 \end{bmatrix}$$

Let's define a generalized case:

$$\epsilon = \epsilon_0 \rho T^\eta$$

$$k = k_0(\alpha) (\rho T^{-7/2})^\alpha$$

The solution is:

$$a_r = \frac{2\eta - 5\alpha - 2}{2\eta - \alpha + 6}$$

$$a_p = -\frac{4\eta + 18\alpha + 20}{2\eta - \alpha + 6}$$

$$a_L = \frac{(6+4\alpha)\eta + 13\alpha + 18}{2\eta - \alpha + 6}$$

$$a_T = \frac{4\alpha + 8}{2\eta - \alpha + 6}$$

Let's define a generalized case:

$$\begin{aligned}\mathcal{E} &= \mathcal{E}_0 \rho T^\eta \\ k &= k_0(\alpha) (\rho T^{-7/2})^\alpha\end{aligned}$$

Let's look at the luminosity relation:

$$a_L = \frac{(6+4\alpha)\eta + 13\alpha + 18}{2\eta - \alpha + 6} = 3 + 2\alpha + \frac{2\alpha(\alpha+2)}{2\eta + 6 - \alpha}$$

$$\frac{L_\star}{L_\odot} = \left(\frac{M_\star}{M_\odot} \right)^{a_L}$$

For low mass stars, $\alpha = 1$

$$a_L = 5 + \frac{6}{2\eta + 5}$$

For low mass stars, $\alpha = 0$

$$a_L = 3 !$$

Independent of the T scaling
of the nuclear reactions !!

Same luminosity, but hotter!

