

①

Helmholtz Problem: $\nabla^2 u + k^2 u = 0$

$\nabla^2 G_i + k^2 G_i = -\delta(\underline{x} - \underline{x}_i)$

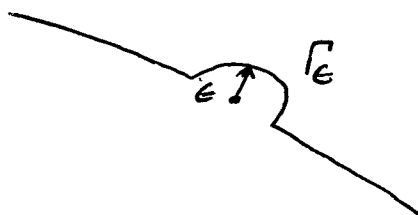
General Boundary expression:

$$\langle \nabla^2 u + k^2 u, G_i \rangle = \langle -\nabla u \cdot \nabla G_i \rangle + \langle k^2 u G_i \rangle + \oint \frac{\partial u}{\partial n} G_i ds$$

$$\langle \nabla^2 G_i + k^2 G_i, u \rangle = \langle \nabla G_i \cdot \nabla u \rangle + \langle k^2 G_i u \rangle + \oint \frac{\partial G_i}{\partial n} u ds$$

$u_i = \oint \frac{\partial u}{\partial n} G_i - \frac{\partial G_i}{\partial n} u ds$ as expected

2D: $G_i = \frac{i}{4} H_0^{(1)}(kr) = \left(J_0(kr) + i Y_0(kr) \right) \frac{i}{4}$



Look at: $\lim_{\epsilon \rightarrow 0} \int_{\Gamma_\epsilon} \frac{\partial u}{\partial n} G_i - \frac{\partial G_i}{\partial n} u ds$

$\Rightarrow G_i = \frac{i}{4} H_0'(k\epsilon)$, so $\lim_{\epsilon \rightarrow 0} \frac{i}{4} H_0'(k\epsilon) = \lim_{\epsilon \rightarrow 0} \frac{i}{4} \left(J_0'(k\epsilon) + i Y_0'(k\epsilon) \right)$
 $= \frac{i}{4} \left(1 + \frac{i2}{\pi} \ln(k\epsilon) \right)$

then

$$\lim_{\epsilon \rightarrow 0} \int_0^\pi \frac{\partial u}{\partial n} G_i \epsilon d\theta = \lim_{\epsilon \rightarrow 0} \frac{i}{4} \int_0^\pi \frac{\partial u}{\partial n} \left(1 + \frac{i2}{\pi} \ln(k\epsilon) \right) \epsilon d\theta = \underline{\underline{0}}$$

$$\frac{\partial G_i}{\partial n} = \frac{\partial G_i}{\partial \epsilon} \underbrace{\frac{\partial \epsilon}{\partial n}}_1 = \frac{i}{4} \left(\frac{i2}{\pi} \frac{1}{\epsilon} \right)$$

(B)

then

$$\lim_{\epsilon \rightarrow 0} \int_0^\pi u_i \frac{\partial G_i}{\partial \theta} \epsilon d\theta = \int_0^\pi -\left(\frac{1}{2\pi}\right) \frac{u_i}{\epsilon} \epsilon d\theta = -\frac{u_i}{2}$$

$$u_i = \oint \frac{2u}{2n} G_i - \frac{2G_i}{2n} u ds + \frac{u_i}{2} \Rightarrow \frac{u_i}{2} = \oint \frac{2u}{2n} G_i - \frac{2G_i}{2n} u ds$$

Also:



$$\lim_{\epsilon \rightarrow 0} \iint_{\Omega_\epsilon} \nabla \cdot \nabla G_i d\Omega_\epsilon + \iint_{\Omega_\epsilon} k^2 G_i d\Omega_\epsilon = -1$$

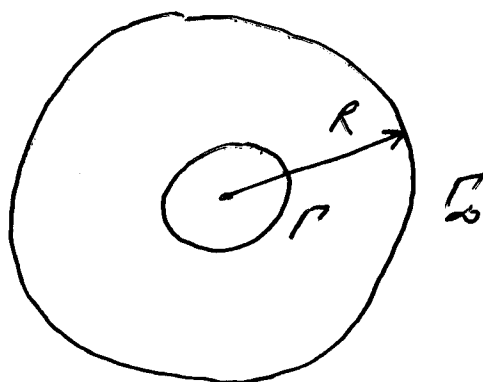
$$\lim_{\epsilon \rightarrow 0} \left[\int_{\Gamma_\epsilon} \frac{\partial G_i}{\partial n} d\Gamma_\epsilon = \int_0^{2\pi} \left(-\frac{1}{2\pi\epsilon}\right) \epsilon d\theta \right] = -1$$

$$\lim_{\epsilon \rightarrow 0} \left[\iint_{\Omega_\epsilon} k^2 G_i d\Omega_\epsilon = \frac{i}{k} \int_0^\epsilon \int_0^{2\pi} k^2 \left[\overset{\frac{i2}{\pi}}{\text{LN}(k\epsilon) + 1} \right] \epsilon d\theta \right] = 0$$

①

Infinite Domains

3D



$$\frac{d\psi_i}{4\pi} = \oint \left(\frac{\partial u}{\partial n} G_i - \frac{\partial G_i}{\partial n} u \right) ds$$

$$\text{Solid angle } \hat{n} = \int_{\Gamma} \left(\frac{\partial u}{\partial n} G_i - \frac{\partial G_i}{\partial n} u \right) d\Omega + \int_{\Gamma_{\infty}} \left(\frac{\partial u}{\partial n} G_i - \frac{\partial G_i}{\partial n} u \right) d\Omega$$

$ds = \hat{n} ds$
 $d\Omega = \frac{ds \cdot \hat{r}}{r^2} = \frac{ds \cdot \hat{r}}{r^3} \Rightarrow \Omega = \int \frac{\hat{r} \cdot ds}{r^3}$

Look at $\lim_{R \rightarrow \infty} \int_{\Gamma_{\infty}} \left(\frac{\partial u}{\partial n} G_i - \frac{\partial G_i}{\partial n} u \right) d\Omega$

3D Helmholtz: $G_i = \frac{e^{ikr}}{4\pi r}$

$$\begin{aligned} \frac{\partial G_i}{\partial n} &= \left(-\frac{e^{ikr}}{4\pi r^2} + \frac{ike^{ikr}}{4\pi r} \right) \hat{r} \cdot \hat{n} \\ &= \frac{e^{ikr}}{4\pi r} \left(ik - \frac{1}{r} \right) \hat{r} \cdot \hat{n} \end{aligned}$$

$$\begin{aligned} \text{So } \int_{\Gamma_{\infty}} \left(\frac{\partial u}{\partial n} G_i - \frac{\partial G_i}{\partial n} u \right) d\Omega &= \int_0^{2\pi} \int_0^{\pi} \left\{ \frac{\partial u}{\partial R} \left(\frac{e^{ikR}}{4\pi R} \right) - \frac{e^{ikR}}{4\pi R} \left(ik - \frac{1}{R} \right) u \right\} R^2 \sin\theta d\theta d\phi \\ &= \int_0^{2\pi} \int_0^{\pi} \frac{e^{ikR}}{4\pi} \left\{ \left(\frac{\partial u}{\partial R} R - ik u R \right) - u \right\} \sin\theta d\theta d\phi \end{aligned}$$

⑦

$$\lim_{R \rightarrow \infty} \iint \frac{e^{ikr}}{4\pi} \left\{ \left(\frac{2u}{2R} - iku \right) R - u \right\} \sin\theta \, d\theta \, d\phi$$

requires $\underbrace{\lim_{R \rightarrow \infty} R \left(\frac{2u}{2R} - iku \right)} = 0 \quad \text{and} \quad \lim_{R \rightarrow \infty} u = 0$

Sommerfeld Radiation Condition

excludes incoming waves and guarantees unique solution to problem ... is satisfied by the Green's function (unbounded space)

e.g. $\lim_{R \rightarrow \infty} R \left(\frac{2G}{2R} - iku \right) = 0 \Rightarrow \lim_{R \rightarrow \infty} -\frac{e^{ikr}}{4\pi R} = 0$

\uparrow \uparrow
 $-\frac{e^{ikr}}{4\pi R^2} + \frac{iku e^{ikr}}{4\pi R}$

Note, in 2D ... Sommerfeld condition is

$$\lim_{R \rightarrow \infty} \sqrt{R} \left(\frac{2u}{2R} - iku \right) = 0$$