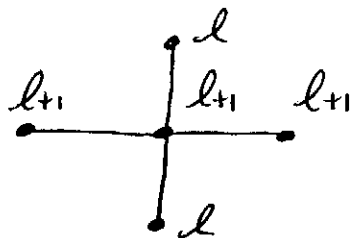


Block Iterative Methods

- "Point" methods... each calculation modifies sol'n at single point... u_{ij}^{l+1} can be computed by itself \Rightarrow "explicit"
- "Block" methods... groups of components of \underline{u}^{l+1} are computed simultaneously... involves sol'n of system of equations \Rightarrow individual components defined in terms of other components of the same Block \Rightarrow "Implicit"
- Basic tradeoff between block size and iteration count... Bigger the block \Rightarrow fewer iterations but cost per block increases... limiting cases
1 Block \Rightarrow 1 "iteration" \Rightarrow This is a direct sol'n !!
- Redefinition of explicit method as implicit one typically increases convergence rate
- Easy to do for Jacobi, Gauss-Seidel, SOR

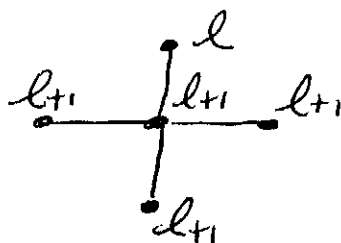
- "Single line" Jacobi



$$\beta_0 u_{ij}^{l+1} - \beta_1 u_{i+1,j}^{l+1} - \beta_2 u_{i-1,j}^{l+1} = \beta_3 u_{i,j+1}^l + \beta_4 u_{i,j-1}^l$$

Tridiagonal system \Rightarrow M eqns in M unknowns
for M nodes in "i" direction

- update sol'n line by line... but must solve system to get $u_{i,j}^{l+1}$ $i=1,2,\dots,M$ for each j
- "Single line" Gauss-Seidel (row wise ordering)



$$\beta_0 u_{ij}^{l+1} - \beta_1 u_{i+1,j}^{l+1} - \beta_2 u_{i-1,j}^{l+1} = \beta_3 u_{i,j+1}^l + \beta_4 u_{i,j-1}^{l+1}$$

SOR: $u_{ij}^{l+1} = \omega \bar{u}_{ij}^{l+1} + (1-\omega) u_{ij}^l$

↖ GS sol'n

$$U_{ij}^{l+1} - \omega \left[\frac{\beta_1}{\beta_0} U_{i+1,j}^{l+1} + \frac{\beta_2}{\beta_0} U_{i-1,j}^{l+1} \right] = (1-\omega) U_{ij}^l + \omega \left[\frac{\beta_3}{\beta_0} U_{i,j+1}^l + \frac{\beta_4}{\beta_0} U_{i,j-1}^{l+1} \right]$$

- Convergence Rates for Block Iterative Methods
... gain must compensate for extra effort to solve tridiagonal... problem dependent
- For single line blocks... Laplace on square
w/ $h = \pi/(N+1)$ i.e. $0 < x < \pi$
 $0 < y < \pi$

Jacobi	$\frac{\text{Point}}{h^2/2}$	$\frac{\text{Line}}{h^2}$
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G-S	h^2	$2h^2$
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SOR(optimal)	$2h$	$\sqrt{2}(2h)$
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$$\omega_{\text{opt}}^{\text{line}} = \frac{2}{1 + \sqrt{1 - \rho_{\text{line}}^2}}$$

- same as $\omega_{\text{opt}}^{\text{Point}}$

- replace ρ_{Point} w/ ρ_{Line}

- Can extend to more lines ... i.e. 2 line etc.

ADI: Alternating Direction Implicit

- Line methods discussed above ... always go in same direction ... Convergence often improved by changing directions!
- We think of a "single iteration" as a 2 step process \Rightarrow first implicit in row direction (i.e. x) followed by implicit in column direction (i.e. y)

eg. Difference form: $\delta_x^2 u_{ij} + \delta_y^2 u_{ij} + fh^2 u_{ij} = gh^2$

- δ_x^2 or δ_y^2 can be tridiagonal, but not simultaneously! $\Rightarrow \delta_x^2 + \delta_y^2 =$

$$\begin{array}{c} \text{Number consecutively across row} \\ \text{or} \\ \text{Number consecutively along a column} \end{array} \quad \delta_x^2 + \delta_y^2 = \begin{bmatrix} \diagup & & \\ & \diagdown & \\ & & \diagup \end{bmatrix} = \begin{bmatrix} \diagup & & \\ & \diagdown & \\ & & \diagup \end{bmatrix}_x + \begin{bmatrix} \diagup & & \\ & \diagdown & \\ & & \diagup \end{bmatrix}_y$$

(5)

- We take advantage of this by writing FD equation

$$\Delta_x^2 u_{ij} + \frac{fh^2}{2} u_{ij} = -\Delta_y^2 u_{ij} - \frac{fh^2}{2} u_{ij} + gh^2$$

- then solve for left side assuming right side is known,
- give ourselves ω tuning factor to speed up
- Repeat in the y direction writing

$$\Delta_y^2 u_{ij} + \frac{fh^2}{2} u_{ij} = -\Delta_x^2 u_{ij} - \frac{fh^2}{2} u_{ij} + gh^2$$

Always dealing w/ Tridiagonal... Complete iteration involves a sweep in both directions

e.g. Step 1: Implicit in x

$$\left(\Delta_x^2 + \frac{fh^2}{2}\right) u_{ij}^{l+1/2} - \omega u_{ij}^{l+1/2} = -\left(\Delta_y^2 + \frac{fh^2}{2}\right) u_{ij}^l - \omega u_{ij}^l + gh^2$$

$\omega \equiv$ Iteration parameter

$$\textcircled{1} - \textcircled{-2 - \omega + \frac{fh^2}{2}} - \textcircled{1} = \begin{array}{c} \textcircled{-1} \\ \textcircled{2 - \omega - \frac{fh^2}{2}} \\ \textcircled{-1} \end{array} + gh^2$$

Step 2: Implicit in y

$$\left(f_y^2 + \frac{fh^2}{2}\right)u_{ij}^{l+1} - \omega u_{ij}^{l+1} = -\left(f_x^2 + \frac{fh^2}{2}\right)u_{ij}^{l+1/2} - \omega u_{ij}^{l+1/2} + gh^2$$

$$\begin{array}{c} \textcircled{1} \\ | \\ \textcircled{-2 - \omega + \frac{fh^2}{2}} \\ | \\ \textcircled{1} \end{array} = \textcircled{-1} - \textcircled{2 - \omega - \frac{fh^2}{2}} - \textcircled{-1} + gh^2$$

- Two Step Procedure ... Solve Tridiagonal Systems
- ... Intermediate Results not considered as valid

Convergence Rate:

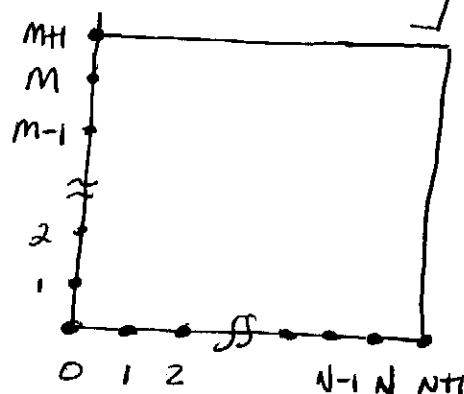
- Key is $\rho(G_{ADI})$... can cast this scheme in form $u^{l+1} = Gu^l + c$ then study spectrum of G
- Complicated but some results are known
 - Convergence only for $\omega > 0$

- Optimal ω for Laplace on square w/
Type I BCs

$$\omega_{opt}^2 = \left[\left(-\frac{fh^2}{2} + 4\sin^2 \frac{\pi}{2R} \right) \left(-\frac{fh^2}{2} + 4\cos^2 \frac{\pi}{2R} \right) \right]$$

$$R = \max(M+1, N+1)$$

\Rightarrow



- ADI convergence \approx SOR when using optimal ω for both
- Can improve significantly by generating a sequence of ω parameters \Rightarrow i.e. ω^L
- optimal sequence difficult, but some approaches for estimates are available