

# Assignment 5

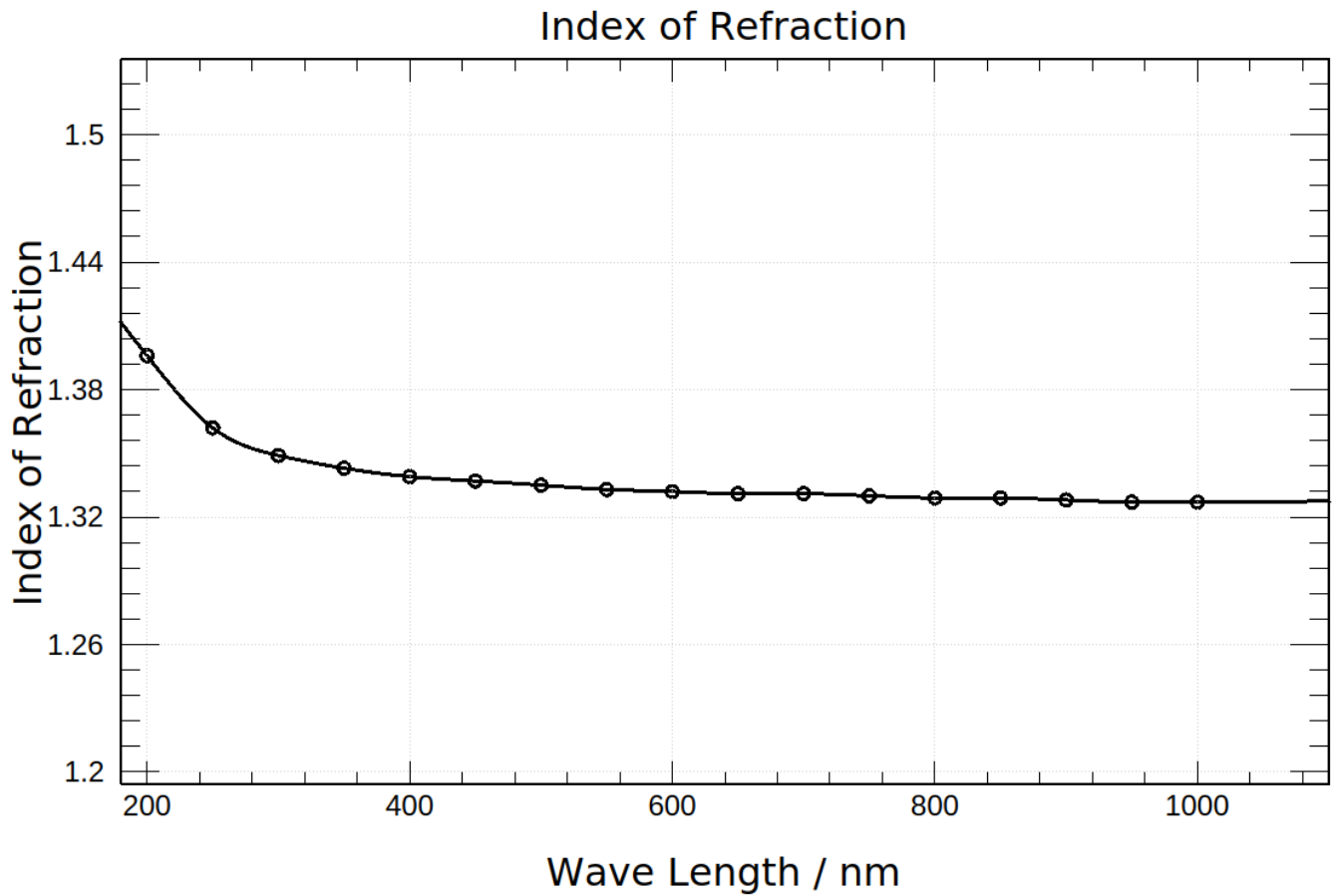
## Shuyang Cao

- [Assignment 5](#)
  - [Shuyang Cao](#)
    - [Chapter 4 Exercise 7](#)
    - [Chapter 5 Exercise 7](#)
    - [Chapter 5 Exercise 9](#)
    - [Chapter 5 Exercise 12](#)
      - [Part a](#)
      - [Part b](#)
      - [Part c](#)

## Chapter 4 Exercise 7

Cubic spline interpolation is used.

```
cat ./data.txt | ./dipersion
```



## Chapter 5 Exercise 7

Use  $M(\hat{A})$  to denote the matrix of operator  $\hat{A}$  and  $\hat{F}(\hat{A})$  to denote the function of operator  $\hat{A}$ . Then, we have

$$M(\hat{F}(\hat{A})) = F(M(\hat{A}))$$

So, for  $\hat{H}$ ,  $\hat{x}$ ,  $\hat{D}$ , we have

$$\begin{aligned} M(\hat{x}^2) &= M(\hat{x})^2 \\ M(\hat{D}^2) &= M(\hat{D})^2 \\ M(\hat{H}) &= -\frac{1}{2}M(\hat{D})^2 + \frac{1}{2}M(\hat{x})^2 \end{aligned}$$

Note that the basis we use the the eigen basis of  $\hat{H}$ . Thus, the matrix of  $\hat{H}$  is diagonal. A value whose magnitude is less than 1e-12 is printed as 0.

\$ ./operator

x

0	0.707107	0	0	0
0.707107	0	1	0	0
0	1	0	1.22474	0
0	0	1.22474	0	1.41421
0	0	0	1.41421	0

x^2

0.5	0	0.707107	0	0
0	1.5	0	1.22474	0
0.707107	0	2.5	0	1.73205
0	1.22474	0	3.5	0
0	0	1.73205	0	4.5

D

0	0.707107	0	0	0
-0.707107	0	1	0	0
0	-1	0	1.22474	0
0	0	-1.22474	0	1.41421
0	0	0	-1.41421	0

D2

-0.5	0	0.707107	0	0
0	-1.5	0	1.22474	0
0.707107	0	-2.5	0	1.73205
0	1.22474	0	-3.5	0
0	0	1.73205	0	-4.5

Hamiltonian

0.5	0	0	0	0
0	1.5	0	0	0
0	0	2.5	0	0
0	0	0	3.5	0
0	0	0	0	4.5

## Chapter 5 Exercise 9

Effective potential is

$$V(r) = -\alpha \frac{4\hbar c}{3r} + \frac{r}{\hbar c a^2} + \frac{l(l+1)\hbar^2}{2\mu r^2}$$

where  $\mu = m_c/2$  is the reduced mass. So, the quantization rule is

$$\int_{r_{\min}}^{r_{\max}} \sqrt{2\mu(E - V(r))} dr = (n + \eta) \pi \hbar$$

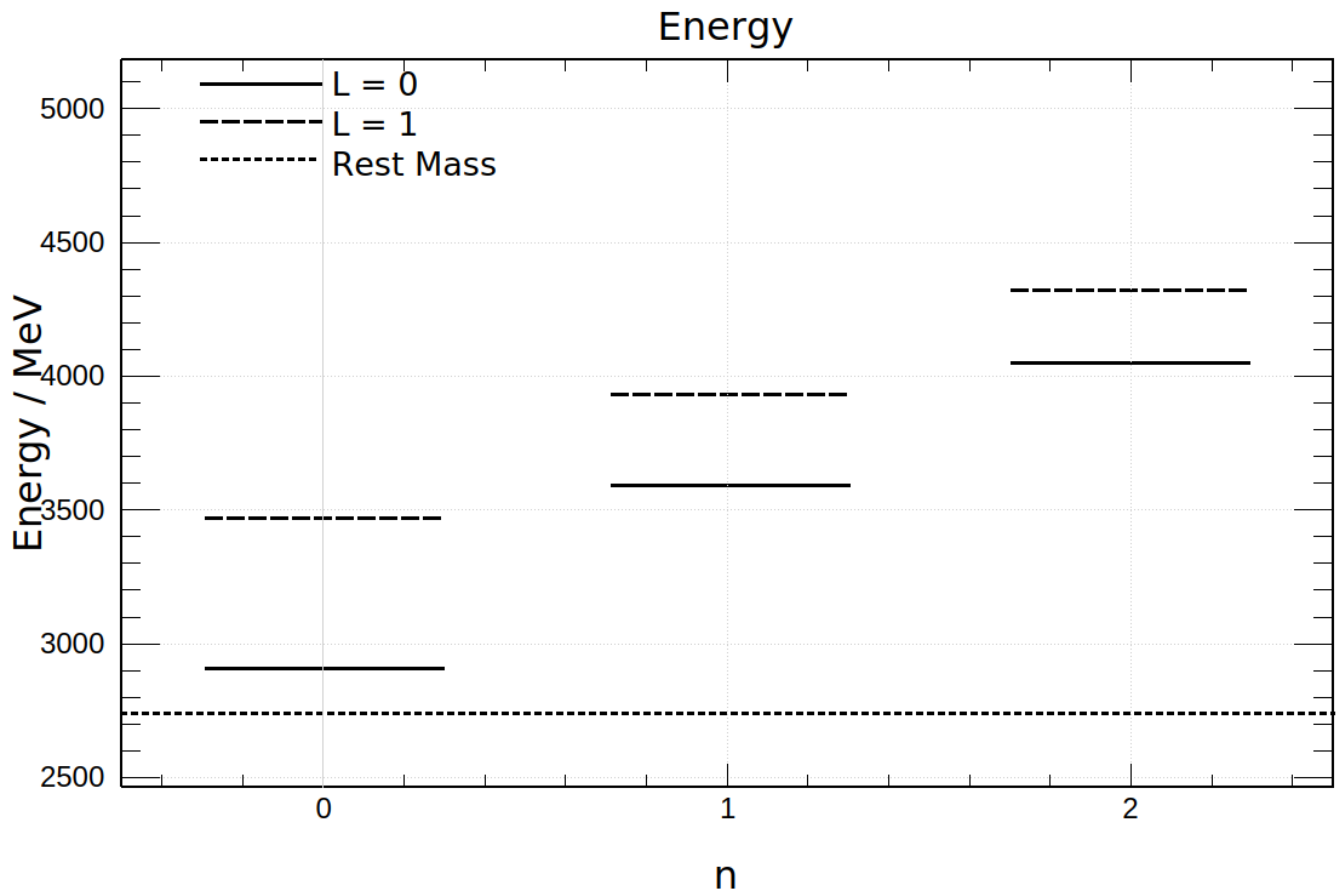
$$\Rightarrow \int_{r_{\min}}^{r_{\max}} \sqrt{\frac{2\mu E}{\hbar^2} + \frac{8\alpha\mu c}{3\hbar r} - \frac{2\mu r}{\hbar^3 c a^2} - \frac{l(l+1)}{r^2}} dr = (n + \eta) \pi$$

where  $\eta = 3/4$  if  $l = 0$ ,  $\eta = 1/3$  if  $l = 1$  because classically particles cannot reach the center of the system if the angular momentum is nonzero. Choose  $b = 10^{-15} m$  as the length unit. The reduced position  $\xi = x/a$ . The quantization rule can be rewritten as

$$\int_{\xi_{\min}}^{\xi_{\max}} \sqrt{\epsilon + \frac{\beta}{\xi} - \gamma\xi - \frac{l(l+1)}{\xi^2}} d\xi = (n + \eta) \pi$$

where

$$\left\{ \begin{array}{l} \epsilon = \frac{E}{E_0} \\ E_0 = \frac{\hbar^2}{2\mu b^2} \approx 28.4219 \text{ MeV} \\ \beta = \frac{8\alpha\mu c b}{3\hbar} \approx 3.51768 \\ \gamma = \frac{2\mu b^3}{\hbar^3 c a^2} \approx 30.1959 \end{array} \right.$$



## Chapter 5 Exercise 12

Suppose we wait for a long enough time so that all N atoms decay.

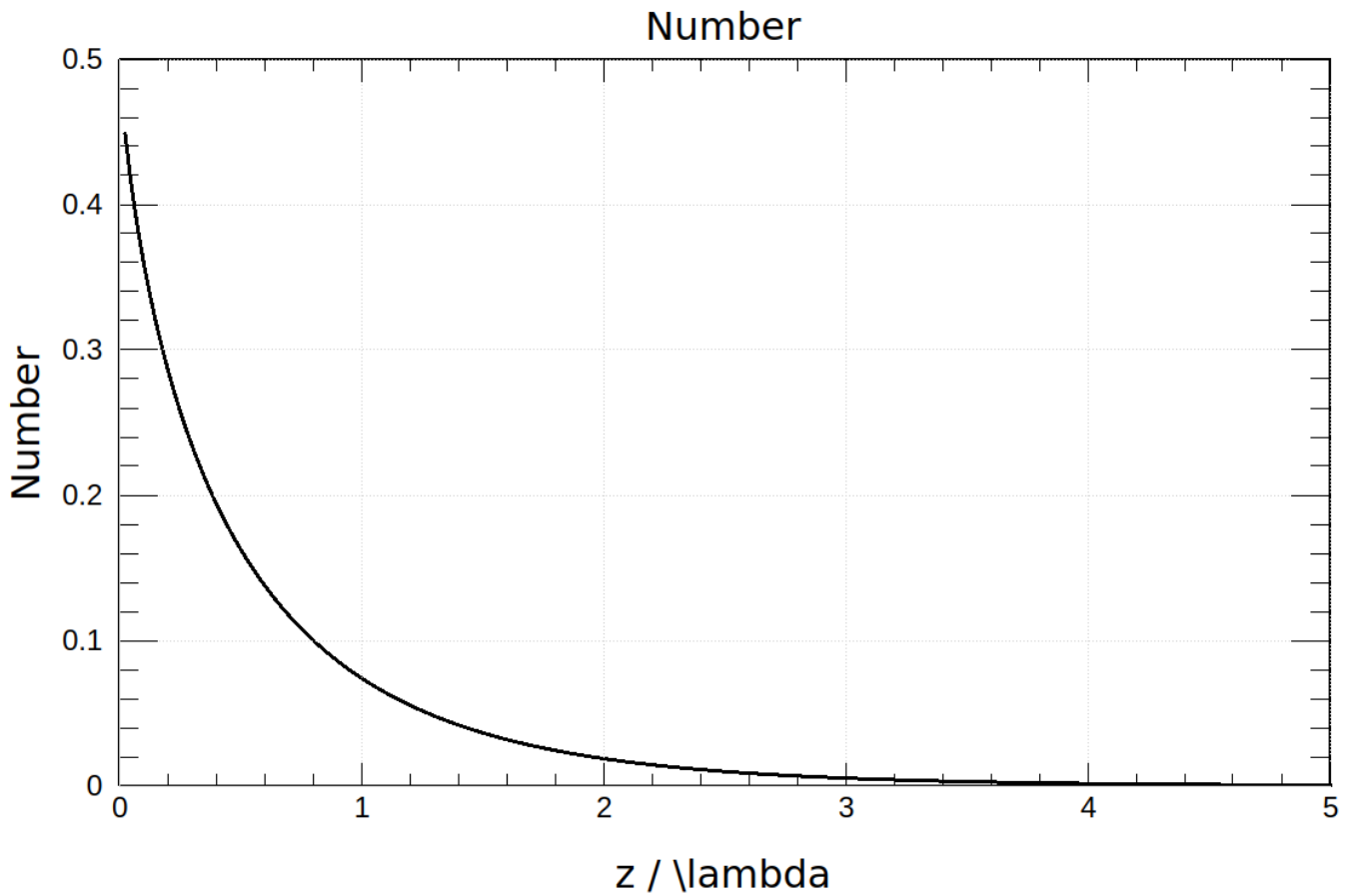
### Part a

$$\begin{aligned}
 F\left(\frac{z}{\lambda}\right) &= \frac{N}{4\pi} \int_0^{2\pi} d\phi \int_0^{\frac{\pi}{2}} \sin\theta d\theta \left(1 - \int_0^{\frac{z}{\cos\theta}} \frac{1}{\lambda} e^{-\frac{x}{\lambda}} dx\right) \\
 &= \frac{N}{2} \left(e^{-\frac{z}{\lambda}} - \frac{z}{\lambda} \Gamma\left(0, \frac{z}{\lambda}\right)\right)
 \end{aligned}$$

where  $\Gamma(a, z)$  is the upper incomplete gamma function. Check asymptotic behaviours of  $\Gamma\left(\frac{z}{\lambda}\right)$

$$\begin{aligned}
 \lim_{z \rightarrow 0} F\left(\frac{z}{\lambda}\right) &= \frac{N}{2} \\
 \lim_{z \rightarrow \infty} F\left(\frac{z}{\lambda}\right) &= 0
 \end{aligned}$$

In the program, N is set to 1.



## Part b

Ignore the limit of the speed of light and assume the gamma ray travels through the material instantly. This should be a reasonable approximation in practice. The surface flux density is

$$f = \int_0^\infty \int_0^{\frac{\pi}{2}} \int_0^{2\pi} (nr^2 \sin \theta dr d\theta d\phi) \left( \frac{1}{\tau} e^{-\frac{t}{\tau}} \right) \left( \frac{\cos \theta}{4\pi r^2} \right) \left( 1 - \int_0^r \frac{1}{\lambda} e^{-\frac{x}{\lambda}} dx \right)$$

$$= \frac{\lambda n}{4} \frac{e^{-\frac{t}{\tau}}}{\tau}$$

## Part c

Define

$$\left\{ \begin{array}{ll} \text{Density of Granite} & \rho = 2.7g/cm^3 \\ \text{Weight fraction of } K_2O & \alpha = 4\% \\ \text{Natural Abundance of } K^{40} & \beta = 0.012\% \\ \text{Gamma ray production efficiency} & \eta = 11\% \\ \text{Half of } K^{40} & t_0 = 10^9 \text{ years} \\ \text{Mass of } K_2O & m = 1.56415 \times 10^{-22} g \end{array} \right.$$

So,

$$\tau = t_0 / \ln(2)$$

$$n = \frac{2\rho\alpha}{m}\beta$$

Surface flux density

$$\begin{aligned} f &= \eta \frac{\lambda n}{4} \frac{e^{-\frac{t}{\tau}}}{\tau} \\ &\approx 0.62602 \left( \frac{1}{2} \right)^{-\frac{t}{1.2 \times 10^9 \text{ yr}}} \text{ cm}^{-2} \text{ s}^{-1} \end{aligned}$$

If we wait long enough time so that all  $K^{40}$  decay. Gamma ray surface number density is

$$N = \int_0^\infty f dt = \frac{n\eta\lambda}{4} \approx 3.41782 \times 10^{16} \text{ cm}^{-2}$$