

Explicit Schemes on FE

- "Mass lumping" sum terms in the mass matrix and place result on diagonal
i.e. "diagonalize" Mass Matrix \Rightarrow Inversion is trivial \Rightarrow explicit scheme
- Controversial... works well, but no strong rationale... Stiffness matrices not treated this way... only mass matrix

eg. 1-D Diffusion: ($\theta = 0$)

$$\frac{h}{6} [u_{i-1} + 4u_i + u_{i+1}]^{k+1} = \frac{h}{6} [u_{i-1} + 4u_i + u_{i+1}]^k + \frac{K\Delta t}{h} [u_{i-1} - 2u_i + u_{i+1}]^k$$

$$\downarrow \quad \quad \quad \downarrow$$

$$u_i^{k+1} = u_i^k + \frac{K\Delta t}{h^2} [u_{i-1} - 2u_i + u_{i+1}]^k$$

No matrices!

Stability? \Rightarrow same as FD on uniform mesh

- More systematic approach "Integral lumping"

Treat all integrations (i.e. matrices) equivalently

View "lumping" as a quadrature approximation

e.g. in 1-D $\int_{x_i}^{x_{i+1}} f(x) dx \approx \frac{f(x_i) + f(x_{i+1})}{2} h$

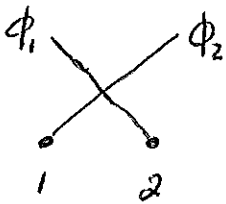
(Trapezoidal Rule)

$$\int_{x_i}^{x_{i+1}} f(x) dx = \int_{-1}^1 f(x(\xi)) \frac{h}{2} d\xi \approx \frac{h}{2} [f(x(-1)) + f(x(1))]$$

$x = \sum_i x_i \phi_i$ ξ x_i x_{i+1} Same!

Nodal quadrature rule \Rightarrow "place gauss points at nodes", sufficient as long as we can at least integrate a constant \Rightarrow i.e. get Area of element

e.g. $\langle \phi_i \phi_j \rangle = \frac{h}{2} [\phi_i(-1) \phi_j(-1) + \phi_i(1) \phi_j(1)]$



$$A^e = \begin{bmatrix} \frac{h}{2} & 0 \\ 0 & \frac{h}{2} \end{bmatrix} \text{ diagonal!}$$

Note: "Gauss pt weights" are unity

e.g. 1-D diffusion eqn

get same thing here
since $\frac{\partial \phi_i}{\partial x}, \frac{\partial \phi_j}{\partial x}$ are constant on an element

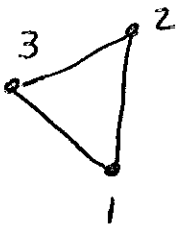
$$h u_i^{k+1} = h u_i^k + \frac{K \Delta t}{h} \left[u_{i-1} - 2u_i + u_{i+1} \right]^k$$

$$u_i^{k+1} = u_i^k + \frac{K \Delta t}{h^2} \left[u_{i-1} - 2u_i + u_{i+1} \right]^k$$

Works in Multi-D situations as well

e.g. Linear triangle:

$$M_{ij}^e = \underbrace{2A^e}_{|K|} \underbrace{\int \phi_i \phi_j dA}_{g_{ij}} \approx 2A^e \sum_{k=1}^3 g_{ij}(z_k) \omega_k$$



ω_k 's such that
if g_{ij} constant
still get A^e
 $\Rightarrow \omega_k = 1/3$

$$[M]^e = \begin{bmatrix} A^e/3 & 0 & 0 \\ 0 & A^e/3 & 0 \\ 0 & 0 & A^e/3 \end{bmatrix}$$

$$2A^e \left[\phi_1(1) \phi_1(1) \omega_1 + \underbrace{\phi_1(2) \phi_1(2)}_0 \omega_2 + \underbrace{\phi_1(3) \phi_1(3)}_0 \omega_3 \right]$$

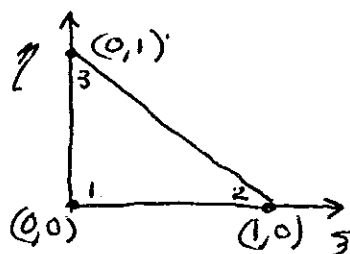
$$2A^e \left[\omega_1 \underbrace{\phi_1(1)}_0 \phi_2(1) + \omega_2 \underbrace{\phi_1(2)}_0 \phi_2(2) + \underbrace{\phi_1(3) \phi_2(3)}_0 \omega_3 \right]$$

etc
⋮

Diagonal!

Jacobian on Triangle:

"parent" element in ξ, η space:



$$\left. \begin{aligned} \phi_1 &= 1 - \xi - \eta \\ \phi_2 &= \xi \\ \phi_3 &= \eta \end{aligned} \right\} \text{check} \Rightarrow \begin{aligned} \phi_1(2) &= 0, \phi_1(3) = 0 \\ \phi_2(1) &= 0, \phi_2(3) = 0 \\ \phi_3(1) &= 0, \phi_3(2) = 0 \end{aligned}$$

$$\frac{\partial \phi_1}{\partial \xi} = -1 \quad \frac{\partial \phi_1}{\partial \eta} = -1$$

$$\frac{\partial \phi_2}{\partial \xi} = 1 \quad \frac{\partial \phi_2}{\partial \eta} = 0$$

$$\frac{\partial \phi_3}{\partial \xi} = 0 \quad \frac{\partial \phi_3}{\partial \eta} = 1$$

$$\sum \phi_i = 1 ; \sum \frac{\partial \phi_i}{\partial \xi} = 0 ; \sum \frac{\partial \phi_i}{\partial \eta} = 0$$

$$|J| = \frac{\partial x}{\partial \xi} \frac{\partial y}{\partial \eta} - \frac{\partial x}{\partial \eta} \frac{\partial y}{\partial \xi}$$

$$\frac{\partial x}{\partial \xi} = x_1 \frac{\partial \phi_1}{\partial \xi} + x_2 \frac{\partial \phi_2}{\partial \xi} + x_3 \frac{\partial \phi_3}{\partial \xi} = x_2 - x_1$$

$$\frac{\partial x}{\partial \eta} = x_1 \frac{\partial \phi_1}{\partial \eta} + x_2 \frac{\partial \phi_2}{\partial \eta} + x_3 \frac{\partial \phi_3}{\partial \eta} = x_3 - x_1$$

$$\text{same for } \frac{\partial y}{\partial \xi} = y_2 - y_1 ; \frac{\partial y}{\partial \eta} = y_3 - y_1$$

$$\begin{aligned}
 \text{Then } |J| &= (x_2 - x_1)(y_3 - y_1) - (y_2 - y_1)(x_3 - x_1) \\
 &= (x_2 - x_1)\Delta y_2 + (x_3 - x_1)\Delta y_3 \\
 &= x_2\Delta y_2 + x_3\Delta y_3 - x_1y_3 + x_1y_1 - x_1y_1 + x_1y_2 \\
 &= x_2\Delta y_2 + x_3\Delta y_3 + x_1\Delta y_1 \\
 &= 2A^e
 \end{aligned}$$

Derivative terms remain the same
 since $\underbrace{\left\langle \frac{\partial \phi_i}{\partial x} \frac{\partial \phi_j}{\partial x} \right\rangle^e}_{\text{constant}}$ and $\langle 1 \rangle^e$ w/

nodal quadrature gives exactly the Area
 of the element (i.e. same result produced
 by the "regular" quadrature

• What about Bilinear element (deformed)?

- easiest to call Basis subroutine

w/ "Gauss pts" $\left. \begin{aligned} (z_1, \eta_1) &= (-1, -1) \\ (z_2, \eta_2) &= (1, -1) \\ (z_3, \eta_3) &= (1, 1) \\ (z_4, \eta_4) &= (-1, 1) \end{aligned} \right\} \begin{array}{l} \text{node} \\ \text{positions} \\ \text{in} \\ (z, \eta) \end{array}$

- Need to determine what the weights will be

- Must be able to show $\langle 1 \rangle = \iint J dz d\eta = \text{Area}$
 exact w/ nodal gauss pts

- What about derivatives? change?

Hyperbolic PDEs ...

- easiest approach:

replace time derivatives w/ centered differences

$$\text{e.g. } \frac{\partial^2 u}{\partial t^2} \rightarrow \frac{u^{k+1} - 2u^k + u^{k-1}}{\Delta t^2}$$

$$u \rightarrow \theta \left(\frac{u^{k+1} + u^{k-1}}{2} \right) + (1-\theta) u^k$$

- Alternately as integration in time... both result in

$$[M] \left(\{u\}^{k+1} - 2\{u\}^k + \{u\}^{k-1} \right) + [\Delta t^2 K] \left(\frac{\theta}{2} \{u\}^{k+1} + \frac{\theta}{2} \{u\}^{k-1} + (1-\theta) \{u\}^k \right) = \{R\}^{k+1/2}$$

$$\underbrace{\left[M - \frac{\theta}{2} \Delta t^2 K \right]}_{[A]} \{u\}^{k+1} = \underbrace{\left[2M - \Delta t^2 (1-\theta) K \right]}_{[B]} \{u\}^k - \underbrace{\left[M + \frac{\Delta t^2 \theta}{2} K \right]}_{[C]} \{u\}^{k-1} + \{R\}^{k+1/2}$$

So:

$$[A]\{u\}^{k+1} = \underbrace{[B]\{u\}^k + [C]\{u\}^{k-1}}_{\text{old info... known!}} + \{R\}^{k+1/2}$$

- Need two-levels of u to proceed

Typically start system at rest

$$u^0 = u^{-1} = 0$$

Formally, must have $\frac{\partial u}{\partial t}(0)$ specified

How to enforce? ... use shadow node approach since are essentially using finite differences in the time dimension

$$\frac{\partial u}{\partial t}(0) = f(x, y) \Rightarrow \frac{u' - u^{-1}}{2\Delta t} = f(x, y)$$

$$u^{-1}(x, y) = -2\Delta t f(x, y) + u'(x, y)$$

$$\text{so } u_i^{-1} = u_i' - 2\Delta t f_i \Rightarrow \{u\}^{-1} = \{u\}' - 2\Delta t \{f\}$$

$$[A - C]\{u\}' = [B]\{u\}^0 - 2\Delta t [C]\{f\} + \{R\}^{k+1/2}$$