

(6)

Analytically we have $\frac{\partial u}{\partial t} = D \frac{\partial^2 u}{\partial x^2}$

$$\text{w/ } u = e^{\alpha t} e^{j\sigma x} \Rightarrow \alpha = -D\sigma^2$$

- all modes decay (i.e. all $\alpha < 0$)
- longest waves decay slowest
- sol'n gets smoother over time

$$\text{Now } \gamma = e^{\alpha \Delta t} = e^{-D\sigma^2 k} = e^{-r(\sigma h)^2}$$

- but as $k \rightarrow 0$, $\gamma_0 \rightarrow \gamma \rightarrow 1$ don't learn much

Common to introduce characteristic time, T

and examine $\left(\frac{\gamma_0}{\gamma}\right)^N$ where $N = \frac{T}{k}$

(i.e. # time-steps
to advance sol'n by T)

- use time constant of T in analytic sol'n

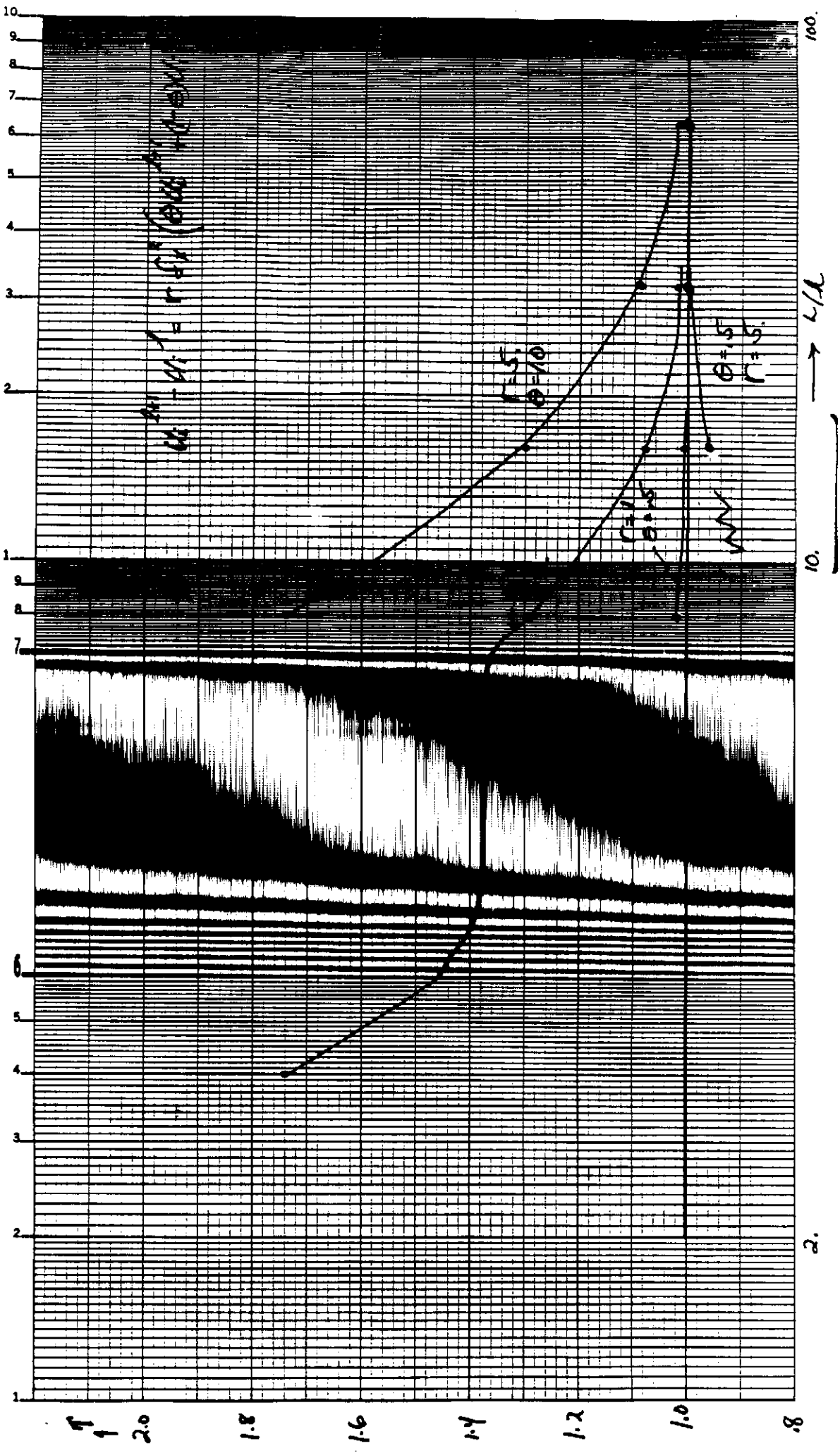
$$\text{i.e. } T = \frac{1}{|\alpha|} = \frac{1}{D\sigma^2}, \therefore N = \frac{1}{D\sigma^2 k} = \frac{1}{r(\sigma h)^2}$$

$$\text{Define } \uparrow T = \left(\frac{\gamma_0}{\gamma}\right)^N = \frac{\gamma_0^{\frac{1}{r(\sigma h)^2}}}{\left(e^{-r(\sigma h)^2}\right)^{\frac{1}{r(\sigma h)^2}}} = \frac{\gamma_0^{\frac{1}{r(\sigma h)^2}}}{e^{-1}}$$

"Propagation Factor"

Plot T vs $\sigma h = \frac{2\pi h}{L}$ for various r

$T=1$ is perfect



Also note...

$$\gamma = e^{-r(\sigma h)^2} = 1 - r(\sigma h)^2 + \frac{(r(\sigma h)^2)^2}{2} - \frac{(r(\sigma h)^2)^3}{3!} + \dots$$

$$\gamma_0 = 1 - 2r(1 - \cos \sigma h)$$

$$= 1 - 2r \left[\frac{(\sigma h)^2}{2!} - \frac{(\sigma h)^4}{4!} + \frac{(\sigma h)^6}{6!} - \dots \right]$$

$$= 1 - r(\sigma h)^2 + \underbrace{\frac{r(\sigma h)^4}{12}}_{\text{leading error term}} - \frac{2r(\sigma h)^6}{6!} + \dots$$

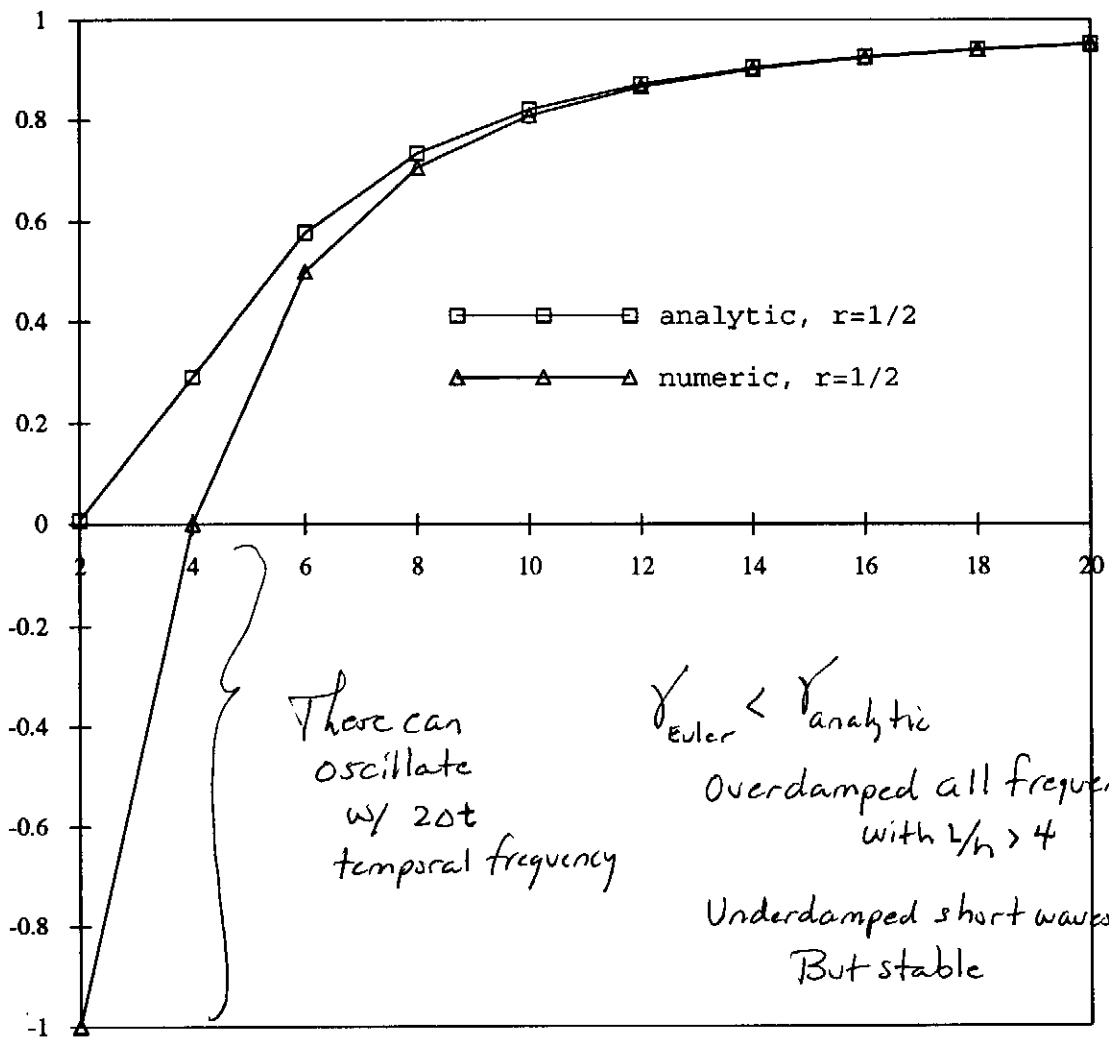
leading error term

But $\frac{r(\sigma h)^4}{12} = \frac{r^2(\sigma h)^4}{2}$ when $r = 1/6$

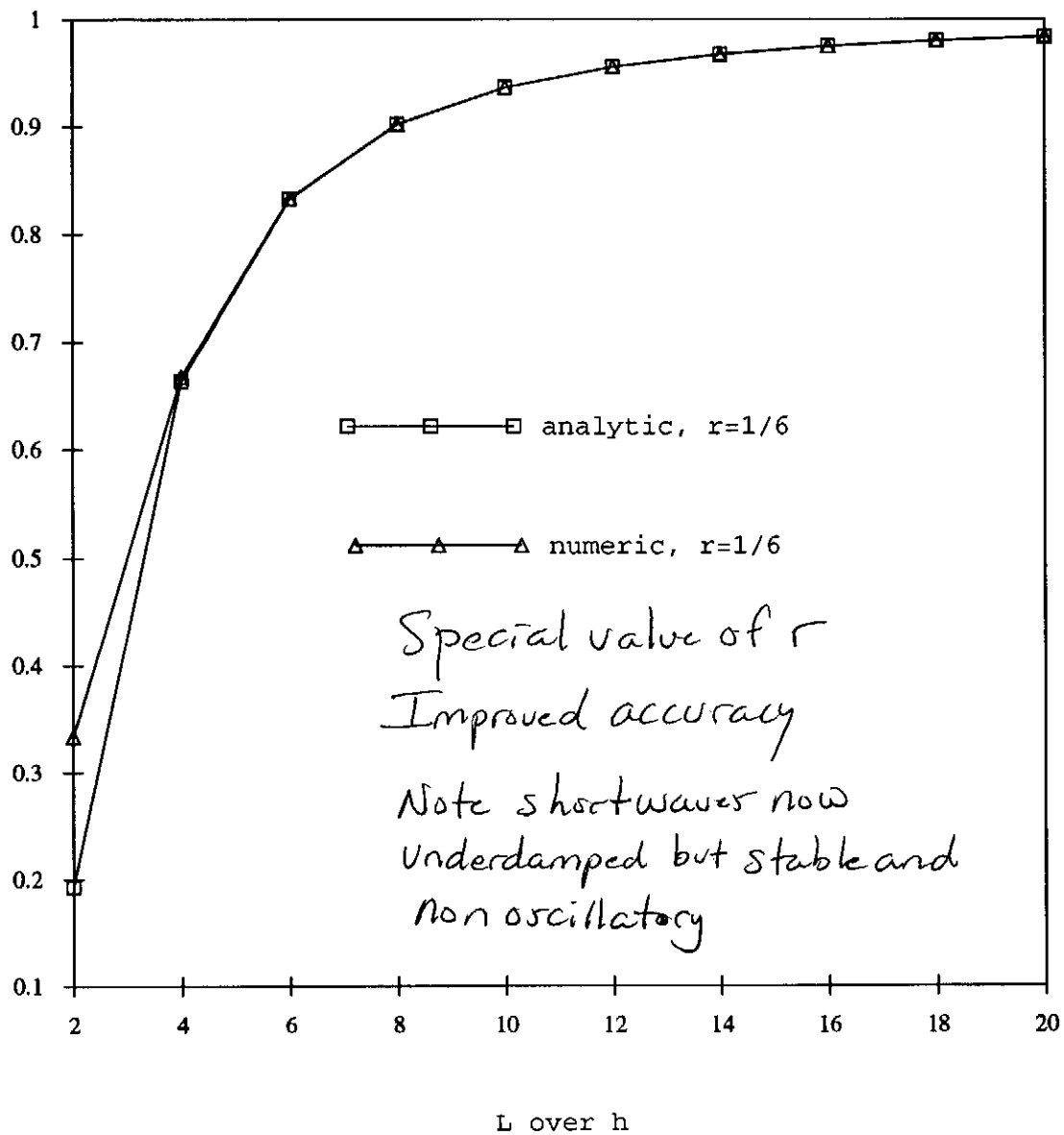
Error is "pushed back" one more term
(we saw this earlier... st error just
cancels leading Δx error)

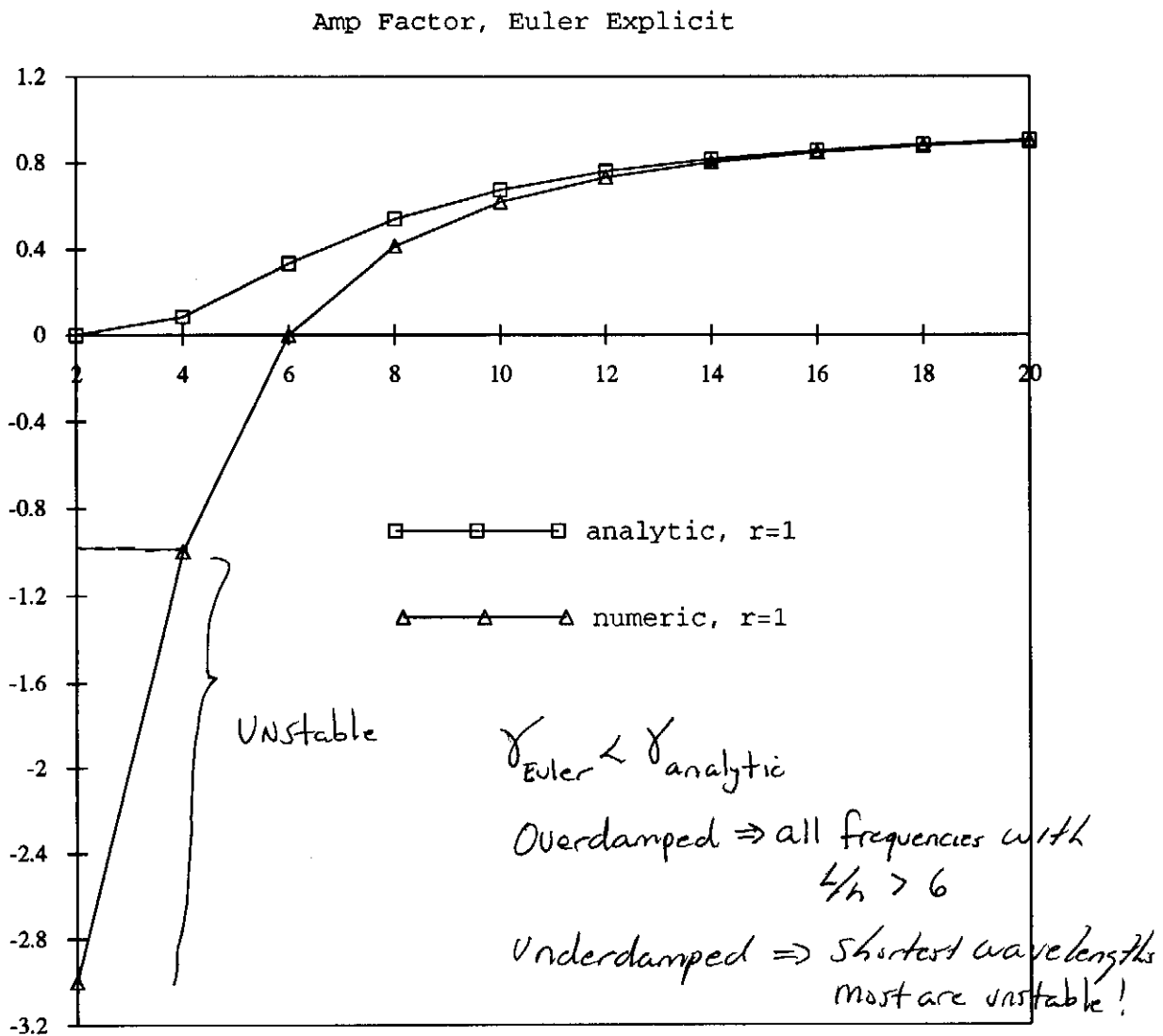
For Euler: $\gamma < \gamma_0 \Rightarrow$ Numerical sol'n underdamped
(generally true... depends on r and σh)

Amp Factor, Euler Explicit

 L over h

Amp Factor, Euler Explicit





(9)

Can examine any scheme in this manner

e.g. "Richardson" $u_i^{l+1} - u_i^{l-1} = r \delta_x^2 u_i^l$

$$\gamma - \frac{1}{\gamma} = 2r(\cosh - 1)$$

$$\gamma^2 + 2r(1 - \cosh)\gamma - 1 = 0 \quad \underline{2 \text{ Roots!}}$$

- Stability for general quadratic

$$a\gamma^2 + b\gamma + c = 0 \Rightarrow |\gamma| \leq 1$$

$$\text{when } \frac{c}{a} \leq 1 \text{ and } |b| \leq a + c$$

In our case... $\frac{c}{a} = -1 \leq 1$ always

$$|b| = 2r(1 - \cosh) \leq 0 \quad \xrightarrow{\text{always positive!}}$$

No value of r satisfies this constraint for all values of σh ... short waves are the biggest offenders as usual!

Unconditionally unstable!!

- Stability Analysis using Matrix Methods

- In Lax-Richtmyer view... if we have the scheme $u^{l+1} = Au^l + c^l$; A grows in size; need to show $\|A\| \leq 1$ guarantees stability (and \therefore convergence for a consistent molecule)
- In practical view of fixed mesh lengths... if have a scheme of form $u^{l+1} = Au^l + c^l$... A has fixed size; sufficient to show $\rho(A) \leq 1$ to ensure boundedness
- Formally must have $\rho(A) \leq \|A\| \leq 1$ as size of $A \rightarrow \infty$ to guarantee convergence for a consistent scheme (possible to have $\rho(A) \leq 1$ w/ $\|A\| > 1$)

e.g. Euler Explicit: $u_i^{l+1} = r u_{i-1}^l + (1-2r) u_i^l + r u_{i+1}^l$
w/ Type I BCs...

$$A = \begin{bmatrix} 1-2r & r & 0 & \dots & 0 \\ r & 1-2r & r & 0 & \dots & 0 \\ 0 & r & 1-2r & r & 0 & \dots \\ 0 & \dots & \dots & \dots & \dots & \dots \end{bmatrix}$$

$$\|A\|_{\infty} = |r| + |1-2r| + |r| ; \text{ Need } \|A\|_{\infty} \leq 1$$

$$\text{if } 1-2r > 0 \text{ then } \|A\|_{\infty} = 1$$

$$1-2r < 0 \quad \|A\|_{\infty} = 4r-1 > 1 \text{ since } r > 1/2$$

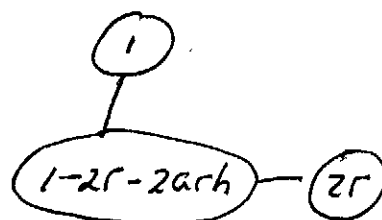
Conclude ... Need $r \leq 1/2$... Same as Von Neumann

(B)

Now if we have derivative BCs... e.g. Type III

$$\frac{\partial u}{\partial x} = au + b \text{ at } x=0 \text{ boundary}$$

... then molecule becomes



$$\text{i.e. } U_i = U_i - 2ahU_0 - bh$$

So A has the structure ...

$$\begin{bmatrix} 1-2r(1+ah) & 2r & & & \\ & r & 1-2r & r & \\ & & r & 1-2r & r \\ & & & \ddots & \ddots \\ & & & & \ddots \end{bmatrix}$$

- all but row 1 require $r \leq 1/2$ for $\|A\|_\infty \leq 1$
- must see if row 1 changes this restriction...

$$\text{we want } |1-2r(1+ah)| + |2r| \leq 1$$

Two cases to consider:

$$(a) 1-2r(1+ah) \geq 0 \text{ (i.e. diagonal term positive)}$$

$$\text{then } |1-2r(1+ah)| + 2r = 1-2r(1+ah) + 2r \leq 1$$

$$1-2rah \leq 1 \text{ always OK}$$

(C)

But for diagonal coefficient to be positive

$$2r(1+ah) \leq 1 \Rightarrow r \leq \frac{1}{2(1+ah)}$$

So if coefficient is positive, problem is stable and we need $r \leq \frac{1}{2(1+ah)}$ to achieve this

(b) $1 - 2r(1+ah) \leq 0$ (diagonal is negative)

$$\text{then } |1 - 2r(1+ah)| + |2r| = 2r(1+ah) - 1 + 2r \leq 1$$

$$2r(2+ah) \leq 2$$

$$r \leq \frac{1}{2+ah}$$

But $\frac{1}{2(1+ah)} < \frac{1}{2+ah} \dots$ So we can maintain stability

when diagonal turns negative provided $r \leq \frac{1}{2+ah}$

IF r gets any bigger... diagonal still negative, but $\|A_h\| > 1$

- Also note $\frac{1}{2+ah} < \frac{1}{2}$ i.e. stability restriction greater w/ Type III than Type I BCG!