

PDE Solution via Spectral Methods

Example: Linear eqn w/ variable coefficient

e.g. $\frac{\partial u}{\partial t} + c(x) \frac{\partial u}{\partial x} = 0$ Periodic BCs
known ICs

$$u = \sum_{k=0}^{N-1} a_k \phi_k(x)$$

Galerkin: $\left\langle \left(\sum_{k=0}^{N-1} \frac{da_k}{dt} \phi_k + c(x) \sum_{k=0}^{N-1} a_k \frac{\partial \phi_k}{\partial x} \right) \phi_m \right\rangle = 0$

$$\sum_{k=0}^{N-1} \frac{da_k}{dt} \langle \phi_k \phi_m \rangle + \left\langle c(x) \frac{\partial \phi_k}{\partial x} \phi_m \right\rangle \sum_{k=0}^{N-1} a_k = 0 \quad m=0, 1, \dots, N-1$$

$$\frac{da_m}{dt} = - \sum_{k=0}^{N-1} a_k \left\langle c(x) \frac{\partial \phi_k}{\partial x} \phi_m \right\rangle \quad m=0, 1, \dots, N-1$$

Center Difference in Time:

$$a_m^{l+1} = a_m^{l-1} - 2\Delta t \underbrace{\sum_{k=0}^{N-1} a_k^l \left\langle c(x) \frac{\partial \phi_k}{\partial x} \phi_m \right\rangle}_{\text{operations/time-step}}$$

$\mathcal{O}(N^2)$ operations/time-step
involves all a_k 's for single
advance of one a_m

Compare w/ FD: $\mathcal{O}(N)$ operations/step

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Note: if $c(x) = c = \text{constant} \dots$ extra overhead goes away ... e.g. Fourier $\Rightarrow \phi_k = e^{ikx}$, $[0, 2\pi]$

$$2\Delta t \sum_{k=0}^{N-1} a_k^l \left\langle c(x) \frac{\partial \phi_k}{\partial x} \phi_m \right\rangle = 2\Delta t c \sum_{k=0}^{N-1} a_k^l \left\langle i k e^{ikx} e^{-imx} \right\rangle$$

$$= 2\Delta t c i m a_m^l$$

$$a_m^{l+1} = a_m^{l-1} - \underbrace{2\Delta t c i m a_m^l}_{O(N) \text{ effort just like FD}} \quad m=0, 1, \dots, N-1$$

Conclude: $c(x)$ "messes-up" orthogonality integrals causing the method to be impractical

What if we use our interpolating polynomial approach with collocation??

$$u(x) = \sum_{k=0}^{N-1} u_k \phi_k(x) \Rightarrow u_j = u(x_j)$$

$$\frac{du_j}{dt} + c(x_j) \sum_{k=0}^{N-1} u_k \frac{\partial \phi_k}{\partial x}(x_j) = 0 \quad j=0, 1, \dots, N-1$$

$$u_j^{l+1} = u_j^{l-1} - 2\Delta t c_j \underbrace{\sum_{k=0}^{N-1} u_k \frac{\partial \phi_k}{\partial x}(x_j)}_{O(N^2) / \text{Time-step}} \quad j=0, 1, \dots, N-1$$

$O(N^2) / \text{Time-step} \Rightarrow$ same problem essentially an N -pt approximation to $\frac{\partial u_j}{\partial x}$

Conclude: Neither Galerkin nor Collocation is very efficient (gain in spatial accuracy offset by computational costs!)

In Galerkin ... orthogonality lost due to $C(x)$, but no difficulty with derivative perse (e.g. derivative of trig function is another trig function)

In Collocation ... derivative evaluation involves an N -pt formula (all grid values of u) but no difficulty with $C(x)$... i.e. that coefficient varies with position

would like to combine the best features of both

$$\text{i.e. } u_j^{l+1} = u_j^{l-1} - 2\Delta t c_j \frac{\partial u^l}{\partial x} \quad j=0, 1, \dots, N-1$$

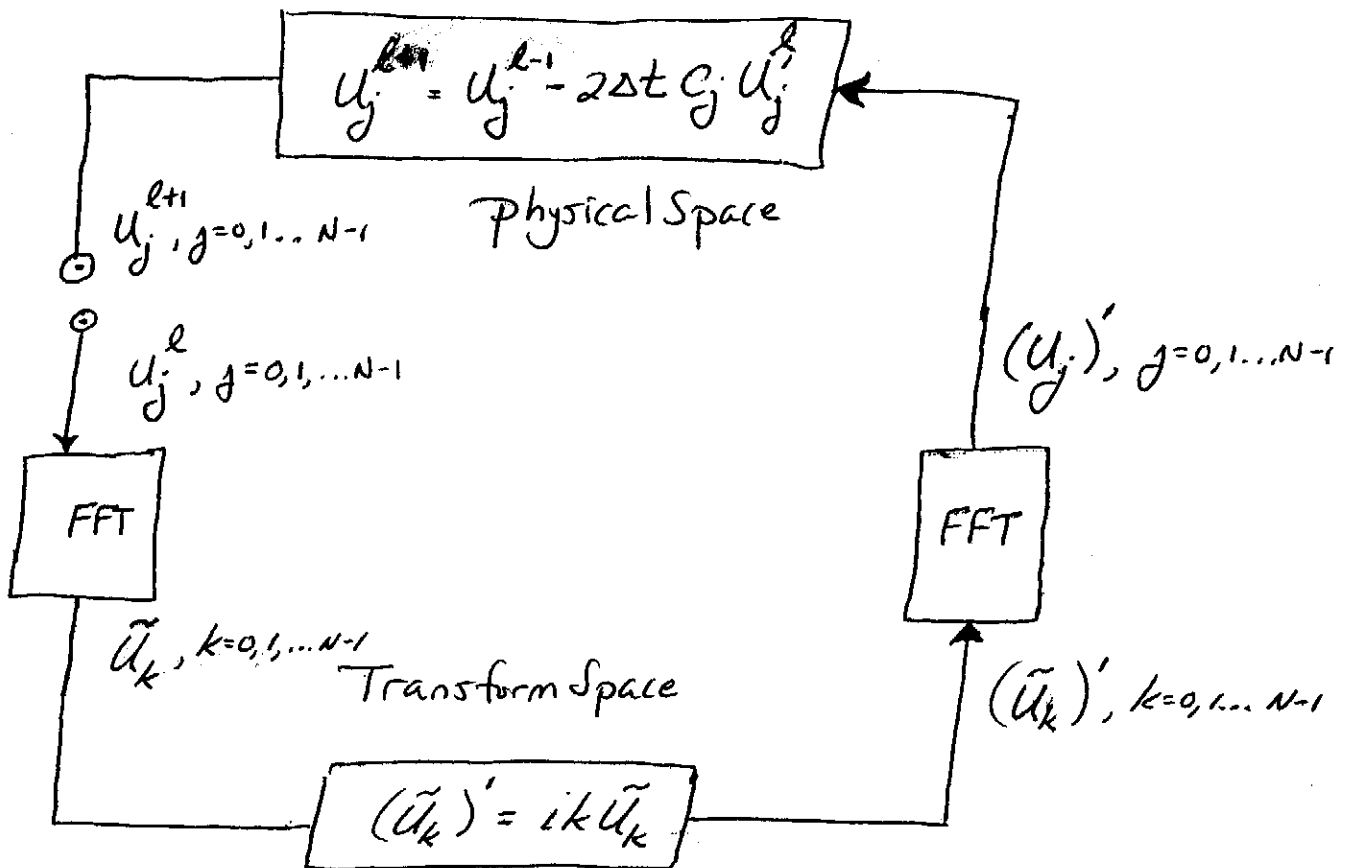
where this can be evaluated efficiently ... hopefully $O(N)/\text{step}$

Can accomplish this by appealing to our collocation derivative concepts but use a fast transform technique to compute $\Rightarrow O(N \log N)$

"pseudospectral method"

So we replace rate-limiting matrix multiply with FFT \Rightarrow limited to Fourier and Chebyshev bases

Algorithm becomes something like:



We "jump back and forth" between the grid point representation of $u(x)$ (in order to multiply by $c(x)$) and the spectral coefficients (in order to differentiate $u(x)$)

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One potential problem with this approach is "aliasing" errors

- IF computational grid supports wavenumbers

$$k \in \left[-\frac{N}{2}, \frac{N}{2}\right]$$

then higher wavenumbers (outside the basis set) appear in the computations as if they were wavenumbers

$$k_A = k \pm mN \quad \text{where } k_A \in \left[-\frac{N}{2}, \frac{N}{2}\right]$$

i.e. wavenumbers outside the truncation are aliased to wavenumbers within

- This is a re-statement of the error occurring from interpolation rather than truncation

$$\text{Recall: } \tilde{u}_k = \hat{u}_k + \sum_{\substack{n=-\infty \\ m \neq 0}}^{\infty} \hat{u}_{k \pm mN}$$

$$\text{e.g. } \cos((k+mN)x) = \cos(kx)\cos(mNx) - \sin(kx)\sin(mNx)$$

$$\sin((k+mN)x) = \sin(kx)\cos(mNx) + \cos(kx)\sin(mNx)$$

$$\text{Now if } x_i = \frac{2\pi i}{N}, \quad i=0, 1, \dots, N-1$$

$$\left. \begin{aligned} \cos(mNx_i) &= \cos\left(mN \frac{2\pi i}{N}\right) = \cos(2\pi mi) = 1 \\ \sin(mNx_i) &= \sin(2\pi mi) = 0 \end{aligned} \right\} \begin{array}{l} \text{all } m, i \\ \text{integer} \end{array}$$

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$$\text{So } \cos(kx) = \cos((k+mN)x) = \cos((-k+mN)x)$$

$$\sin(kx) = \sin((k+mN)x) = -\sin((-k+mN)x)$$

at the N grid points ... so we have linearly independent functions that are point-by-point equal on the grid \Rightarrow aliasing error

- Can generate higher wavenumbers outside the basis set truncation through variable coefficients, non linear terms, etc.; these can lead to numerical instabilities (or inaccuracies)

$$\begin{aligned} \text{e.g. } u \frac{\partial u}{\partial x} &= \left(\sum_{l=-\frac{N}{2}}^{\frac{N}{2}-1} \tilde{u}_l e^{ilx} \right) \left(\sum_{m=-\frac{N}{2}}^{\frac{N}{2}-1} im \tilde{u}_m e^{imx} \right) \\ &= \sum_{l=-\frac{N}{2}}^{\frac{N}{2}-1} \sum_{m=-\frac{N}{2}}^{\frac{N}{2}-1} im \tilde{u}_l \tilde{u}_m e^{i(l+m)x} \\ &= \sum_{k=-N}^{N-1} a_k e^{ikx} \end{aligned}$$

Non linear interaction produces wavenumbers on the range $k \in [-N, N]$ which will get aliased onto $k \in [-\frac{N}{2}, \frac{N}{2}]$

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- Recall... highest frequency on a discrete mesh
"2Δx"-waves ()

Also recall Nyquist Sampling Criterion...
need to sample at twice the highest frequency
for perfect reconstruction

- Obvious remedy... apply spatial filter
so as to eliminate all waves with
wavelengths less than 4Δx ... i.e.
eliminate/remove all wavenumbers with
 $|k| > \frac{N}{4} \Rightarrow \frac{N}{4} < |k| \leq \frac{N}{2}$

Turns out can do better than this... obvious
remedy is wasteful

"3/2 Rule": accurate, alias-free computation
of a total of N wavenumbers for a quadratically
nonlinear interaction, must use $\frac{3}{2} N$ samples

Also known as "padding" because we add $N/2$ zeros to the spectral coefficients before taking coefficient to grid transforms

- Alternately ... if use a total of N grid points, then only $2/3 N$ wavenumbers should be retained in the truncation \Rightarrow only necessary to filter waves w/ wavelengths between $2\Delta x$ and $3\Delta x$ (not $4\Delta x$) \Rightarrow "2/3 Rule"

Example: $U \cdot V$

Transform \tilde{U}_k, \tilde{V}_k on N_{pts} to U_j, V_j on M_{pts}

Product $U_j V_j$ on M_{pts}

Transform Back $(\tilde{U}\tilde{V})_k$ on M_{pts}

Truncate $(\tilde{U}\tilde{V})_k$ to N (padding)

$$y_j = \frac{2\pi j}{M} \quad j=0, 1, \dots, M-1$$

$$U_j = \sum_{k=-M/2}^{M/2-1} \tilde{U}_k e^{iky_j}, \quad V_j = \sum_{k=-M/2}^{M/2-1} \tilde{V}_k e^{iky_j}$$

$$\text{where } \tilde{U}_k = \begin{cases} \tilde{U}_k & k \leq \frac{N}{2}-1 \\ 0 & \text{otherwise} \end{cases}$$

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$$\tilde{V}_k = \begin{cases} \tilde{V}_k & k \leq \lfloor \frac{N}{2} \rfloor \\ 0 & \text{otherwise} \end{cases}$$

then $W_j = U_j - V_j \quad j=0, 1, \dots, M-1$

and $\tilde{W}_k = \frac{1}{M} \sum_{j=0}^{M-1} W_j e^{-iky_j} \quad k = -\frac{M}{2}, \dots, \frac{M}{2}-1$

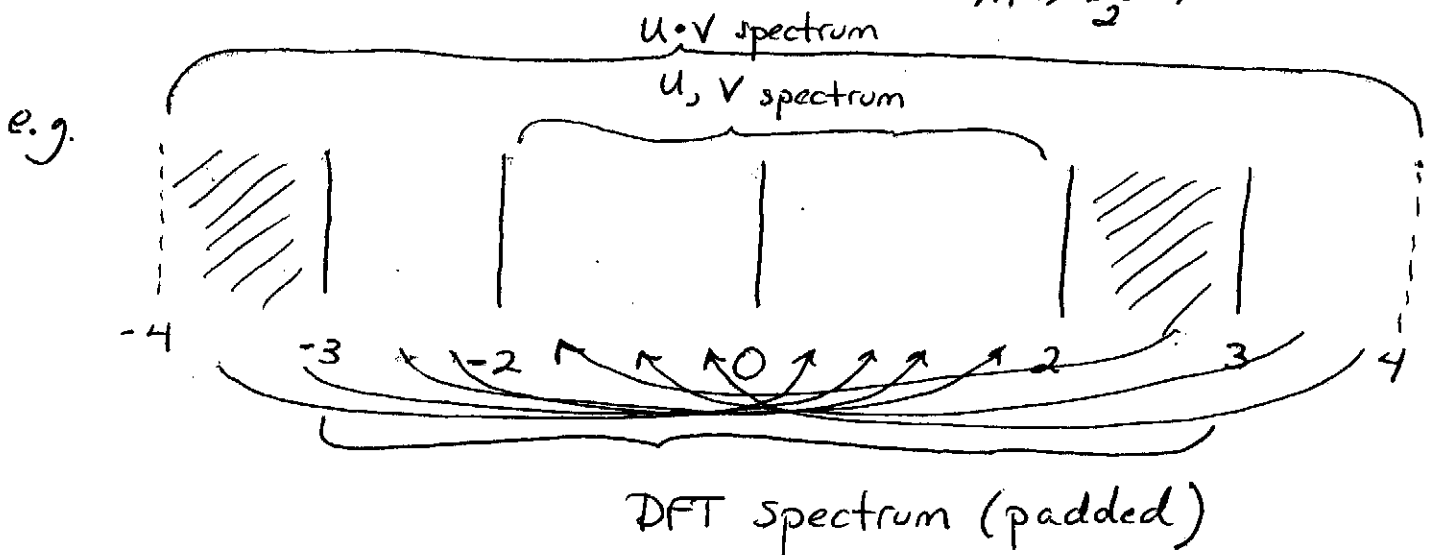
But $\tilde{W}_k = \hat{W}_k + \underbrace{\sum_{m+n=k+M} \hat{U}_m \hat{V}_n}_{\text{Choose } M \text{ so that this term vanishes for } k \in [-\frac{N}{2}, \frac{N}{2}]}$

Choose M so that this term vanishes for $k \in [-\frac{N}{2}, \frac{N}{2}]$

i.e. when $|m|/M > N/2$

\Rightarrow worst case $-\frac{N}{2} - \frac{N}{2} \leq \frac{N}{2} - 1 - M$

$M \geq \frac{3N}{2} - 1$



DFT (aliased) get contamination from the U-V spectrum outside $[-2, 2]$ on $[-2, 2]$

Another Strategy: Dealiasing by phase shift

Aliased: $\tilde{u}_k \Rightarrow u_j \xrightarrow{\quad} w_j \Rightarrow \tilde{w}_k = \sum_{m+n=k} \tilde{u}_m \tilde{v}_m + \sum_{m+n=k \pm N} \tilde{u}_m \tilde{v}_m$

$\tilde{v}_k \Rightarrow v_j \xrightarrow{\quad}$

Correct Error

Shifted:

"Shifted Transforms"

$$u'_j = \sum_{k=-\frac{N}{2}}^{\frac{N}{2}-1} \tilde{u}_k e^{ik(x_j + \Delta)} = \sum (\tilde{u}_k e^{ik\Delta}) e^{ikx_j}$$

$$\tilde{u}'_k = \frac{1}{N} \sum_{j=0}^{N-1} u'_j e^{-ik(x_j + \Delta)} = e^{-ik\Delta} \frac{1}{N} \sum_{j=0}^{N-1} u'_j e^{-ikx_j}$$

$$\Rightarrow e^{ik\Delta} \tilde{u}'_k = \frac{1}{N} \sum_{j=0}^{N-1} u'_j e^{-ikx_j}$$

$$\tilde{u}_k \xRightarrow{\Delta} u'_j \xrightarrow{\quad} w'_j \xRightarrow{\Delta} \tilde{w}'_k = \sum_{m+n=k} \tilde{u}_m \tilde{v}_m + e^{\pm iN\Delta} \sum_{m+n=k \pm N} \tilde{u}_m \tilde{v}_m$$

$$\tilde{v}_k \xRightarrow{\Delta} v'_j \xrightarrow{\quad}$$

(i.e. $N\Delta = \pi$)
 $e^{\pm i\pi} = -1$

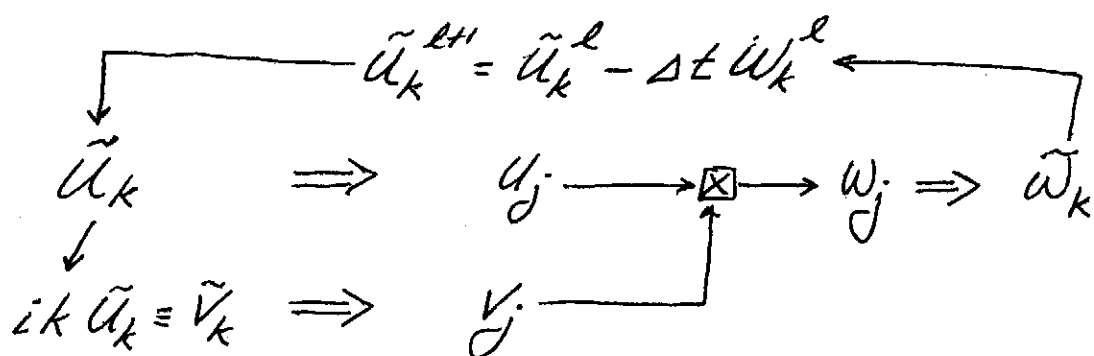
dealiased $\tilde{w}_k = \frac{\tilde{w}_k + \tilde{w}'_k}{2}$

Cost turns out to be greater than padding technique
 so not used that often

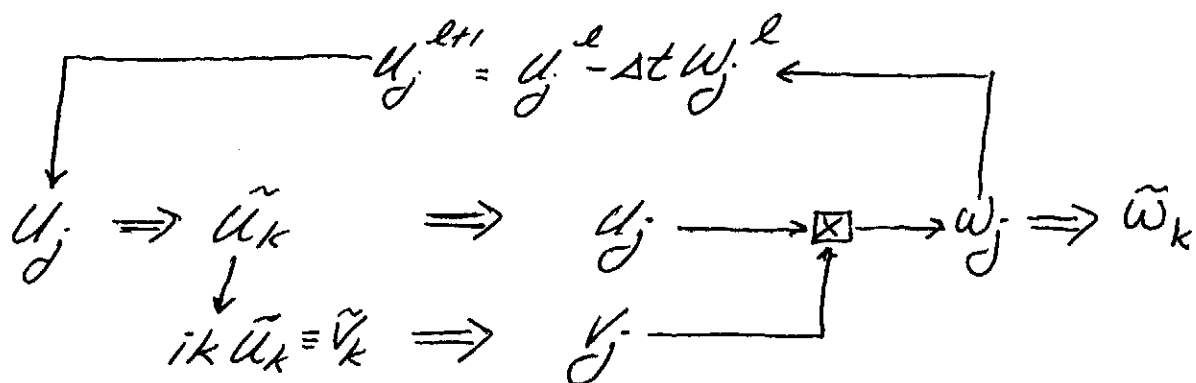
Rule of thumb: dealiasing not really necessary
for well-resolved simulations; may be helpful
when simulation is marginally resolved
(this can become important in large-scale problems)

- Look at some possible strategies (and their relationship)

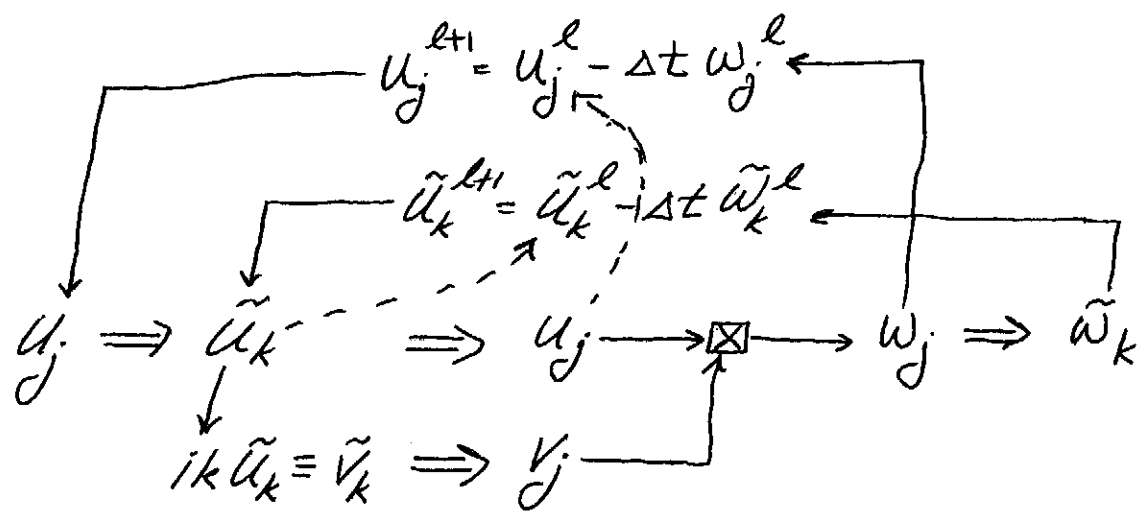
for: $\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = 0$



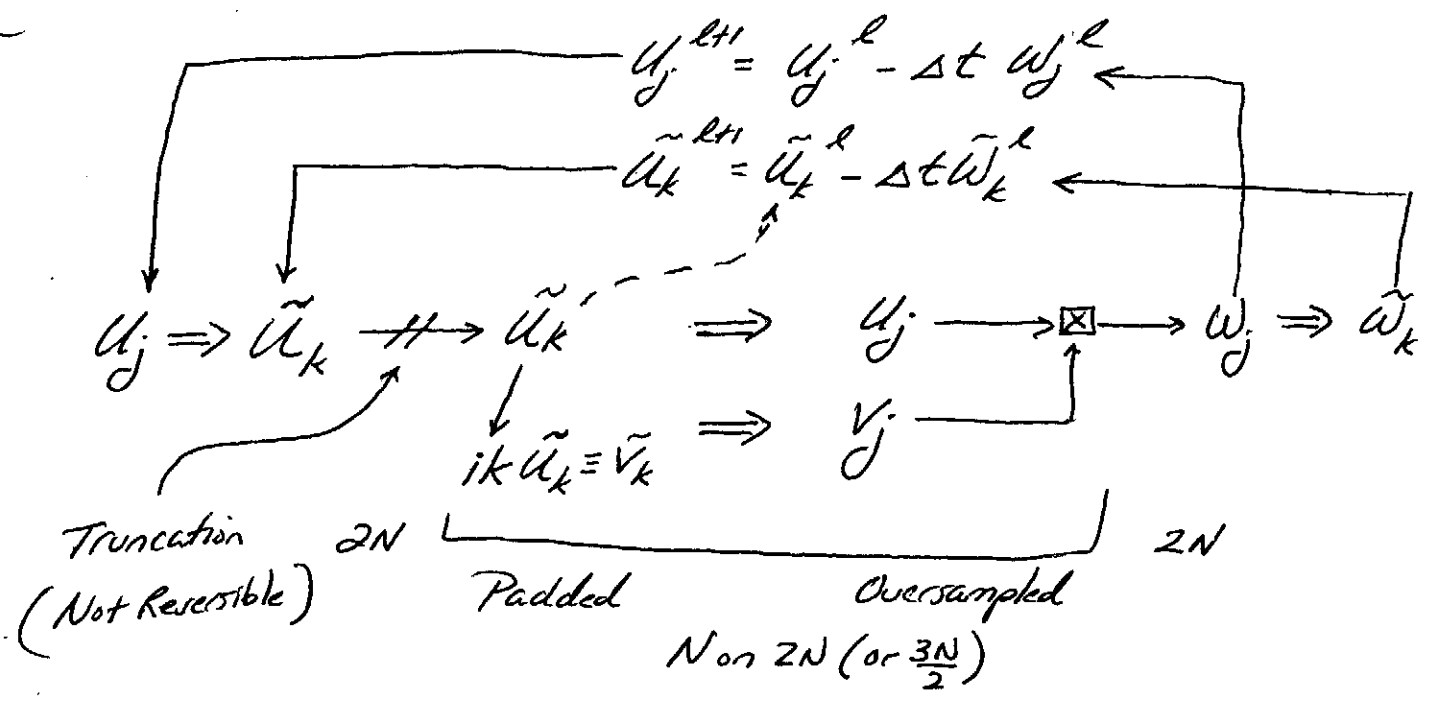
"Pseudospectral"



"Collocation"



Collocation = Pseudospectral (aliased)
(reversible!!)



Pseudospectral with dealiasing
(collocation)

Rules of Thumb on Accuracy:

To achieve 5% accuracy for wave-like sol'n

20 samples/wavelength FD $O(\Delta x^2)$
 (40 samples/wavelength \Rightarrow 1% error)

10 samples/wavelength FD $O(\Delta x^4)$
 (15 samples/wavelength \Rightarrow 1% error)

3.5 polys/wavelength Chebyshev
 pseudospectral
 (3.5 polys/wavelength \Rightarrow 1% error)

if Δx is smallest scale to resolve on $x \in [-1, 1]$
 then an estimate of number of wavelengths

$$M \approx \frac{1}{\pi \Delta x}$$

Errors are based on smooth solutions (i.e.
 no discontinuities on the interior or boundary)

$\Rightarrow N \approx 3/\epsilon$ polynomials for 1% accuracy
 for boundary layer width ϵ at $x = \pm 1$

Simple Time-stepping

e.g. $\frac{d\hat{u}_k}{dt} + \nu k^2 \hat{u}_k = 0 \quad -\frac{N}{2} \leq k \leq \frac{N}{2}-1$
 $\Delta x = \frac{2\pi}{N}$

Euler: $\frac{\hat{u}_k^{l+1} - \hat{u}_k^l}{\Delta t} = -\nu k^2 \hat{u}_k^l$

let $\hat{u}_k^l = \lambda \hat{u}_k^{l-1}$

$$\frac{\lambda - 1}{\Delta t} = -\nu k^2 \Rightarrow \lambda = 1 - \nu k^2 \Delta t$$

want $-1 \leq \lambda \leq 1$ for stability

$$1 - \nu k^2 \Delta t \geq -1$$

$$2 \geq \nu k^2 \Delta t \Rightarrow \text{worst case } k^2 = \left(\frac{N}{2}\right)^2 = \left(\frac{\pi}{\Delta x}\right)^2$$

$$2 \geq \frac{\nu \pi^2 \Delta t}{\Delta x^2}$$

so $\boxed{\frac{\nu \Delta t}{\Delta x^2} \leq \frac{2}{\pi^2}}$ (compare w/ FD: $\frac{\nu \Delta t}{\Delta x^2} \leq \frac{1}{2}$)

2-level Implicit

$$\frac{\hat{u}_k^{l+1} - \hat{u}_k^l}{\Delta t} = -vk^2 (\theta \hat{u}_k^{l+1} + (1-\theta) \hat{u}_k^l)$$

$$\lambda - 1 = -vk^2 \Delta t (\theta \lambda + (1-\theta))$$

$$\lambda [1 + \theta vk^2 \Delta t] = 1 - (1-\theta) vk^2 \Delta t$$

$$\lambda = \frac{1 - (1-\theta) vk^2 \Delta t}{1 + \theta vk^2 \Delta t} \Rightarrow \text{always } < 1$$

can become < -1
(unstable)

$\theta > 1/2$ unconditionally stable!

Conditionally stable if: $\frac{1 - (1-\theta) vk^2 \Delta t}{1 + \theta vk^2 \Delta t} > -1$

$$1 - (1-\theta) vk^2 \Delta t > -1 - \theta vk^2 \Delta t$$

$$2 > (1-2\theta) vk^2 \Delta t$$

$$\frac{2}{1-2\theta} > vk^2 \Delta t \Rightarrow \underline{\underline{\frac{v \Delta t}{\Delta x^2} < \frac{2}{\pi^2 (1-2\theta)}}}$$

①

Prototype Problem: Burger's Equation

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} - v \frac{\partial^2 u}{\partial x^2} = 0$$

$$0 \leq x \leq 2\pi$$

IC's given
BC's given

$v = \text{positive constant}$

Fourier Galerkin:

$$u^N(x, t) = \sum_{k=-\frac{N}{2}}^{\frac{N}{2}-1} \hat{u}_k(t) e^{ikx} \quad \text{Fourier}$$

$$\int_0^{2\pi} \left(\frac{\partial u^N}{\partial t} + u^N \frac{\partial u^N}{\partial x} - v \frac{\partial^2 u^N}{\partial x^2} \right) e^{-ikx} dx = 0 \quad \text{Galerkin}$$

$$\frac{\partial \hat{u}_k}{\partial t} + \widehat{\left(u \frac{\partial u}{\partial x} \right)}_k + k^2 v \hat{u}_k = 0, \quad k = -\frac{N}{2} \dots \frac{N}{2} - 1$$

$$\text{Product Term: } \widehat{(uv)}_k = \frac{1}{2\pi} \int_0^{2\pi} uv e^{-ikx} dx$$

$$= \sum_{p+q=k} \hat{u}_p \hat{v}_q \quad \text{Convolution Sum}$$

$$\left(\frac{1}{2\pi} \int_0^{2\pi} \underbrace{\left(\sum_p \hat{u}_p e^{ipx} \right) \left(\sum_q \hat{v}_q e^{iqx} \right)}_{\sum \sum \hat{u}_p \hat{v}_q e^{i(p+q)x}} e^{-ikx} dx \right)$$

Implied creation
of info outside
the basis

Fourier Collocation:

- carried out in x -domain
- retain values of u at $x_j = \frac{2\pi j}{N}$, $j=0, 1, \dots, N-1$

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} - v \frac{\partial^2 u}{\partial x^2} = 0 \quad j=0, 1, \dots, N-1$$

$$u^N D_N^2 u^N$$

Fourier Collocation Derivative

$$D_N^2 u^N \quad (\text{Full matrix})$$

Pseudospectral Transform

$$\left. \begin{aligned} w(x) &= u(x)v(x) \\ \hat{w}_k &= \sum_{m+n=k} \hat{u}_m \hat{v}_n \end{aligned} \right\} \text{Convolution}$$

$$\text{or } u_j = \sum_{k=-N/2}^{N/2-1} \hat{u}_k e^{ikx_j}$$

$$v_j = \sum_{k=-N/2}^{N/2-1} \hat{v}_k e^{ikx_j}$$

$$\Rightarrow w_j = u_j v_j$$

$$\hat{w}_k = \frac{1}{N} \sum_{j=0}^{N-1} w_j e^{-ikx_j}$$

$$k = -\frac{N}{2}, \dots, \frac{N}{2}-1$$

produces aliasing error:

$$\hat{w}_k = \hat{w}_k + \underbrace{\sum_{m+n=k \pm N} \hat{u}_m \hat{v}_n}_{\text{Error}}$$

So we have (pseudospectral)

$$\frac{d\hat{u}_k}{dt} + \sum_{m+n=k} \hat{u}_m \hat{v}_n + \sum_{m+n=k \pm N} \hat{u}_m \hat{v}_n + \nu k^2 \hat{u}_k = 0 \quad k = -\frac{N}{2}, \dots, \frac{N}{2}-1$$

Same as Collocation under DFT

$$\left(\frac{\partial u^N}{\partial t} + u^N v^N - \nu \frac{\partial^2 u^N}{\partial x^2} \right) \Big|_{x=x_j} = 0 \quad \text{all } x_j$$

$$\frac{d\hat{u}_k}{dt} + \sum_{m+n=k} \hat{u}_m \hat{v}_n + \underbrace{\sum_{m+n=k \pm N} \hat{u}_m \hat{v}_n}_{\text{eliminate aliasing error}} + \nu k^2 \hat{u}_k = 0 \quad \text{same!!}$$

IF can dealias Collocation, then have Galerkin

